

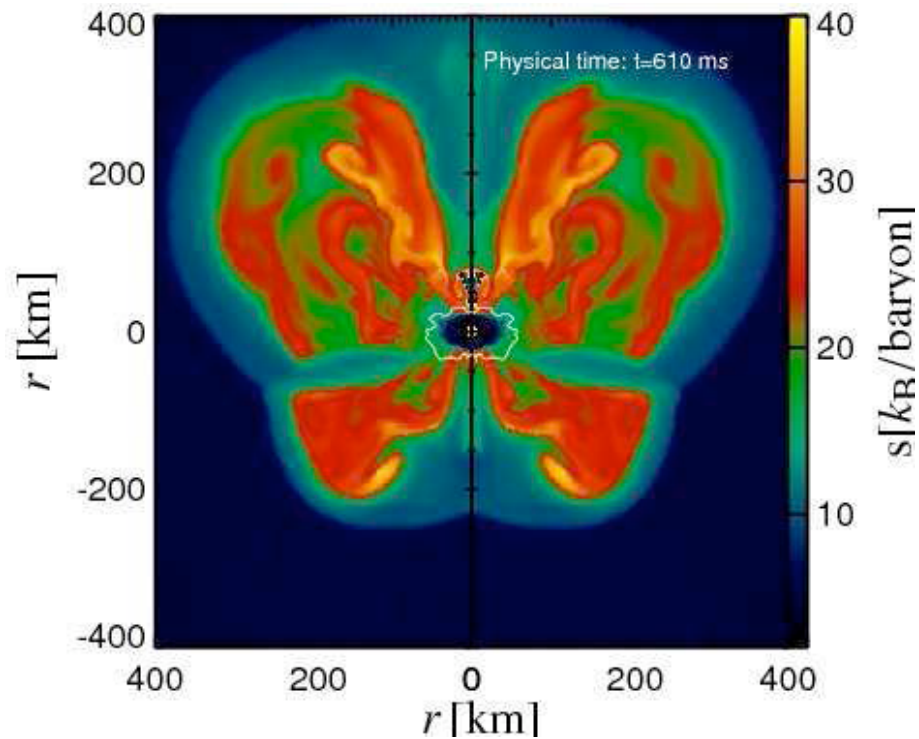
TOWARDS A MHD VERSION OF SASI

JÉRÔME GUILLET
&
THIERRY FOGLIZZO

Outline

- Core collapse with a moderate magnetic field
- Studying MHD SASI : The set up and formalism
- Effect of the magnetic field in different geometries :
 - B perpendicular to the shock
 - B parallel to the shock
- Conclusion and discussion

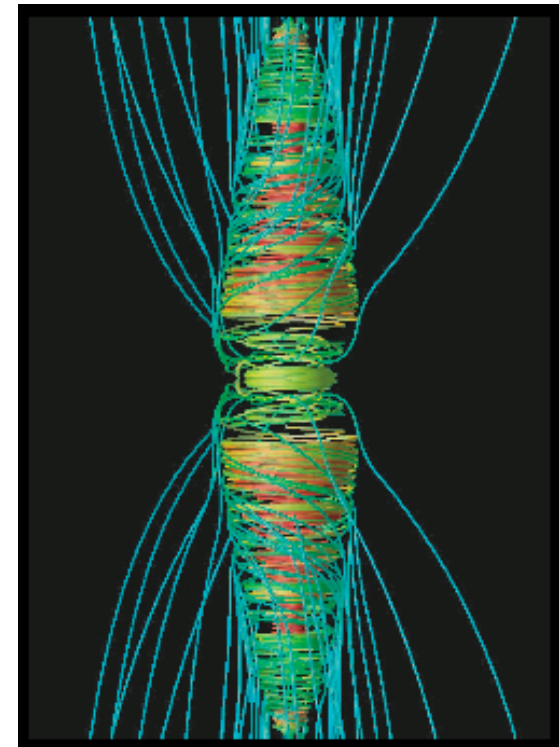
Non magnetic SN



Marek & Janka 2008

SASI plays a crucial role

Magnetic explosion



Burrows *et al.* 2007

Jet explosion with no time
for SASI to develop

A few questions

- What happens with a moderate rotation & magnetic field? Mixed features?
- Could a moderate magnetic field help the neutrino explosion?
- What is the effect of the magnetic field on the SASI?
- How strong a magnetic field should be to have a significant effect on SASI?

A few numbers

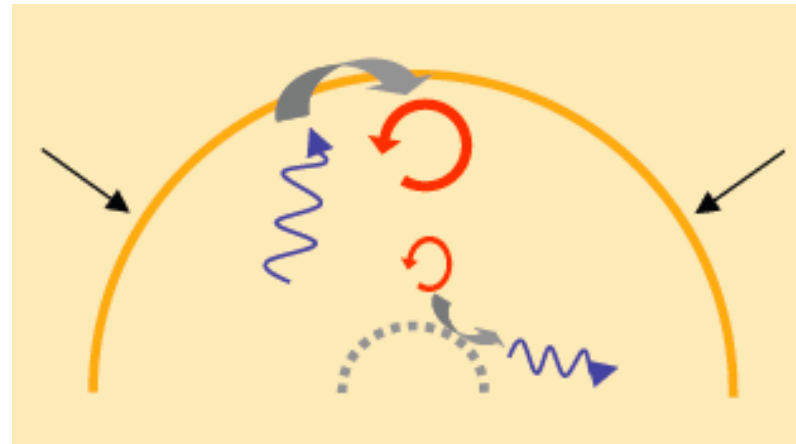
- Strength of B needed to have $v_{\text{Alfven}} = v_{\text{flow}}$?
 - $r_{\text{sh}} = 150\text{km} : B \sim 3 \cdot 10^{13} \text{ G}$
 - $r \sim 50 \text{ km} : B \sim 10^{14} \text{ G}$
- Magnetic field from flux conservation :
 - Heger *et al.* 2005 + flux conservation :
 - At $r=150\text{km} : B_{\phi} \sim 3 \cdot 10^{11} \text{ G}$ too small
 - at $r=50\text{km} : B_{\phi} \sim 2 \cdot 10^{12} \text{ G}$
- Scaling of the saturation of MRI (from Burrows *et al.* 2007) :
 - Supposing a period of the core $P_{\text{core}} \sim 20\text{s}$
 - $R = 150\text{km} : \text{time scale too long } (>500\text{ms})$
 - $R = 50\text{km} : t_{\text{growth}} \sim 50\text{ms}, B_{\text{sat}} \sim 3 \cdot 10^{14} \text{ G}$

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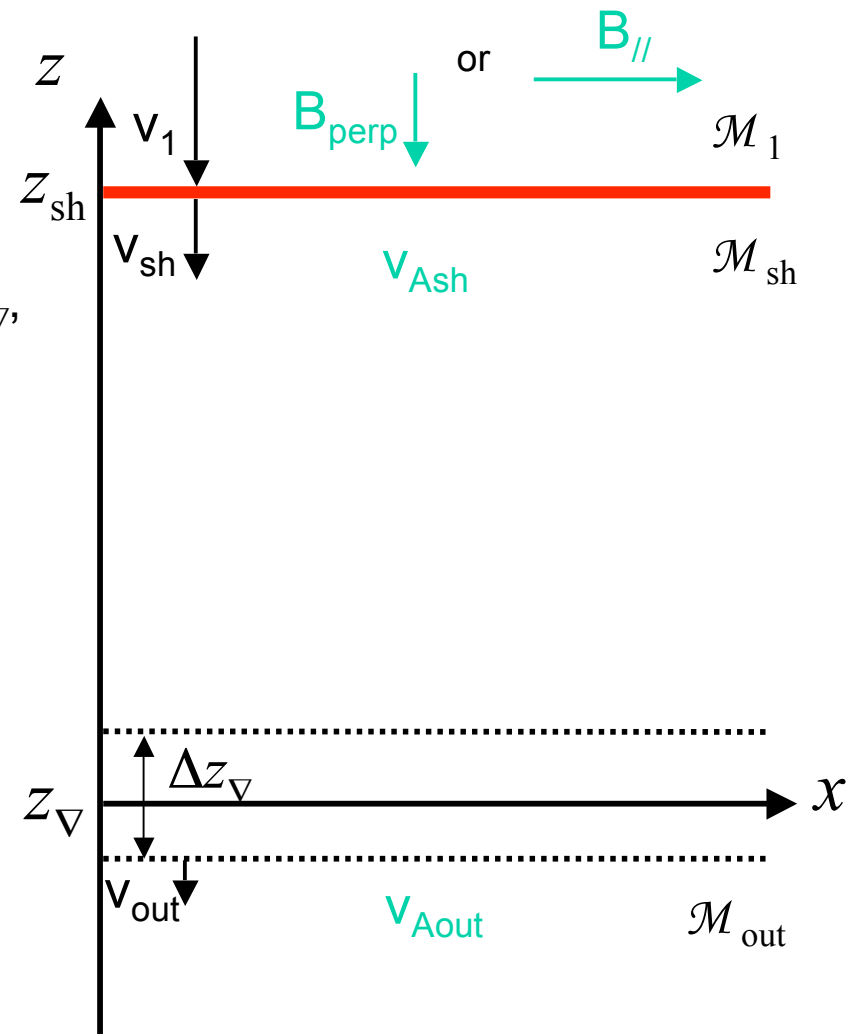
The physical mechanism behind SASI

- Two proposed mechanisms :
 - Advective-acoustic cycle : Foglizzo 2000, 2002, 2007, Blondin *et al.* 2003
 - Purely acoustic mechanism : Blondin & Mezzacappa 2006, Laming 2007
- We will focus on the advective-acoustic interpretation
- The mechanism relies on the decomposition of the mode into waves
- Cycles between the shock and a region of gradient which couples advected waves to acoustic waves




A (very) simplified set up

- Based on the hydro toy model of Foglizzo 2006 (Talk in KITP)
- Parallel adiabatic flow ($\gamma=4/3$)
- Localized coupling with a potential jump (Δz_{∇} , c_{in}/c_{out})
- Lower boundary condition : no wave propagating up from below
- Periodic box of aspect ratio 4
- 2 geometries of magnetic field are studied : parallel and perpendicular to the shock
- 3D perturbations
- Perturbative approach :
 → only the linear phase...



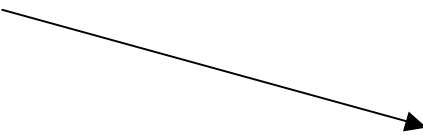


What does B change?

Decomposition of the perturbations : 3 waves -----> **7 waves!**

2 acoustic waves :  2 fast magnetosonic waves

1 advected wave :

entropy		1 advected entropy wave
Vorticity		2 slow magnetosonic waves
		2 Alfvén waves

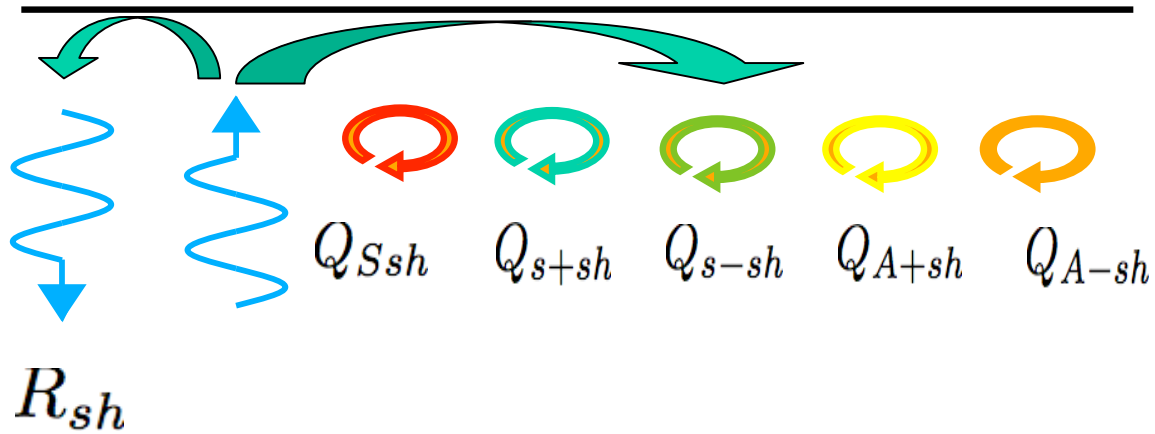
Vorticity and entropy decouple : **vorticity can propagate!**

But : slow or Alfvén waves cannot propagate up to the shock...

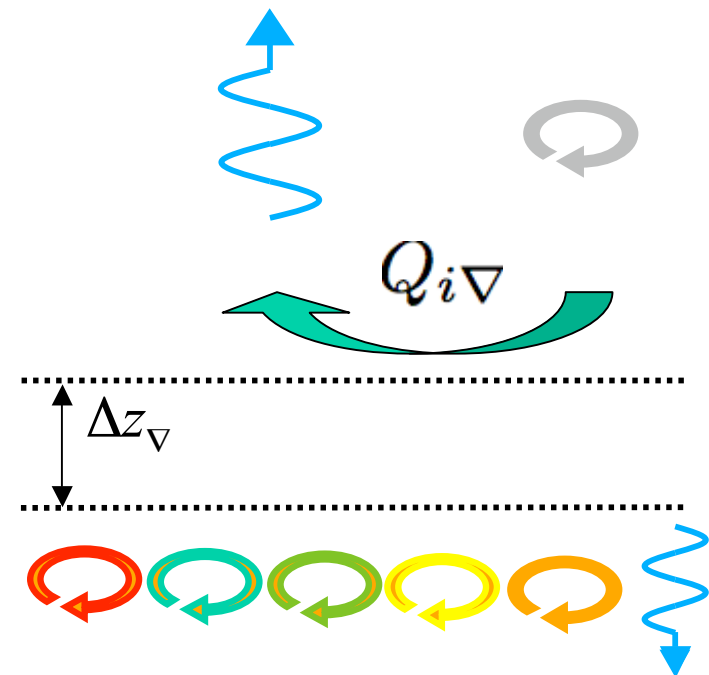
(Condition of evolutionarity of the fast MHD shock wave : Akhiezer *et al.* 1959)

The 6 cycles

Coupling at the shock



Coupling at the potential jump



Cycle constants :

$$Q_i = Q_{ish} \times Q_{i\Delta}$$

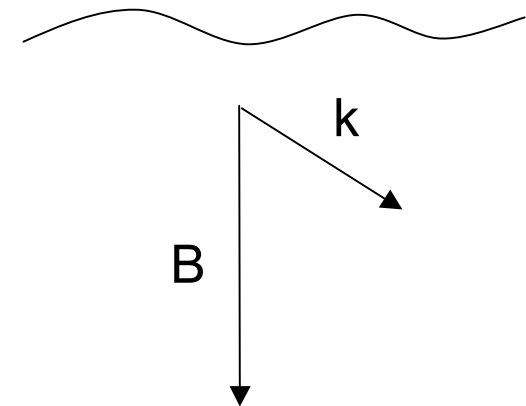
$$R = R_{sh} \times R_{\Delta}$$

The non magnetic limit

- Which one of slow or Alfvén waves contains the vorticity of the hydro cycle?

It depends on the geometry... can be understood by geometric arguments.

- If B perp shock : slow waves
- If B // shock :
 - If k_{tr} // B : slow waves
 - If k_{tr} perp B : Alfvén waves
 - If k_{tr} oblique : both
- What is the relative contribution of the vortical and entropic cycles ?
 - Depends on the details of the coupling region :
 $c_{in}/c_{out}, \Delta z_{\nabla}$
 - Usually in this set up : $Q_S/Q_{vort} \sim 2-5$



Formalism for studying the 6 cycles

$$Qe^{i\omega\tau_Q} + Re^{i\omega\tau_R} = 1$$



$$Qse^{i\omega\tau_S} + Q_{s+}e^{i\omega\tau_{s+}} + Q_{s-}e^{i\omega\tau_{s-}} + Q_{A+}e^{i\omega\tau_{A+}} + Q_{A-}e^{i\omega\tau_{A-}} + Re^{i\omega\tau_R} = 1$$

entropy

Slow waves

Alfven waves

fast waves

Usefull quantities :

$$Q_{tot} \equiv \sum_i Q_i e^{i\omega_r \tau_i} + Re^{i\omega_r \tau_R}$$

$$\omega_i \sim \frac{\log Q_{tot}}{\tau}$$

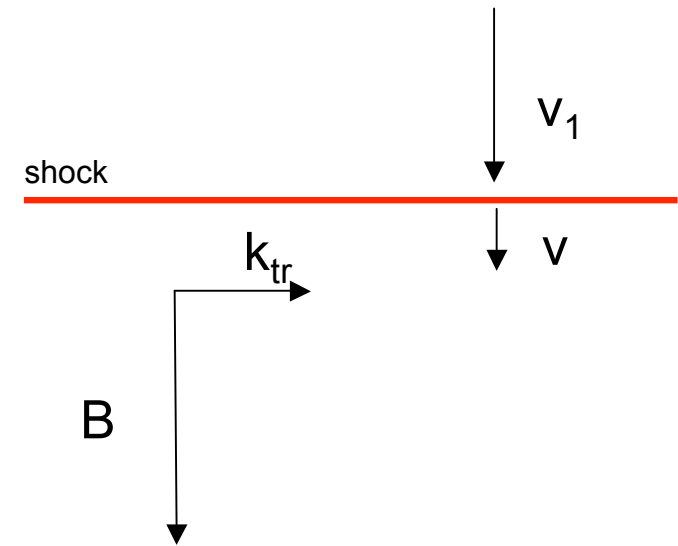
$$Q_{vort} \equiv Q_{s+}e^{i\omega_r \tau_{s+}} + Q_{s-}e^{i\omega_r \tau_{s-}} + Q_{A+}e^{i\omega_r \tau_{A+}} + Q_{A-}e^{i\omega_r \tau_{A-}}$$

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Magnetic field perpendicular to the shock

- Corresponds to a radial field in core collapse (split monopole)
- If $v=v_A$: a singular point appears which is unstable (Williams 1975)
 - We restrict our study to $v_A < v$ everywhere
- 2D perturbations are enough to describe the flow : Alfvén waves are not produced at the shock
- Vorticity propagates through slow waves along the magnetic field direction

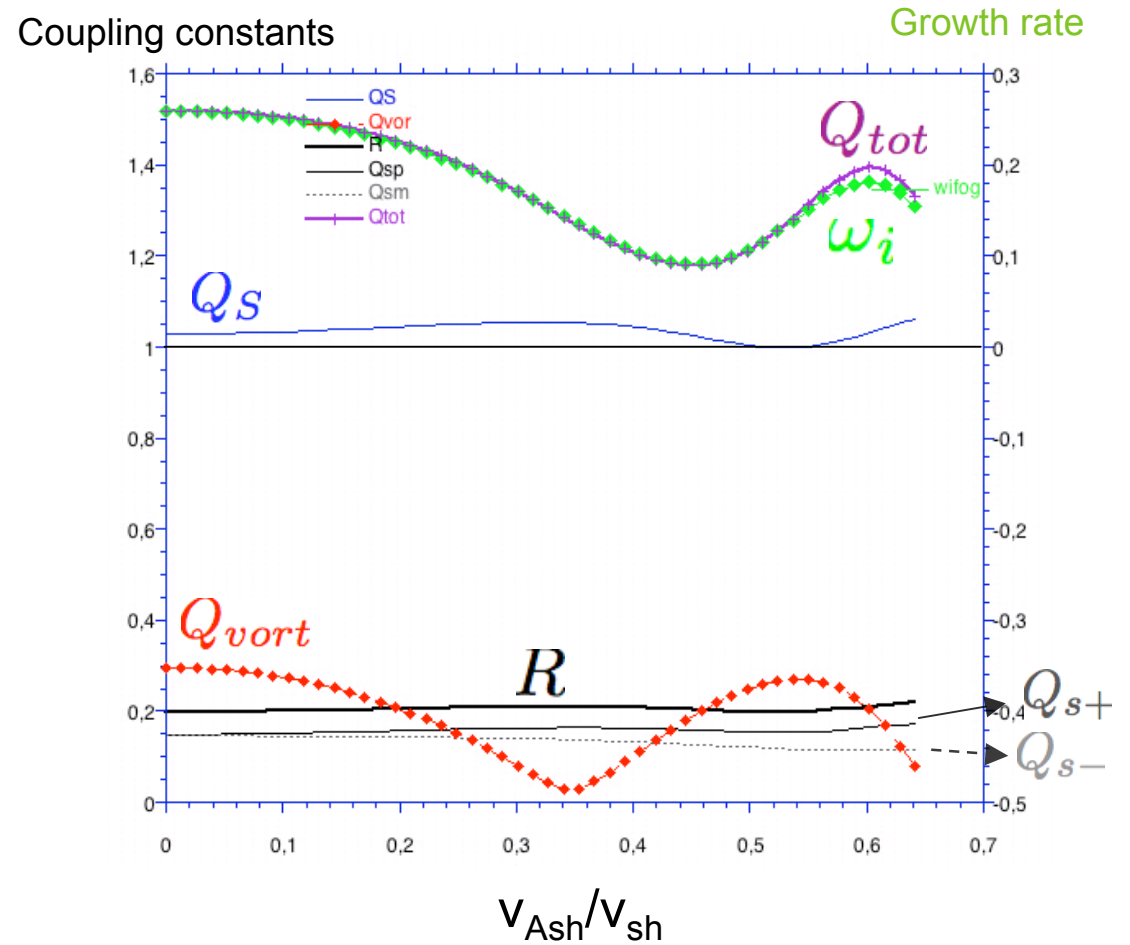


n=1

- Effect on the coupling : not much...
 - Q_{S+} and Q_{S-} change very little

- Dominant effect : interferences between the 3 cycles : Q_S , Q_{S+} , Q_{S-}

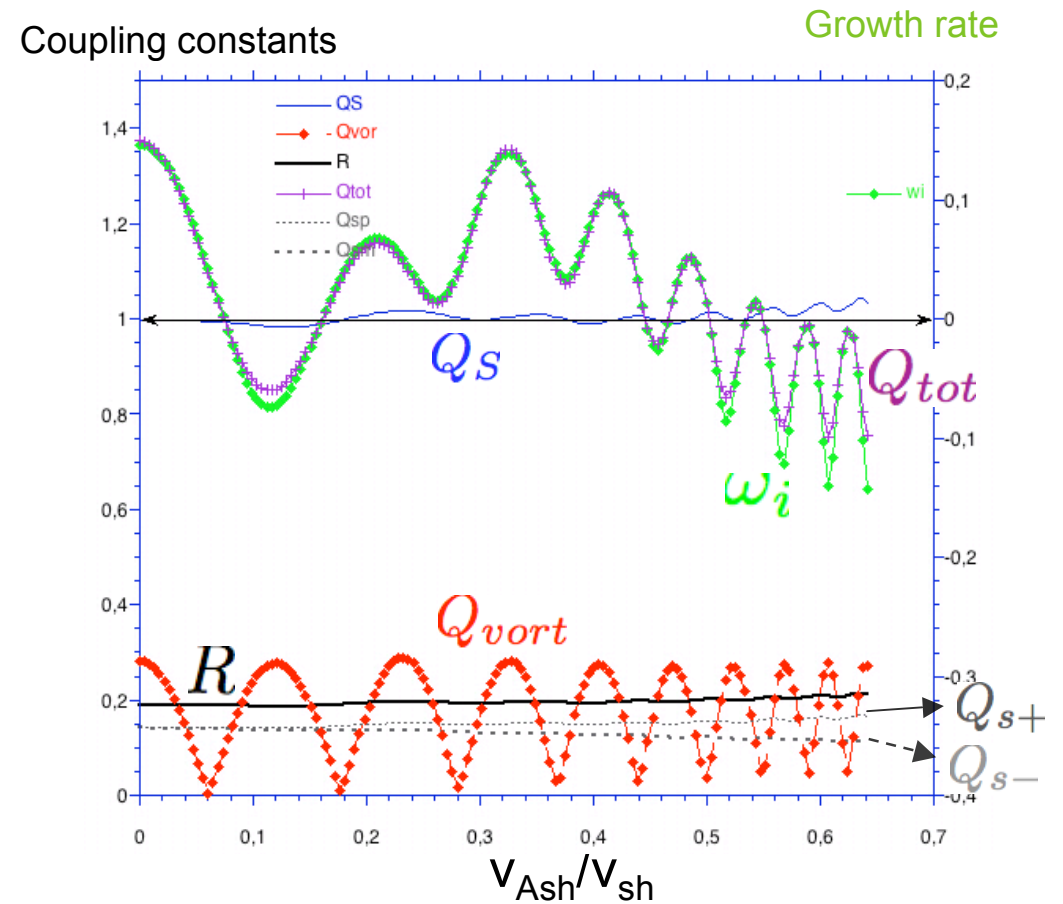
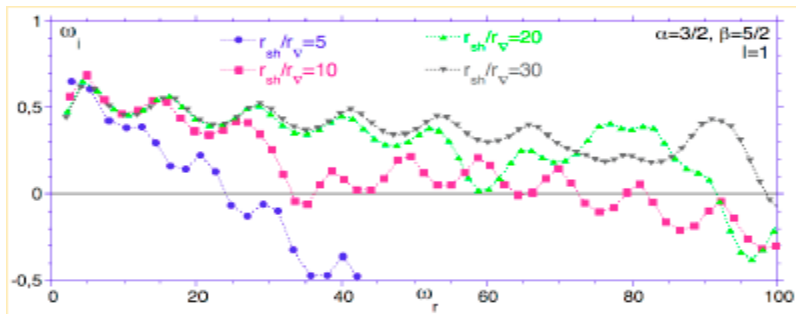
→ Stabilization...



$n=6, h=6$

- Slow « + » : $v-v_A$ $k_z \sim \frac{\omega}{v - v_A}$
- Slow oscillations
- Slow « - » : $v+v_A$ $k_z \sim \frac{\omega}{v + v_A}$
- Fast oscillations

Similar to the oscillations in the eigenspectrum :
due to the acoustic interferences (Foglizzo et al 2007)



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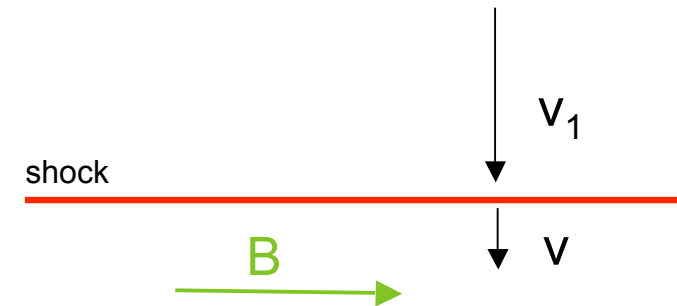
Magnetic field parallel to the shock

- Transverse magnetic field is amplified by compression through the shock and in a decelerating region
- Winding by rotation creates this geometry (cf magnetic explosions)

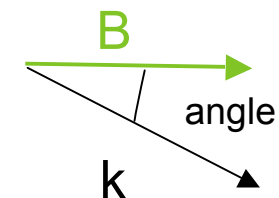
→ **More realistic than a perpendicular field**

- Alfvén waves propagate only in a transverse direction
- Slow waves propagate mainly along the field line : transverse direction as well

$$k_z \sim \frac{\omega}{v} \pm k_{//B} \frac{v_A}{v}$$

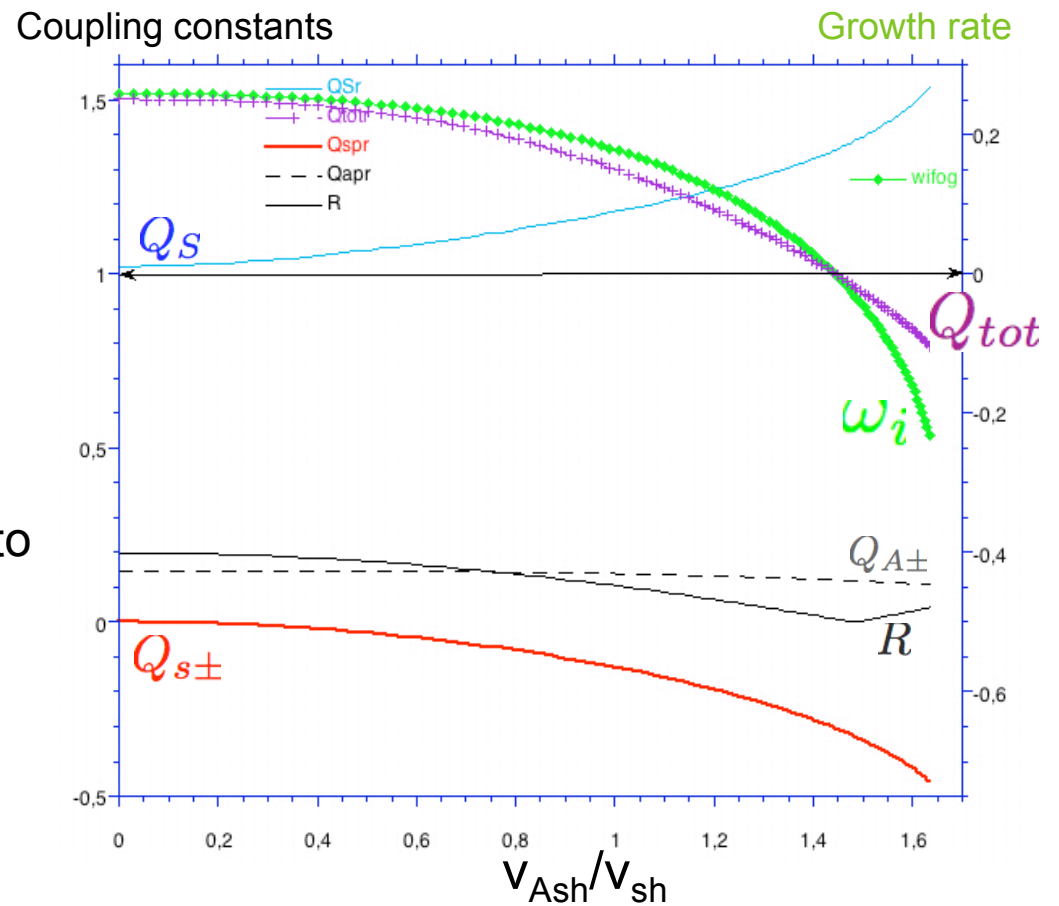


Viewed from above :

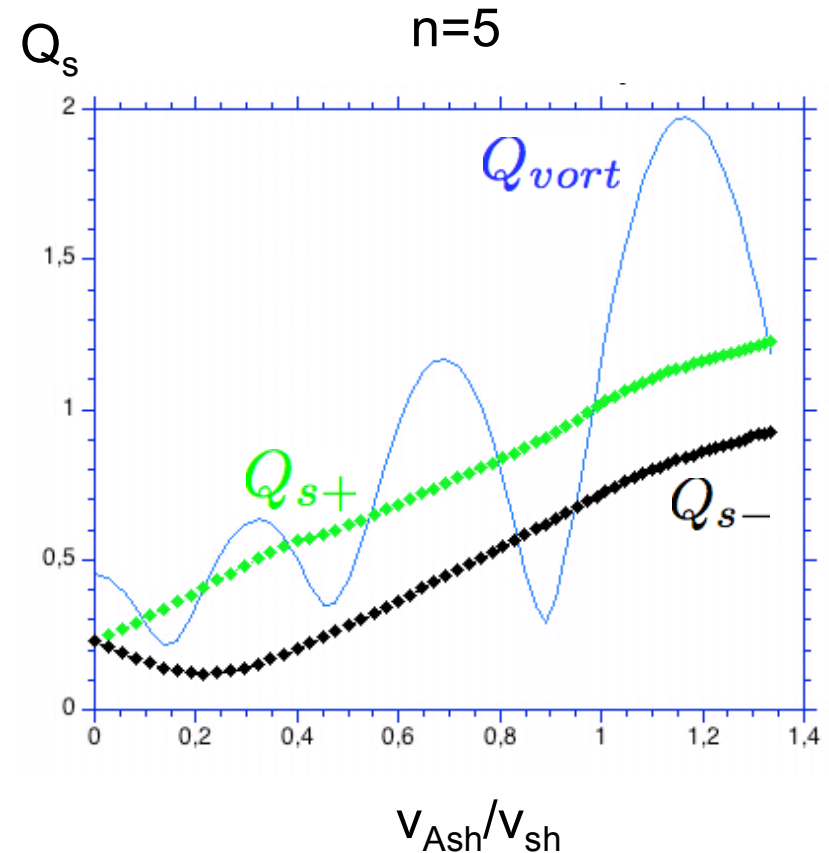
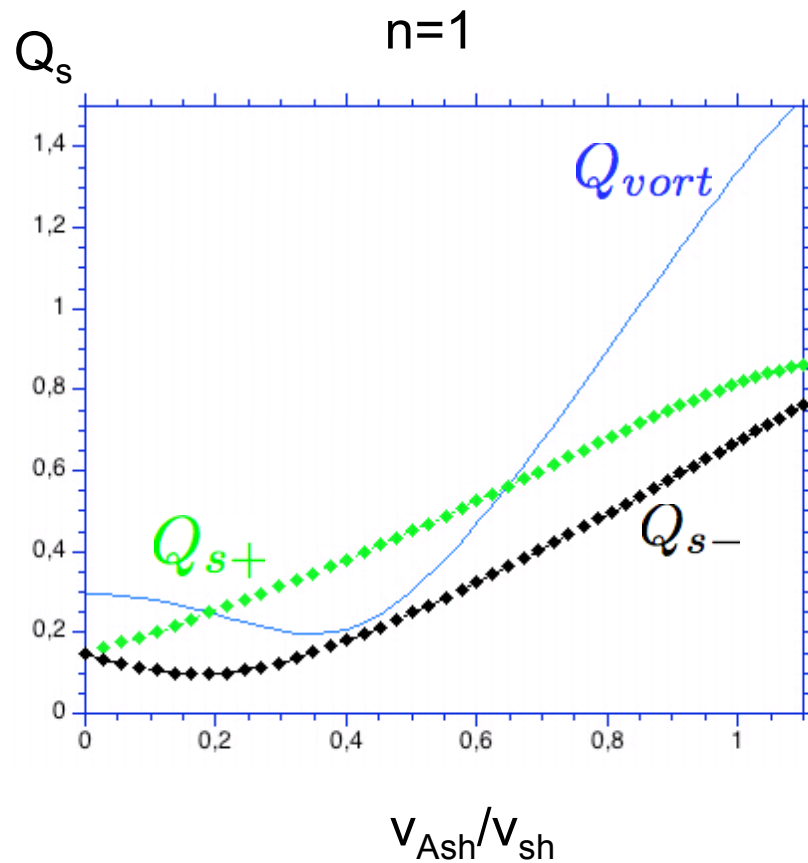


K perpendicular to B

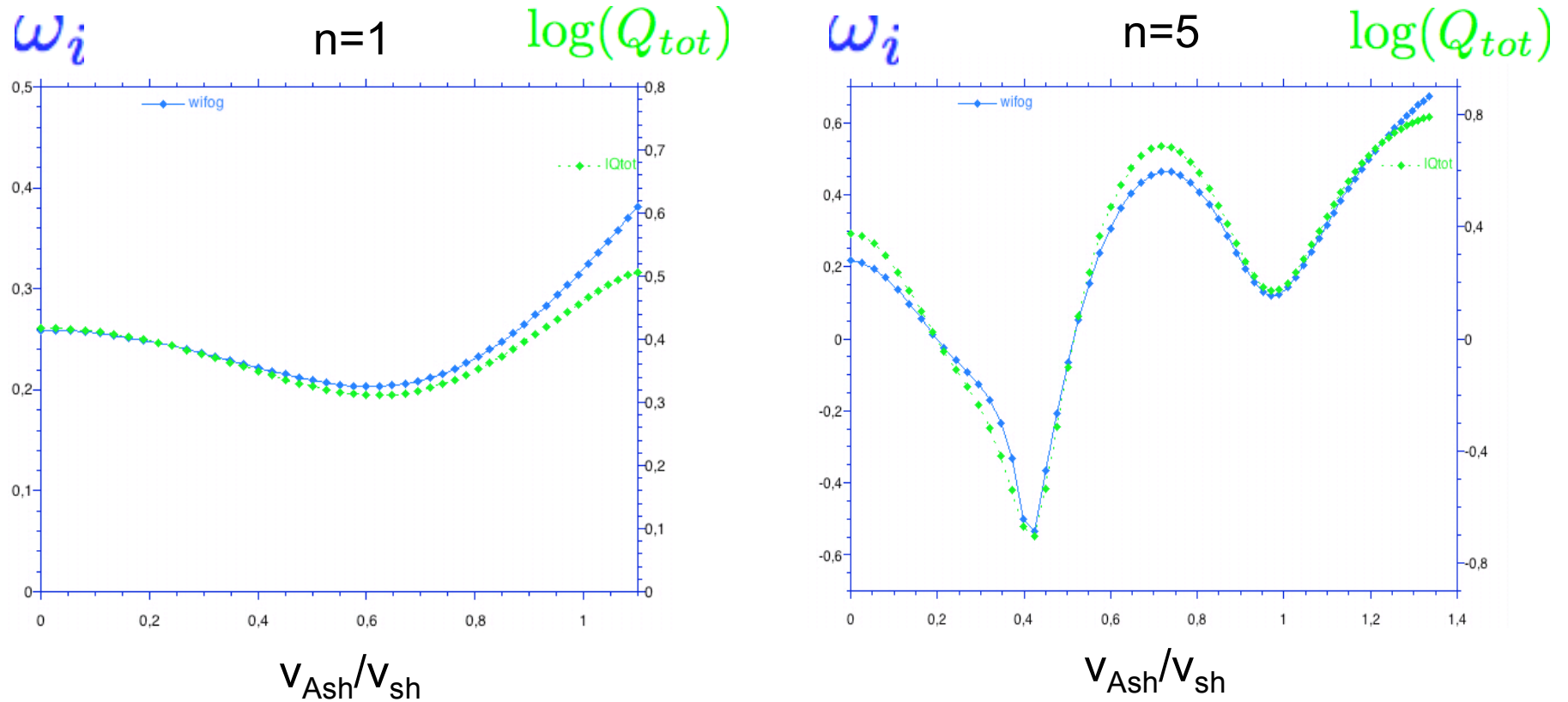
- Slow and Alfvén waves simply advected
- If $v_A=0$: only Alfvén waves
- With a magnetic field : slow waves excited but transverse velocity associated to it = 0
- Stabilizing effect due to the slow waves which have an opposite effect to entropy...



K parallel to B : the 'vortical' cycles

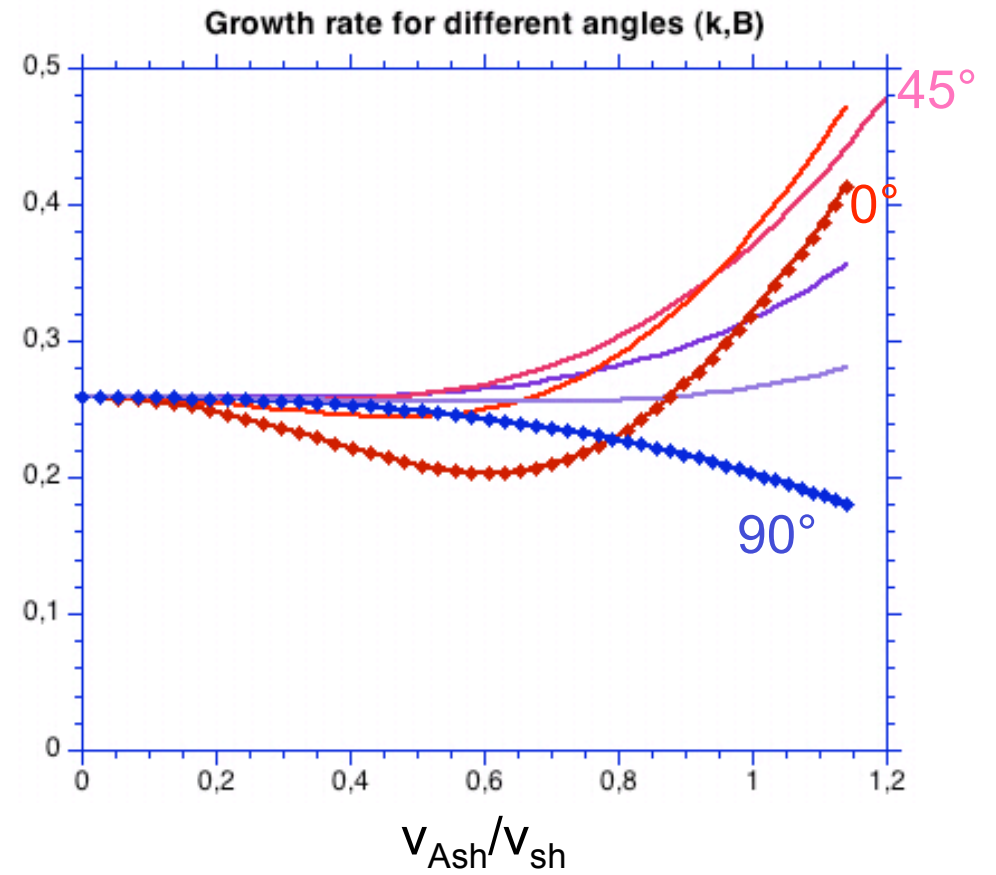
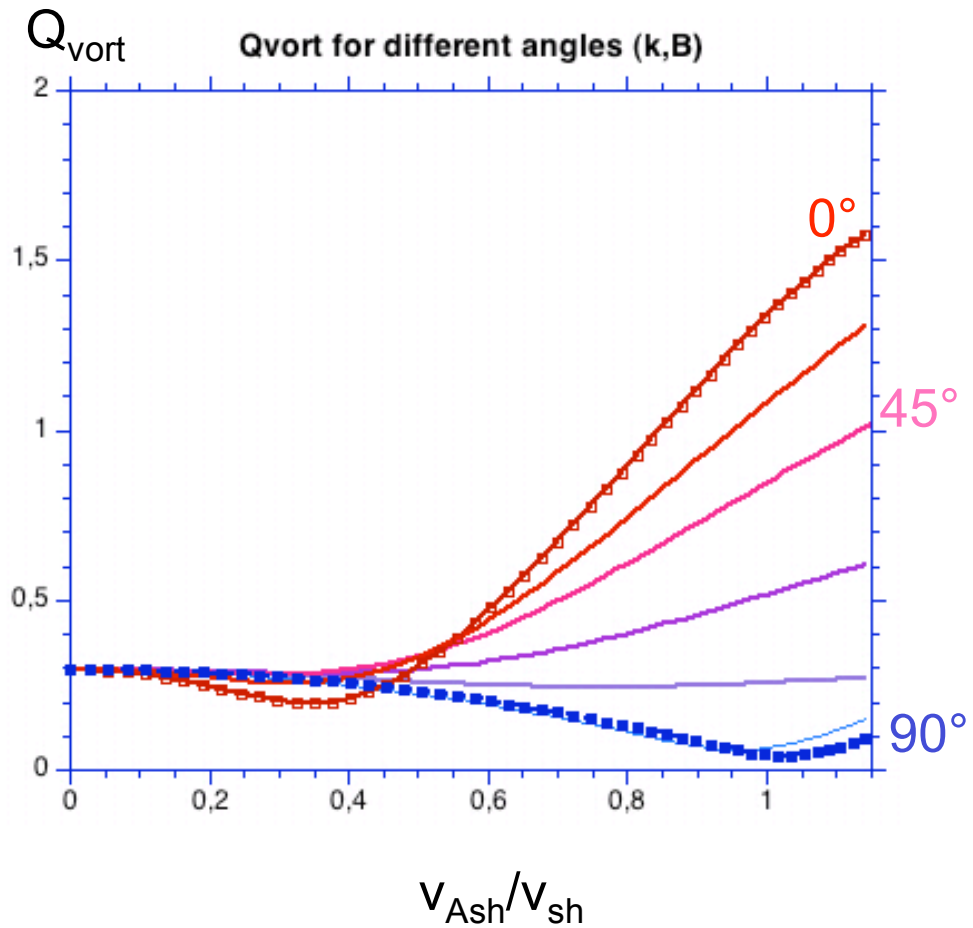


Growth rate : stabilization or destabilization



Stabilization at small B, destabilization for strong B

Oblique wave vector



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Summary

- Effect always of the second order ($\sim v_A^2$)
- No significant effects on the acoustic cycle : remains stable.
- Vorticity can propagate through up to 4 different waves : not in phase with entropy or with each other...
 - Interference between the different cycles : stabilizing effect...
- If B perpendicular to the shock :
 - no strong effect on the coupling, only the interference effect
- If B parallel to the shock :
 - K perp B : magnetic perturbation associated to entropy
 - stabilization...
 - K // B : strong amplification of vortical cycles
 - Possible destabilization

Towards more realism

- Missing ingredients : cooling and convergence (and gravity).
- Effect of cooling : Favor the vortical cycle compared to the entropic cycle (?)
 - Effects due to the vortical cycle are accentuated :
 - If B perpendicular : more stabilized
 - If $B //$:
 - If $k // B$: the destabilization could be much more efficient
 - If $k \cdot B$: effect independant of vorticity (associated to entropy) : less stabilized?
 - New effect?
- Effect of convergence :
 - Strength needed to have a strong effect? v_A/v at the shock or deep inside are very different...
- Future works :
 - Include cooling, convergence in cylindrical geometry
 - Magnetic field + rotation...