

Polarization effects in $e^+ + e^- \rightarrow \bar{d} + d$ and determination of time like deuteron form factors

G. I. Gakh

*National Science Centre "Kharkov Institute of Physics and Technology",
61108 Akademicheskaya 1, Kharkov, Ukraine*

E. Tomasi-Gustafsson

DAPNIA/SPhN, CEA/Saclay, 91191 Gif-sur-Yvette Cedex, France

C. Adamušćin and S. Dubnička

Institute of Physics, Slovak Academy of Sciences, Bratislava, Slovakia

A. Z. Dubničková

Department of Theoretical Physics, Comenius University, Bratislava, Slovakia

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Abstract

Polarization effects in the reaction $e^+ + e^- \rightarrow \bar{d} + d$ have been investigated for the case of longitudinally polarized electron beam and arbitrary polarization of the produced deuteron, with the aim of a determination of the time-like complex deuteron electromagnetic form factors. General expressions of polarization observables are derived and numerical estimations have been carried out by means of various models of deuteron electromagnetic form factors, for kinematical conditions near threshold.

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I. INTRODUCTION

The electromagnetic form factors (FFs) of hadrons and nuclei provide important information about the structure and internal dynamics of these systems. Recent progress in electron–scattering experiments allowed to measure not only the cross sections but also various polarization observables in the region of the momentum transfers where these data can help to discriminate between different theoretical predictions.

The deuteron, the only bound two–nucleon system, is one of the fundamental systems of nuclear physics. Accordingly, many studies, both experimental and theoretical, have been devoted to it. Of particular interest today is the degree to which the deuteron can be understood as a system of two nucleons interacting via the known nucleon–nucleon interaction.

When addressing, more specifically, to the electromagnetic properties of the deuteron, the main question concerns the reliability to predict the three deuteron FFs starting from the calculated deuteron wave function and nucleon FFs known from electron–nucleon scattering. At low momentum transfers, predictions and data agree quite well when accounting for one–body terms only, whereas at the higher momentum transfers, two–body contributions are known to be important. Whether quark degrees of freedom do need to be taken explicitly into account, is still a matter of debate. A status of the experimental and theoretical research of the deuteron can be found in recent reviews [1, 2].

Elastic electron–deuteron scattering has been investigated in many experiments, and cross section data today covers a large range of momentum transfers (see review [2]). Some of these data obviously are not very precise, other data, mainly of more recent origin, have reached accuracies down to the 1 % level. During the last years, it has become possible to measure not only cross sections, but also spin observables, due to the developments of polarized electron beams, polarized deuteron targets and polarimeters. The knowledge of these spin observables is unavoidable, if one wants to separate the contributions of the different multipolarities to the $A(Q^2)$ structure function. On the side of experiment, good progress has been made. In particular, recent polarization data for electron–deuteron elastic scattering allowed the individual determination of the deuteron charge and quadrupole FFs up to a value of the momentum transfer squared $Q^2 = 1.8 \text{ GeV}^2$.

The deuteron charge FF G_C is particularly interesting for the understanding the deuteron

structure, beyond the impulse approximation. G_C displays a node at $Q^2=0.7 \text{ GeV}^2$, and the position of this node is especially sensitive to the ingredients of the models, in particular meson-exchange currents.

The experimental investigation of deuteron FFs should help to determine the region where it is necessary to introduce explicitly quark and gluon degrees of freedom, for a correct description of the deuteron. At present, as it was shown in Ref. [3], the overall experimental results on elastic electron-deuteron cross sections are not consistent with pQCD predictions. The best global descriptions of the existing deuteron data are based on impulse approximation (including eventually relativistic corrections, meson exchange currents, Δ isobars...).

The interaction of electrons with deuterons is usually assumed to occur through the exchange of a virtual photon (one-photon exchange approximation) due to the smallness of the electromagnetic fine structure constant, which suppress two-or more- photon exchange. However, a few decades ago it was suggested [4] that the two-photon exchange mechanism may be significant in the region of large momentum transfer. More recently, the possible contribution of two-photon exchange to the elastic electron-deuteron scattering was discussed in Ref. [3].

As for the nucleon, the knowledge of electromagnetic FFs in the time-like (TL) region of momentum transfer can give additional important information about the internal composite structure of the hadron. Measurements are certainly more difficult in the deuteron case, as shown in Ref. [5], where the total cross section of the reaction $e^- + e^+ \rightarrow d + \bar{d}$ was predicted up to $q^2 = 30 \text{ GeV}^2$, using a model of deuteron FFs based on an extension of the vector-meson-dominance model (VMD) of the electromagnetic hadron interactions. However, other mechanisms, as the presence of a two-photon contribution, could favor a larger cross section.

After the challenging discovery of antideuteron [6], which established the existence of nuclear antimatter, the production of antideuteron was recorded in different reactions. Very recently, the production of deuterons and antideuterons in $Au + Au$ collisions has been reported by the PHENIX experiment at RHIC [7] and interpreted in terms of coalescence model. It was found that the spectra of d and \bar{d} decrease less steeply than p (\bar{p}) spectra. The cross section for \bar{d} photoproduction was also measured at HERA at $W_{\gamma p} = 200 \text{ GeV}$ [8]. The production of \bar{d} in e^+e^- -annihilation at $W = 10 \text{ GeV}$ was also measured at DORIS II storage ring [9].

In the present paper we calculate the polarization observables in the reaction

$$e^-(k_1) + e^+(k_2) \rightarrow d(p_1) + \bar{d}(p_2). \quad (1)$$

where the momenta of the particles are indicated in brackets.

We consider the case of unpolarized and longitudinally-polarized electron beam with production of vector- and tensor-polarized deuterons. The expressions of polarization observables are given in terms of the deuteron electromagnetic FFs. Due to final state interaction, FFs are complex functions of the variable q^2 . Nevertheless, not all models of FFs, which are mainly parametrizations built for the space-like region, can be consistently applied to the TL region and give origin to an imaginary part. However, numerical estimations for the cross section and polarization observables are tentatively given, on the base of the analytic continuation of existing parameterizations of the deuteron FFs, similarly to the case of nucleon FFs in Ref. [10].

II. POLARIZATION OBSERVABLES

In the one-photon approximation, the differential cross section of the reaction (1) in terms of the leptonic $L_{\mu\nu}$ and hadronic $W_{\mu\nu}$ tensors contraction (in the Born approximation we can neglect the electron mass) is written as

$$\frac{d\sigma}{d\Omega} = \frac{\alpha^2 \beta}{4q^2} \frac{L_{\mu\nu} W_{\mu\nu}}{q^4}, \quad (2)$$

where $\alpha = 1/137$ is the electromagnetic constant, $\beta = \sqrt{1 - 4M^2/q^2}$ is the deuteron velocity in the reaction center of mass system (CMS), M is the deuteron mass and q is the four momentum of the virtual photon, $q = k_1 + k_2 = p_1 + p_2$ (note that the cross section is not averaged over the spins of the initial beams).

The leptonic tensor (for the case of longitudinally polarized electron beam) is

$$L_{\mu\nu} = -q^2 g_{\mu\nu} + 2(k_{1\mu} k_{2\nu} + k_{2\mu} k_{1\nu}) + 2i\lambda \varepsilon_{\mu\nu\sigma\rho} k_{1\sigma} k_{2\rho}, \quad (3)$$

where λ is the degree of the beam polarization (further we assume that the electron beam is completely polarized and consequently $\lambda = 1$).

The hadronic tensor can be expressed via the nucleon electromagnetic current J_μ , describing the transition $\gamma^* \rightarrow \bar{d}d$, as

$$W_{\mu\nu} = J_\mu J_\nu^*. \quad (4)$$

As the deuteron is a spin-one nucleus, its electromagnetic current is completely described by three FFs. Assuming the P- and C-invariance of the hadron electromagnetic interaction this current can be written as [11]

$$J_\mu = (p_1 - p_2)_\mu \left[-G_1(q^2) U_1^* \cdot U_2^* + \frac{G_3(q^2)}{M^2} (U_1^* \cdot q U_2^* \cdot q - \frac{q^2}{2} U_1^* \cdot U_2^*) \right] - G_2(q^2) (U_{1\mu}^* U_2^* \cdot q - U_{2\mu}^* U_1^* \cdot q), \quad (5)$$

where $U_{1\mu}$ ($U_{2\mu}$) is the polarization four-vector describing the spin one deuteron (antideuteron), and $G_i(q^2)$ ($i = 1, 2, 3$) are the deuteron electromagnetic FFs. The FFs $G_i(q^2)$ are complex functions of the variable q^2 in the region of the TL momentum transfer ($q^2 > 0$). They are related to the standard deuteron electromagnetic FFs: G_C (charge monopole), G_M (magnetic dipole) and G_Q (charge quadrupole) by

$$G_M = -G_2, \quad G_Q = G_1 + G_2 + 2G_3, \quad G_C = -\frac{2}{3}\tau(G_2 - G_3) + (1 - \frac{2}{3}\tau)G_1, \quad \tau = \frac{q^2}{4M^2}. \quad (6)$$

The standard FFs have the following normalizations:

$$G_C(0) = 1, \quad G_M(0) = (M/m_n)\mu_d, \quad G_Q(0) = M^2 Q_d, \quad (7)$$

where m_n is the nucleon mass, $\mu_d = 0.857(Q_d = 0.2859 \text{ fm}^2)$ is deuteron magnetic (quadrupole) moment.

When calculating the expression for the hadron tensor $W_{\mu\nu}$ in terms of the deuteron electromagnetic FFs, using the explicit form of the electromagnetic current (5), the spin-density matrices of the deuteron and antideuteron are

$$U_{1\mu} U_{1\nu}^* = -\left(g_{\mu\nu} - \frac{p_{1\mu} p_{1\nu}}{M^2} \right) + \frac{3i}{2M} \varepsilon_{\mu\nu\rho\sigma} s_\rho p_{1\sigma} + 3Q_{\mu\nu}, \quad U_{2\mu} U_{2\nu}^* = -\left(g_{\mu\nu} - \frac{p_{2\mu} p_{2\nu}}{M^2} \right), \quad (8)$$

if the deuteron polarization is measured and the antideuteron polarization is not measured. Here s_μ and $Q_{\mu\nu}$ are the deuteron polarization four vector and quadrupole tensor, respectively. The four vector of the deuteron vector polarization s_μ and the deuteron quadrupole-polarization tensor $Q_{\mu\nu}$ satisfy the following conditions:

$$s^2 = -1, \quad s p_1 = 0, \quad Q_{\mu\nu} = Q_{\nu\mu}, \quad Q_{\mu\mu} = 0, \quad p_{1\mu} Q_{\mu\nu} = 0.$$

Taking into account Eqs. (4), (5) and (8), the hadronic tensor in the general case can be written as the sum of three terms

$$W_{\mu\nu} = W_{\mu\nu}(0) + W_{\mu\nu}(V) + W_{\mu\nu}(T), \quad (9)$$

where $W_{\mu\nu}(0)$ corresponds to the case of unpolarized deuteron and $W_{\mu\nu}(V)(W_{\mu\nu}(T))$ corresponds to the case of the vector (tensor) polarized deuteron. The explicit form of these terms is:

- the unpolarized term $W_{\mu\nu}(0)$:

$$W_{\mu\nu}(0) = W_1(q^2)\tilde{g}_{\mu\nu} + \frac{W_2(q^2)}{M^2}\tilde{p}_{1\mu}\tilde{p}_{1\nu}, \quad \tilde{g}_{\mu\nu} = g_{\mu\nu} - \frac{q_\mu q_\nu}{q^2}, \quad \tilde{p}_{1\mu} = p_{1\mu} - \frac{p_1 q}{q^2}q_\mu, \\ W_1(q^2) = 8M^2\tau(1-\tau)|G_M|^2, \quad W_2(q^2) = 12M^2(|G_C|^2 - \frac{2}{3}\tau|G_M|^2 + \frac{8}{9}\tau^2|G_Q|^2). \quad (10)$$

- the term for vector polarization $W_{\mu\nu}(V)$:

$$W_{\mu\nu}(V) = \frac{i}{M}S_1(q^2)\varepsilon_{\mu\nu\sigma\rho}s_\sigma q_\rho + \frac{i}{M^3}S_2(q^2)[\tilde{p}_{1\mu}\varepsilon_{\nu\alpha\sigma\rho}s_\alpha q_\sigma p_{1\rho} - \tilde{p}_{1\nu}\varepsilon_{\mu\alpha\sigma\rho}s_\alpha q_\sigma p_{1\rho}] + \\ + \frac{1}{M^3}S_3(q^2)[\tilde{p}_{1\mu}\varepsilon_{\nu\alpha\sigma\rho}s_\alpha q_\sigma p_{1\rho} + \tilde{p}_{1\nu}\varepsilon_{\mu\alpha\sigma\rho}s_\alpha q_\sigma p_{1\rho}], \\ S_1(q^2) = -3M^2(\tau-1)|G_M|^2, \quad S_2(q^2) = 3M^2[|G_M|^2 - 2Re(G_C - \frac{\tau}{3}G_Q)G_M^*], \\ S_3(q^2) = 6M^2Im(G_C - \frac{\tau}{3}G_Q)G_M^*. \quad (11)$$

- the term for tensor polarization $W_{\mu\nu}(T)$:

$$W_{\mu\nu}(T) = V_1(q^2)\tilde{Q}\tilde{g}_{\mu\nu} + V_2(q^2)\frac{\tilde{Q}}{M^2}\tilde{p}_{1\mu}\tilde{p}_{1\nu} + \\ + V_3(q^2)(\tilde{p}_{1\mu}\tilde{Q}_\nu + \tilde{p}_{1\nu}\tilde{Q}_\mu) + V_4(q^2)\tilde{Q}_{\mu\nu} + iV_5(q^2)(\tilde{p}_{1\mu}\tilde{Q}_\nu - \tilde{p}_{1\nu}\tilde{Q}_\mu), \quad (12)$$

where

$$\tilde{Q}_\mu = Q_{\mu\nu}q_\nu - \frac{q_\mu}{q^2}\tilde{Q}, \quad \tilde{Q}_\mu q_\mu = 0, \\ \tilde{Q}_{\mu\nu} = Q_{\mu\nu} + \frac{q_\mu q_\nu}{q^4}\tilde{Q} - \frac{q_\nu q_\alpha}{q^2}Q_{\mu\alpha} - \frac{q_\mu q_\alpha}{q^2}Q_{\nu\alpha}, \quad \tilde{Q}_{\mu\nu}q_\nu = 0, \quad \tilde{Q} = Q_{\alpha\beta}q_\alpha q_\beta. \quad (13)$$

The tensor structure functions $V_i(q^2)$ are combinations of deuteron FFs as follows:

$$V_1(q^2) = -3|G_M|^2, \quad V_2(q^2) = 3\left[|G_M|^2 + \frac{4}{1-\tau}Re(G_C - \frac{\tau}{3}G_Q - \tau G_M)G_Q^*\right], \quad (14)$$

$$V_3(q^2) = -6\tau\left[|G_M|^2 + 2ReG_Q G_M^*\right], \quad V_4(q^2) = -12M^2\tau(1-\tau)|G_M|^2, \quad V_5(q^2) = -12\tau Im(G_Q G_M^*).$$

Using the definitions of the cross-section (2), leptonic (3) and hadronic (9) tensors, one can easily derive the expression for the unpolarized differential cross section in terms of the structure functions $W_{1,2}$ (after averaging over the spins of the initial particles)

$$\frac{d\sigma^{un}}{d\Omega} = \frac{\alpha^2\beta}{4q^4} \left\{ -W_1(q^2) + \frac{1}{2}W_2(q^2) \left[\tau - 1 - \frac{(u-t)^2}{4M^2q^2} \right] \right\}, \quad (15)$$

where $t = (k_1 - p_1)^2$, $u = (k_1 - p_2)^2$.

In the reaction CMS this expression can be written as

$$\frac{d\sigma^{un}}{d\Omega} = \frac{\alpha^2\beta^3}{4q^2} D, \quad D = \tau(1 + \cos^2\theta)|G_M|^2 + \frac{3}{2}\sin^2\theta \left(|G_C|^2 + \frac{8}{9}\tau^2|G_Q|^2 \right), \quad (16)$$

where θ is the angle between the momenta of the deuteron (\vec{p}) and the electron beam (\vec{k}). Integrating the expression (16) with respect to the deuteron angular variables one obtains the following formula for the total cross section of the reaction (1)

$$\sigma_{tot}(e^+e^- \rightarrow \bar{d}d) = \frac{\pi\alpha^2\beta^3}{3q^2} \left[3|G_C|^2 + 4\tau(|G_M|^2 + \frac{2}{3}\tau|G_Q|^2) \right]. \quad (17)$$

One can define also an angular asymmetry, R , with respect to the differential cross section measured at $\theta = \pi/2, \sigma_0$

$$\frac{d\sigma^{un}}{d\Omega} = \sigma_0(1 + R\cos^2\theta), \quad (18)$$

where R can be expressed as a function of the deuteron FFs

$$R = \frac{2\tau(|G_M|^2 - \frac{4}{3}\tau|G_Q|^2) - 3|G_C|^2}{2\tau(|G_M|^2 + \frac{4}{3}\tau|G_Q|^2) + 3|G_C|^2}. \quad (19)$$

This observable should be sensitive to the different underlying assumptions on deuteron FFs; therefore, a precise measurement of this quantity, which does not require polarized particles, would be very interesting.

One can see that, as in the space-like (SL) region, the measurement of the angular distribution of the outgoing deuteron determines the modulus of the magnetic form factor, but the separation of the charge and quadrupole form factors requires the measurement of polarization observables [12]. The outgoing-deuteron polarization can be measured in a secondary analyzing scattering [12]. For vector polarization up to a few GeV, an inclusive measurement on a carbone target as $d + C \rightarrow \text{one charged particle} + X$ is sufficient, when the charged protons from deuteron break up are eliminated with help of an absorber [13]. For tensor polarization, however, only exclusive reactions as elastic $d + p$ scattering [14] or charge exchange [15] give sufficient efficiency and analyzing powers.

As it was shown in Ref. [16], a nonzero phase difference between FFs of two baryons (with 1/2 spins) leads to the T-odd single-spin asymmetry normal to the scattering plane in the baryon-antibaryon production $e^+e^- \rightarrow B\bar{B}$. It is more convenient to derive polarization observables in CMS. When considering the polarization of the final particle, we choose a reference system with the z axis along the momentum of this particle (in our case it is \vec{p}). The y axis is normal to the reaction plane in the direction of $\vec{k} \times \vec{p}$; x , y and z form a right-handed coordinate system.

The cross section can be written, in the general case, as the sum of unpolarized and polarized terms, corresponding to the different polarization states and polarization directions of the incident and scattered particles:

$$\frac{d\sigma}{d\Omega} = \frac{d\sigma^{un}}{d\Omega} [1 + P_y + \lambda P_x + \lambda P_z + P_{zz}R_{zz} + P_{xz}R_{xz} + P_{xx}(R_{xx} - R_{yy}) + \lambda P_{yz}R_{yz}], \quad (20)$$

where P_i (P_{ij}), $i, j = x, y, z$ are the components of the polarization vector (tensor) of the outgoing deuteron, R_{ij} , $i, j = x, y, z$ the components of the quadrupole polarization tensor of the outgoing deuteron $Q_{\mu\nu}$, in its rest system and $\frac{d\sigma^{un}}{d\Omega}$ is the differential cross section for the unpolarized case.

The degree of longitudinal polarization of the electron beam, λ , is explicitly indicated, in order to stress the origin of the specific polarization observables.

Let us consider the different polarization observables and give their expression in terms of the deuteron FFs.

- The vector polarization of the outgoing deuteron, P_y , which does not require polarization in the initial state is

$$P_y = -\frac{3}{2}\sqrt{\tau}\sin(2\theta)\text{Im}\left[\left(G_C - \frac{\tau}{3}G_Q\right)G_M^*\right]/D. \quad (21)$$

- The part of the differential cross section that depends on the tensor polarization can be written as follows

$$\frac{d\sigma_T}{d\Omega} = \frac{d\sigma_{zz}}{d\Omega}R_{zz} + \frac{d\sigma_{xz}}{d\Omega}R_{xz} + \frac{d\sigma_{xx}}{d\Omega}(R_{xx} - R_{yy}), \quad (22)$$

$$\frac{d\sigma_{zz}}{d\Omega} = \frac{\alpha^2\beta^3}{4q^2}\frac{3\tau}{4}\left[(1 + \cos^2\theta)|G_M|^2 + 8\sin^2\theta\left(\frac{\tau}{3}|G_Q|^2 - \text{Re}(G_C G_Q^*)\right)\right], \quad (23)$$

$$\frac{d\sigma_{xz}}{d\Omega} = -\frac{\alpha^2\beta^3}{4q^2}3\tau^{3/2}\sin(2\theta)\text{Re}(G_Q G_M^*), \quad (24)$$

$$\frac{d\sigma_{xx}}{d\Omega} = -\frac{\alpha^2\beta^3}{4q^2}\frac{3\tau}{4}\sin^2\theta|G_M|^2, \quad (25)$$

- Let us consider now the case of a longitudinally polarized electron beam. The other two components of the deuteron vector polarization (P_x, P_z) require the initial particle polarization and are

$$P_x = -3\frac{\sqrt{\tau}}{D} \sin\theta \operatorname{Re}\left(G_C - \frac{\tau}{3}G_Q\right) G_M^*, \quad P_z = \frac{3\tau}{2D} \cos\theta |G_M|^2. \quad (26)$$

From angular momentum and helicity conservations it follows that the sign of the deuteron polarization component P_z in the forward direction ($\theta = 0$) must coincide with the sign of the electron beam polarization. This requirement is satisfied by Eq. (26).

A possible nonzero phase difference between the deuteron FFs leads to another T-odd polarization observable proportional to the R_{yz} component of the tensor polarization of the deuteron. The part of the differential cross section that depends on the correlation between the longitudinal polarization of the electron beam and the deuteron tensor polarization can be written as follows

$$\frac{d\sigma_{\lambda T}}{d\Omega} = \frac{\alpha^2\beta^3}{4q^2} 6\tau^{3/2} \sin\theta \operatorname{Im}(G_M G_Q^*) R_{yz}. \quad (27)$$

The deuteron FFs in the TL region are complex functions. In the case of unpolarized initial and final particles, the differential cross section depends only on the squared modulus $|G_M|^2$ and on the combination $G = |G_C|^2 + \frac{8}{9}\tau^2|G_Q|^2$. So, the measurement of the angular distribution allows one to determine $|G_M|$ and the quantity G , as in the elastic electron-deuteron scattering.

Let us discuss which information can be obtained by measuring the polarization observables derived above. Three relative phases exist for three FFs, which we note as follows: $\alpha_1 = \alpha_M - \alpha_Q$, $\alpha_2 = \alpha_M - \alpha_C$, and $\alpha_3 = \alpha_Q - \alpha_C$, where $\alpha_M = \operatorname{Arg}G_M$, $\alpha_C = \operatorname{Arg}G_C$, and $\alpha_Q = \operatorname{Arg}G_Q$. These phases are important characteristics of FFs in the TL region since they result from the strong interaction between final particles.

Let us consider the ratio of the polarizations P_{yz} (let us remind that it requires a longitudinally polarized electron beam) and P_{xz} (when the electron beam is unpolarized). One finds:

$$R_1 = \frac{P_{xz}}{P_{yz}} = -\cos\theta \cot\alpha_1. \quad (28)$$

So, the measurement of this ratio gives us information about the relative phase α_1 . The measurement of another ratio of polarizations, $R_2 = P_{xz}/P_{xx}$ gives us information about the

quantity $|G_Q|$:

$$R_2 = \frac{P_{xz}}{P_{xx}} = 8\sqrt{\tau} \cot \theta \cos \alpha_1 \frac{|G_Q|}{|G_M|}. \quad (29)$$

This allows one to obtain the modulus of the charge FF, $|G_C|$, from the quantity G , known from the measurement of the differential cross section. The measurement of a third ratio

$$R_3 = \frac{P_y}{P_x} = -\cos \theta \frac{\sin \alpha_2 - r \sin \alpha_1}{\cos \alpha_2 - r \cos \alpha_1}, \quad r = \frac{\tau |G_Q|}{3 |G_C|} \quad (30)$$

allows to determine the phase difference α_2 . And at last, if we measure the ratio of the polarizations P_{zz} and P_{xx}

$$R_4 = \frac{P_{zz}}{P_{xx}} = -\frac{1}{\sin^2 \theta} \left[1 + \cos^2 \theta + 8 \sin^2 \theta \frac{|G_C| |G_Q|}{|G_M|^2} (r - \cos \alpha_3) \right] \quad (31)$$

we can obtain information about the third phase difference α_3 . Moreover, one can verify the relation:

$$\alpha_3 = \alpha_2 - \alpha_1.$$

Thus, the measurement of these polarization observables allows to fully determine the deuteron FFs in TL region.

Note that using the ratio of two polarization components that are simultaneously measured, greatly reduces systematic uncertainties. It is not necessary to know neither the beam polarization or the polarimeter analyzing power, since both of these quantities cancel in the ratio.

This procedure can be considered as the generalization of the polarization method proposed almost four decades ago [17], which could be applied only recently to elastic electron proton scattering [18].

Let us note here that, in principle, one should take into account the problem of the two-photon-exchange contribution, which, as mentioned in the Introduction, may become important at large momentum transfer, as it is expected that the reactions mechanisms are similar for the crossed channel (1). As it was shown in Ref. [19], if the detection of the final particles does not distinguish between deuteron and antideuteron, then the interference between one-photon and two-photon amplitudes does not contribute to the cross section of the reaction (1).

III. NUMERICAL ESTIMATIONS

In the previous section, the expressions of cross section and polarization observables have been given, in terms of the deuteron FFs. Numerical estimations require the knowledge of such FFs, in TL region. Due to the hermiticity of the electromagnetic current, FFs are real in the SL region, and complex in the TL region. At our knowledge, most of the existing parametrizations of these FFs are phenomenological fits to SL data, and are useful for different estimations and to plan corresponding experiments in that kinematical region. However, their analytical expressions were not built to obey fundamental properties of FFs. For example, their extension to the TL region does not induce any phase (i.e., the imaginary part of FFs is equal to zero).

Recent work in this direction [21] describes three different parametrizations of deuteron FFs describing the world data. The first one (Parametrization I) is a sum of inverse polynomial terms, where the first node of the corresponding FFs is introduced as a global multiplicative term. The number of free parameters, necessary to obtain $\chi^2/ndf = 1.5$, was 18.

The second parametrization is based on a previous work [22]. It is an attempt to find a global description based on the vector dominance model, satisfying the asymptotic conditions predicted by QCD at large momentum transfer, and leads to a 12 parameters fit.

The third parametrization is a sum of gaussians, with some physical constraints on the parameters, which are the width and the position of the maximum of the gaussians. In total the parametrization contains 33 parameters for $\chi^2/ndf = 1.5$.

In Ref. [23] a generalization of the nucleon model from Ref. [24] has been successfully applied to the deuteron case. Besides the fact that the VMD model [24] satisfies by construction some of the basic properties of FFs, its extension to the TL region is straightforward [25].

The basic idea of this parametrization is the presence of two components in the hadron structure: an intrinsic structure, very compact, characterized by a dipole (monopole) q^2 dependence and a meson cloud, which contains only the ρ , ϕ and ω (not the ρ) contributions, in the nucleon (deuteron) case. A very good description of all known data on deuteron electromagnetic FFs has been obtained, with as few as six free parameters and few evident physical constraints.

In principle, all these parametrizations are not predictive outside the kinematical domain where the experimental data have been fitted. Therefore, the extrapolation to TL region is just given for illustrative purposes. We give the predictions from one of the parametrizations from Ref. [21] (Parametrization I), and of the model from [23]. Note that the analytical form of all three parametrizations in [21] is such that only real terms are present in TL region. An imaginary part arises naturally from the analytical continuation of model [23], (for $q^2 \rightarrow -q^2$) due to the not integer nature of the exponent of the intrinsic part. Finite widths for the ϕ and ω meson contributions would also give rise to complexity, but it was not necessary to introduce them, for obtaining a good description of data in SL region.

We also consider an updated version of the model [5], based on unitarity and analyticity [26].

The q^2 dependence of these models is illustrated in Fig. 1 for the moduli and in Fig. 2 for the real and imaginary parts of the model from Ref. [23]. One can see that the three models coincide in the SL region, where they are constrained by the experimental data, but, outside this kinematical region, they show very different behavior. In particular, Parametrization I differs by few order of magnitude. Parametrization I does not show any singularity in TL region. Two poles coincide in TL region, for the models [26] and [23], as they correspond to the ω and ϕ contributions. More resonances are built, by construction, in the model [26], and occur in the unphysical region. These two models show a similar trend, near the threshold, for the moduli of FFs, however the sign, which is reflected in the relevant polarization observables, may differ.

From Fig. 2 one can see that in TL region, FFs from Ref. [23] display an imaginary part which is an order of magnitude smaller than the real part, as a consequence of the exponent of the term corresponding to the intrinsic part. As the model [26] fulfills by construction the unitary condition, its imaginary part starts at the deuteron anomalous threshold, $q^2 = 1.73m_\pi^2 \simeq 0.02 \text{ GeV}^2$. Concerning the model [23], the imaginary part is different from zero for $q^2 > 0.08 \text{ GeV}^2$.

The predictions for the different observables are shown in Fig. 3, for $E=1.9 \text{ GeV}$, not far above threshold. The three parametrizations, as expected, give very different results, especially concerning the predictions for the cross section (Fig. 3a), which just reflects the differences in the moduli of FFs. In spite of this, the angular distributions are very similar, as it appears from Fig. 3a, as it is driven by the underlying assumption of the one-photon

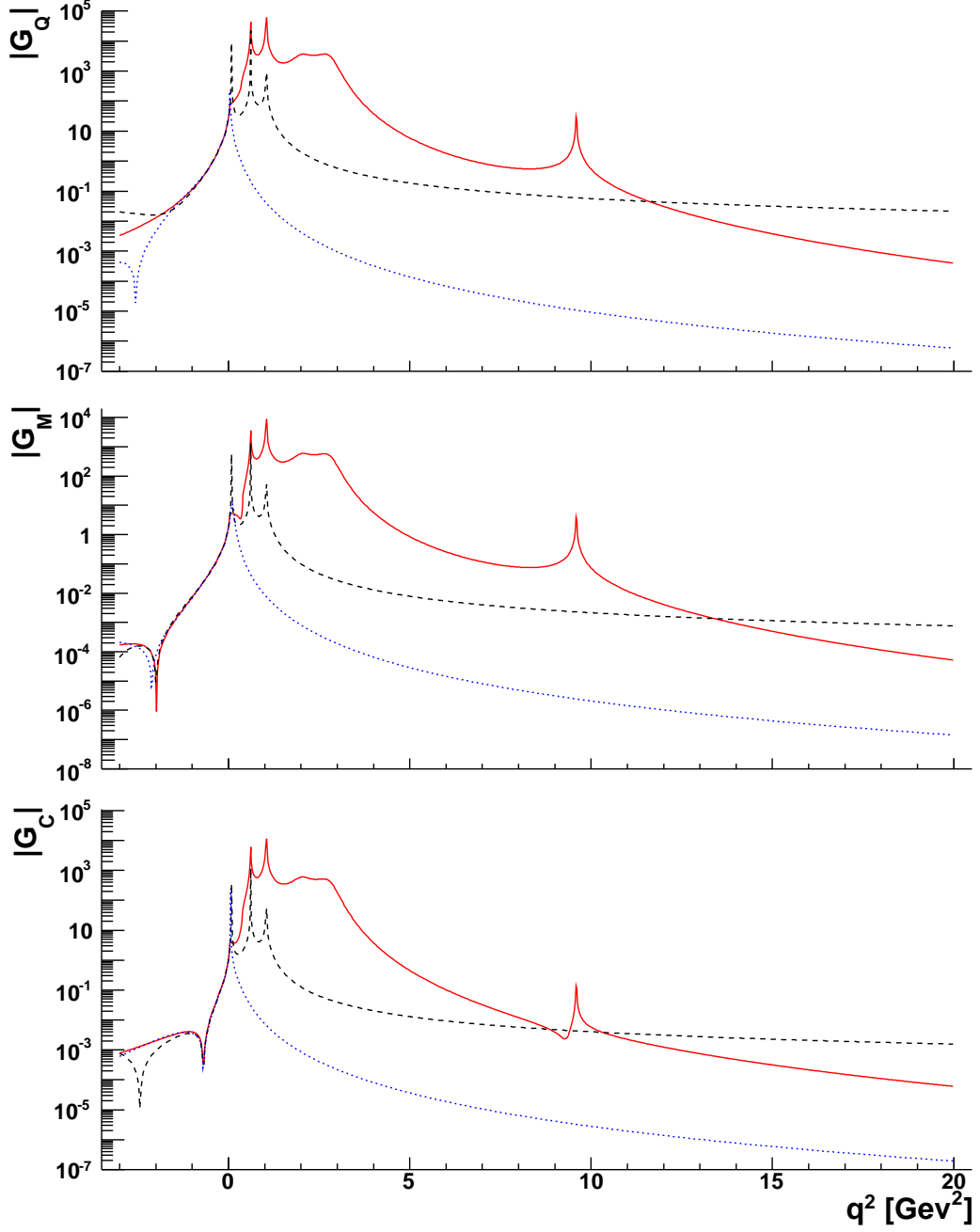


FIG. 1: q^2 -dependence of the G_Q , G_M , G_C from top to bottom (moduli): from Ref. [26] (solid line), from Ref. [23] (dashed line), and from Parametrization I from Ref. [21] (dotted line).

exchange mechanism.

Evidently, the observables such as P_y (Fig. 3d) and P_{yz} (Fig. 3i) vanish, for parametrization I, as they depend only on the imaginary part, see Eqs. (21) and (27), respectively.

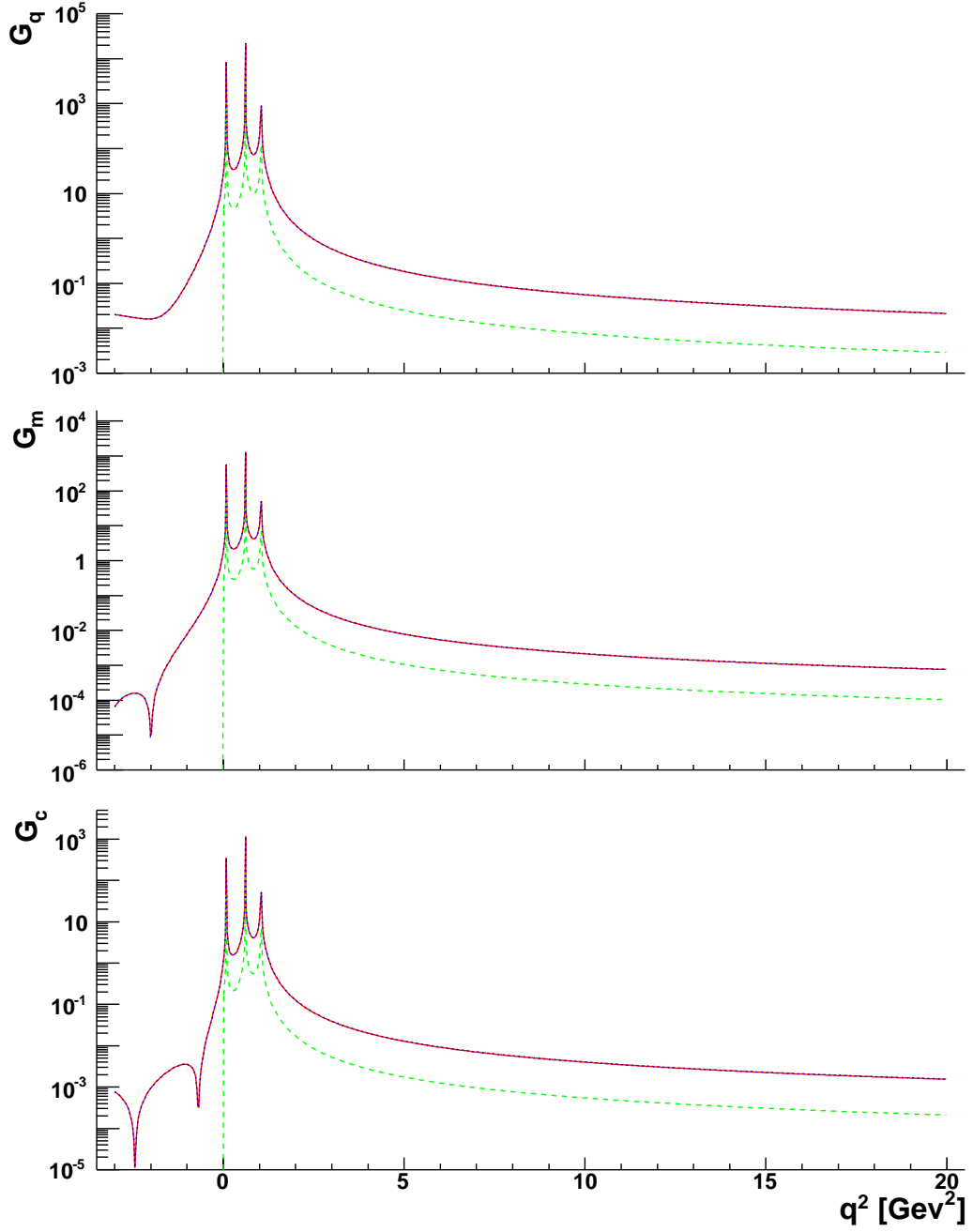


FIG. 2: q^2 -dependence of the G_Q , G_M , G_C from top to bottom: from Ref. [23]: real part (solid line), imaginary part (dashed line).

In the physical region, the angular asymmetry, Eq. (19), is very large in absolute value (over 90%) and negative, for all the considered models, due to the fact that one FF, G_Q , is dominant.

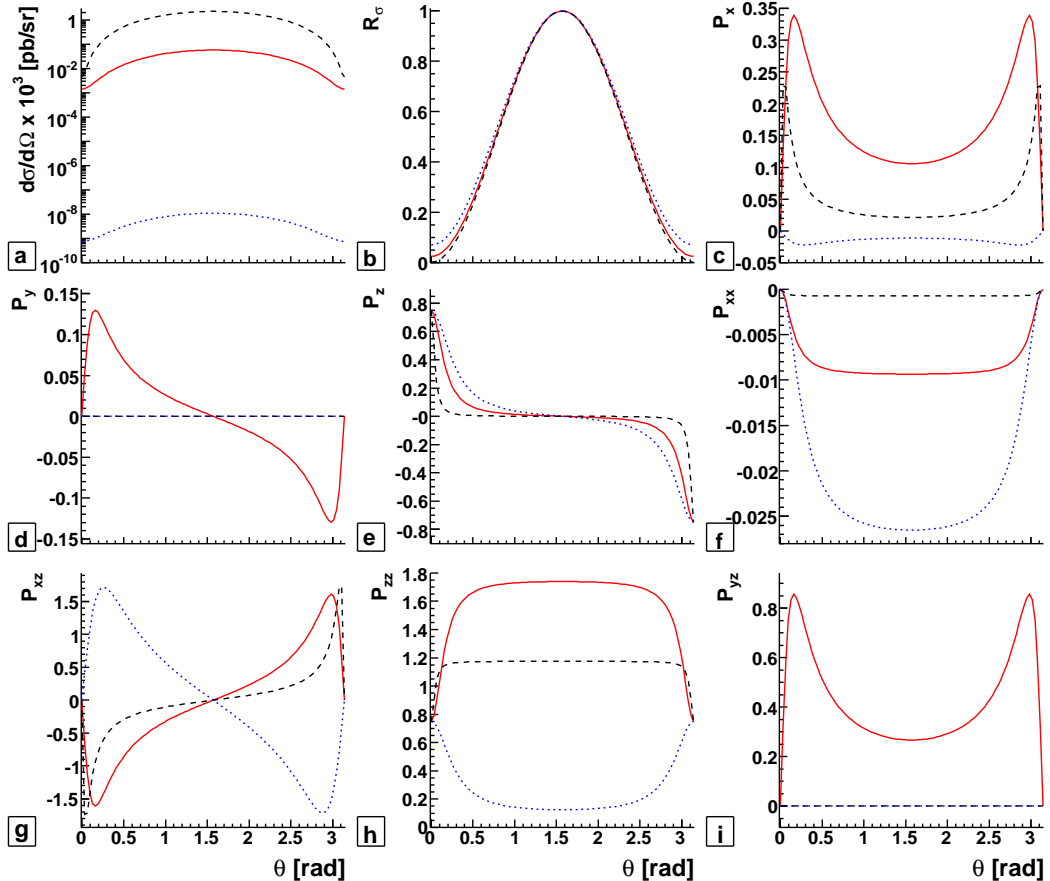


FIG. 3: Predictions of the different observables, for the considered parametrizations of deuteron FFs, extrapolated to the TL region. Notations as in Fig. 1.

It should be noted that the CMS threshold energy of the reaction $e^+ + e^- \rightarrow \bar{d} + d$ is quite large, $E_T = 2M \simeq 3.75$ GeV, which corresponds to $q^2 \simeq 14$ GeV². There are no data in this momentum range in SL region, which could better constrain models and parametrizations.

IV. CONCLUSIONS

Polarization observables have been derived for the production of a deuteron antideuteron pair in electron-positron annihilation. Although the cross section of this process is expected to be very small, the search for the corresponding events it is not excluded in future at high luminosity e^+e^- rings.

In TL region, the electromagnetic structure of the deuteron is characterized by three complex FFs. Generalizing the polarization method, successfully applied to ep elastic scat-

tering, we derive the expressions for the relevant observables in terms of the deuteron FFs and indicate the measurements which are necessary for the full determination of the deuteron structure.

Quantitative estimations require the knowledge of the deuteron FFs, in the corresponding kinematical region. Data are absent in the whole TL region, and also in SL region, at large momentum transfer squared. Therefore, we used the analytical continuations from the SL region of few existing parametrizations and models, keeping in mind that they are poorly constrained in the corresponding SL kinematical region. The results show that polarization effects either vanish or are large and measurable.

The formalism developed here is model independent and based on symmetry properties of electromagnetic and strong interactions. It allows to establish properties of observables that should be satisfied by any model calculation. Moreover, it applies as well to the annihilation reactions involving the production of spin one particles in the final state, such as $e^+ + e^- \rightarrow \rho^+ + \rho^-$, $e^+ + e^- \rightarrow \omega^+ + \omega^-$. The study of these reactions will be the object of a future work.

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VI. APPENDIX

In this Appendix we give useful formulae describing the polarization state of the deuteron for different cases. For the case of arbitrary polarization, the deuteron is described by the spin-density matrix (defined, in the general case, by 8 parameters) which, in the coordinate representation, has the form

$$\rho_{\mu\nu} = -\frac{1}{3}\left(g_{\mu\nu} - \frac{p_\mu p_\nu}{M^2}\right) + \frac{i}{2M}\varepsilon_{\mu\nu\lambda\rho} s_\lambda p_\rho + Q_{\mu\nu}, \quad Q_{\mu\nu} = Q_{\nu\mu}, \quad Q_{\mu\mu} = 0, \quad p_\mu Q_{\mu\nu} = 0, \quad (\text{C.1})$$

where p_μ is the deuteron four momentum, s_μ and $Q_{\mu\nu}$ are the deuteron polarization four vector and the deuteron quadrupole polarization tensor.

In the deuteron rest frame the above formula is written as

$$\rho_{ij} = \frac{1}{3}\delta_{ij} - \frac{i}{2}\varepsilon_{ijk}s_k + Q_{ij}, \quad ij = x, y, z. \quad (\text{C.2})$$

This spin–density matrix can be written in the helicity representation using the following relation

$$\rho_{\lambda\lambda'} = \rho_{ij}e_i^{(\lambda)*}e_j^{(\lambda')}, \quad \lambda, \lambda' = +, -, 0, \quad (\text{C.3})$$

where $e_i^{(\lambda)}$ are the deuteron spin functions which have the deuteron spin projection λ on to the quantization axis (z axis). They are

$$e^{(\pm)} = \mp \frac{1}{\sqrt{2}}(1, \pm i, 0), \quad e^{(0)} = (0, 0, 1). \quad (\text{C.4})$$

The elements of the spin–density matrix in the helicity representation are related to the ones in the coordinate representation by such a way

$$\begin{aligned} \rho_{\pm\pm} &= \frac{1}{3} \pm \frac{1}{2}s_z - \frac{1}{2}Q_{zz}, \quad \rho_{00} = \frac{1}{3} + Q_{zz}, \quad \rho_{+-} = -\frac{1}{2}(Q_{xx} - Q_{yy}) + iQ_{xy}, \\ \rho_{+0} &= \frac{1}{2\sqrt{2}}(s_x - is_y) - \frac{1}{\sqrt{2}}(Q_{xz} - iQ_{yz}), \quad \rho_{-0} = \frac{1}{2\sqrt{2}}(s_x + is_y) + \frac{1}{\sqrt{2}}(Q_{xz} + iQ_{yz}), \\ \rho_{\lambda\lambda'} &= (\rho_{\lambda'\lambda})^*. \end{aligned} \quad (\text{C.5})$$

To obtain these relations we use $Q_{xx} + Q_{yy} + Q_{zz} = 0$.

When the deuteron is used as a target, the spin matrix is diagonal, and the polarization state is described by the population numbers n_+ , n_- and n_0 . Here n_+ , n_- and n_0 are the fractions of the atoms with the nuclear spin projection on to the quantization axis $m = +1$, $m = -1$ and $m = 0$, respectively. If the spin–density matrix is normalized to 1, i.e. $\text{Tr}\rho = 1$, then we have $n_+ + n_- + n_0 = 1$. Thus, the polarization state of the deuteron target is defined in this case by two parameters, called V (vector) and T (tensor) polarizations

$$V = n_+ - n_-, \quad T = 1 - 3n_0. \quad (\text{C.6})$$

Using the definitions for the quantities $n_{\pm,0}$

$$n_{\pm} = \rho_{ij}e_i^{(\pm)*}e_j^{(\pm)}, \quad n_0 = \rho_{ij}e_i^{(0)*}e_j^{(0)}, \quad (\text{C.7})$$

we have the following relation between V and T parameters and parameters of the spin-density matrix in the coordinate representation (with the quantization axis directed along the z axis)

$$n_0 = \frac{1}{3} + Q_{zz}, \quad n_{\pm} = \frac{1}{3} \pm \frac{1}{2}s_z - \frac{1}{2}Q_{zz}, \quad (\text{C.8})$$

or

$$T = -3Q_{zz}, \quad V = s_z. \quad (\text{C.9})$$

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