

# Search for a neutral heavy Higgs boson in the $H \rightarrow ZZ$ decay channels with ATLAS

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IRFU DDays, 1/7/2015



# About me

Enrolled at ED 517 (Université Paris-Sud)

Supervisor: Rosy Nikolaidou

Doing my PhD on the ATLAS experiment at the LHC at CERN

My thesis has two components:

## 1) Physics analysis

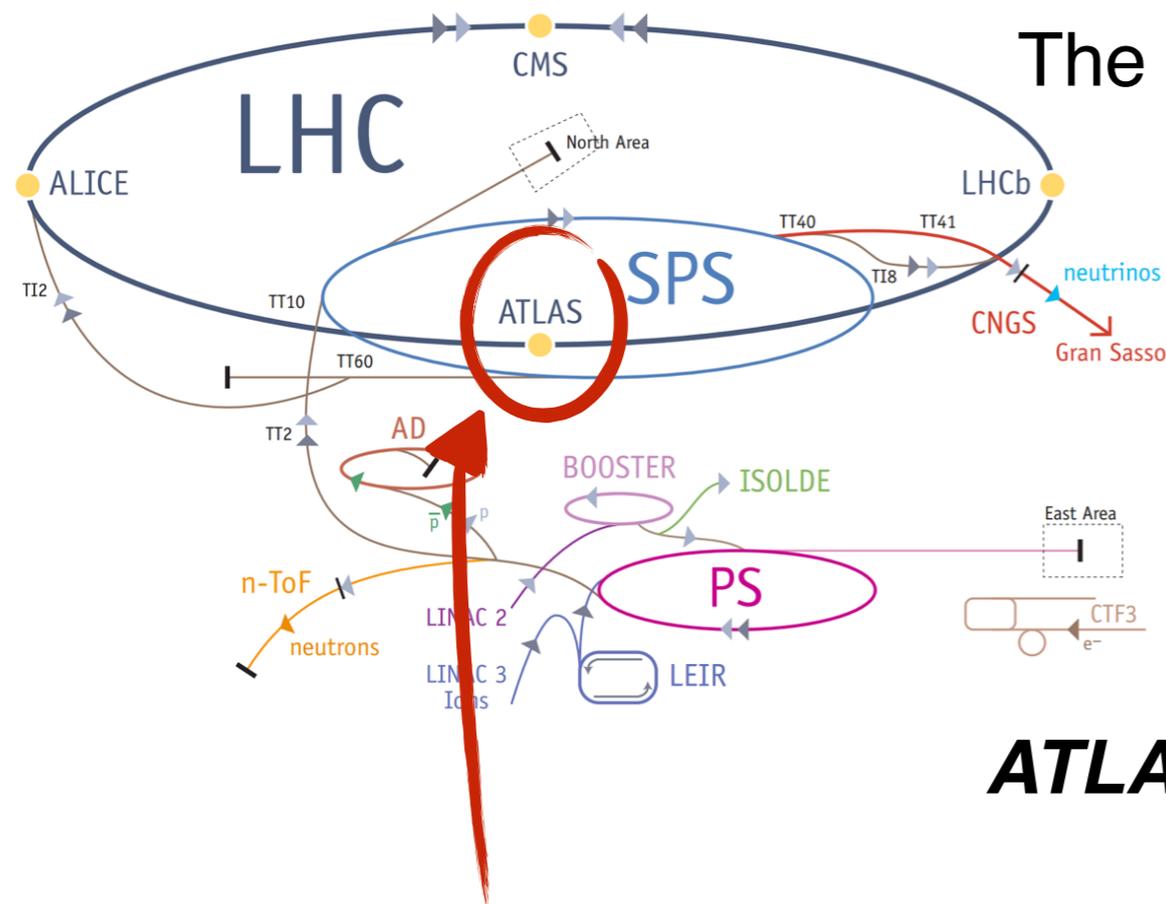
- Search for a beyond the Standard Model heavy Higgs boson in  $H \rightarrow ZZ$  decay channels

## 2) Detector physics

- Development and commissioning of gas and scintillation detectors
- Prototyping of MicroMegas detectors for the upgrade of the ATLAS Muon Spectrometer

Today's presentation will focus on the heavy Higgs search

# The LHC and the ATLAS Experiment

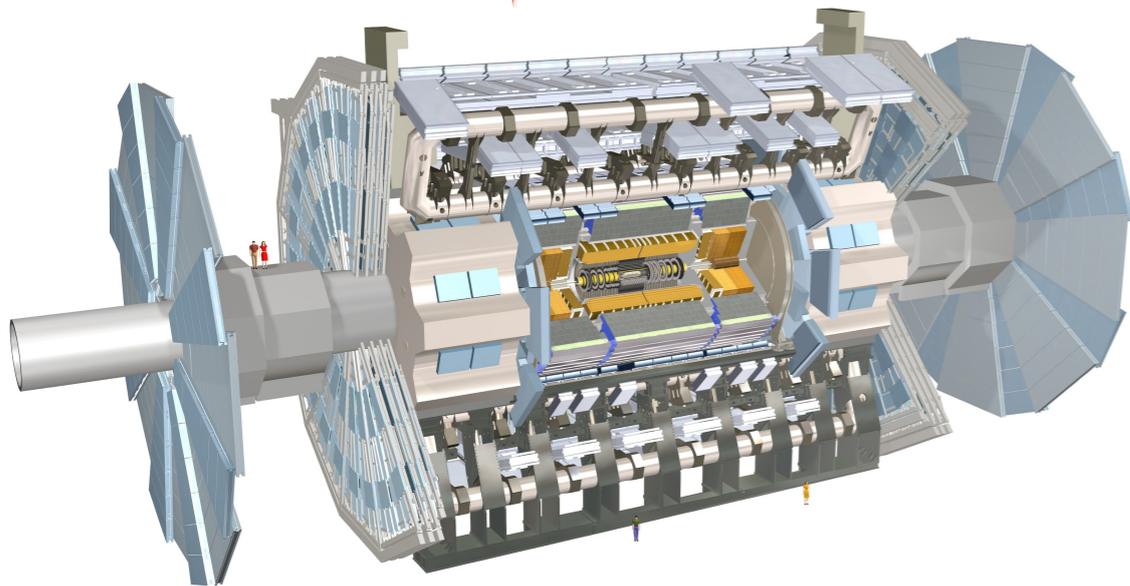


The Large Hadron Collider (LHC) is a ~27 km proton-proton collider situated at CERN

The LHC is the largest and most powerful particle collider ever built

- Able to probe a previously inaccessible kinematic regime

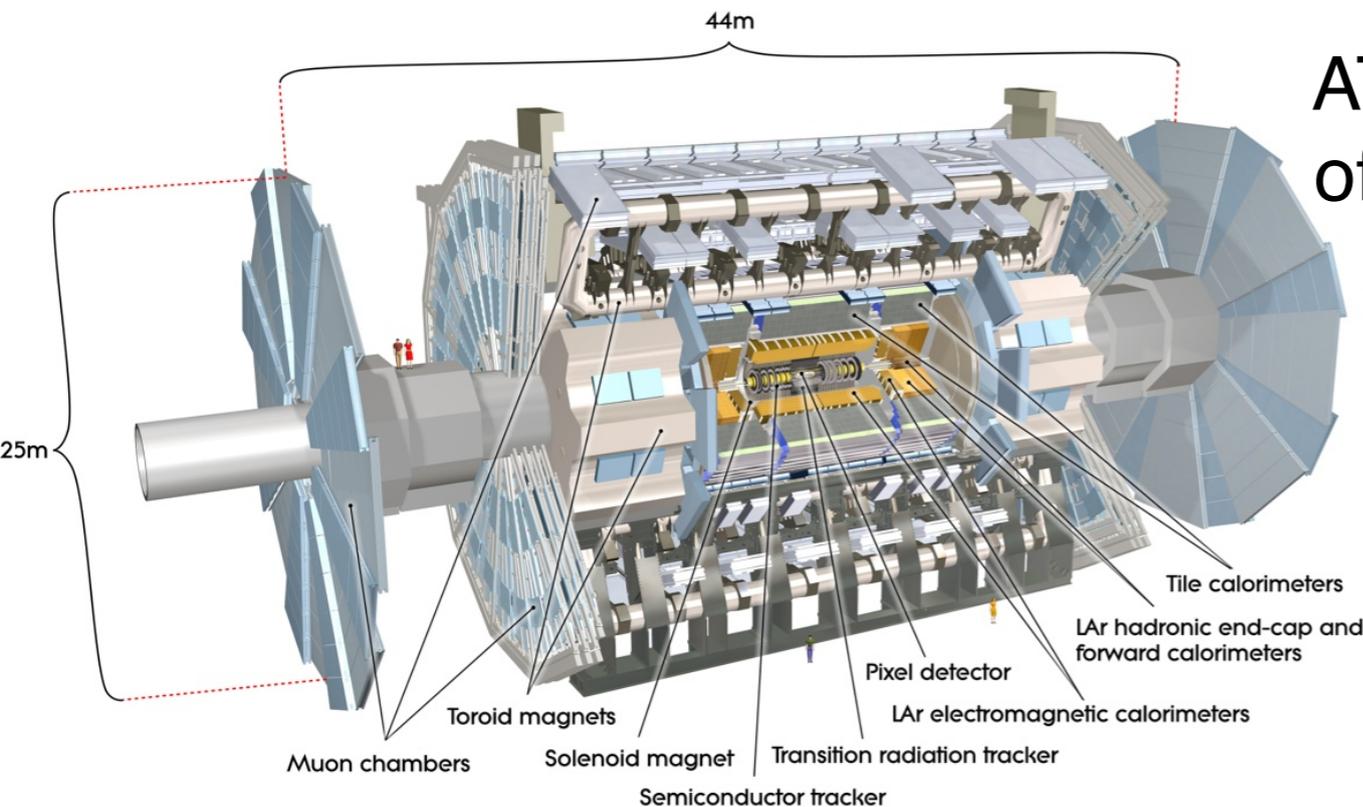
**ATLAS** is one of four experiments at the LHC



ATLAS is designed as a *general purpose* experiment to perform a broad range of physics measurements and searches

- Higgs boson
- Super Symmetric (SUSY) particles
- Heavy, exotic resonances
- Many more!

# The LHC and the ATLAS Experiment



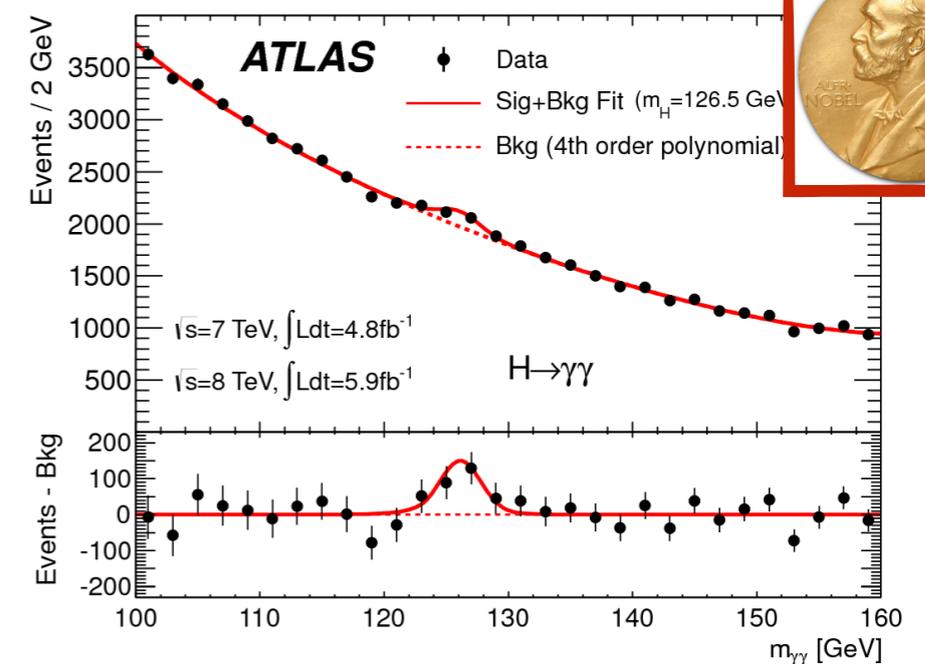
ATLAS works by reconstructing a variety of final state particles with a composition of different detector technologies

- Photons
- Electrons
- Muons
- Jets
- Missing transverse energy

ATLAS successfully took data in **Run-1**, the period from 2010 - 2013

- $5.1 \text{ fb}^{-1}$  recorded at  $\sqrt{s} = 7 \text{ TeV}$
- $21.3 \text{ fb}^{-1}$  recorded at  $\sqrt{s} = 8 \text{ TeV}$

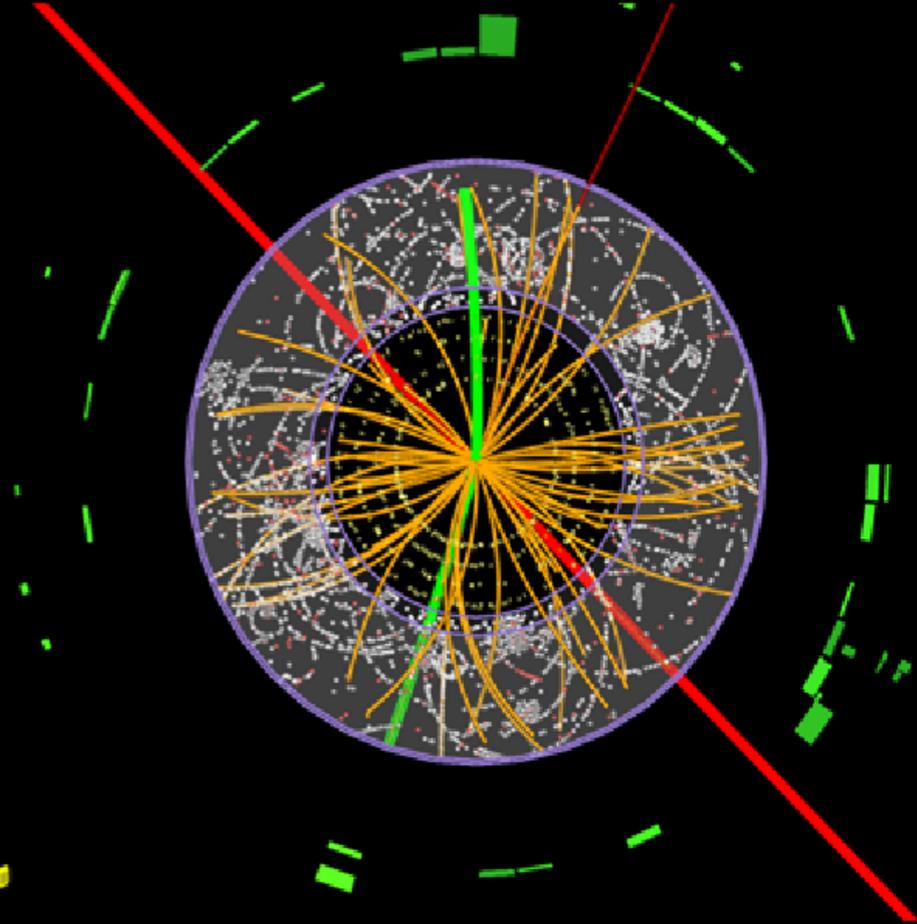
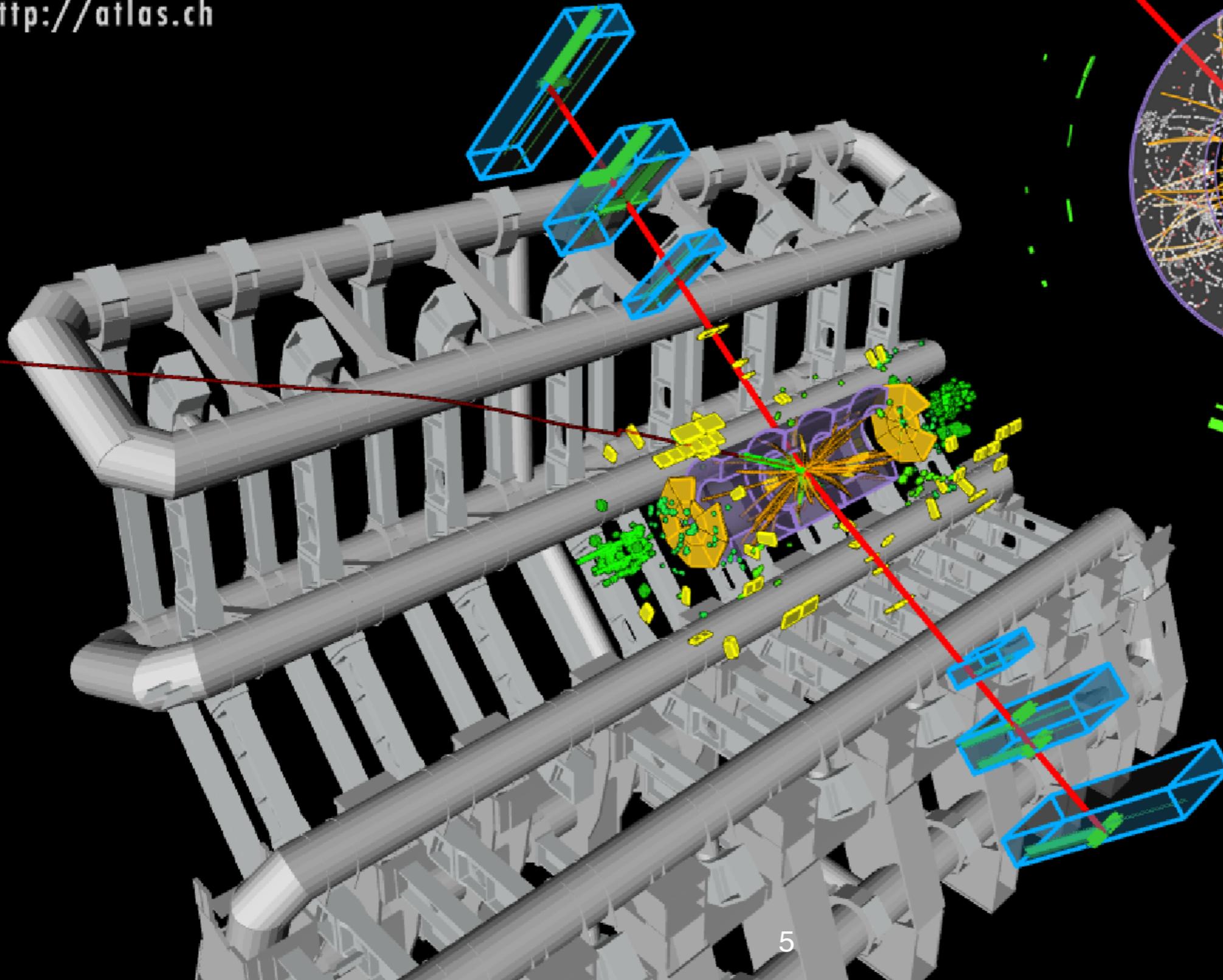
**The Run-1 dataset constituted the foundation for the Higgs boson discovery in 2012**



# The LHC and the ATLAS Experiment

 **ATLAS**  
EXPERIMENT  
<http://atlas.ch>

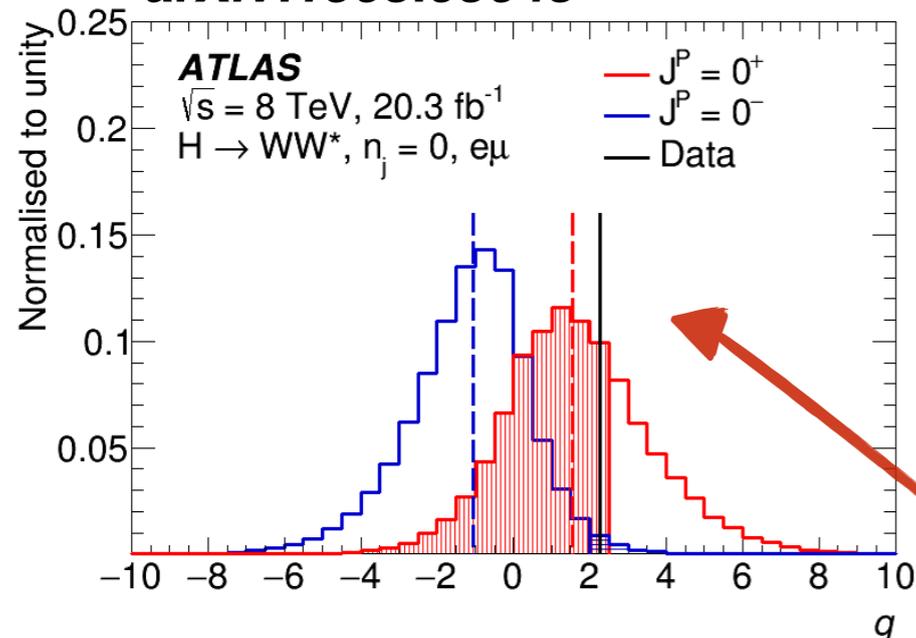
**$h \rightarrow ZZ \rightarrow 2e2\mu$  event**



Run: 186877  
Event: 84622334  
2011-08-05  
15:03:21 CEST

# Introduction to heavy Higgs search

arXiv:1503.03643



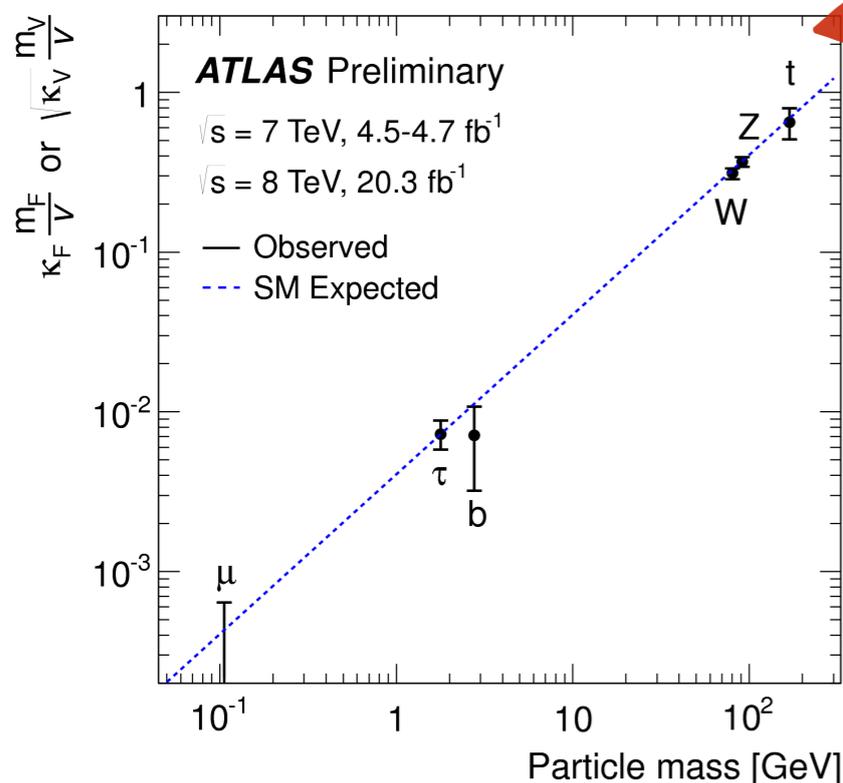
The new boson strongly resembles the Standard Model (SM) Higgs boson

- The last undiscovered particle in the SM

Its compatibility with the SM Higgs boson was determined by measuring its *properties*

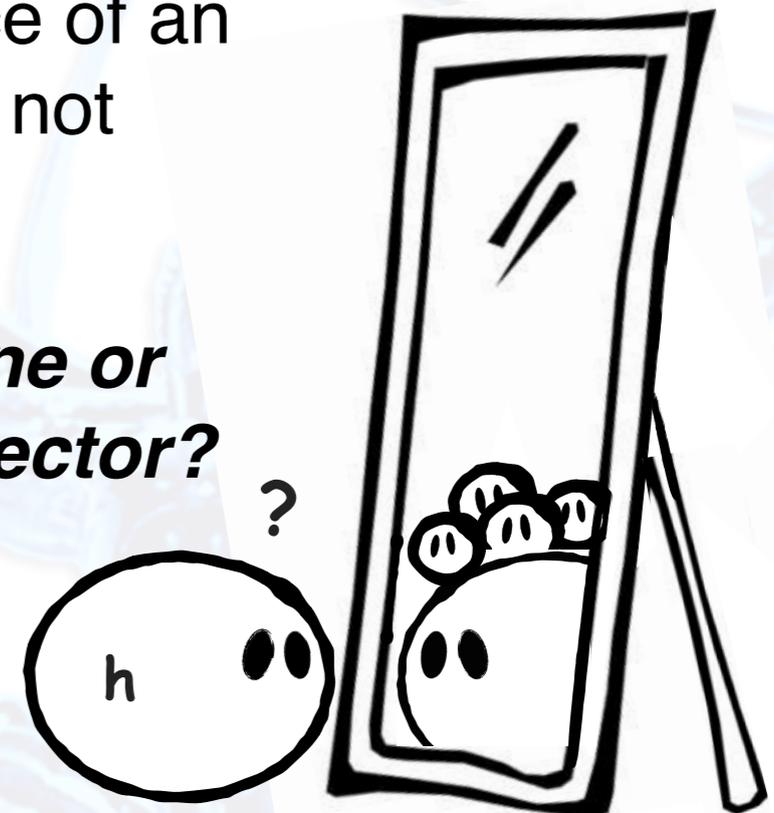
- Mass, spin, coupling strength to other particles

ATLAS-CONF-2015-007



Although this new particle is compatible with the SM, the existence of an extended Higgs sector is not ruled out

***Is our Higgs boson alone or part of a larger Higgs sector?***



# Introduction to heavy Higgs search

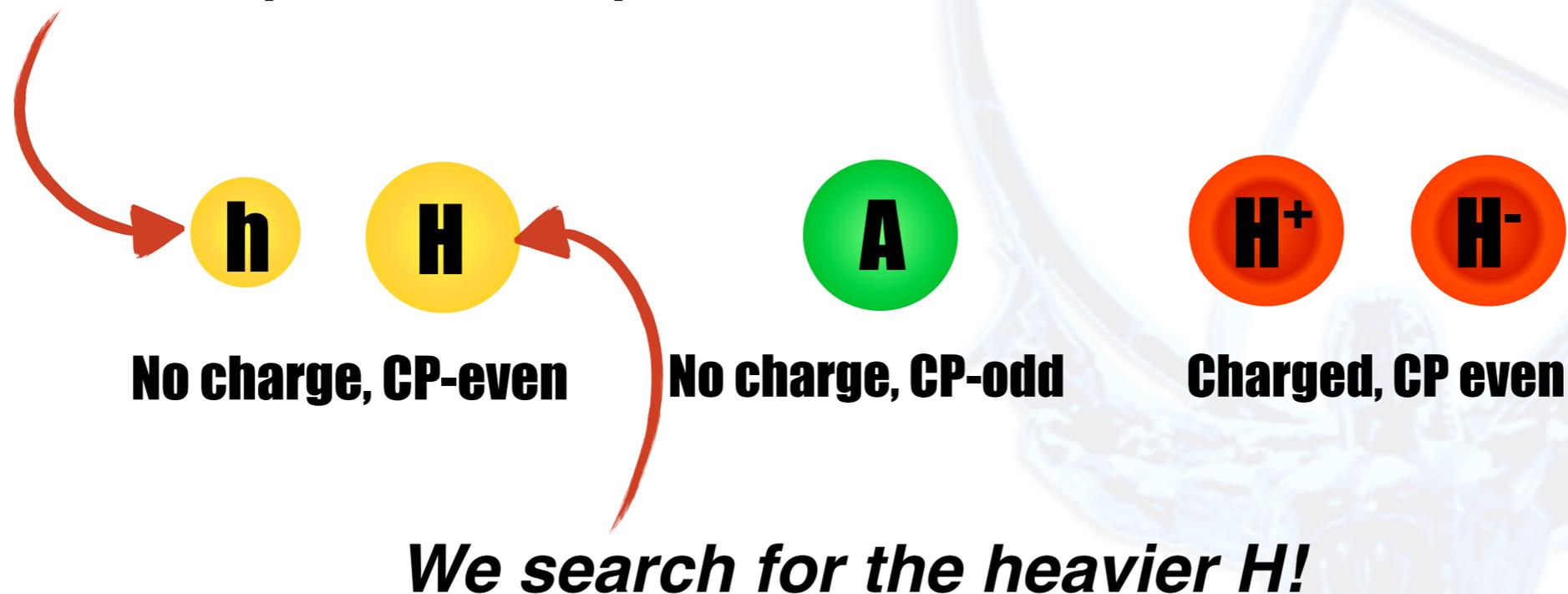
In the SM one Higgs boson appears after spontaneous symmetry breaking

One of the simplest extensions to the SM is obtained by adding a second electroweak doublet: **2-Higgs Doublet Model (2HDM)**

- A very generic model that constitutes a basis for more advanced theories
  - The Higgs sector in Super Symmetry (SUSY) is a 2HDM variety

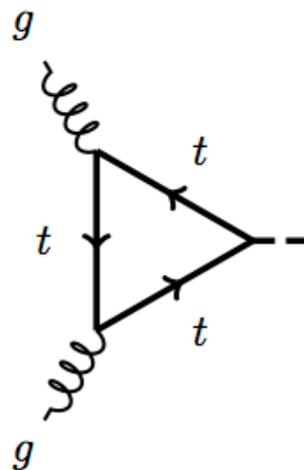
In 2HDM, multiple Higgses appear after spontaneous symmetry breaking

- ***Our 125 GeV particle interpreted as  $h$***



# Main Higgs production modes at the LHC

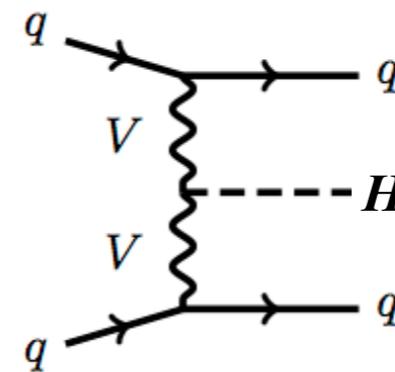
## Gluon fusion (ggF)



### Signature

- Higgs produced with no by-products

## Vector boson fusion (VBF)

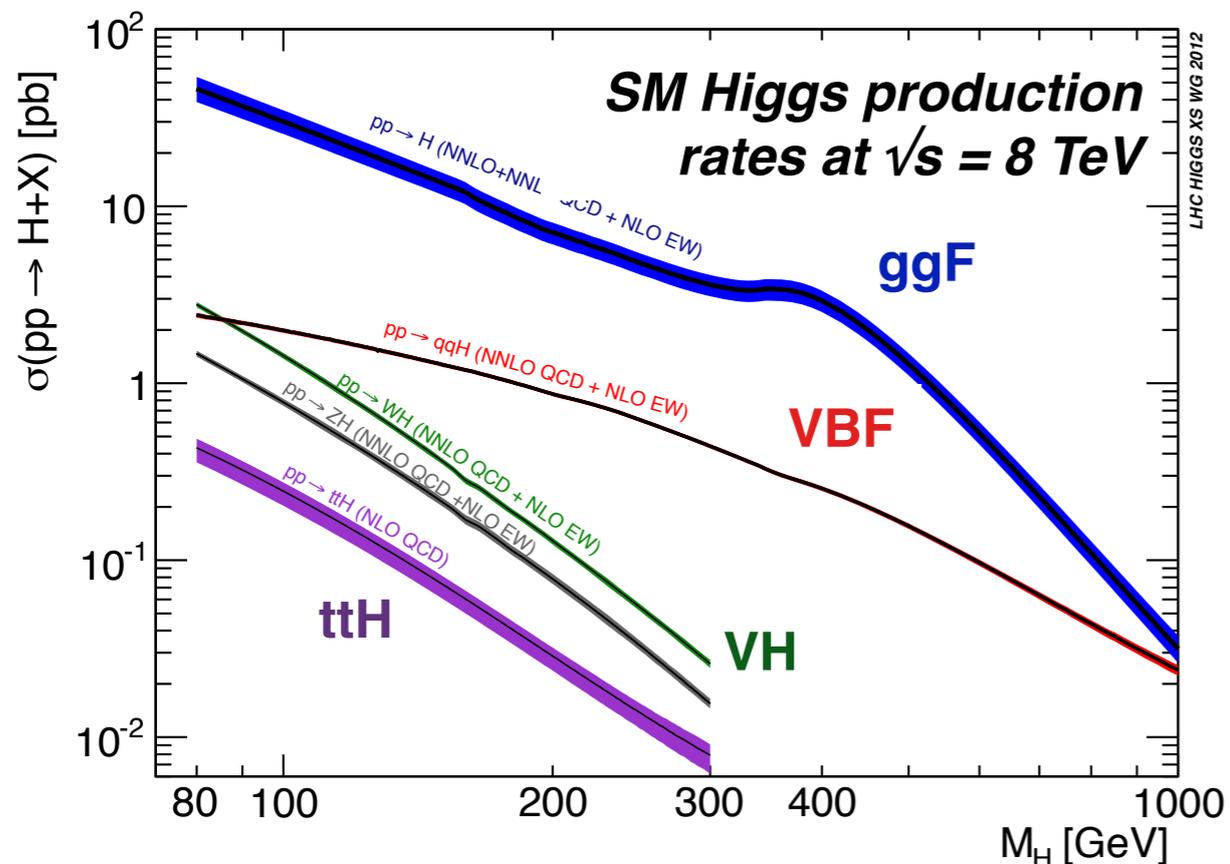


### Signature

- Higgs produced with two energetic jets in opposite, forward region

**Assume  $H$  will also be produced through these mechanisms**

- Production rate heavily dependent on its mass



Number of Higgs bosons produced via ggF with  $21.3 \text{ fb}^{-1}$  at  $\sqrt{s} = 8$  TeV

- $m_H = 125 \text{ GeV}$ : 410k
- $m_H = 300 \text{ GeV}$ : 77k
- $m_H = 600 \text{ GeV}$ : 11k
- $m_H = 900 \text{ GeV}$ : 1k

- $H \rightarrow ZZ$  branching ratio and experimental acceptance should be multiplied to these numbers

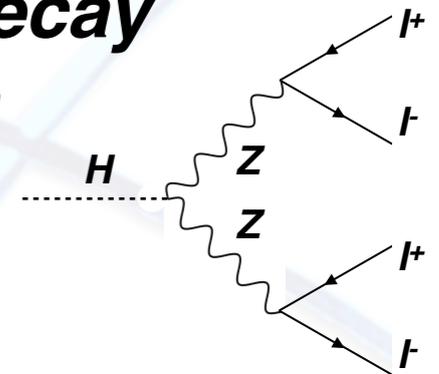
# Overview of search

We search for a heavy Higgs in the range 140 GeV - 1000 GeV

The Run-1 dataset (20.3 fb<sup>-1</sup> @  $\sqrt{s} = 8$  TeV) is used

**We search independently in four different  $H \rightarrow ZZ$  decay channels and afterwards statistically combine them**

- First analysis on ATLAS to combine these channels
- The full ATLAS detector is used to reconstruct the many different final state objects  $\rightarrow$  *complex analysis!*

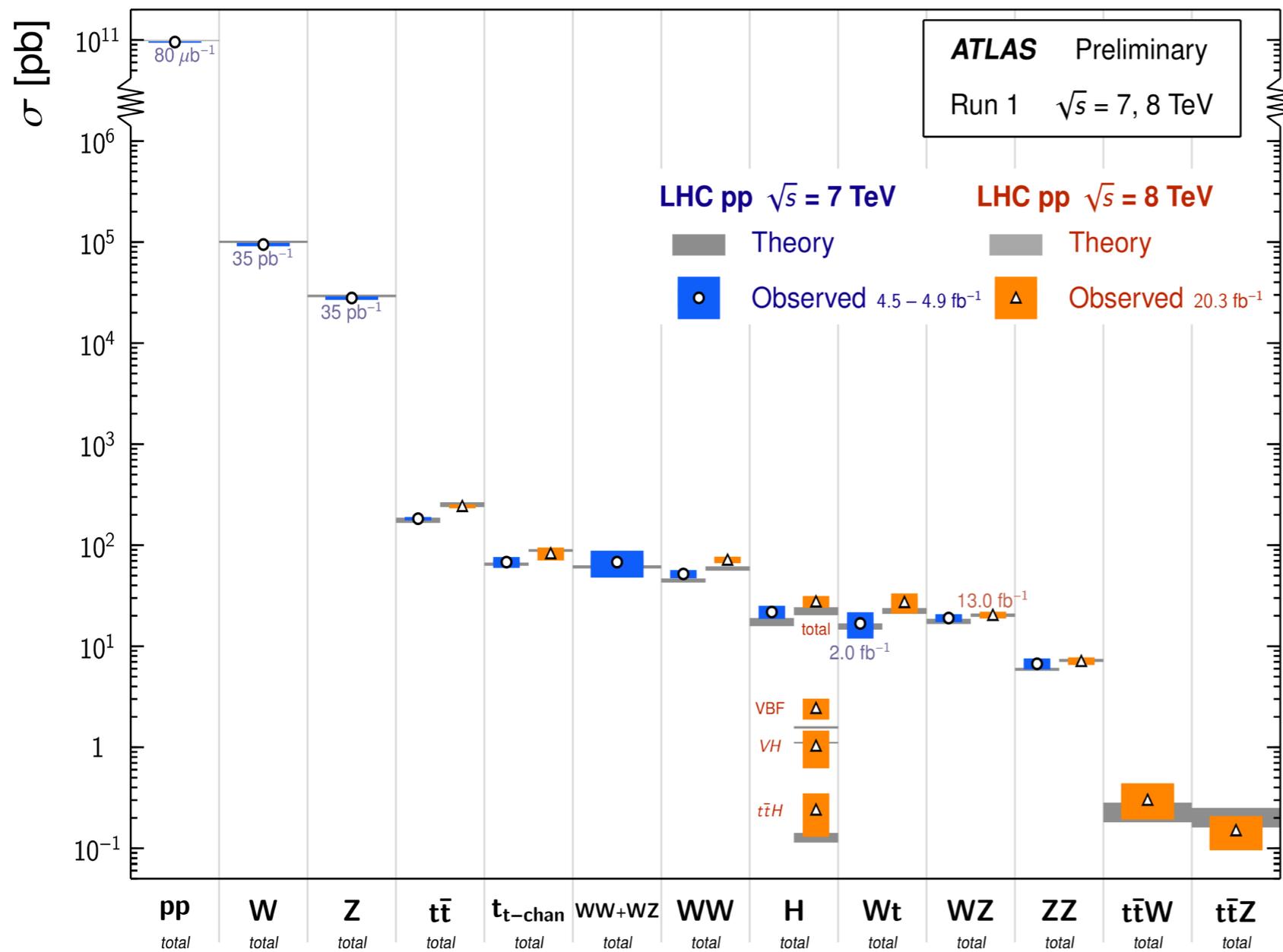


Channel	Sensitive mass range	Reconstructed final state
$H \rightarrow ZZ \rightarrow 4l$	Low	2 lepton pairs ( $\mu^+\mu^-$ or $e^+e^-$ )
$H \rightarrow ZZ \rightarrow 2l2\nu$	Intermediate	1 lepton pair ( $\mu^+\mu^-$ or $e^+e^-$ ) Missing transverse energy
$H \rightarrow ZZ \rightarrow 2l2q$	High	1 lepton pair ( $\mu^+\mu^-$ or $e^+e^-$ ) 2 jets
$H \rightarrow ZZ \rightarrow 2\nu2q$	High	Missing transverse energy 2 jets

Channels have different mass resolution (O(1) - O(100) GeV)

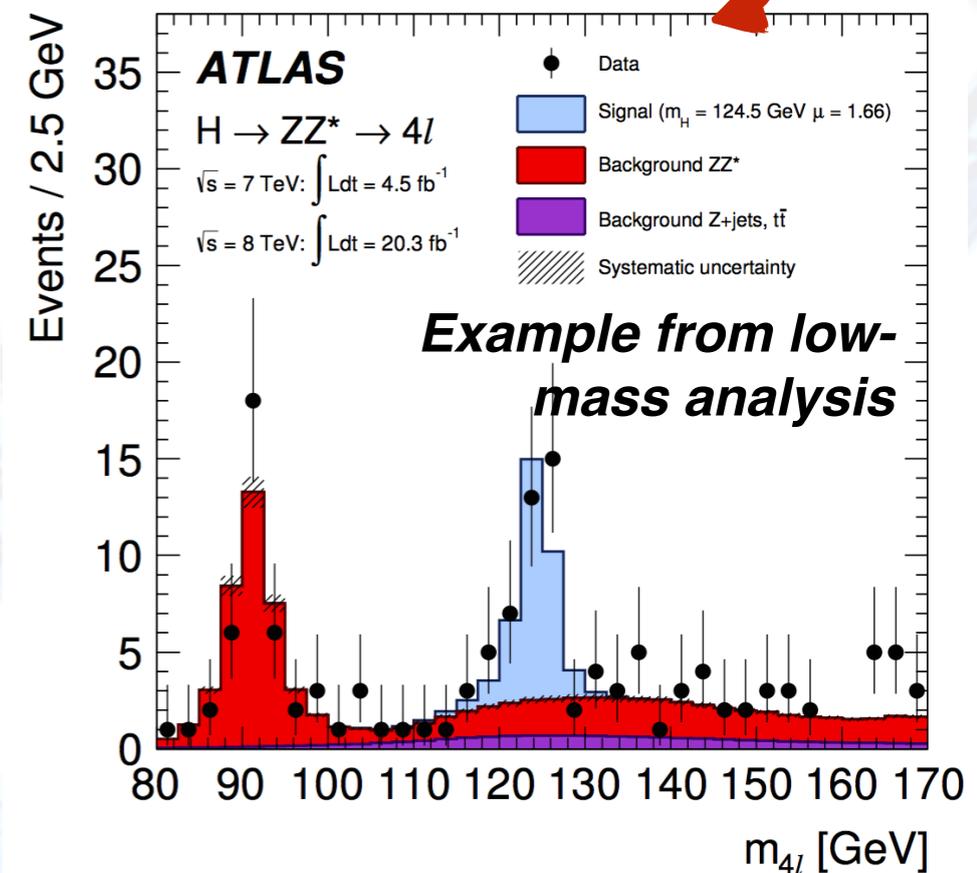
# Overview of search

Standard Model Total Production Cross Section Measurements Status: March 2015



The principle behind each analysis is to separate signal-like events from the many background events

- At first events are sorted with a sequence of requirements (**cuts**) imposed to data
- Afterwards the shape of a **discriminant** is used

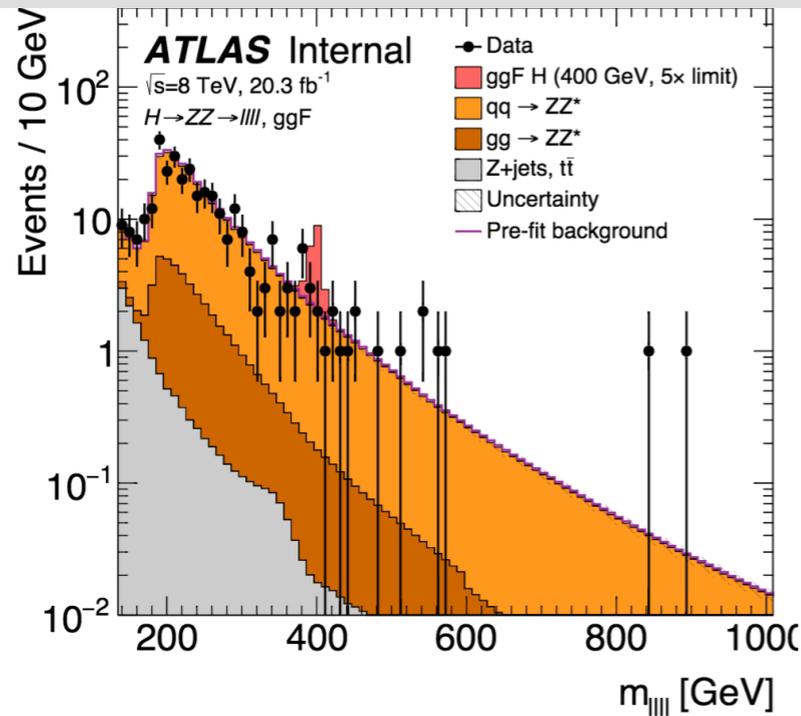


**My contribution was to the  $H \rightarrow ZZ \rightarrow 4l$  analysis, to the combination of the decay channels and to the statistical treatment**

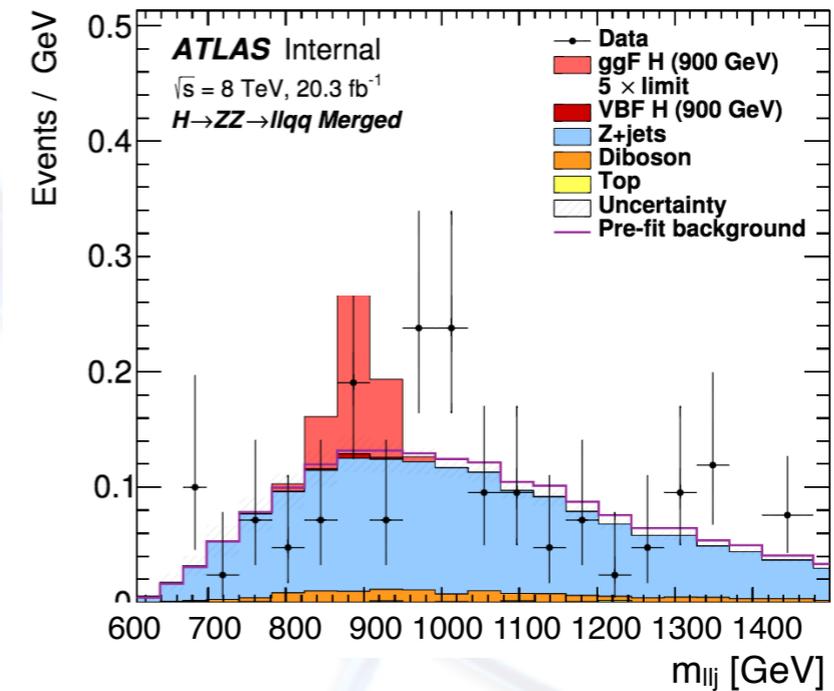
# Overview of search

## Discriminants of the four different $H \rightarrow ZZ$ decay channels

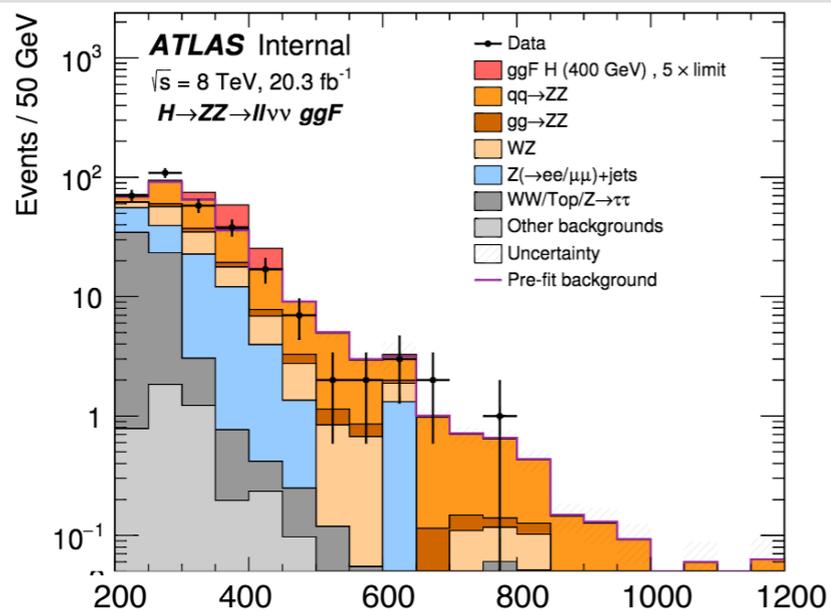
### 4l: inv. mass of 4 leptons



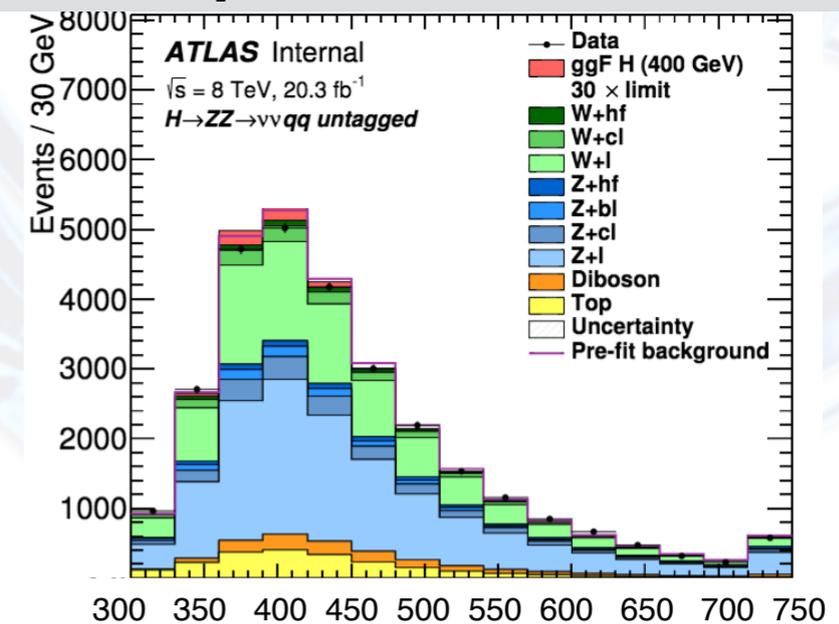
### 2l2q: inv. mass of jets and leptons



### 2l2v: transverse mass



### 2v2q: transverse mass



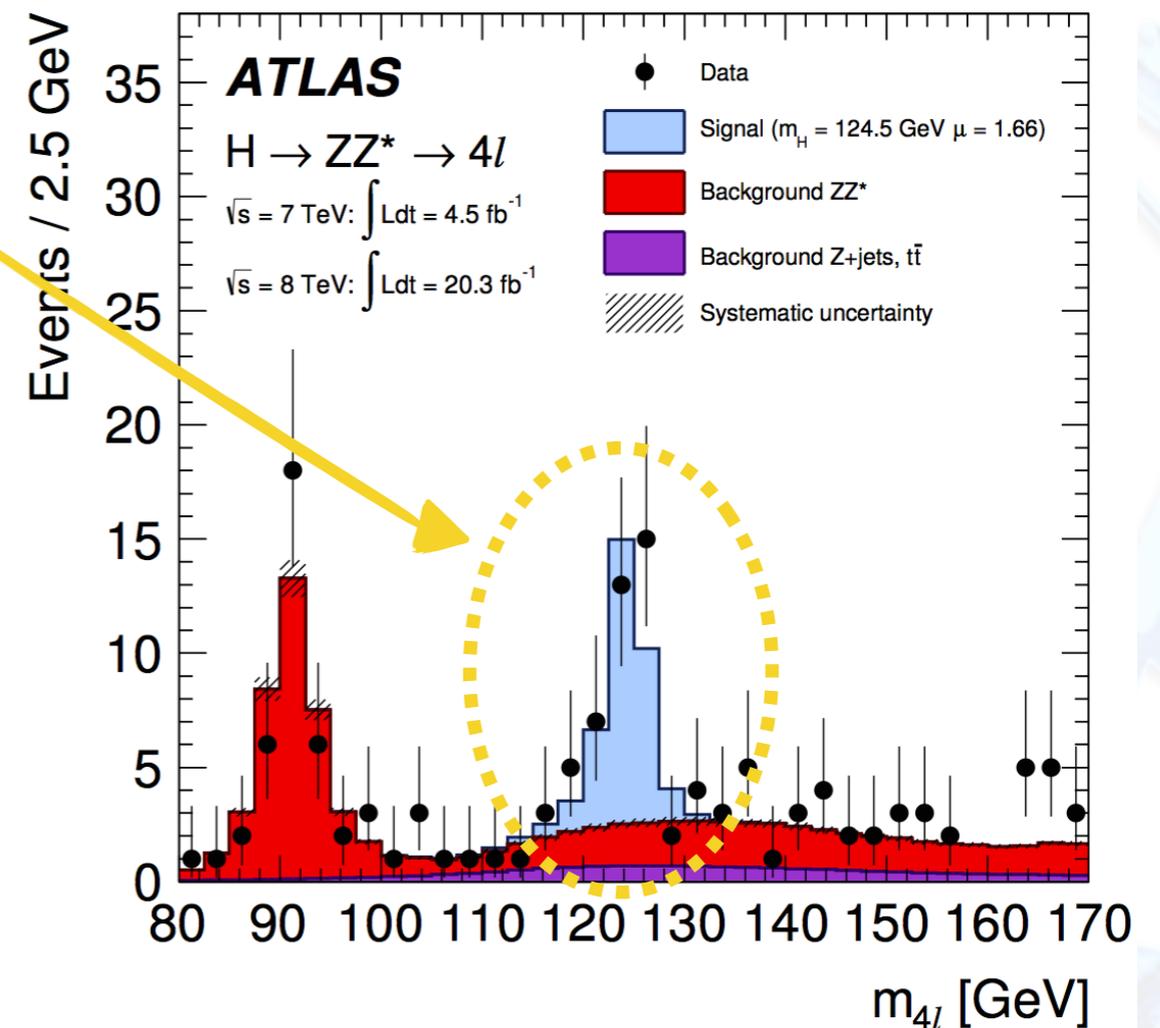
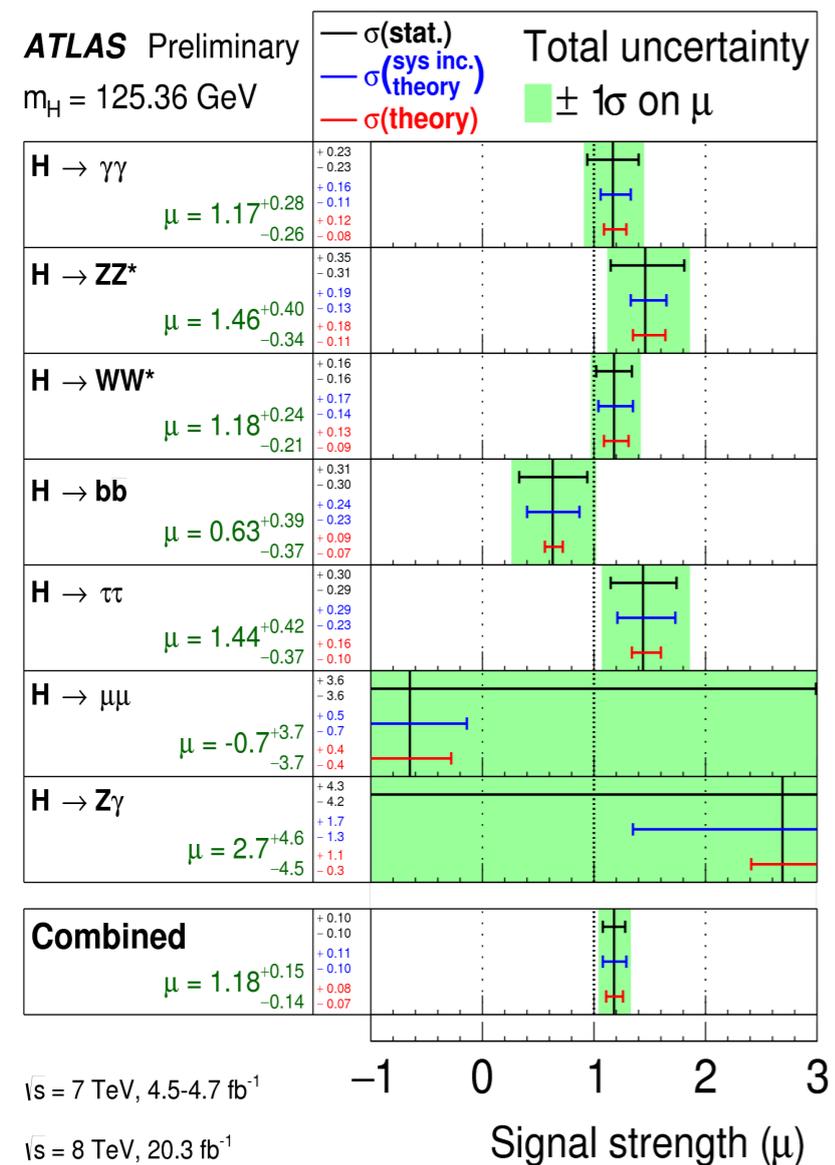
$$(m_T^{ZZ})^2 \equiv \left( \sqrt{m_Z^2 + |\vec{p}_T^{\ell\ell}|^2} + \sqrt{m_Z^2 + |E_T^{\text{miss}}|^2} \right)^2 - |\vec{p}_T^{\ell\ell} + \vec{E}_T^{\text{miss}}|^2$$

$$(m_T^{ZZ})^2 \equiv \left( \sqrt{m_Z^2 + |\vec{p}_T^{jj}|^2} + \sqrt{m_Z^2 + |E_T^{\text{miss}}|^2} \right)^2 - |\vec{p}_T^{jj} + \vec{E}_T^{\text{miss}}|^2$$

# Statistical treatment and combination

The signal production rate is quantified with the **signal strength  $\mu$**

- Acts as signal scale factor



$\mu = \text{observed rate/rate predicted by theory}$

- $\mu = 0$ : data matches background only prediction
- $\mu = 1$ : data matches theoretical signal prediction

Our final results rely on the determination of  $\mu$

# Statistical treatment and combination

$\mu$  is determined by the use of a **likelihood function**

- Parametrisation of our assumptions about signal and background
- Represents the probability for our dataset to originate from our model with a signal scaled by  $\mu$

$$\mathcal{L}(\mu) = \text{Pois}(x_1 \dots x_n | \mu S + B) \left[ \prod_{e=1}^n \frac{\mu S f_s(x_e) + B f_B(x_e)}{\mu S + B} \right]$$

The power of the search is enhanced by **combining the  $H \rightarrow ZZ$  channels**, done by merging the  $i$  individual likelihood functions

$$L(\mu) = \prod_i L_i(\mu)$$

- Assume channels are statistically independent
- $\mu$  is common between channels
- Define which systematic uncertainties are correlated between channels

The combined likelihood is used to fit the data from all channels

- Increased statistics enhances sensitivity

# Results

Our final results are under circulation in ATLAS, therefore, material for this section is limited

- Examples on the next slide borrowed from an earlier analysis

The combined likelihood is used to derive two types of results

***Model-independent*** limits on the heavy Higgs boson production

- Reduce assumptions about underlying model to make search general

***Model-specific*** interpretation

- Limits on the 2HDM parameter-space in certain benchmark scenarios
- Detailed information about this still confined to ATLAS

# Results

arXiv:1207.7214

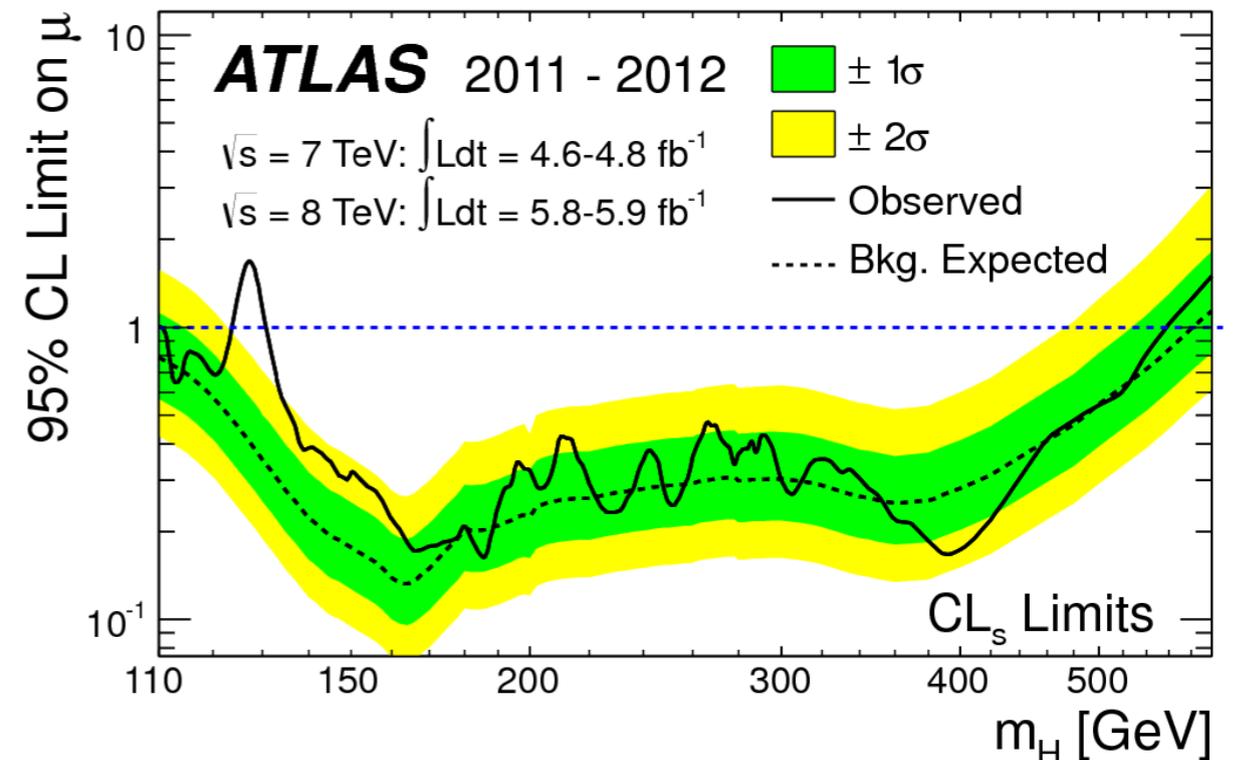
Our *model-independent* results are shown as upper limit on the heavy Higgs boson cross-section  $\times$  the  $H \rightarrow ZZ$  branching ratio

- Limit converted from  $\mu$  to fb
- Computed with the CLs technique separately for each mass point

The *expected limit* represents the maximum possible signal rate (at 95% CL) assuming the data is modelled by the background only hypothesis ( $\mu = 0$ )

- The more data, the smaller rates can be probed
- Background fluctuation represented by the yellow/green bands

New physics would show as *data excess above the expected limits*



# Summary and outlook

***Presented an overview of my main thesis topic: the search for a heavy Higgs boson in the  $H \rightarrow ZZ$  channels using  $\sqrt{s} = 8$  TeV data recorded with ATLAS***

Our analysis is complete at this point and going through the final internal ATLAS review

The LHC Run-2 data-taking period just began (in May this year)

- Collisions at  $\sqrt{s} = 13$  TeV
- Expect  $\sim 5 \text{ fb}^{-1}$  towards end of year

***A corresponding search is foreseen with the early Run-2 data***

- Sensitivity of Run-1 search expected to be exceeded in Run-2 with as little as  $< 10 \text{ fb}^{-1}$

At the moment I work on finishing the final details with the Run-1 paper and in parallel prepare for the Run-2 search

**Thanks for your attention**



# Backup

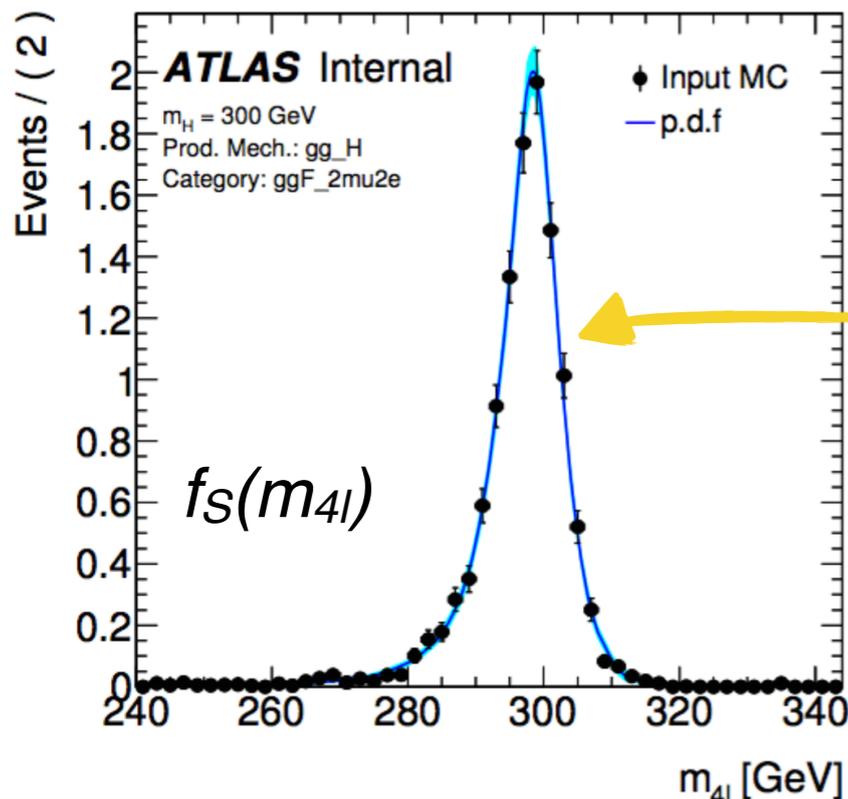


# Statistical treatment and combination

$\mu$  is determined by the use of a **likelihood function**\*

- Parametrisation of our assumptions about signal and background
- Represents the probability for our dataset to originate from our model with a signal scaled by  $\mu$

$$\mathcal{L}(\mu) = \text{Pois}(x_1 \dots x_n | \mu S + B) \left[ \prod_{e=1}^n \frac{\mu S f_S(x_e) + B f_B(x_e)}{\mu S + B} \right]$$



The likelihood consists of

- $\mathbf{x}_e$ : dataset of  $n$  events in units of the discriminant
- $\mathbf{S/B}$ : expected signal/background events
- $\mathbf{f_{S/B}(x_e)}$ : function describing signal/background

**By minimising the likelihood, the value of  $\mu$  favoured by data is found**

# Statistical treatment and combination

A crucial (and time-consuming) aspect of the analyses is the determination of **systematic uncertainties**

- For example, the uncertainty on the jet energy scale is 1-3% (depending on the jet momentum)

Three types of systematics are defined, affecting

- Experimental factors
- Signal modelling
- Background estimation

The systematic uncertainties affecting our search are integrated into the statistical procedure as **nuisance parameters (NP)**

Each systematic is represented with a NP, which models it as a distribution with mean and width representing its value and uncertainty

- Usually a normal distribution or Poissonian

A systematic with the measured value  $\tilde{\theta} \pm \sigma_{\theta}$  would be constrained by the normal distribution  $G(\theta|\tilde{\theta}, \sigma_{\theta})$

# Statistical treatment and combination

The NPs are added to the likelihood function as **constraint terms**

$$\mathcal{L}(\mu, \theta) = \text{Pois}(x_1 \dots x_n | \mu S + B) \left[ \prod_{e=1}^n \frac{\mu S f_s(x_e) + B f_B(x_e)}{\mu S + B} \right] \prod_i G(\theta_i | \tilde{\theta}_i, \sigma_{\theta_i})$$

The NP values favoured by data is found by minimising the likelihood w.r.t. each NP

- Profile likelihood approach
- The data constrains the NPs

The presence of nuisance parameters broadens the profile likelihood scan

- Reflects the loss of information due to the systematic uncertainties

