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Gyrokinetic Turbulence in Laboratory and Astrophysical Plasmas

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## **Research topics**

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Plasma astrophysics (incl. dynamo physics) Computational fluid dynamics

Computational plasma physics

Turbulence in magnetized plasmas

- Code development
- Basic physics issues
- Applications to experiments

Colloidal (dusty) plasmas

Non-Gaussian diffusion in complex systems

## Approach

Computer simulations (Virtual experiments)

Analytical theories

GENE (kinetic continuum code)

DPS (molecular dynamics code)

DYNAMO (nonlinear MHD code)

TURB3D (nonlinear fluid code)

XGC-0 (kinetic PIC code)

Kinetic theory

Nonequilibrium statistics

Nonlinear dynamics

etc.

**Closure theories** 

Experiments or observations

Fusion experiments (AUG, JET etc.)

Smaller lab experiments

Astrophysical data

etc.

etc.

# Plasma turbulence

## Plasma turbulence in astrophysics and in the laboratory









Key issues: Turbulent transport, Anomalous heating/dissipation

## Fusion energy in the stars



This points the way to a  $CO_2$  free energy source for the 21<sup>st</sup> century...

## Fusion energy in the laboratory



Deuterium-Tritium fusion has by far the highest probability under experimentally accessible conditions.

Still, temperatures of about 100 million degrees are required! Thus, we are dealing with a fully ionized gas (plasma).

Two main approaches: Inertial fusion and magnetic fusion.

## ITER and plasma turbulence

**ITER** is an extremely challenging scientific project, aimed at investigating burning **D-T** plasmas

It is currently being built in Cadarache

The success of ITER depends on our understanding of plasma turbulence



www.iter.org

## Plasma turbulence: GENE simulations



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#### www.ipp.mpg.de/~fsj/gene

## Key idea: Applications of insights, theories, and tools to astrophysics

Some issues under investigation:

- Role of microturbulence in astrophysics?
- How does MHD turbulence dissipate energy?
- How do fast particles interact with turbulence?

## Gyrokinetic turbulence

## What is gyrokinetic theory?

Dilute and/or hot plasmas are almost collisionless.

Thus, if kinetic effects (finite Larmor radius, Landau damping, magnetic trapping etc.) play a role, MHD is not applicable, and one has to use a (reduced) kinetic description!

 $\left[\frac{\partial}{\partial t}\right]$ 

Vlasov-Maxwell equations

$$+\mathbf{v}\cdot\frac{\partial}{\partial\mathbf{x}} + \frac{q}{m}\left(\mathbf{E} + \frac{\mathbf{v}}{c}\times\mathbf{B}\right)\cdot\frac{\partial}{\partial\mathbf{v}}\left[f(\mathbf{x},\mathbf{v},t) = \mathbf{0}\right]$$

Removing the fast gyromotion  $\omega$  (gyrokinetic ordering of scales / amplitudes)

[Frieman, Chen, Lee, Hahm, Brizard et al., 1980s]



Charged rings as quasiparticles; gyrocenter coordinates; keep FLR effects



Brizard & Hahm, Rev. Mod. Phys. **79**, 421 (2007)

## The nonlinear gyrokinetic equations

 $f = f(\mathbf{X}, v_{\parallel}, \mu; t)$ 

Advection/Conservation equation

$$\frac{\partial f}{\partial t} + \dot{\mathbf{X}} \cdot \frac{\partial f}{\partial \mathbf{X}} + \dot{v}_{\parallel} \frac{\partial f}{\partial v_{\parallel}} = 0$$

$$\dot{\mathbf{X}} = v_{\parallel} \mathbf{b} + \frac{B}{B_{\parallel}^*} \left( \frac{v_{\parallel}}{B} \bar{\mathbf{B}}_{1\perp} + \mathbf{v}_{\perp} \right)$$

X = gyrocenter position  $\lor_{II} =$  parallel velocity  $\mu =$  magnetic moment

Appropriate field equations

$$\frac{n_1}{n_0} = \frac{\bar{n}_1}{n_0} - \left(1 - \|I_0^2\|\right) \frac{e\phi_1}{T} + \|xI_0I_1\| \frac{B_{1\|}}{B}$$

$$\mathbf{v}_{\perp} \equiv \frac{c}{B^2} \bar{\mathbf{E}}_1 \times \mathbf{B} + \frac{\mu}{m\Omega} \mathbf{b} \times \nabla (B + \bar{B}_{1\parallel}) + \frac{v_{\parallel}}{\Omega} (\nabla \times \mathbf{b})_{\perp}$$

..2

$$\nabla_{\perp}^2 A_{1\parallel} = -\frac{4\pi}{c} \sum \bar{J_{1\parallel}}$$

$$\dot{v}_{\parallel} = \frac{\dot{\mathbf{X}}}{mv_{\parallel}} \cdot \left( e\bar{\mathbf{E}}_{1} - \mu \nabla (B + \bar{B}_{1\parallel}) \right) \qquad \qquad \frac{B_{1\parallel}}{B} = -\sum \epsilon_{\beta} \left( \frac{\bar{p}_{1\parallel}}{n_{0}T} + \|xI_{1}I_{0}\| \frac{e\phi_{1}}{T} + \|x^{2}I_{1}^{2}\| \frac{B_{1\parallel}}{B} \right)$$

Nonlinear integro-differential equations in **5 dimensions**...

## Gradient-driven microinstabilities



## Anisotropy of turbulent fluctuations

Strong background magnetic field causes turbulence to be quasi-2D



Possible simulation volume: flux tube, annulus, full (or fractional) torus

## Turbulent mixing in a tokamak

ExB drift velocity

$$\tilde{\mathbf{v}}_E = \frac{c}{B^2} \mathbf{B} \times \nabla \tilde{\phi}$$

$$\mathbf{Q} \equiv \frac{3}{2} \langle \tilde{p} \, \tilde{\mathbf{v}}_E \rangle = -n \chi \nabla T$$



potential contours streamlines of ExB velocity **Gradient-driven instabilities**  $\rightarrow$  fluctuations  $\rightarrow$  transport

Typical heat and particle diffusivities are of the order of 1 m<sup>2</sup>/s.

### Major theoretical speedups

relative to original Vlasov/pre-Maxwell system on a naïve grid, for ITER  $1/\rho_* = a/\rho \sim 1000$ 

- Nonlinear gyrokinetic equations
  - □ eliminate plasma frequency:  $\omega_{pe}/\Omega_i \sim m_i/m_e$  x10<sup>3</sup>
  - □ eliminate Debye length scale:  $(\rho_i / \lambda_{De})^3 \sim (m_i / m_e)^{3/2}$  x10<sup>5</sup>
  - □ average over fast ion gyration:  $\Omega_i / \omega \sim 1 / \rho_*$  x10<sup>3</sup>

#### Field-aligned coordinates

□ adapt to elongated structure of turbulent eddies:  $\Delta_{\mu}/\Delta_{\perp} \sim 1/\rho_{*}$  x10<sup>3</sup>

#### Reduced simulation volume

- $\Box$  reduce toroidal mode numbers (i.e., 1/15 of toroidal direction) x15
- $\Box$  L<sub>r</sub> ~ a/6 ~ 160  $\rho$  ~ 10 correlation lengths x6

#### Total speedup

For comparison: Massively parallel computers (1984-2009) x10<sup>7</sup>

**x10**<sup>16</sup>

## The simulation code GENE

- GENE is a physically comprehensive and well benchmarked code
- GENE is publicly available (<u>www.ipp.mpg.de/~fsj/gene</u>)
- various applications in fusion research (tokamaks *and* stellarators) and astrophysics; can be run as a (radially) local or global code

• the differential operators are discretized via a combination of spectral, finite difference, and finite volume methods; the time stepping is done via a (non-standard) 4<sup>th</sup>-order explicit Runge-Kutta method

- in addition, GENE can also be run as a linear eigenvalue solver
- two main goals: deeper understanding of fundamental physics issues and direct comparisons with experiments (interfaces to MHD codes)

## GENE is a massively parallel code





- Parallelization due to high-dimensional domain decomposition (either pure MPI or mixed MPI/OpenMP paradigm)
- GENE runs very efficiently on a large number of parallel platforms (including IBM BlueGene, IBM Power6, Cray XT4, SGI Altix etc.)
- GENE is part of the European DEISA benchmark suite and the EU-Japanese IFERC benchmark suite

## GENE on BlueGene/L (strong scaling)



measurements in co-processor mode at IBM Watson Research Center)

Role of microturbulence in astrophysics?

## Broad-line regions in AGNs

Measured electromagnetic spectra from AGNs suggest the existence of...

...cold, dense clouds in a hot, dilute, magnetized medium in the central region of AGNs.

How can those cold clouds survive?

Standard model (e.g. Kuncic et al., MNRAS 1996): Cold clouds are magnetically confined and form filaments; perpendicular transport is negligible.

What if microturbulence is taken into account?



## Microturbulence effects?

 $\nabla T \neq 0, \ \nabla n \neq 0$ , magnetic confinement  $\Rightarrow$  evaporating clouds prone to **plasma microinstabilities** 

Taking **typical parameters** (e.g., Kuncic 1996):  $n_{\rm cloud} \sim 10^{17} \, {\rm cm}^{-3}$ ,  $T_{\rm medium} \sim 10^9 \, {\rm K}$  and assuming gyrokinetic transport (ITG) levels:

$$t_{
m evap}^{
m GK} \sim 10^{9-10}~
m s$$

significantly larger than unhindered evaporation
 however: too large for the standard model?

More detailed investigations are currently underway.

## How does MHD turbulence dissipate energy?

## NS turbulence: Richardson cascade

Turbulence as a local cascade in wave number space...



"Big whorls have little whorls, little whorls have smaller whorls that feed on their velocity, and so on to viscosity"

The Richardson cascade is a nonlinear route to dissipation.

## Solar wind turbulence: Dissipation?



High-k MHD turbulence satisfies the gyrokinetic ordering!

## High-k Alfvén wave turbulence: Observations versus simulations



## How do fast particles interact with turbulence?

## **Diffusivities and correlations**

Test particle approach

- Diffusion equation:  $\partial_t n(\mathbf{x},t) = D(t) \nabla^2 n(\mathbf{x},t)$
- Einstein (1905):  $D_x(t) = \frac{1}{2} \frac{d}{dt} \langle \delta x^2(t) \rangle$
- Taylor (1920):  $D_x(t) = \frac{1}{2} \frac{d}{dt} \langle \delta x^2(t) \rangle = \int_0^t d\xi \langle v_x(0) v_x(\xi) \rangle \equiv \int_0^t d\xi L_{v_x}(\xi)$



## Validity of gyro/orbit averaging



Magnetic transport along fluctuating field lines:  $v_B \equiv v_{\parallel}(\tilde{B}_r/B_0)$ 

## Scaling laws: Diffusivities vs energy



Magnetic transport of beam ions is independent of energy!

### Cosmic rays in astrophysical turbulence

Example: Solar wind

(other candidates: interstellar medium, intergalactic medium, jets etc.)



Cross-field transport:

- Gyro averaging is not valid for highly energetic particles
- Particle diffusivity may clearly exceed naïve estimates

## Conclusions

## Some key points of this talk

Fusion-oriented plasma physics has developed sophisticated theories / tools, some of which can be applied to astrophysics

Questions under investigation (partial list):

- Role of microturbulence in astrophysics?
- How does MHD turbulence dissipate energy?
- How do fast particles interact with turbulence?
- Kinetic effects in the magnetorotational instability?
- Heating of the solar corona?

Collaborative efforts are underway (e.g. MPI for Solar Systems Research & MPI for Plasma Physics)