



The role of tidal dissipation on the evolution of planetary systems

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Tidal interactions in exoplanetary systems

- No exomoons discovered yet!
- Tidal torque $\propto a^{-6}$
- Close-in planets ($a < 0.15$ AU), around MS late-type stars

- Close-in giant planets cannot form *in situ*

- How did they migrate?

- planet-disc interactions
- planet-planet interactions
- planet-planetesimal disc interactions
- planet-distant star companion (Kozai-Lidov)

- The end of migration is the beginning of tidal interactions

Tidal evolution outcome

Barker & Ogilvie 2009

Tidal circularization time (for co-planar orbit)

$$\tau_e \approx 16.8 \text{ Myr} \left(\frac{Q'_\star}{10^6} \right) \left(\frac{m_\star}{M_\odot} \right)^{\frac{10}{3}} \left(\frac{M_J}{m_p} \right) \left(\frac{R_\odot}{R_\star} \right)^5 \left(\frac{P_{\text{orb}}}{1 \text{ d}} \right)^{\frac{13}{3}} \\ \times \left[\left(f_1(e) - \frac{11}{18} \frac{P_{\text{orb}}}{P_\star} f_2(e) \right) + \frac{Q'_p}{Q'_\star} \left(\frac{m_\star}{m_p} \right)^2 \left(\frac{R_p}{R_\star} \right)^5 \left(f_1(e) - \frac{11}{18} f_2(e) \right) \right]^{-1}$$

$$\tau_e \approx 4 \text{ Myr, for } e = 0.4 \text{ and } P_{\text{orb}} = 3 \text{ d}$$

Tidal alignment time (for circular orbit and small inclination)

$$\tau_i \approx 70 \text{ Myr} \left(\frac{Q'_\star}{10^6} \right) \left(\frac{m_\star}{M_\odot} \right) \left(\frac{M_J}{m_p} \right)^2 \left(\frac{R_\odot}{R_\star} \right)^3 \left(\frac{P_{\text{orb}}}{1 \text{ d}} \right)^4 \frac{12.5 \text{ d}}{P_\star} \left[1 - \frac{P_{\text{orb}}}{2P_\star} \left(1 - \frac{I\Omega}{h\mu} \right) \right]^{-1}$$

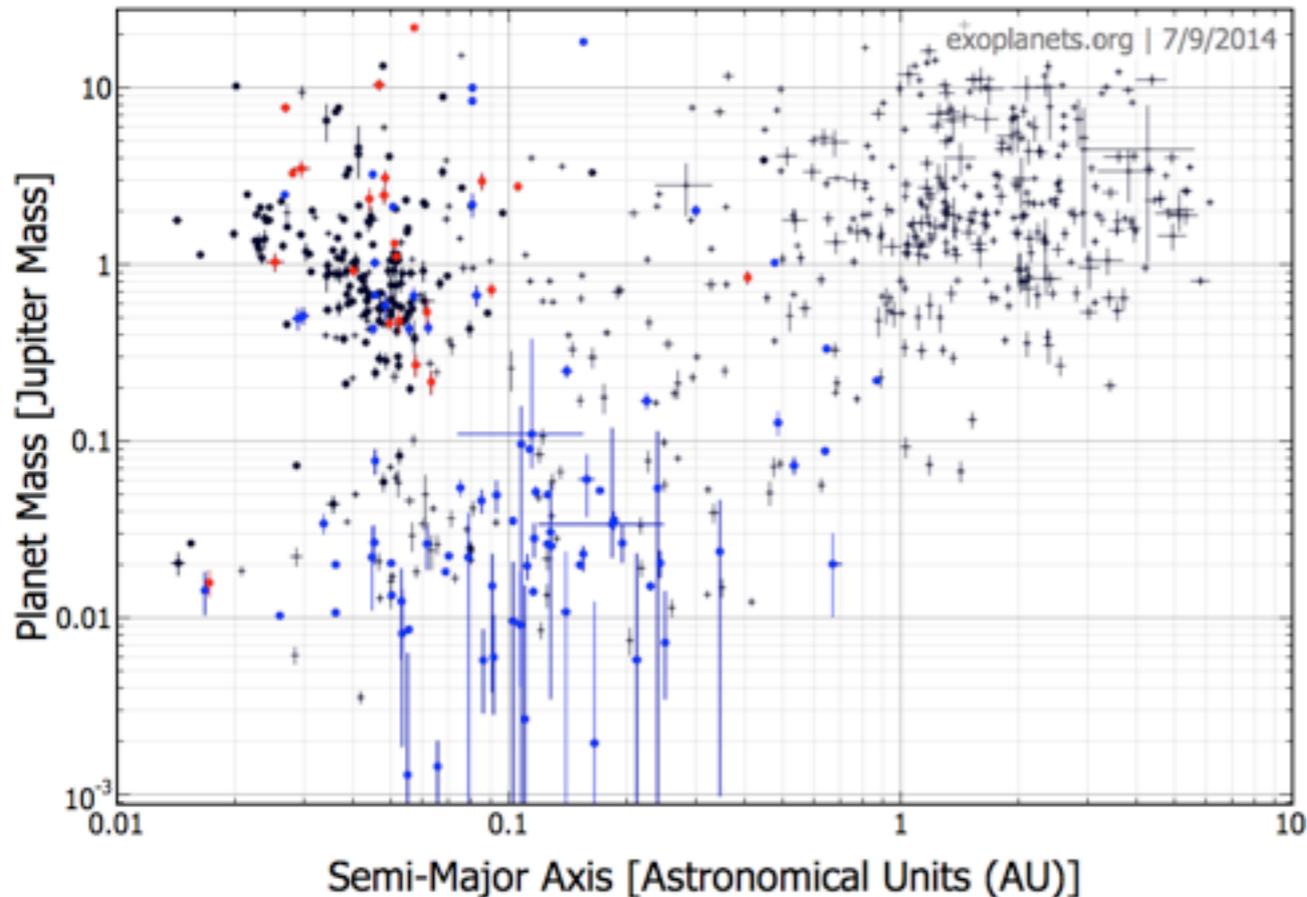
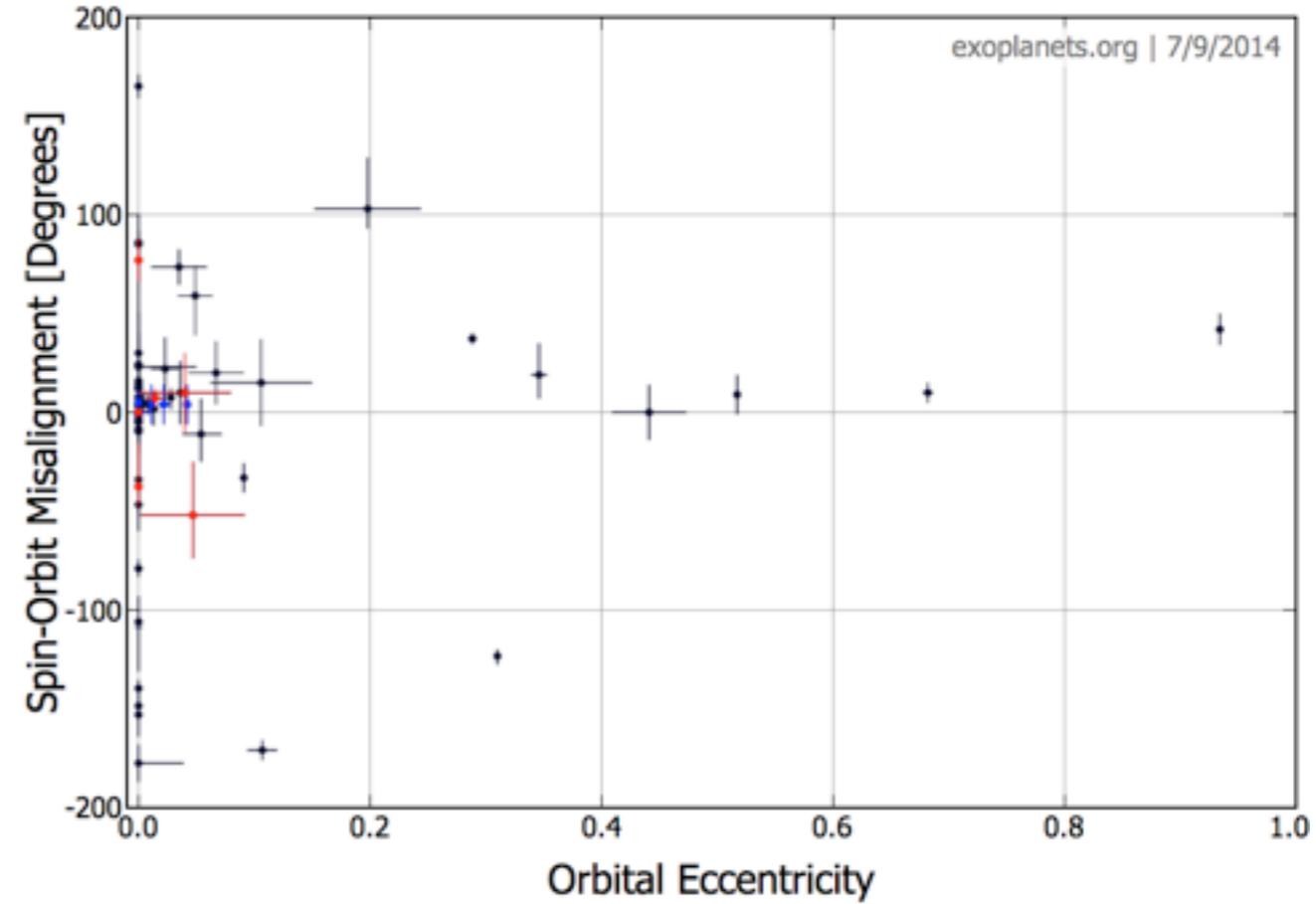
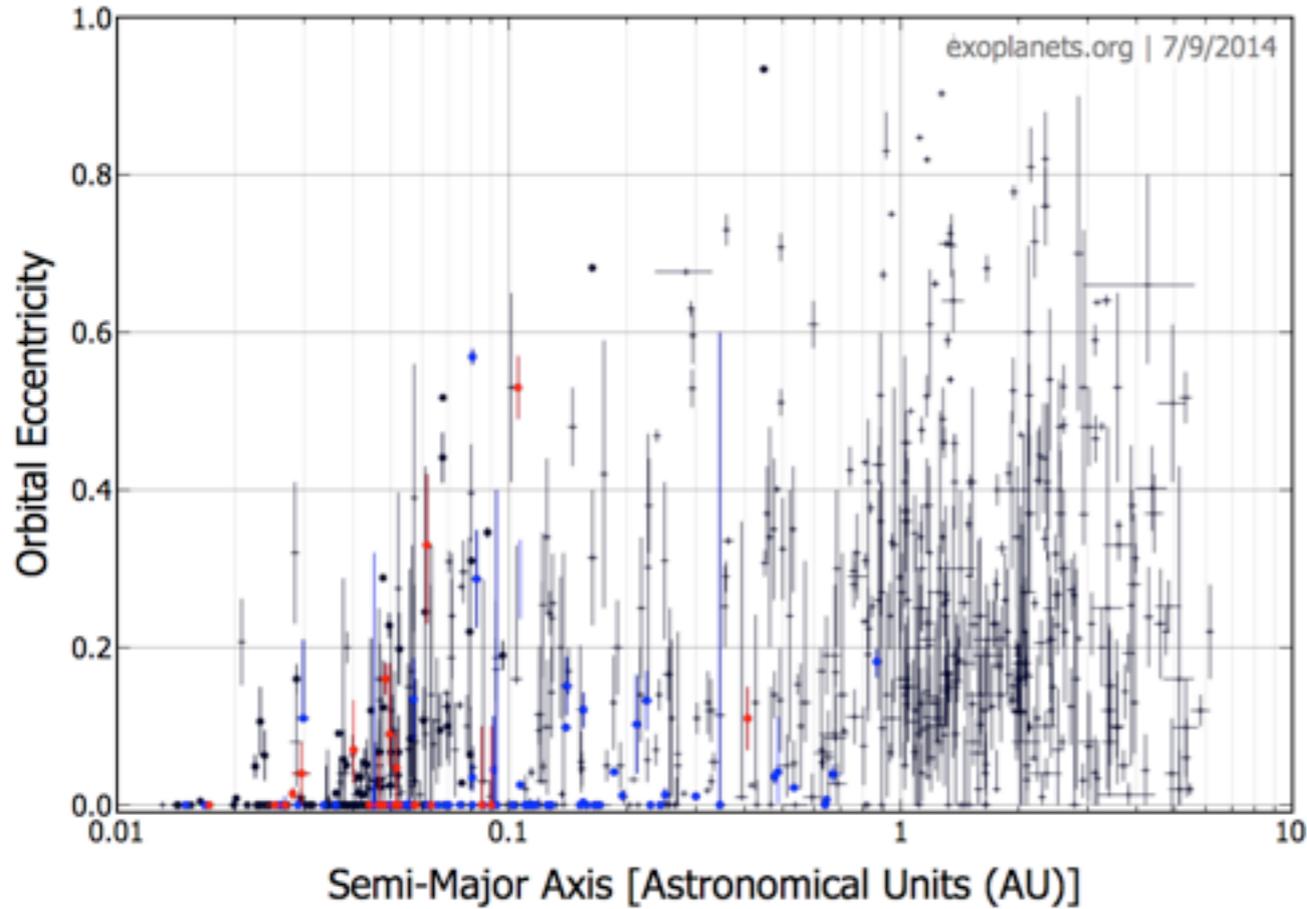
$$\tau_i \approx 6 \text{ Gyr, for } P_{\text{orb}} = 3 \text{ d and } P_\star = 12.5 \text{ d}$$

Tidal inspiral time (neglecting tides in the planet and for circular and co-planar orbit)

$$\tau_a \approx 12.0 \text{ Myr} \left(\frac{Q'_\star}{10^6} \right) \left(\frac{m_\star}{M_\odot} \right) \left(\frac{M_J}{m_p} \right) \left(\frac{P_{\text{orb}}}{1 \text{ d}} \right)^{\frac{13}{3}} \left(1 - \frac{P_{\text{orb}}}{P_\star} \right)^{-1}$$

$$\tau_a \approx 2 \text{ Gyr, for } P_{\text{orb}} = 3 \text{ d and } P_\star = 12.5 \text{ d}$$

Observations



- How efficient is tidal dissipation?
- observational constraints: Jackson et al 2008, Matsumura et al 2008, Deleuil et al 2012, Carone & Patzold 2007, Lanza et al 2011...

$$10^4 \leq Q'_p \leq 10^7, 10^5 \leq Q'_s \leq 10^9$$

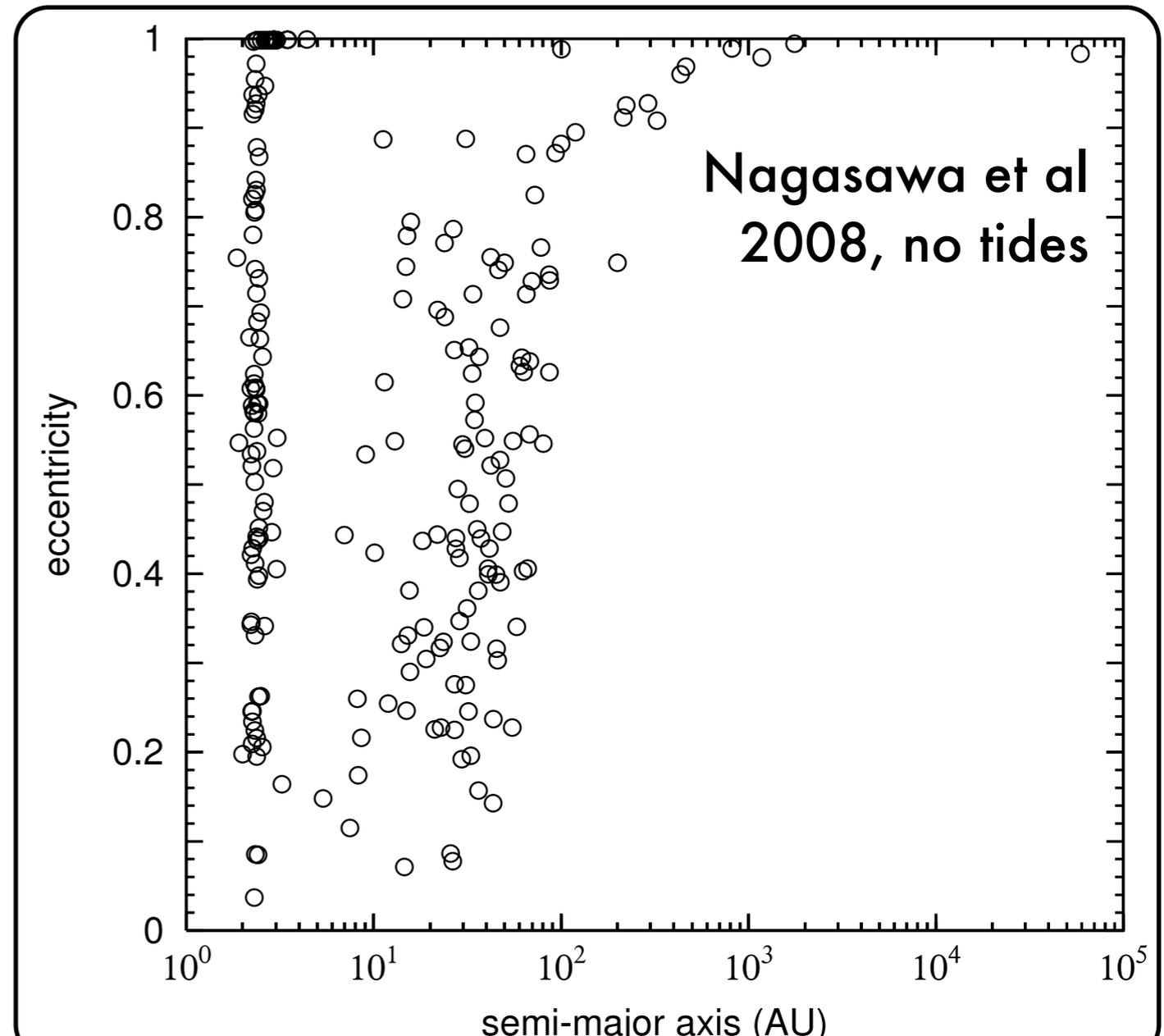
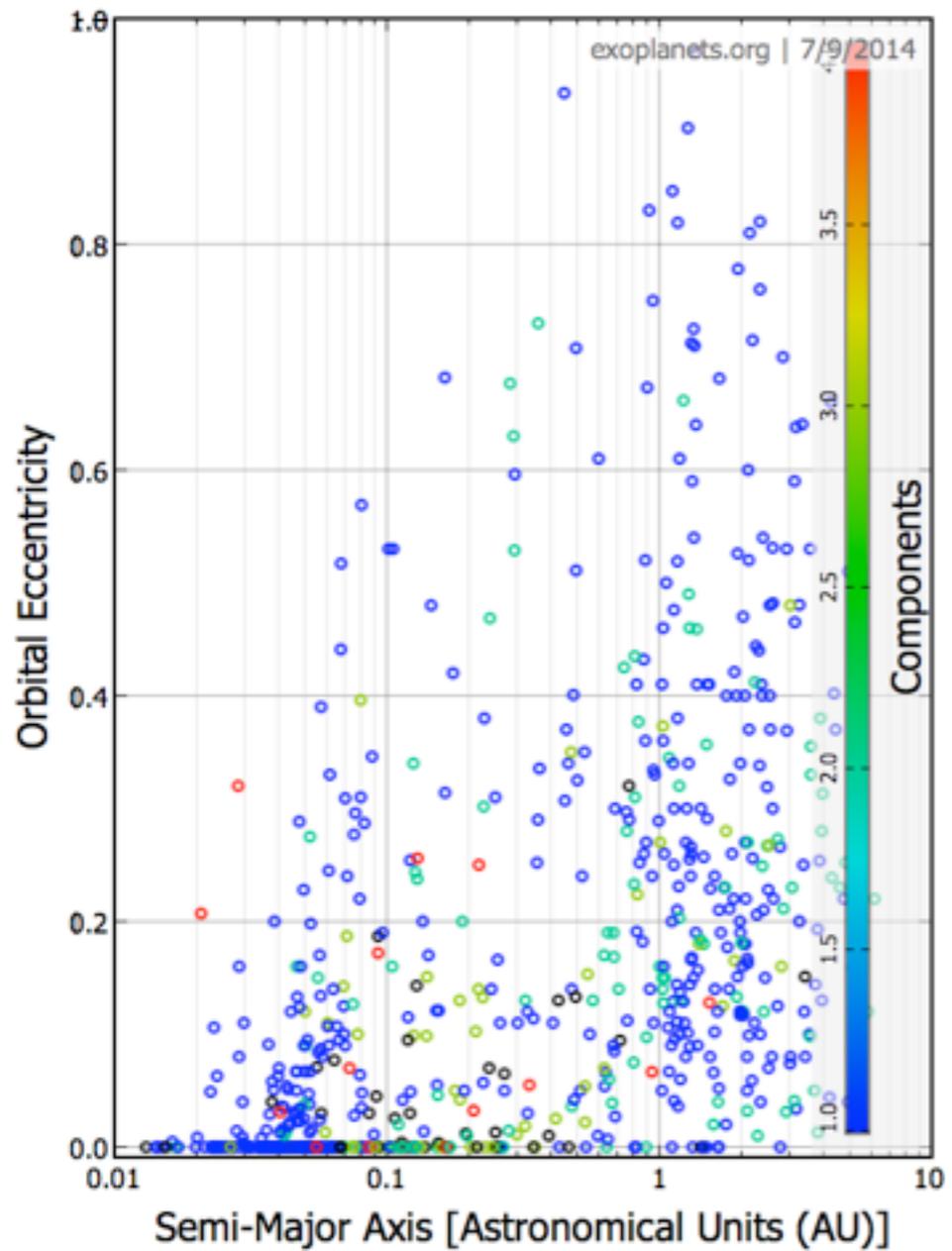
➔ See next talks

Excentricity

Origin?

Planet-planet scattering

(Rasio & Ford 1996, Weidenschilling & Marzari 1996)

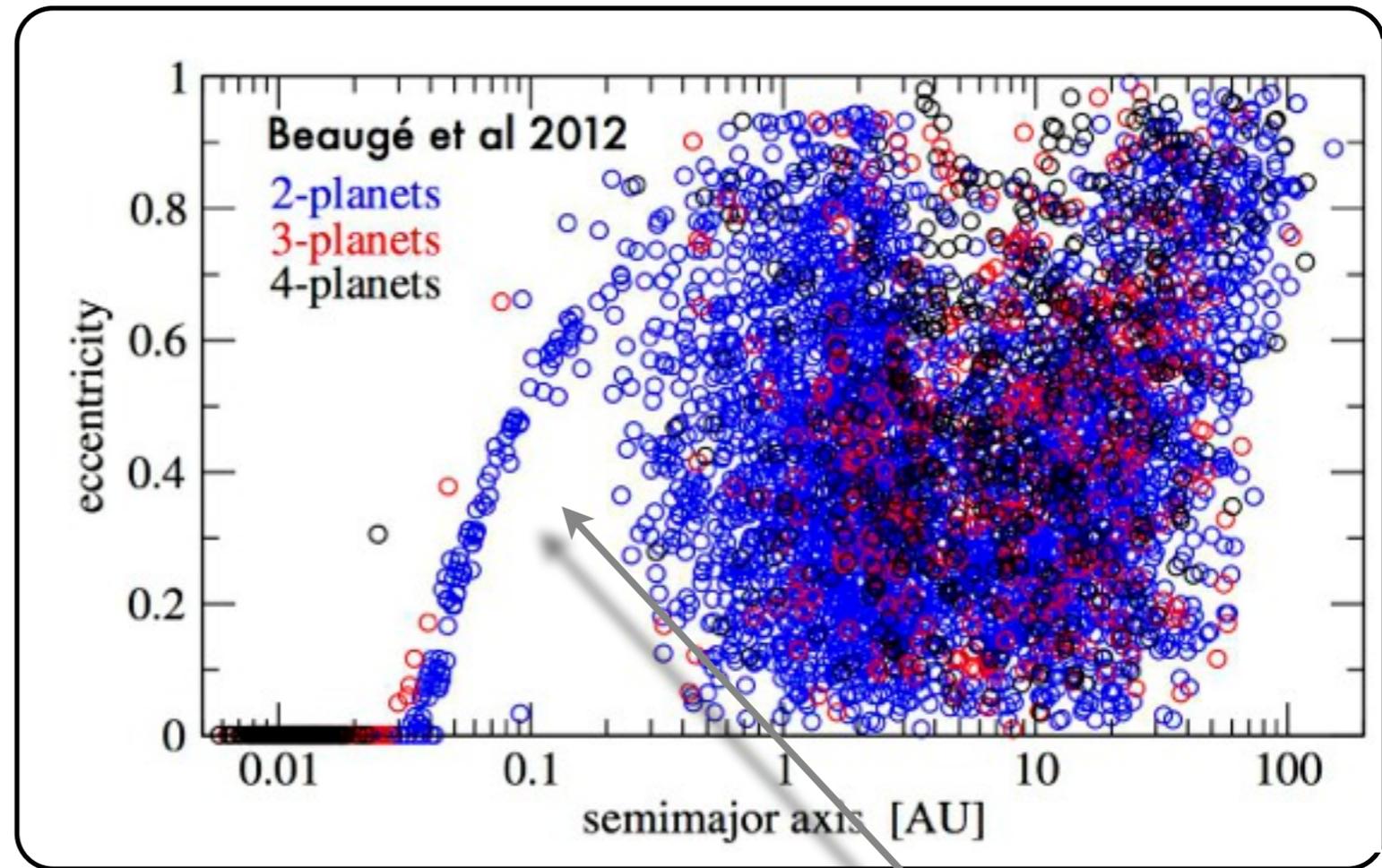
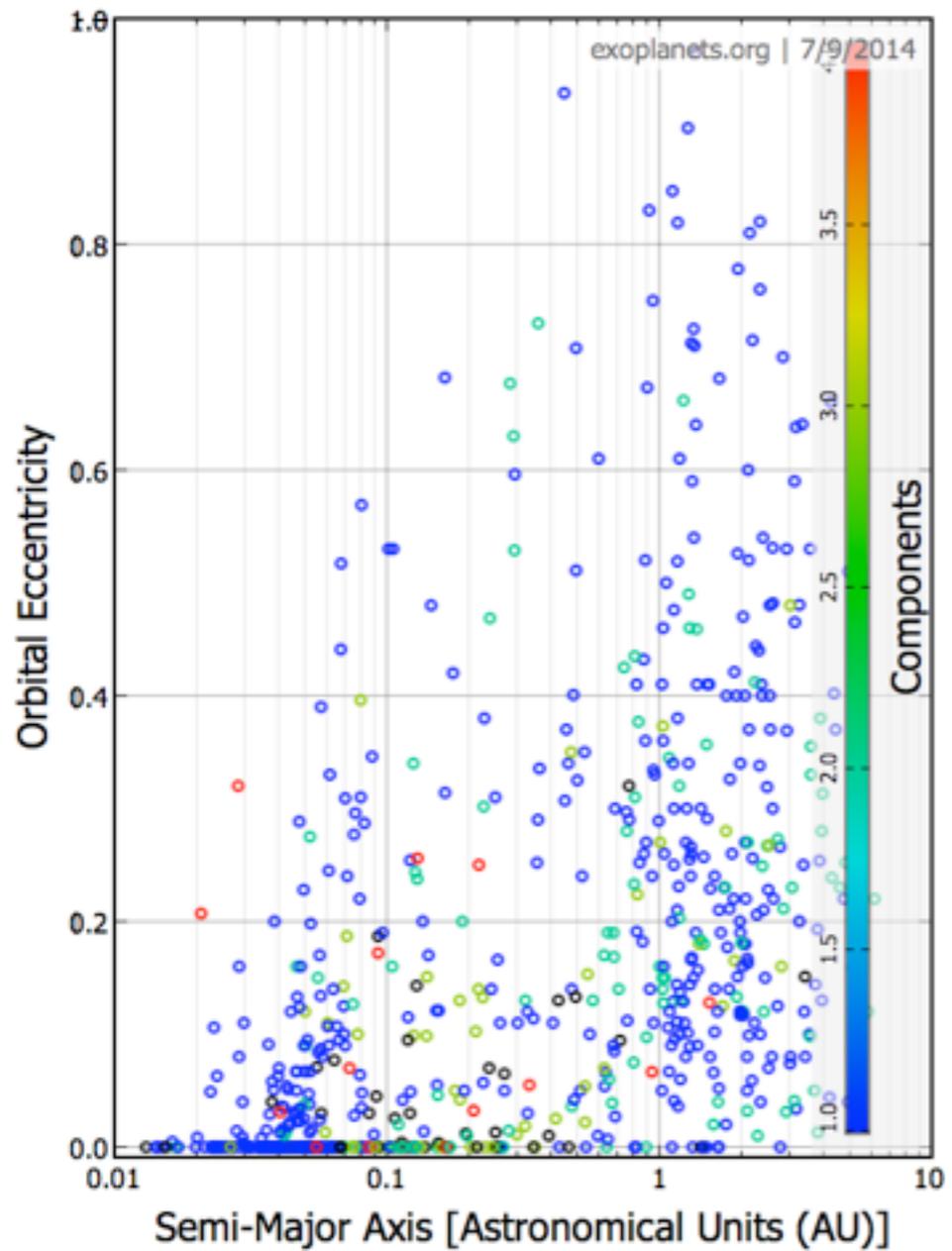


Excentricity

Origin?

Planet-planet scattering &
tidal circularization

(Nagasawa et al 2008)



$$Q'_p \sim 10^6 - 10^7$$

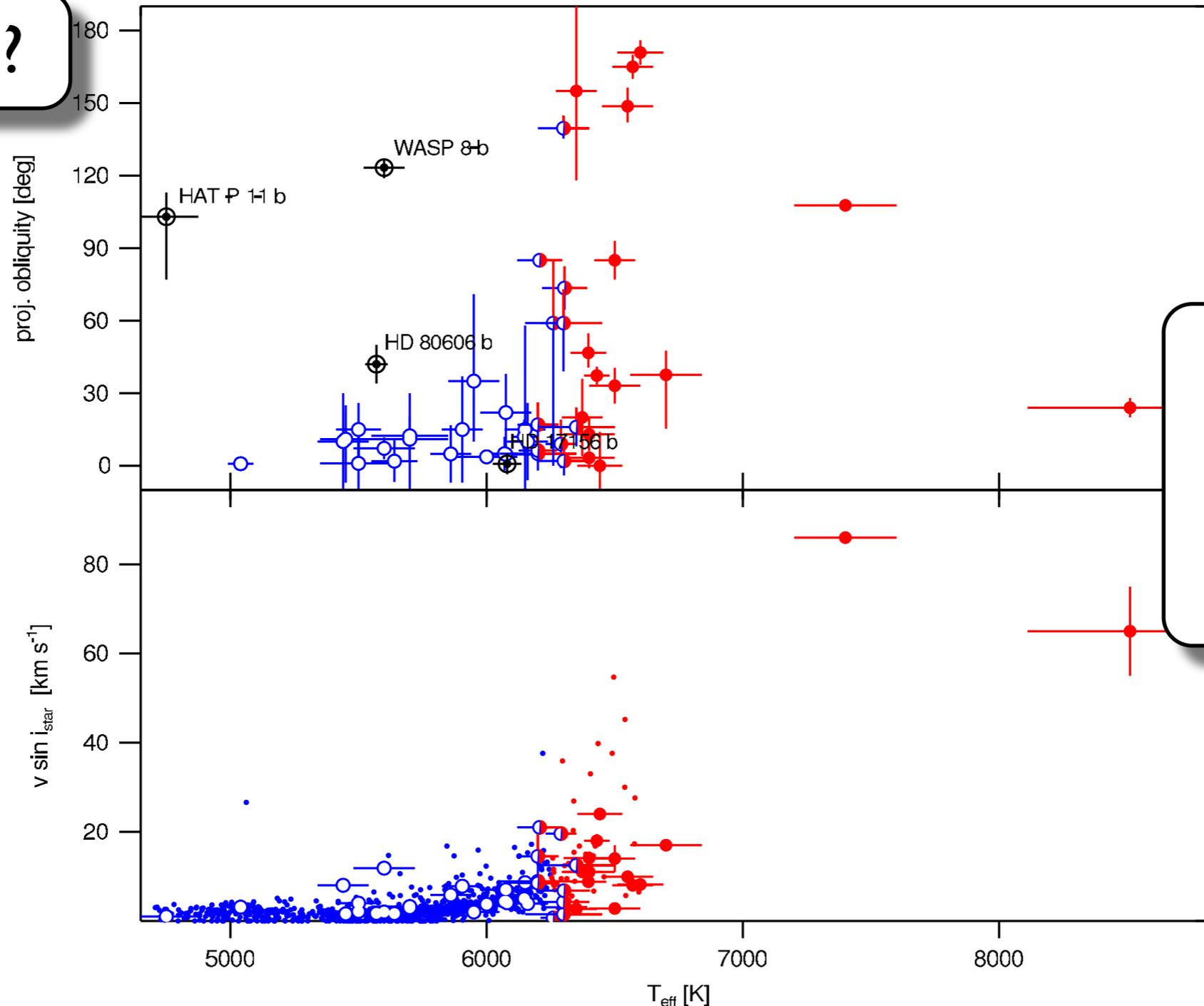
Spread?

Obliquity

THE ASTROPHYSICAL JOURNAL, 757:18 (25pp), 2012 September 20

ALBRECHT ET AL.

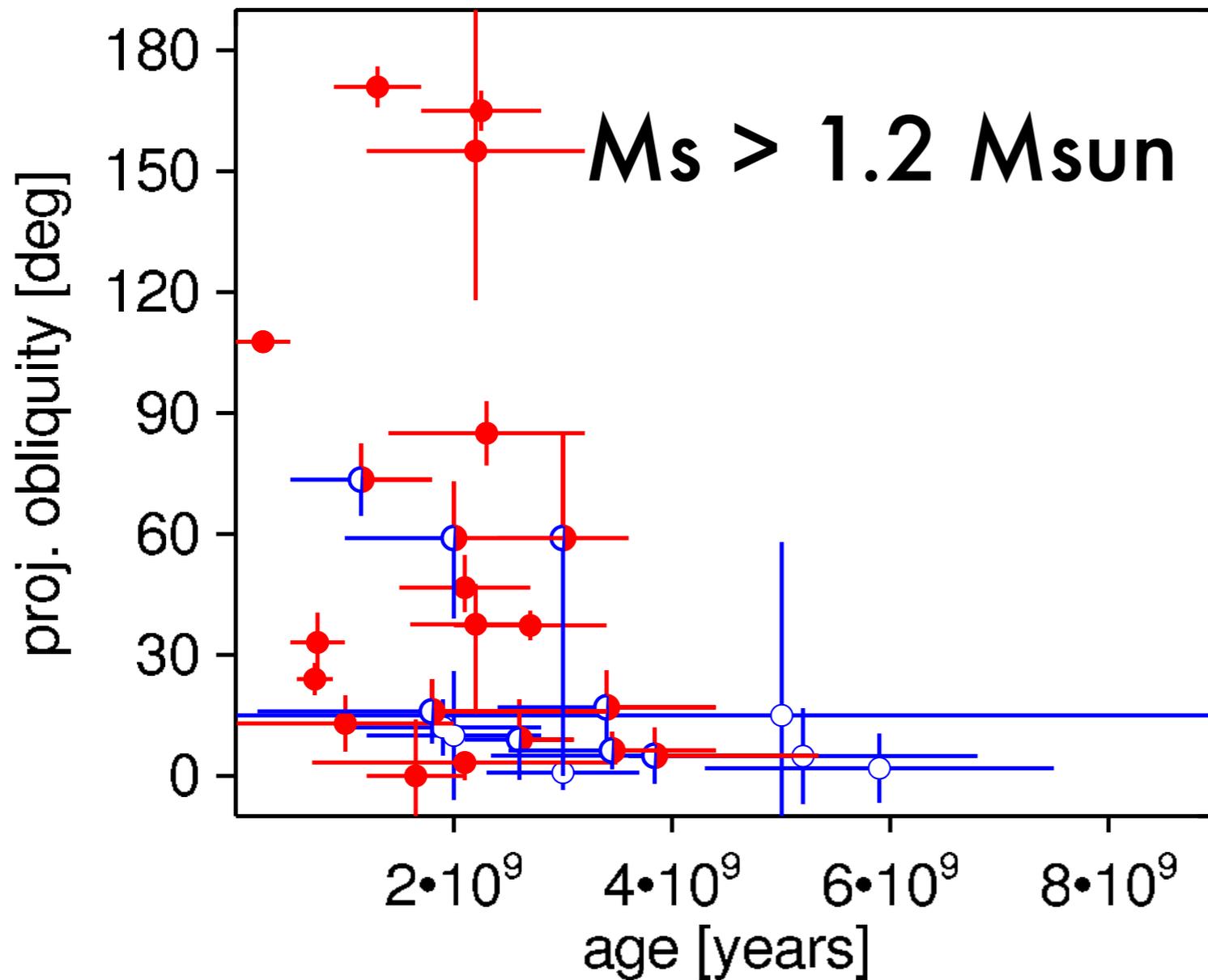
Primordial?



Hot star
=> high
obliquity?

Figure 20. Projected obliquities and projected stellar rotation speeds as a function of the stellar effective temperature. Upper panel: measurements of projected obliquities as a function of the effective temperature of the host star. Stars which have temperatures higher than 6250 K are shown with red filled symbols. Blue open symbols show stars with temperatures lower than 6250 K. Stars which measured effective temperature include 6250 K in their 1σ interval are shown by split symbols. Systems which harbor planets with mass $< 0.2 M_{\text{Jup}}$ or have an orbital period more than 7 days are marked by a black filled circle with a ring. Lower panel: projected stellar rotation speeds $v \sin i_*$ of the stars in our sample. In addition, $v \sin i_*$ measurements of stars in the catalog by Valenti & Fischer (2005) are shown as small dots.

Obliquity



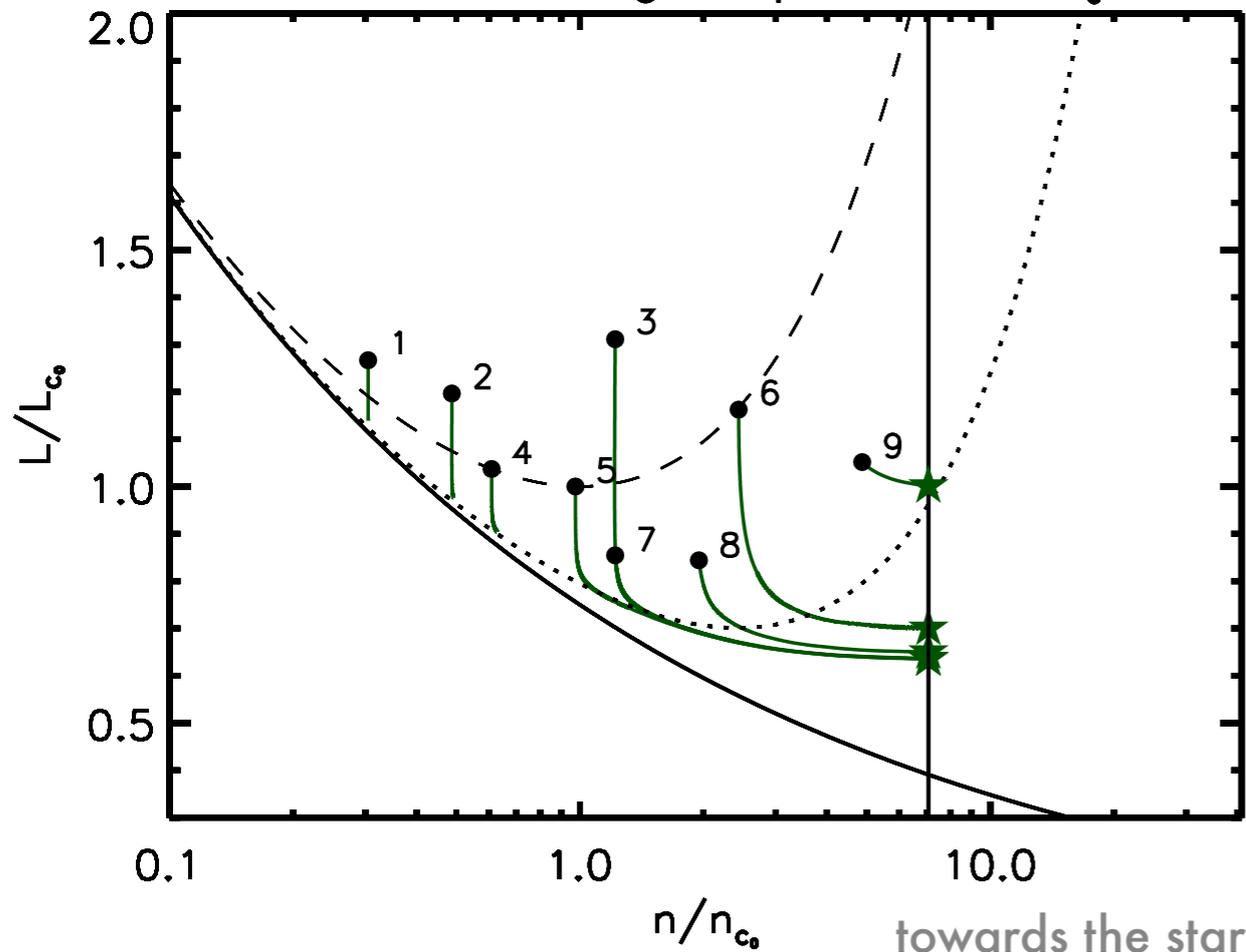
Albrecht et al.2012

- Development of the convective envelope with time
- Initially random
- Tidally damped

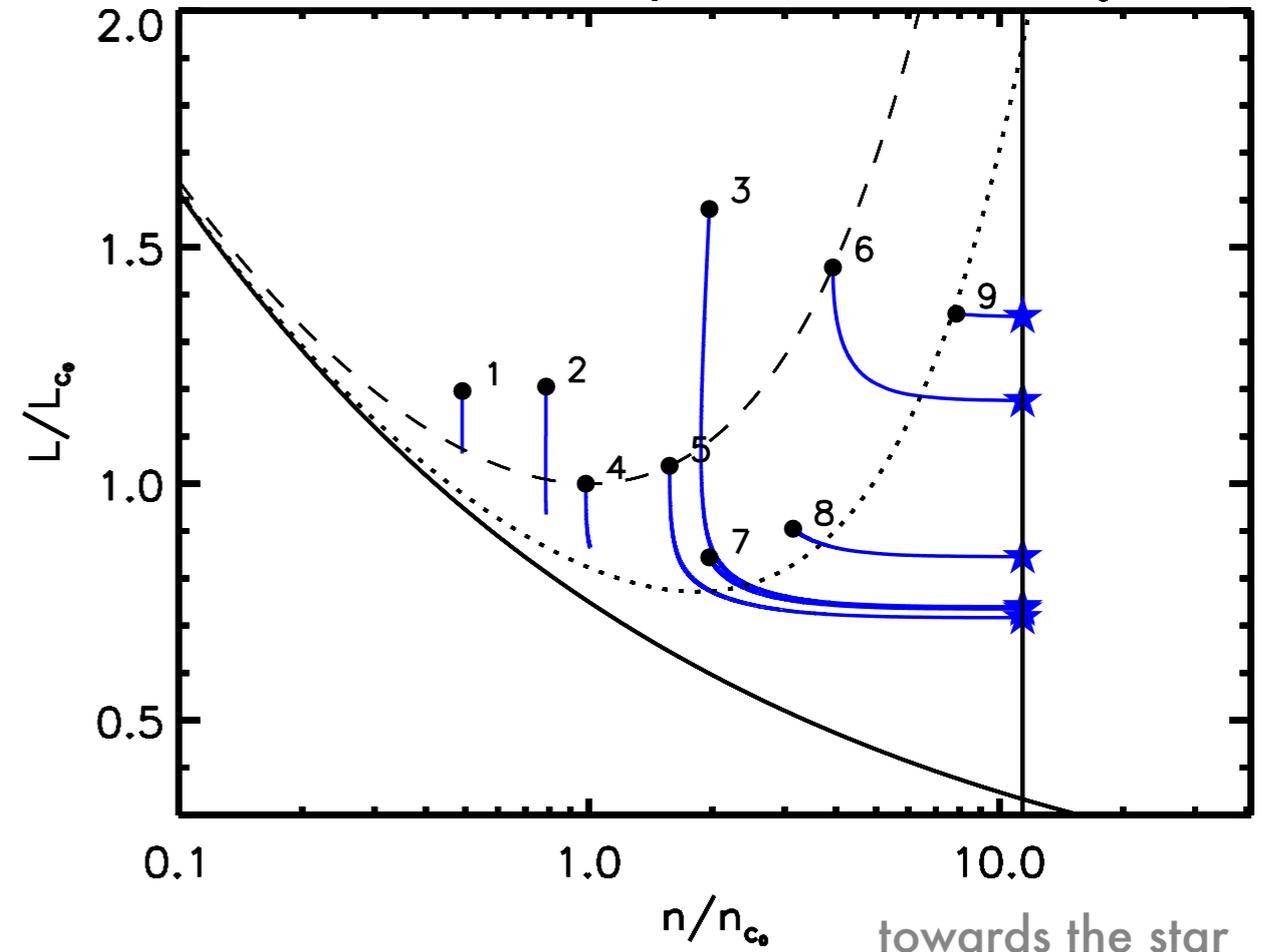
Role of angular momentum loss

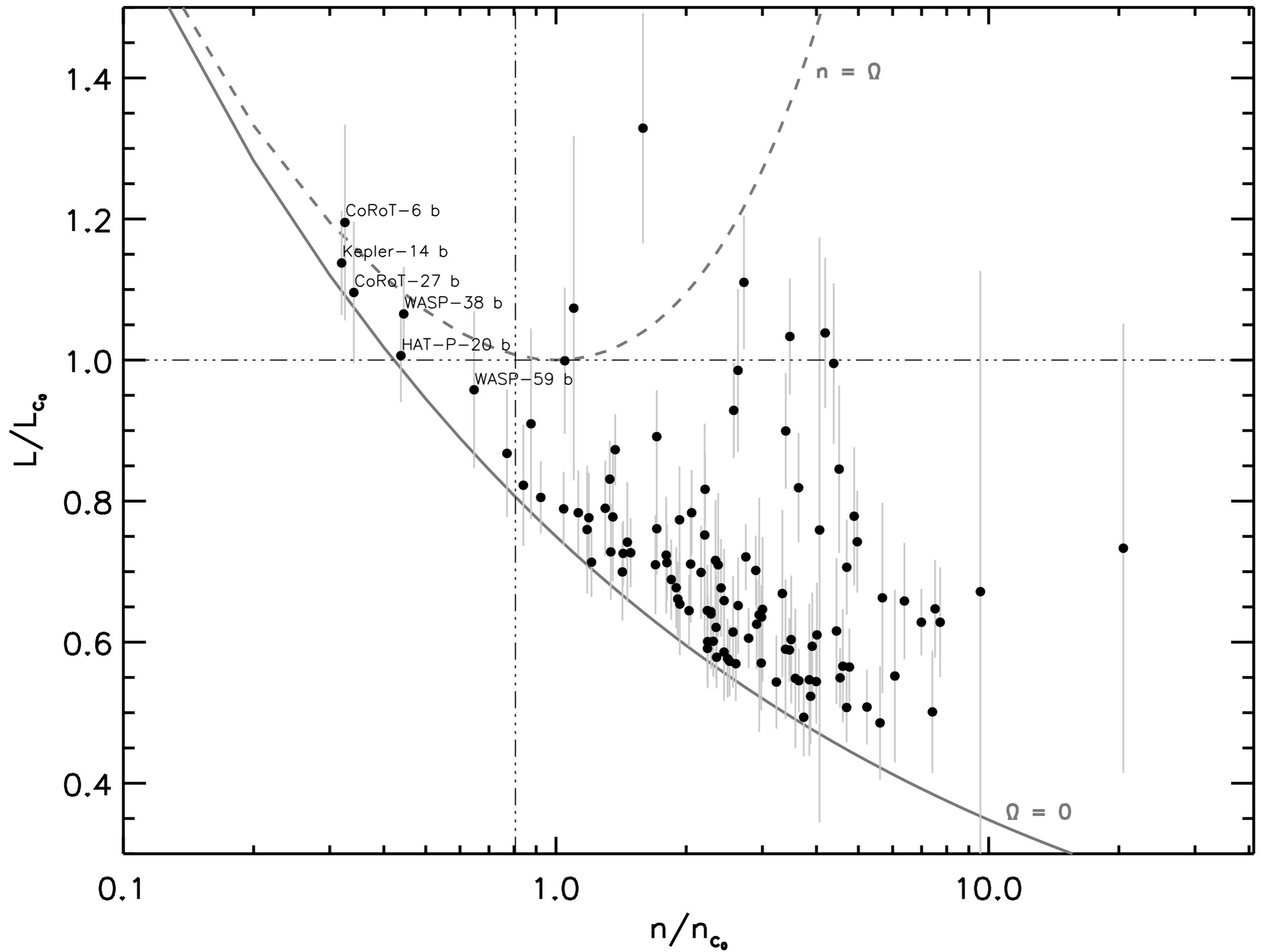
- G-type stars lose angular momentum from their magnetized wind, F-type stars too but less so
- The dynamical evolution of orbital elements is driven by the resultant of the wind torque and the tidal torque

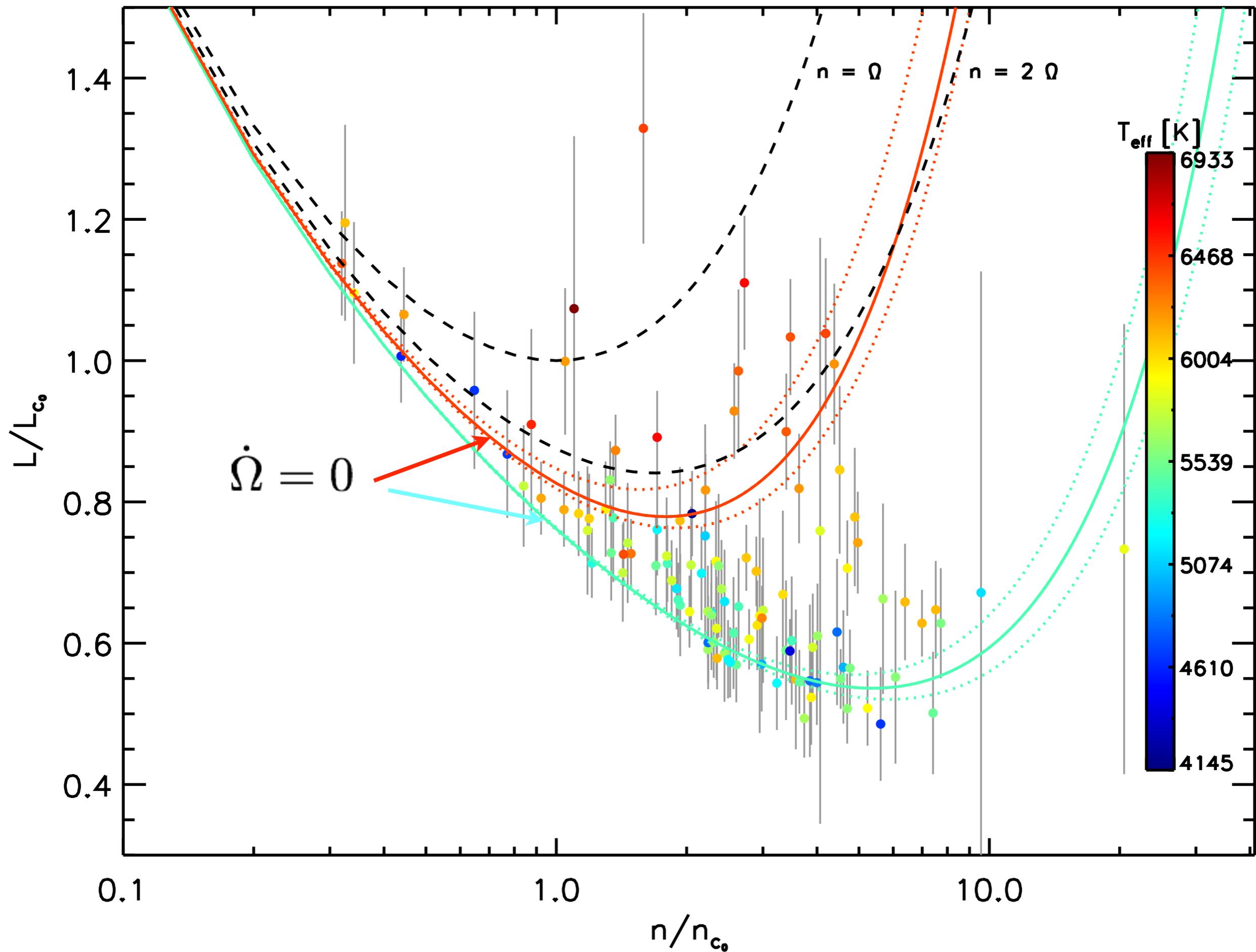
$M_s = 1.0 M_\odot$, $M_p = 1.0 M_J$



$M_s = 1.4 M_\odot$, $M_p = 1.0 M_J$







Role of angular momentum loss

- G-type stars lose angular momentum from their magnetized wind, F-type stars too but less so
- The dynamical evolution of orbital elements is driven by the resultant of the wind torque and the tidal torque
- The wind efficiency dependance on stellar parameters is not well known but
 - ➔ Could explain the spread in excentricity (Dobbs-Dixon et al 2004)
 - ➔ Could explain the spin/orbit misalignement (Dawson 2014)
 - ➔ Could explain the delay of the tidal decay (Damiani et al 2014)

Conclusion

- Understanding star-planet interaction is a necessary step to confront observations and predictions of formation/migration models
- For hot-Jupiters around late-type stars the magnetized wind torque can be comparable (and opposite) to the tidal torque
- Dynamics and environmental effect in multi-planet systems are fundamental to assess the habitability
- Better ages and stellar physics are essential to understand exoplanetary systems dynamics (we need PLATO)

Thank you!

