

# **Recent origin of Saturn's moonlets and F-ring from rings viscous spreading**

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The regular satellites of the giant planets are believed to have finished their accretion concurrent with the planets at about 4.5 Gy ago<sup>1,2,3,4</sup>. A population of Saturn's small moons, that orbit just exterior to the main rings, challenges this conception because they are dynamically young<sup>5,6</sup> (less than  $10^7$  years). They are also under-dense<sup>7</sup> ( $\sim 600 \text{ kg/m}^3$ ) and show spectral characteristics similar to the rings<sup>8,9</sup>. Some authors have suggested they might have accreted at the rings' edge<sup>7,10,11</sup>. However their formation and coupled evolution with the rings was never simulated because of the too-short orbital timescales, casting doubts on this hypothesis. Using a hybrid simulation we unveil a mechanism by which the viscous spreading of Saturn's rings beyond the Roche Limit (the distance beyond which the rings are gravitationally unstable) gives birth to the today's small-moons: masses distribution and orbital architecture are reproduced. The current confinement of the main rings and the presence of the dusty F-ring are shown to be direct consequences of this mechanism that couples viscous evolution and satellite formation. Saturn's rings, like a mini-protoplanetary disk, may be the last place where accretion was recently active in the Solar System, some  $10^6$  to  $10^7$  years ago.

The low density of Saturn's small moons, their icy composition, closeness to the rings and rapid tidal timescales have suggested for long that a genetic link may exist with Saturn's icy rings. On the one hand, the population of small moons exterior to the Roche Limit (Atlas, Prometheus, Pandora, Janus and Epimetheus) have a mass/distance relation remarkably different from Saturn's main satellites (Fig. 1). On the other hand, N-body simulations<sup>10,12</sup> show that ring material spreading viscously beyond 142,000 km would be gravitationally unstable and that  $\sim 100 \text{ m}$  aggregates would form in less than 10 orbits. Below 138,000 km, accretion becomes inefficient<sup>12,13,14</sup>, so the outer edge of the region where accretion is

prevented (the ‘‘Roche Region’’) is  $R_L \approx 140000 \pm 2000 \text{ km}$ .  $R_L$  differs slightly from the classical definition of the Roche Limit since its precise location depends on a diversity of factors like the bodies’ relative size, spin<sup>14,15</sup> etc. Today, the ring’s outer edge is sharply confined at 136,775 km by Janus’ gravitational torque<sup>5,16</sup>. However this configuration must be transient (<10 My) because the rings repel Janus and, in turn, the ring is pushed inward to conserve angular momentum. As the moon migrates outward, the rings’ edge confinement location moves accordingly and may pass the nearby Roche Limit. So ring material may have spread across Saturn’s Roche Limit in the past, or will do in the future.

To simulate on long-timescales ( $\sim 10^9$  years) the coupled evolution of aggregates with the rings we designed a hybrid model in which two codes are self-consistently coupled<sup>17</sup>: (1) a 1D-hydrodynamical model to track the ring’s viscous evolution and (2) an analytical orbital model to track the moonlets. The evolution of the ring’s surface density  $\sigma(r)$  is computed using a finite-element scheme, on a staggered mesh, solving for the surface-density equation of a keplerian disk under the effect of viscous torque and moonlets’ gravitational-torques<sup>18</sup>:

$$\frac{d\sigma}{dt} = \frac{3}{r} \frac{\partial}{\partial r} \left[ r^{1/2} \frac{\partial}{\partial r} \left( \nu(r) \sigma r^{1/2} \right) - \frac{1}{3\pi\sqrt{GM}} r^{1/2} T(r) \right], \quad \text{Eq. 1}$$

where  $r$ ,  $\nu(r)$ ,  $G$  and  $M$  are the distance, local viscosity, gravitational constant, and Saturn’s mass respectively.  $T(r)$  is the sum, for all the satellites, of the Lindblad resonances torque densities located at distance  $r$  from Saturn (Supplementary Information, section 1). A realistic viscosity model, which is an increasing function of the surface density, is implemented<sup>19</sup>. Because aggregates are formed at  $r \geq R_L$  after only a few orbital periods (negligible compared to viscous timescales), the accretion of ring material into moonlets is considered to occur instantaneously: at each time step, all ring mass located beyond  $R_L$  is removed from the hydrodynamical simulation and transformed into one additional satellite.

Every individual satellite is tracked and its mass  $m_s$ , semi-major axis  $a_s$  and eccentricity  $e_s$  are tabulated (SI, section 2). The resultant eccentricities are so small ( $<10^{-4}$ ) that they do not influence the system's evolution. The semi-major axes evolve under the effects of planet's tides and ring torque according to:

$$\frac{da_s}{dt} = \frac{3k_{2p}m_sG^{1/2}R_p^5}{Q_pM^{1/2}a_s^{11/2}} + \frac{2a_s^{1/2}\Gamma_s}{m_s(GM)^{1/2}} \quad \text{Eq.2}$$

where  $Q_p$ ,  $k_{2p}$ ,  $R_p$  stand for the planet's dissipation factor, Love number, and radius respectively (SI, section 1.2).  $\Gamma_s$  is the sum of the ring torques for all first-order Lindblad resonances of the satellite<sup>6,18,20</sup>. When the orbital separation of two satellites is smaller than 2.2 mutual Hill radii, they are merged<sup>12</sup>. To test the full procedure, the Moon formation from a circumterrestrial disk was successfully reproduced<sup>21</sup> (SI, section 3).

The initial conditions are as follows : the A ring is initially represented by a disk extending from 122,000 km to 136,000km with a constant surface density  $\sigma_0$ . The initial state of Saturn's rings is unknown and they could have been denser in the past<sup>22</sup>. Cases with  $\sigma_0=400$ , 1000, 5000 and  $10^4$  kg/m<sup>2</sup> have been considered. The case with  $\sigma_0=400$  kg/m<sup>2</sup> is presented (about the present surface density of Saturn's A ring. Fig. 2).  $R_L$  is set to 140000 km.

The simulation starts with no satellite so that the ring spreads freely initially and reaches  $R_L$  in  $\sim 10^6$  years (Fig. 2). Then, moonlets accrete at  $R_L$  from the ring material and grow through mutual encounters. They induce step-like structures in the ring's surface density near moonlets' Lindblad resonances (Fig 2.a-d). At these locations the ring angular-momentum is directly transferred to the satellites, inducing an orbital decay of the ring material and an expansion of the satellites' orbit. As a consequence, the ring material moves inward and

accumulates just interior to the resonant location, resulting in the visible step-like structures in the surface-density function. While  $\Gamma_s$  increases with the moonlet's mass, the disk's surface density decreases due to spreading and thus so do the viscosity and the viscous torque ( $\Gamma_v$ )<sup>16</sup>. When the magnitude of  $\Gamma_s$  becomes larger than that of  $\Gamma_v$ , the disk's outer edge is confined and the disk stops spreading (Fig.2.e). This happens when the largest moon reaches  $\sim 10^{17}$  kg (and  $\sim 10^{18}$  kg) for  $\sigma_0 \leq 1000$  kg/m<sup>2</sup> ( and  $\sigma_0 \geq 5 \times 10^3$  kg/m<sup>2</sup>). As the moonlets' masses increase, the ring's outer edge shifts inward down to the position of the first Lindblad resonance for which the torque is strong enough to counterbalance  $\Gamma_v$ . On long timescales, due to the decrease of the surface density, confinement is increasingly easier and, in this model, the outer edge is around  $\sim 135,000$  km at  $4 \times 10^9$  yrs for  $\sigma_0 = 400$  kg/m<sup>2</sup>. For higher values of  $\sigma_0$ , the disk spreads viscously more rapidly, more massive moonlets are formed; this ultimately shifts the disk's edge even below 130,000km. In conclusion, the current confinement of Saturn's rings outer edge seems the consequence of on-going satellite accretion occurring at the Roche Limit. The outer-edge jumps from one resonance to another, depending on the local balance between the positive viscous torque and negative torque induced by the population of small moons.

Concerning the moonlets, they migrate outward due to the positive induced by the rings and the planet. Since both torques are increasing functions of the satellite's mass<sup>20</sup> ( $\Gamma_s \propto m_s^2$ ) their migration rate increases accordingly and more massive satellites migrate more rapidly. Different migration rates lead to crossings and merging. Since all satellites appear at the same location ( $R_L$ ) a simple orbital architecture appears in which satellites are radially sorted: their distance to Saturn is an increasing function of their mass, in agreement with observations (Fig. 1). So the actual orbital architecture of the small moons may be the direct consequence of ring/satellites interactions.

These results imply also that Saturn's small moons are gravitational aggregates made of icy ring particles which would explain naturally their very low densities. When they form, they should be initially elongated like Hill spheres<sup>7,10</sup> like observed for some of the moons<sup>7</sup>. Pandora or Epimetheus seem to deviate somewhat from this shape<sup>7</sup>, but this could result from post-accretional restructuration. The moons' spectral similarities with Saturn's rings<sup>9,24</sup> may be the direct result of their formation within the rings, as well as their apparent absence of silicates<sup>7</sup> since Saturn's rings seem devoid of such material<sup>8,24</sup>.

The coupling of the ring-confinement with moonlet's migration induces a feedback on the formation rate of the moonlets. When the ring's edge is repelled below  $R_L$ , the moonlet production stops (compare Fig. 2.e and 2.f). Conversely, as satellites recede away, the location of the ring's edge follows and satellite production can set-in again for any ring material that reaches  $R_L$ . This stop-and-go mechanism is apparent in the mass history of satellites (Fig. 2.f): it limits their mass to about the smallest mass necessary to confine the ring and can be analytically estimated (SI, section 4). The mass-distributions of moonlets obtained in our simulations matches well the observed one with similar cumulative indexes and overall shape (Fig. 3). However, these results may be affected by non-linearities that could appear in the torque of the most massive satellites (SI, section 3.4). A direct consequence of this feedback mechanism is that the mass of the largest moon must be of the order of the mass necessary to confine the ring: this is actually the case with Janus, the most massive of the small moons, that confines the A ring's outer edge<sup>5,16</sup>. The current mass of Janus implies that the A ring surface density was between  $10^3$  kg/m<sup>2</sup> and  $5 \times 10^3$  kg/m<sup>2</sup> at the time Janus formed (whereas the ring's total mass remains undetermined<sup>23</sup>).

Saturn's main rings are encircled by a dusty and dynamically active ringlet, called F-ring, located at 140,500 km, whose origin is still debated. In the present work, its presence may have a simple explanation: due to the spreading of the A ring, aggregates are formed around 140,000 km and suffer subsequent collisional evolution while migrating outward. Since accretion should not be 100% efficient below 142,000km<sup>14</sup> colliding aggregates will release dust and produce a dusty ring with a non-negligible mass, like today's F-ring<sup>25,26,27</sup>. In the current orbital configuration, the F ring is not feed with new aggregates coming from the A ring because of the confinement induced by Janus. However, on longer timescales, the dusty F ring will be fed again when the A ring viscously spreads again. Thus, the F ring would have been always present because of the regular renewing of its material. In this scenario, it is envisioned as the dusty signature of ring material crossing the Roche Limit due to the global viscous spreading of the rings. The age of the F-ring material would be about the same as the nearby moonlets<sup>25</sup>,  $10^{6-7}$  years, even though the main rings could be older.

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## LEGENDS OF FIGURES

**Figure 1:** Mass of Saturn's inner moons versus distance. Their names and average diameters are indicated in thumbnails images. Mimas, Enceladus and Tethys are also plotted for comparison. The vertical dashed line shows the location of the outer edge of Saturn's A ring at 136,750km, the vertical dashed-dotted lines stands for the F-ring location (a ~1000 km wide ringlet located between Prometheus and Pandora). Blue and red thick lines show simple log-fits to the mass/distance relations for small moons and main moons, respectively. Images credit: Cassini (NASA/JPL/SSI).

**Figure 2:** The time evolution of a model with  $\sigma_0=400 \text{ Kg/m}^2$ . (a) to (d): the time evolution of the ring's surface density (solid line, left scale) and the satellite system as a function of distance (black circles, right scale). As the ring spreads inward and outward, its surface density decreases. Ring material crossing the Roche Limit ( $R_L=140000 \text{ km}$  here) is transformed into moonlets. At the end of the simulation only  $1.5 \times 10^{18} \text{ kg}$  of material remains in the A ring whereas  $3.5 \times 10^{18} \text{ kg}$  spread below 120,000 Km and  $1.8 \times 10^{17} \text{ kg}$  were transformed into moonlets; (e) The time evolution of the location of the ring's outer edge, defined as the place where the surface density drops below  $1 \text{ kg/m}^2$ . When the edge is confined at the location of a satellite's resonance, the disk's viscous torque (that tends to push material beyond the ring's edge by transferring angular momentum) is perfectly balanced by the satellite's gravitational torque, thus the ring material is prevented from spreading further outward; (f) solid line: total mass into satellites as a function of time; dashed line: mass of the largest satellite.

**Figure 3:** Comparing the mass-distribution of the moonlets obtained in our simulation with observations. The cumulative mass-distributions of moonlets obtained in four simulations with different initial surface-densities ( $\sigma_0$ ) are fitted with single power-law functions  $N(>m) \propto m^{-\alpha}$ : Case A (green) :  $\sigma_0 = 400 \text{ kg/m}^2$  ( $\alpha=0.31\pm0.06$ ); Case B (yellow):  $\sigma_0 = 1000 \text{ kg/m}^2$  ( $\alpha=0.19\pm0.01$ ); Case C (red):  $\sigma_0 = 5000 \text{ kg/m}^2$  ( $\alpha=0.22\pm0.03$ ); Case D (blue):  $\sigma_0 = 10000 \text{ kg/m}^2$  ( $\alpha=0.17\pm0.04$ ). The actual population of Saturn's moonlets (black solid line, with label Sat.,  $\alpha=0.27\pm0.07$ ) is well in the range of masses and number of bodies obtained in the simulations. Note that some of our distributions (A,B and C) show a knee and a shallower slope at smaller sizes, as observed for Saturn's small moons.







