## Covariant extension of Generalized Parton Distributions

From an Overlap of Light-cone Wave-functions to a Double Distribution

Nabil Chouika<br>Irfu/SPhN, CEA Saclay - Université Paris-Saclay

Café SPhN, 20 février 2017



## Outline

(1) Introduction to Generalized Parton Distributions

- Definition and properties
(2) Overlap and Double Distribution representations of GPDs
- Overlap of Light-cone wave functions
- Double Distributions
(3) From an Overlap of LCWFs to a Double Distribution
- Inversion of Incomplete Radon Transform
- Results

4 Conclusion

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## Definition of GPDs

- Quark GPD (twist-2, spin-0 hadron): (Müller et al., 1994; Radyushkin, 1996; Ji, 1997)

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\begin{equation*}
H^{q}(x, \xi, t)=\left.\frac{1}{2} \int \frac{\mathrm{~d} z^{-}}{2 \pi} e^{i x P^{+} z^{-}}\left\langle P+\frac{\Delta}{2}\right| \bar{q}(-z) \gamma^{+} q(z)\left|P-\frac{\Delta}{2}\right\rangle\right|_{z^{+}=0, z_{\perp}=0} \tag{1}
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with:


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- Impact parameter space GPD (at $\xi=0$ ): (Burkardt, 2000)

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q\left(x, \overrightarrow{b_{\perp}}\right)=\int \frac{\mathrm{d}^{2} \overrightarrow{\Delta_{\perp}}}{(2 \pi)^{2}} e^{-i \overrightarrow{b_{\perp}} \cdot \overrightarrow{\Delta_{\perp}}} H^{q}\left(x, 0,-{\overrightarrow{\Delta_{\perp}}}^{2}\right) . \tag{2}
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\begin{gather*}
\int \mathrm{d} x H^{q}(x, \xi, t)=F^{q}(t)  \tag{3}\\
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- Cauchy-Schwarz theorem in Hilbert space.


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- A given hadronic state is decomposed in a Fock basis: (Brodsky and Lepage, 1989)

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\begin{equation*}
|H ; P, \lambda\rangle=\sum_{N, \beta} \int[\mathrm{~d} x]_{N}\left[\mathrm{~d}^{2} \mathbf{k}_{\perp}\right]_{N} \Psi_{N, \beta}^{\lambda}\left(x_{1}, \mathbf{k}_{\perp 1}, \ldots\right)\left|N, \beta ; k_{1}, \ldots, k_{N}\right\rangle \tag{7}
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\begin{align*}
H^{q}(x, \xi, t) & =\sum_{N, \beta}{\sqrt{1-\xi^{2}}}^{2-N} \sqrt{1+\xi^{2-N}} \sum_{a} \delta_{a, q}  \tag{9}\\
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H^{q}(x, \xi, t)=\int_{\Omega} \mathrm{d} \beta \mathrm{~d} \alpha\left(F^{q}(\beta, \alpha, t)+\xi G^{q}(\beta, \alpha, t)\right) \delta(x-\beta-\alpha \xi) \cdot(10)
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- Pobylitsa gauge (One Component DD): (Pobylitsa, 2003)

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\begin{equation*}
H(x, \xi, t)=(1-x) \int_{\Omega} \mathrm{d} \beta \mathrm{~d} \alpha f(\beta, \alpha, t) \delta(x-\beta-\alpha \xi) \tag{12}
\end{equation*}
$$

with

$$
\left\{\begin{array}{l}
F(\beta, \alpha)=(1-\beta) f(\beta, \alpha)  \tag{13}\\
G(\beta, \alpha)=-\alpha f(\beta, \alpha)
\end{array}\right.
$$

## Radon transform




- Radon Transform:

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\mathcal{R} f(x, \xi) \propto \int \mathrm{d} \beta \mathrm{~d} \alpha f(\beta, \alpha) \delta(x-\beta-\alpha \xi) \tag{14}
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- ERBL region: $|x|<|\xi|$.


## Outline

(1) Introduction to Generalized Parton Distributions

- Definition and properties
(2) Overlap and Double Distribution representations of GPDs
- Overlap of Light-cone wave functions
- Double Distributions
(3) From an Overlap of LCWFs to a Double Distribution
- Inversion of Incomplete Radon Transform
- Results

4) Conclusion

## From DGLAP GPD to a DD

- In Overlap representation: DGLAP region only (e.g. two-body LCWFs).


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Find $f(\beta, \alpha)$ on square $\{|\alpha|+|\beta| \leq 1\}$ such that

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- DD $\boldsymbol{f}(\beta, \alpha)$ exists (if the GPD behaves well) and is unique.
- We can reconstruct the GPD everywhere.


## Support properties




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- Domains $\beta<0$ and $\beta>0$ are uncorrelated in the DGLAP region.
- Divide and conquer:
- Better numerical stability.
- Lesser complexity: $O\left(N^{p}+N^{p}\right) \ll O\left((N+N)^{p}\right)$.


## Domain for the inversion



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- Rotated square $\left[-\frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}}\right] \times\left[-\frac{1}{\sqrt{2}}, \frac{1}{\sqrt{2}}\right]$ :

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- $\alpha$-parity of the DD:

$$
\begin{equation*}
f(\beta,-\alpha)=f(\beta, \alpha) \tag{16}
\end{equation*}
$$

## First result



- Real application to a DSE toy model.

$$
\begin{gathered}
f(\beta, \alpha)= \begin{cases}? & \beta>0 \\
0 & \beta<0\end{cases} \\
\downarrow \\
\left.H(x, \xi)\right|_{x>|\xi|}=30 \frac{(1-x)^{2}\left(x^{2}-\xi^{2}\right)}{\left(1-\xi^{2}\right)^{2}}
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- Goal: extend the DGLAP GPD of Ref. (Mezrag, 2015; Mezrag et al., 2016).

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- Gauge introduced for positivity.


## Quantitative comparison of DDs (DSE model)



Figure: Quantitative comparison for the DSE toy model DD. Left: Theoretical discretized DD. Middle: Numerical solution at tolerance $10^{-6}$. Right: Absolute difference.

$$
f(\beta, \alpha)= \begin{cases}\frac{30}{4}\left(1-3 \alpha^{2}-2 \beta+3 \beta^{2}\right) & \beta>0 \\ 0 & \beta<0\end{cases}
$$

## Quantitative comparison of GPDs (DSE model)





Figure: Quantitative comparison for the DSE toy model GPD obtained from the numerical DD solution. Left: Theoretical GPD. Middle: Numerical solution at tolerance $10^{-6}$. Right: Absolute difference.

$$
H(x, \xi)= \begin{cases}30 \frac{(1-x)^{2}\left(x^{2}-\xi^{2}\right)}{\left(1-\xi^{2}\right)^{2}} & x>|\xi| \\ \frac{15(x-1)\left(x^{2}-\xi^{2}\right)\left(\xi^{2}+2|\xi| x+x\right)}{2|\xi|^{3}(|\xi|+1)^{2}} & |x|<|\xi|\end{cases}
$$

## DDs \& GPDs (DSE model)



Figure: DDs and GPDs for the DSE model. Left: $t=0 \mathrm{GeV}^{2}$. Right: $t=1 \mathrm{GeV}^{2}$.

## DDs \& GPDs (Gaussian model)



Figure: DDs and GPDs for a Gaussian toy model (similar to AdS/QCD). Left: $t=0 \mathrm{GeV}^{2}$. Right: $t=1 \mathrm{GeV}^{2}$.

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- Any questions?



## Bibliography I

D. Müller, D. Robaschik, B. Geyer, F. M. Dittes, and J. Hořejši, "Wave functions, evolution equations and evolution kernels from light ray operators of QCD", Fortsch. Phys. 42 (1994) 101-141, arXiv:hep-ph/9812448 [hep-ph].
A. V. Radyushkin, "Scaling limit of deeply virtual Compton scattering", Phys. Lett. B380 (1996) 417-425, arXiv:hep-ph/9604317 [hep-ph].
X.-D. Ji, "Deeply virtual Compton scattering", Phys. Rev. D55 (1997) 7114-7125, arXiv:hep-ph/9609381 [hep-ph].
M. Burkardt, "Impact parameter dependent parton distributions and off forward parton distributions for zeta —> 0", Phys. Rev. D62 (2000) 071503, arXiv:hep-ph/0005108 [hep-ph], [Erratum: Phys. Rev.D66,119903(2002)].
B. Pire, J. Soffer, and O. Teryaev, "Positivity constraints for off - forward parton distributions", Eur. Phys. J. C8 (1999) 103-106, arXiv:hep-ph/9804284 [hep-ph].
S. J. Brodsky and G. P. Lepage, "Exclusive Processes in Quantum Chromodynamics", Adv. Ser. Direct. High Energy Phys. 5 (1989) 93-240.
M. Diehl, T. Feldmann, R. Jakob, and P. Kroll, "The overlap representation of skewed quark and gluon distributions", Nucl. Phys. B596 (2001) 33-65, arXiv:hep-ph/0009255 [hep-ph], [Erratum: Nucl. Phys.B605,647(2001)].
M. Diehl, "Generalized parton distributions", Phys. Rept. 388 (2003) 41-277, arXiv:hep-ph/0307382 [hep-ph].

## Bibliography II

P. V. Pobylitsa, "Solution of polynomiality and positivity constraints on generalized parton distributions", Phys. Rev. D67 (2003) 034009, arXiv:hep-ph/0210150 [hep-ph].
H. Moutarde, "Nucleon Reverse Engineering: Structuring hadrons with colored degrees of freedom", 2015.
C. Mezrag, "Generalised Parton Distributions : from phenomenological approaches to Dyson-Schwinger equations", PhD thesis, IRFU, SPhN, Saclay, 2015.
C. Mezrag, H. Moutarde, and J. Rodriguez-Quintero, "From Bethe-Salpeter Wave functions to Generalised Parton Distributions", Few Body Syst. 57 (2016), no. 9, 729-772, arXiv:1602.07722 [nucl-th].
P. C. Hansen, "Regularization Tools version 4.0 for Matlab 7.3", Numerical Algorithms 46 (2007), no. 2, 189-194.

## Discrete ill-posed problem



Theoretical "L-curve": curve parameterized by the regularization factor.
(fig. taken from Ref. (Hansen, 2007))

L-curve


L-curve with the iteration number as regularization factor.

