

# *Current-Driven Instabilities in the Crab Nebula Jet: Results from Numerical Simulations*

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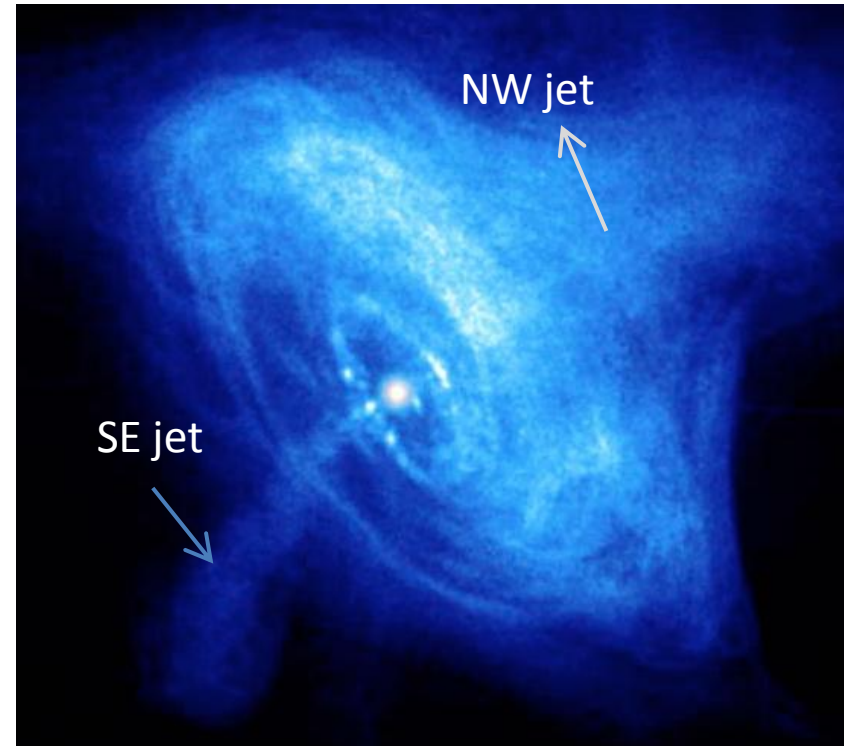
# Outline

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1. *Observational Evidence*
2. *Numerical Models of Relativistic MHD Jets*
  - *2D Axisymmetric models*
  - *3D models → Kink instabilities*
3. *Results*
4. *Summary*

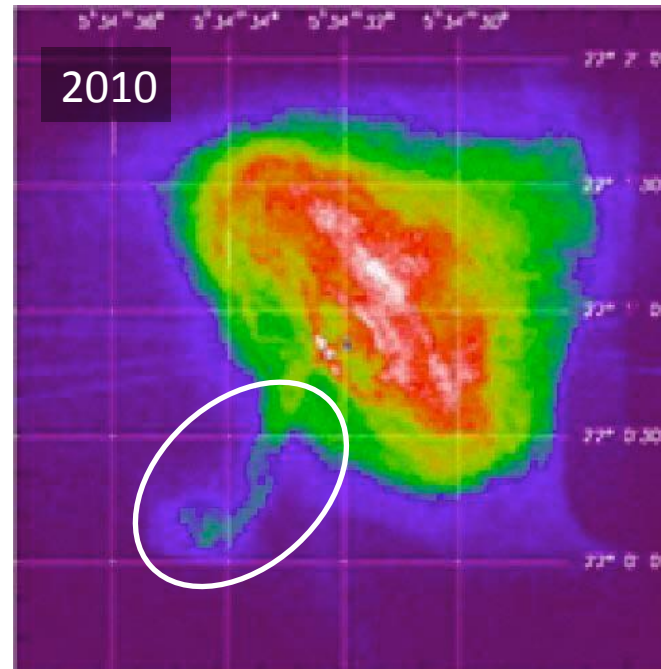
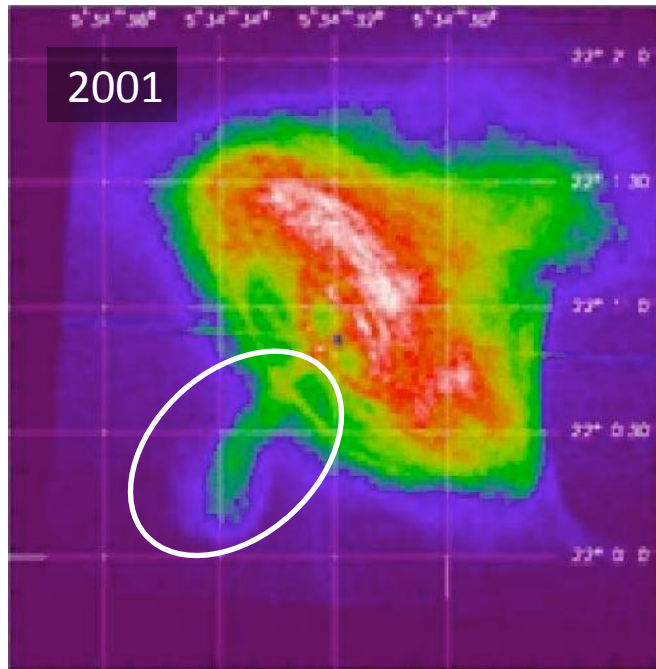
# Observational Evidence

- X-ray observations (Chandra) show the emergence of bipolar jets extending to the SE and NW of the pulsar;
- A region of diffuse emission (Anvil) may be associated with shocks and marks the base of the X-ray and optical jet;
- Knots of emission are seen along the jets;
- In the SE jet material flows with  $v/c \sim 0.4$  slowing down to  $\sim 0.02$  into the nebula;



# Jet Wiggling

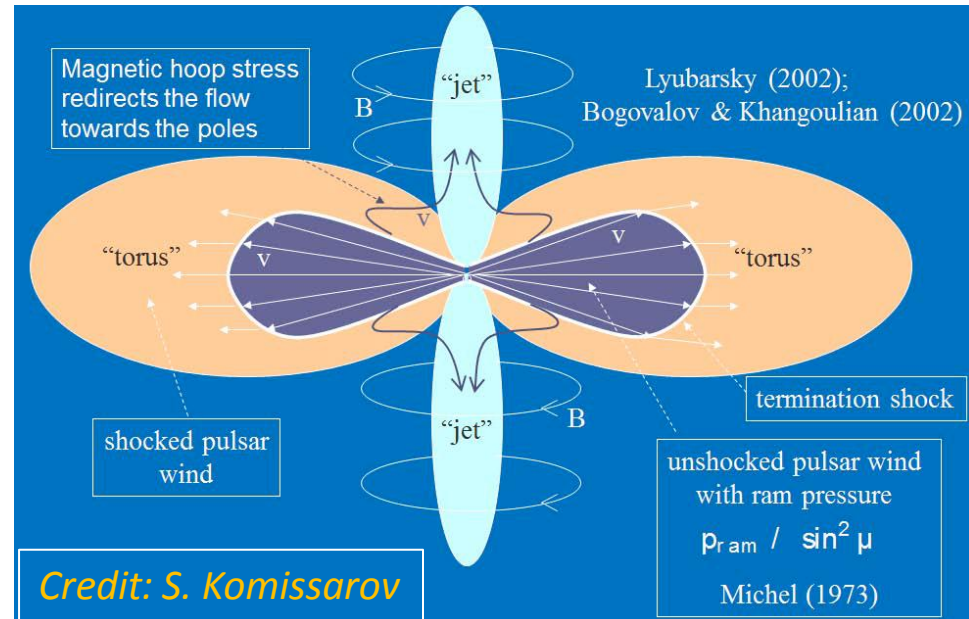
- SE jet morphology is “S” shaped and show remarkable time variability:



- → evidence for some kind of flow instability (Current Driven ?)

# On the Origin of the Jet

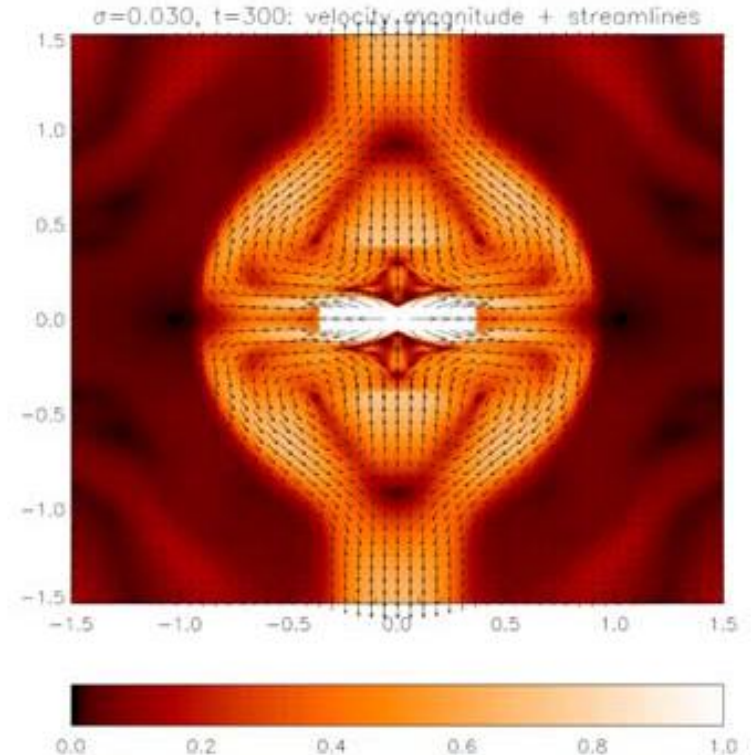
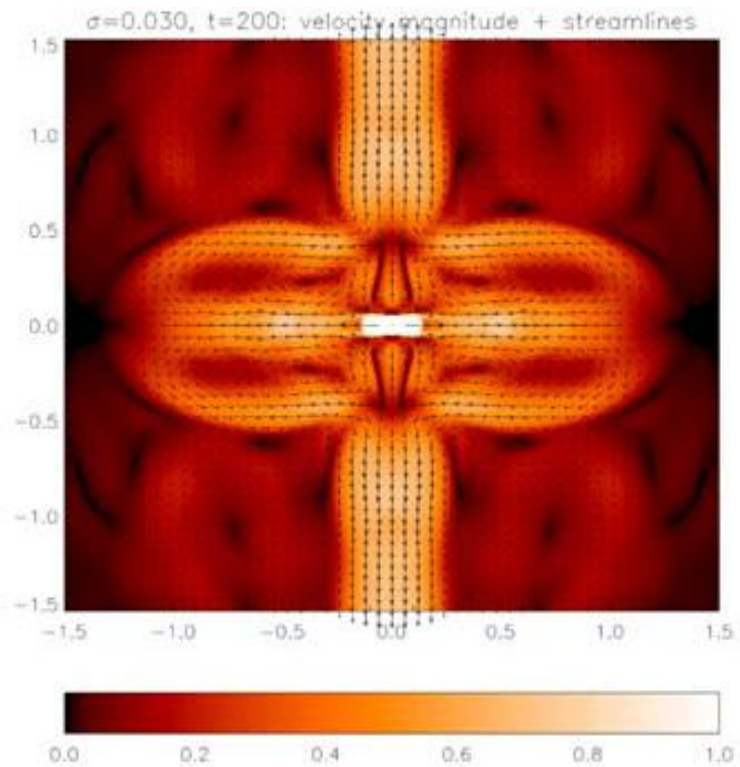
- Jet forms downstream of the wind termination shock;
- Magnetic fields confine matter towards polar axis;  
→ **“tooth-paste”** effect: hoop stress of the azimuthal magnetic field carried by the wind (Lyubarsky 2002).



- Models confirmed by 2D axisymmetric numerical simulations (Komissarov & Lyubarski 2003,2004, Del Zanna et al. 2004, Bogovalov et al. 2005)

# Jet Origin: Axisymmetric Models

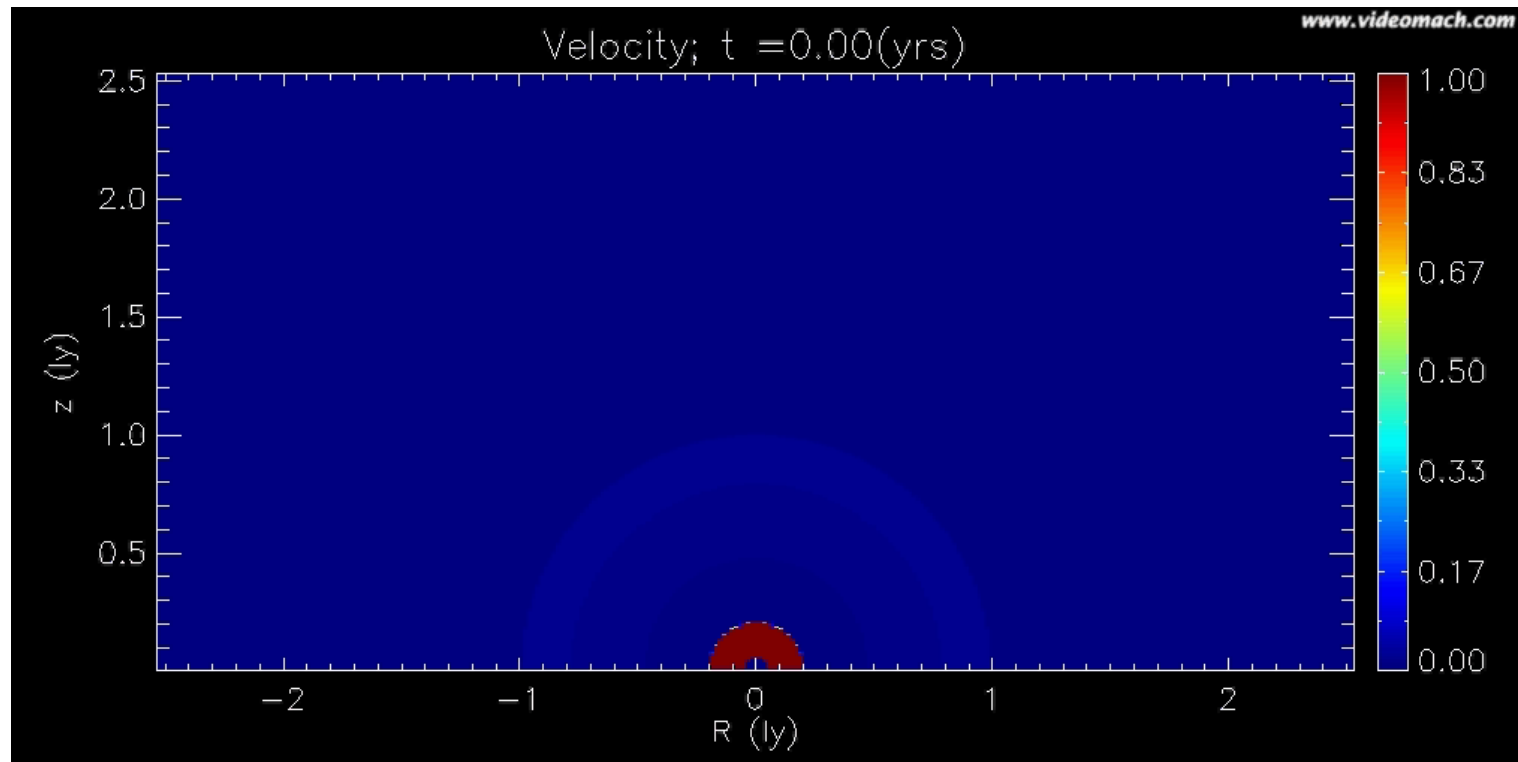
- For moderate/large  $\sigma = B^2/(4\pi\rho c^2\gamma^2)$  magnetic hoop stress suppresses high velocity outflows in the equatorial plane and divert them towards the polar axis partially driving the super-fast jet<sup>1</sup>



<sup>1</sup>Del Zanna et al, A&A (2004) 421,1063

# Axisymmetric PWN Models

- Results from 2D axisymmetric simulations predict hollow and hot jets initially carrying purely axial current ( $B_\phi \neq 0, B_z = B_R = 0$ );



- $B_z = 0 \rightarrow$  Pitch = 0;  $1.3 \lesssim Ms \lesssim 2$  (hot jet);  $\rho_j / \rho_e \lesssim 10^{-6}$
- Two free parameters:  **$0.1 \lesssim \sigma \lesssim 10$**  and  **$2 \lesssim \gamma \lesssim 4$** ;

# Jet Models

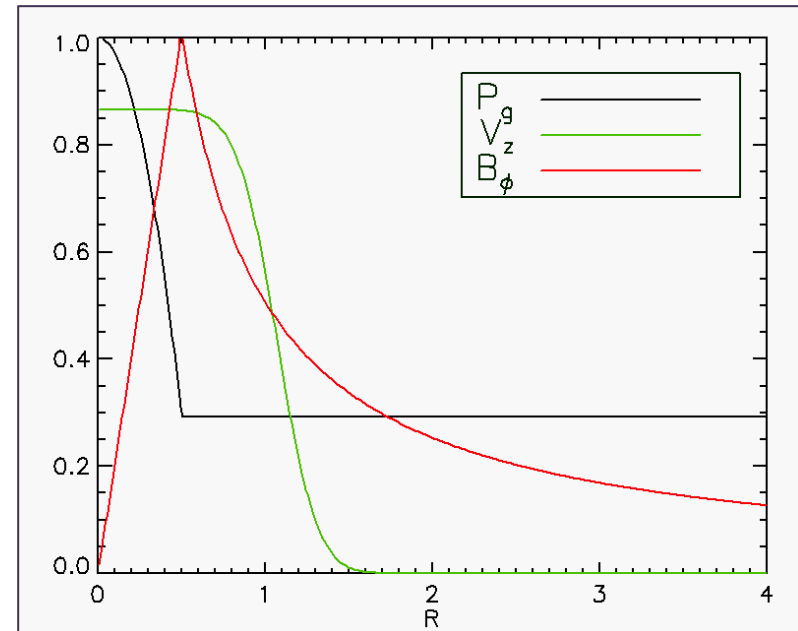
- We consider a 2-parameter ( $\gamma$ ,  $\sigma$ ) family of light, hot jets with  $\rho_j/\rho_e=10^{-6}$ ;  $M_s = 1.7$ ;

$$v_z(R) = \sqrt{1 - \frac{1}{\gamma_j^2} \psi \left( \frac{R^8}{R_j^8} \right)} \quad B_\phi(R) = \begin{cases} B_m \frac{R}{a} & \text{for } R \leq a, \\ B_m \frac{a}{R} \psi \left( \frac{R^6}{R_j^6} \right) & \text{for } R > a, \end{cases}$$

with  $(B_m)^2 \propto \sigma$ .

- Radial momentum balance holds across the beam

$$\left( \frac{\partial p}{\partial r} - \frac{w\gamma^2 v_\phi^2}{r} \right) \hat{r} = (\nabla \cdot \mathbf{E}) \mathbf{E} + \mathbf{J} \times \mathbf{B}$$





# Equations and Numerical Method

- We solve the equations of a relativistic perfectly conducting fluid describing energy/momentum and particle conservation (relativistic MHD equations)

$$\frac{\partial(\rho\gamma)}{\partial t} + \nabla \cdot (\rho\gamma \mathbf{v}) = 0,$$

$$\frac{\partial \mathbf{m}}{\partial t} + \nabla \cdot [w\gamma^2 \mathbf{v}\mathbf{v} - \mathbf{B}\mathbf{B} - \mathbf{E}\mathbf{E}] + \nabla p_t = 0,$$

$$\frac{\partial \mathbf{B}}{\partial t} - \nabla \times (\mathbf{v} \times \mathbf{B}) = 0,$$

$$\frac{\partial \mathcal{E}}{\partial t} + \nabla \cdot (\mathbf{m} - \rho\gamma \mathbf{v}) = 0,$$

$$\mathbf{m} = w\gamma^2 \mathbf{v} + \mathbf{E} \times \mathbf{B}$$

$$\mathcal{E} = w\gamma^2 - p + \frac{\mathbf{B}^2 + \mathbf{E}^2}{2} - \rho\gamma$$

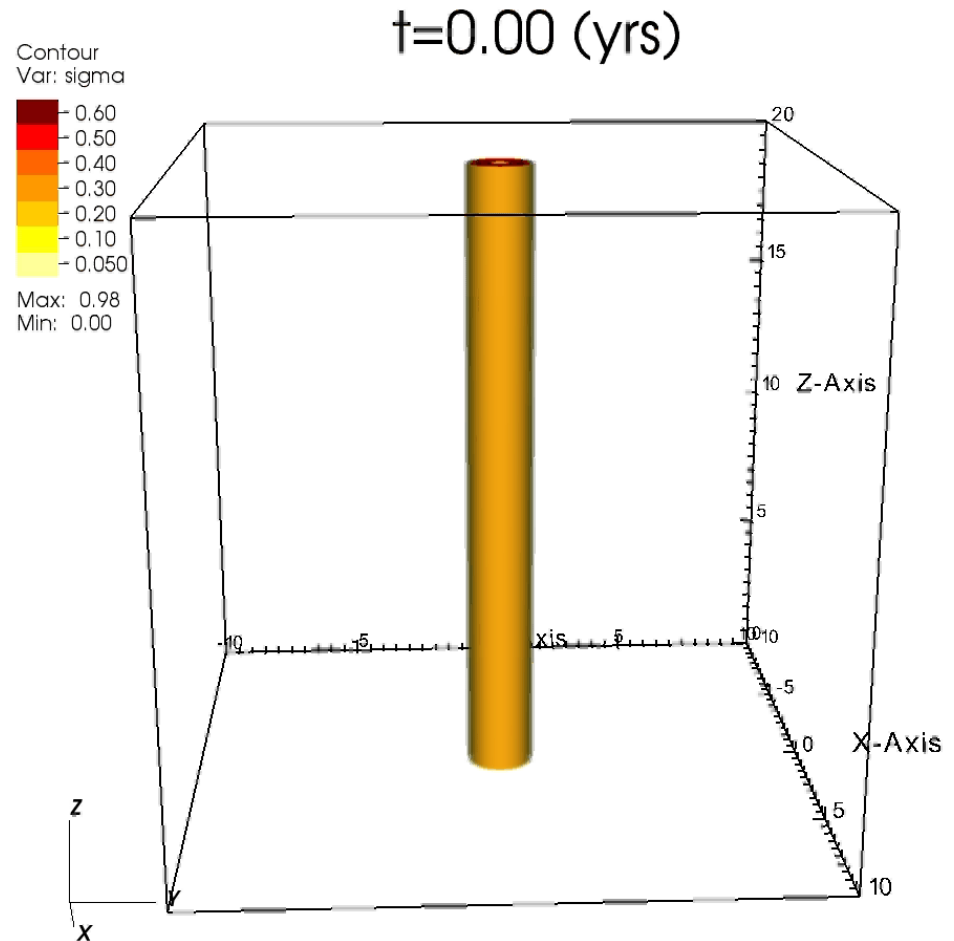
$$w = \rho + \frac{\Gamma p}{\Gamma - 1}$$

- We use the PLUTO<sup>1,2</sup> code for astrophysical fluid dynamics (<http://plutocode.ph.unito.it>);
- Linear reconstruction + HLLD Riemann solver;
- Numerical resolution 320 x 320 x 768 zones (≈ 20 zones on the jet).

<sup>1</sup>Mignone et al, *ApJS* (2007) 170, 228; <sup>2</sup>Mignone et al, *ApJS* (2012) 198, 7

# Instabilities in Periodic Jets

- These jet configurations are unstable to a variety of modes, mainly KH and CD;
- For non-zero velocities KH and CD modes mix up<sup>1</sup>.
- At large magnetizations, the  $m=1$  CD mode (kink) prevails.
- At large velocities KH modes prevails.

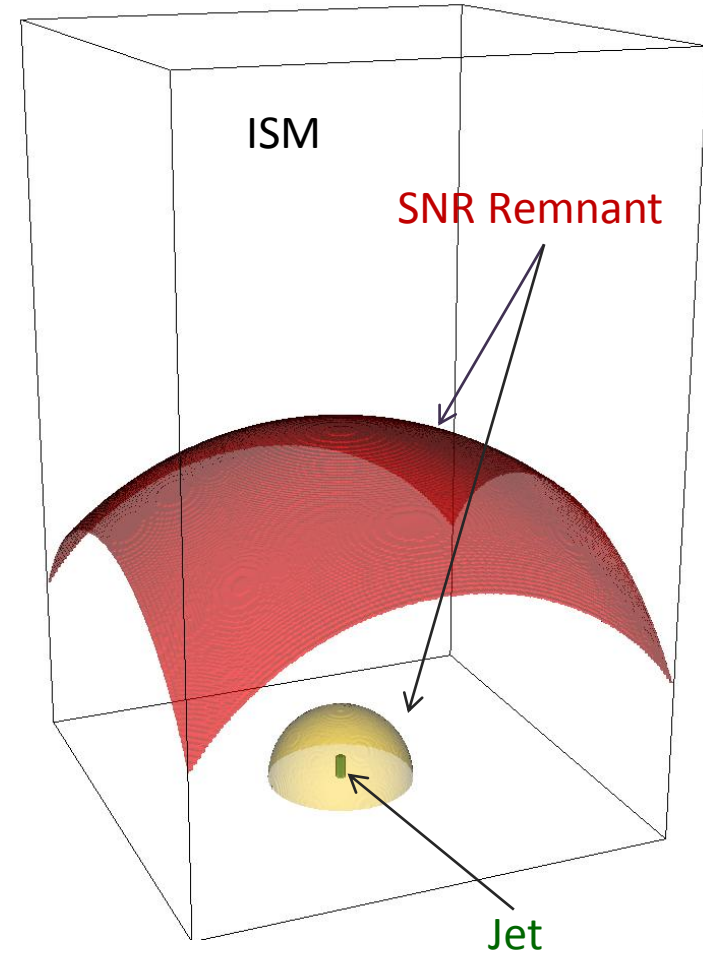


$$\sigma = 1; \gamma = 2$$

<sup>1</sup>Bodo et al. *MNRAS* (2013, accepted)

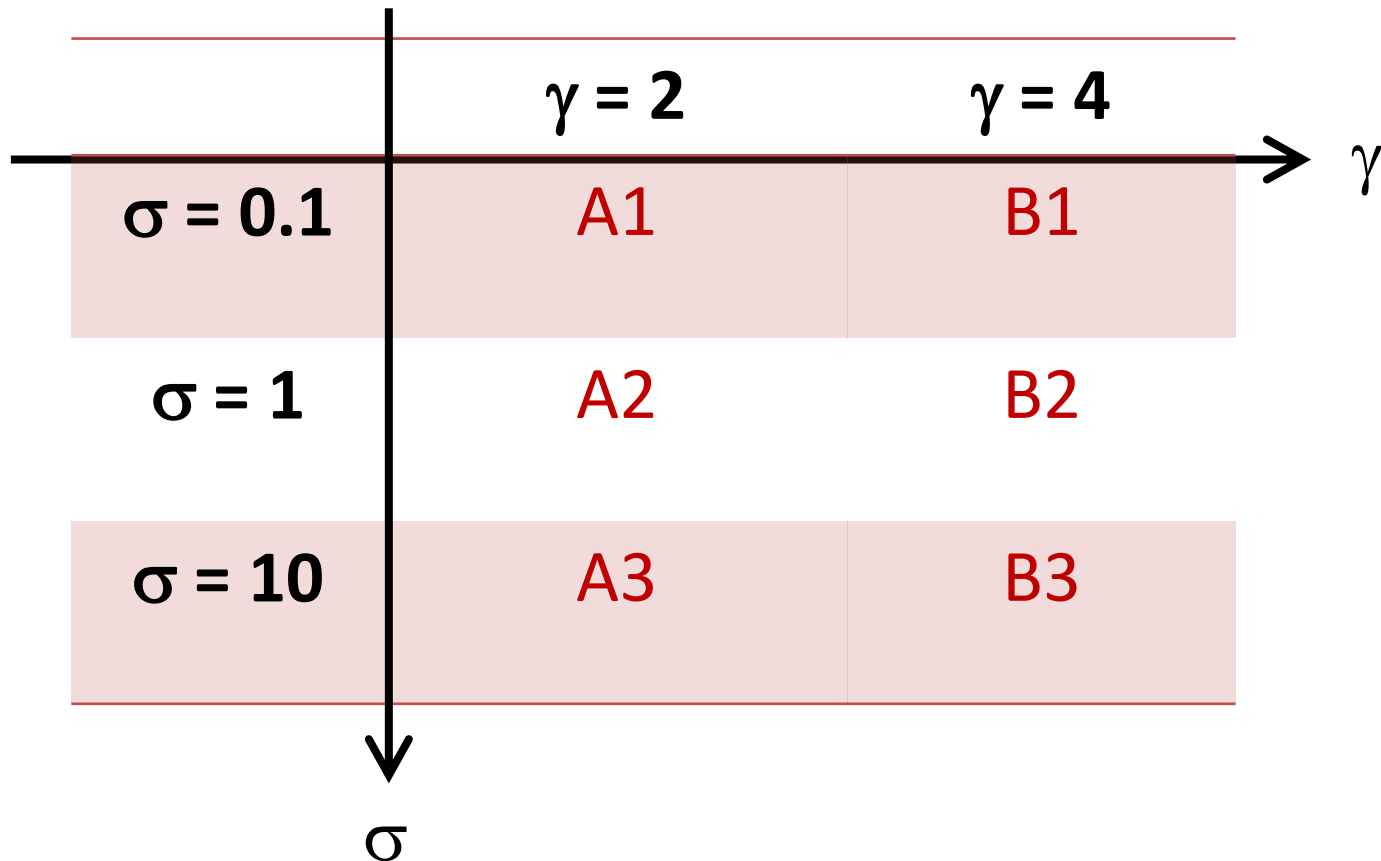
# A More Realistic 3D Scenario

- We consider a 3D Cartesian domain with  $x, y \in [-0.8, 0.8]$  (ly),  $z \in [0, 2.5]$  (ly).
- Freely expanding supernova ejecta ( $3 M_{\text{sun}}$ ,  $E = 10^{51}$  erg) for  $0.2 < r < 1$  (ly)
- Pulsar wind structure not considered: jet already formed as the result of the collimation process;
- Supersonic injection nozzle at the lower z-boundary.



# Simulation Cases

- $\gamma$  and  $\sigma$  are free parameters. We consider slow and fast jets with weak, moderate and strong magnetic fields (6 cases)



# Results: Case A2

$\gamma = 2$

$\gamma = 4$

$\sigma = 0.1$

A1

B1

$\sigma = 1$

A2

B2

$\sigma = 10$

A3

B3

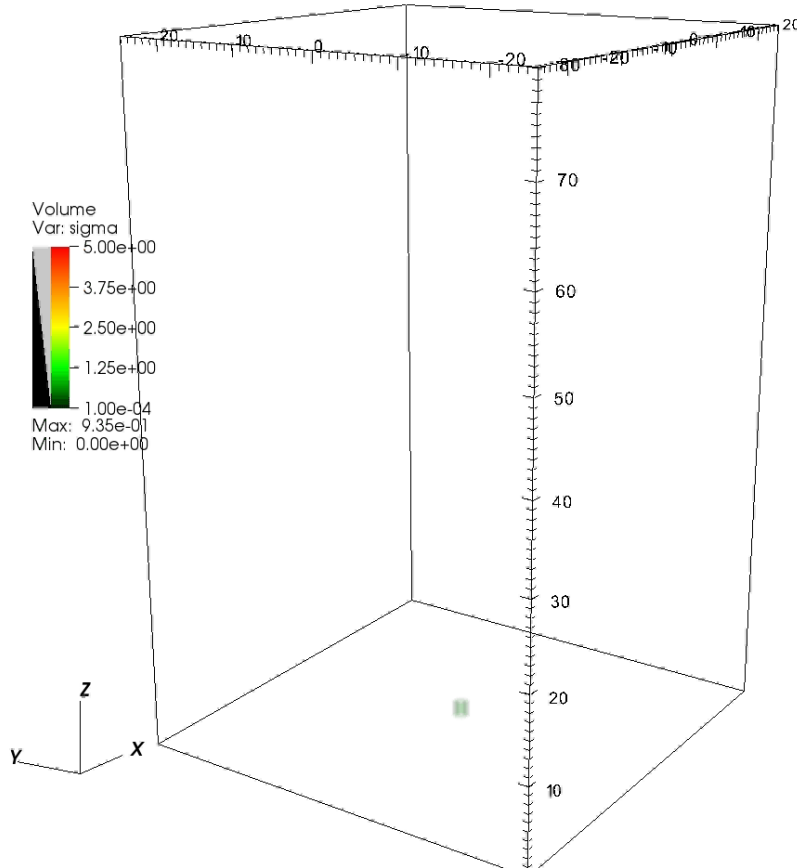
[www.videomach.com](http://www.videomach.com)

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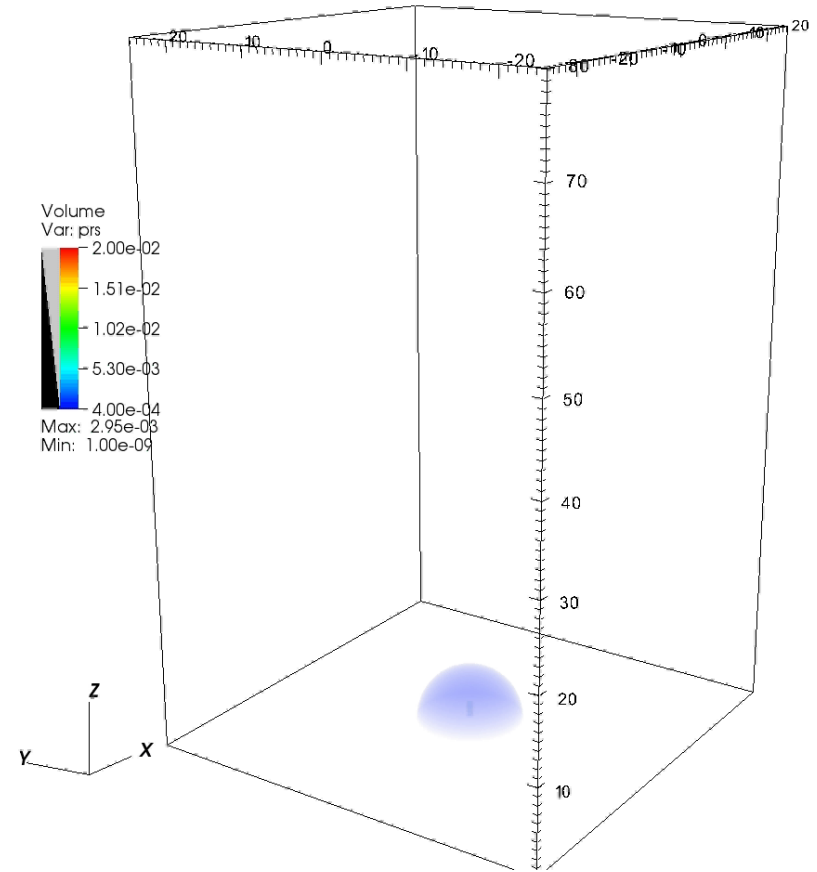
$\sigma$

$\rho$

Case A2, t=0.00 (yrs)



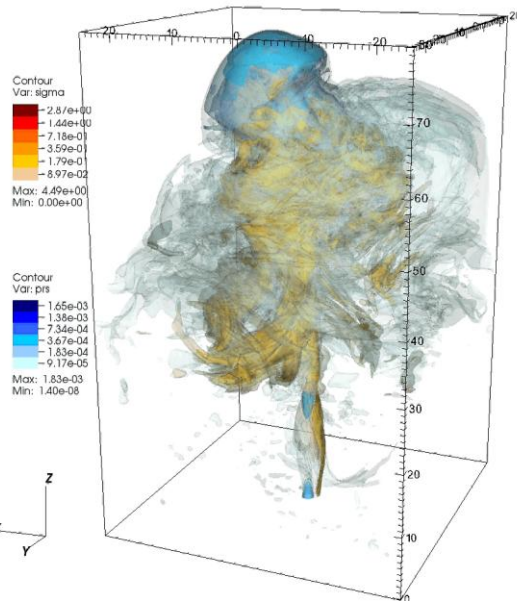
Case A2, t=0.00 (yrs)



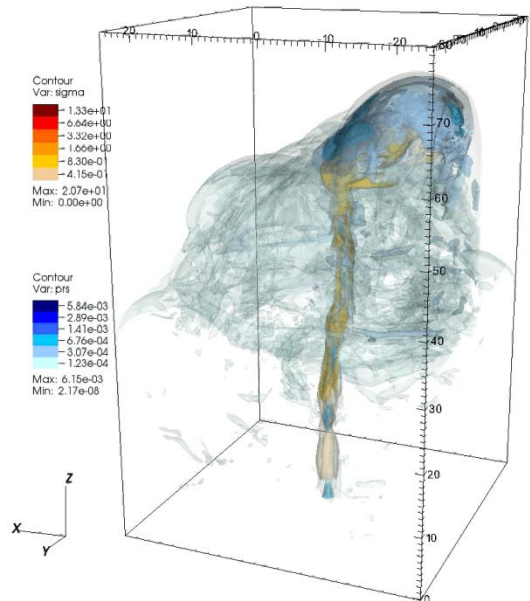
# General Features: low speed jets

- Low speed jets advance slowly ( $v_{\text{head}} < 0.02$ ) ← large density contrast;
- Evolve entirely inside the remnant;
- Larger  $\sigma$  drive magnetically supported jets and show the largest deflections;

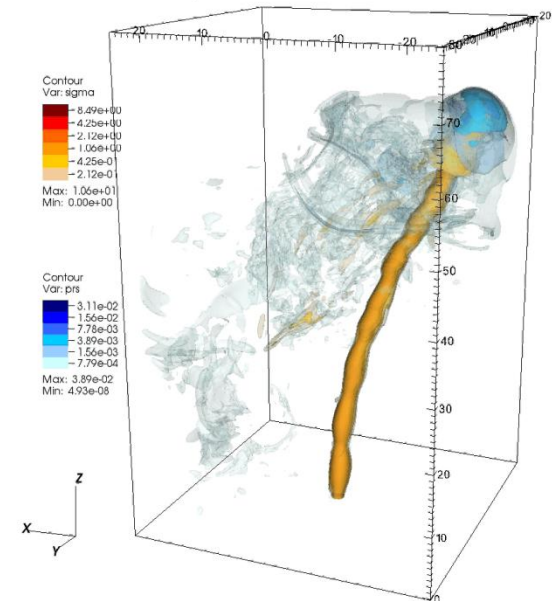
Case A1,  $t=151.67$  (yrs)



Case A2,  $t=130.09$  (yrs)



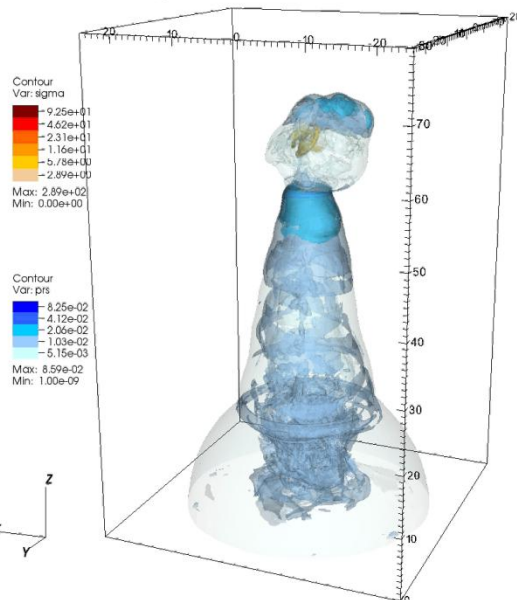
Case A3,  $t=95.19$  (yrs)



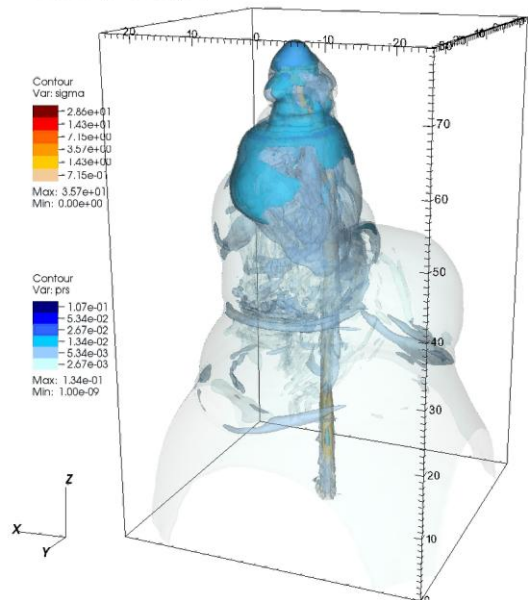
# General Features: high speed jets

- High-speed jets propagate faster ( $v_{\text{head}} < 0.05$ );
- Reach the outer supernova remnant after  $\approx 50$  years;
- For large  $\sigma$  deflections are present but smaller than low speed jets  $\rightarrow$  Lorentz factor has a stabilizing effect.

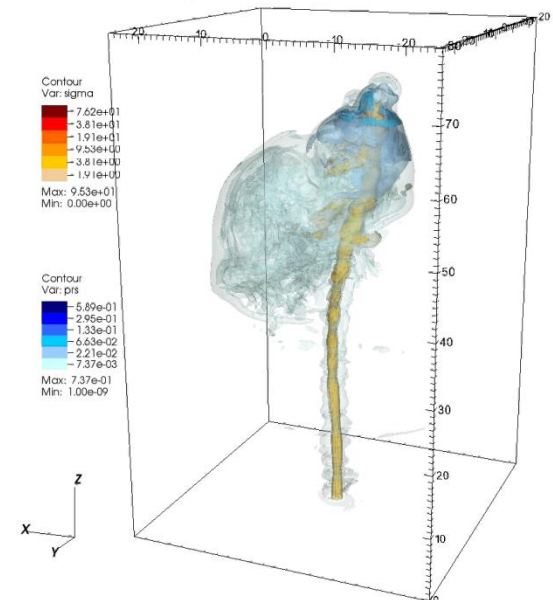
Case B1,  $t=30.14$  (yrs)



Case B2,  $t=47.60$  (yrs)

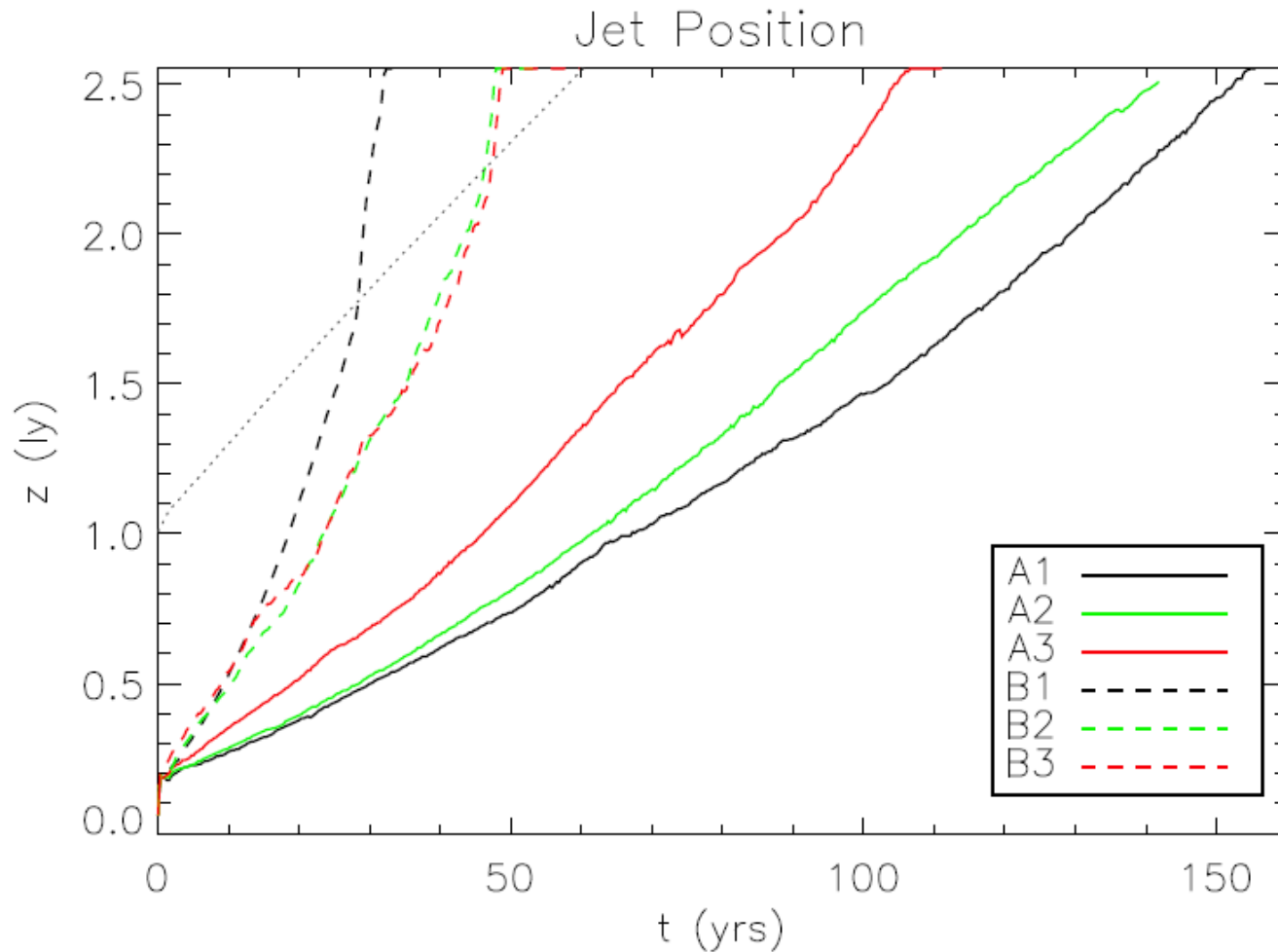


Case B3,  $t=47.60$  (yrs)



# Propagation Speed

- Jets with  $\gamma=4$  “drill out” of the remnant in less than 50 years...



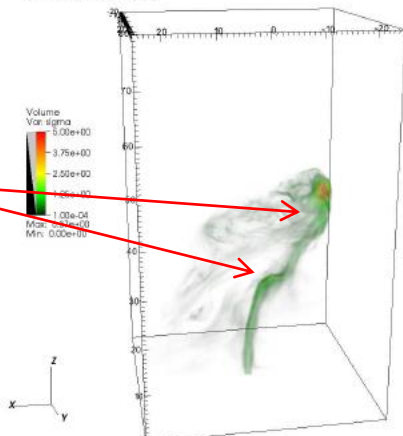


# Jet Structure

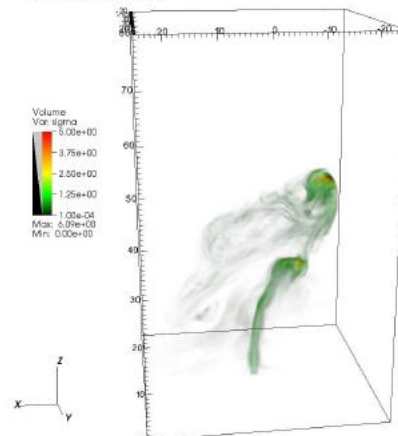
- *Back-end regions*: quasi-periodic stationary pinch ( $m=0$ ) shocks;
- *Front-end regions*: jet fragmentation at deflection sites forming short-lived unstable structures;

Kinked deflections

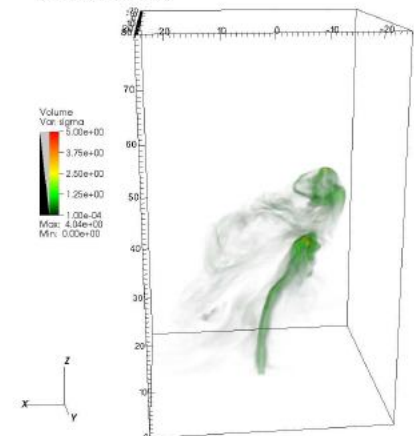
Case A2, t=92.97 (yrs)



Case A2, t=93.29 (yrs)

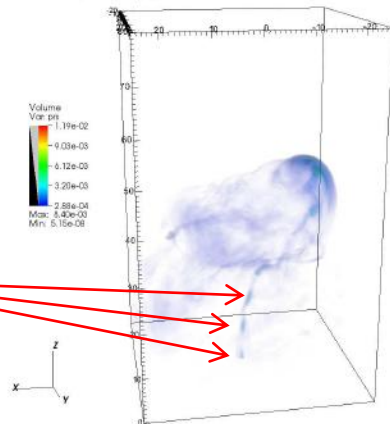


Case A2, t=93.60 (yrs)

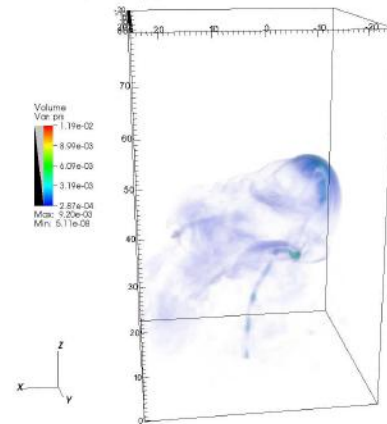


$\sigma$

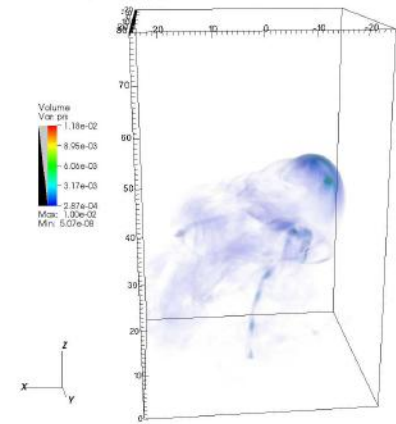
Case A2, t=92.97 (yrs)



Case A2, t=93.29 (yrs)



Case A2, t=93.60 (yrs)



$\psi$

Pinching shocks

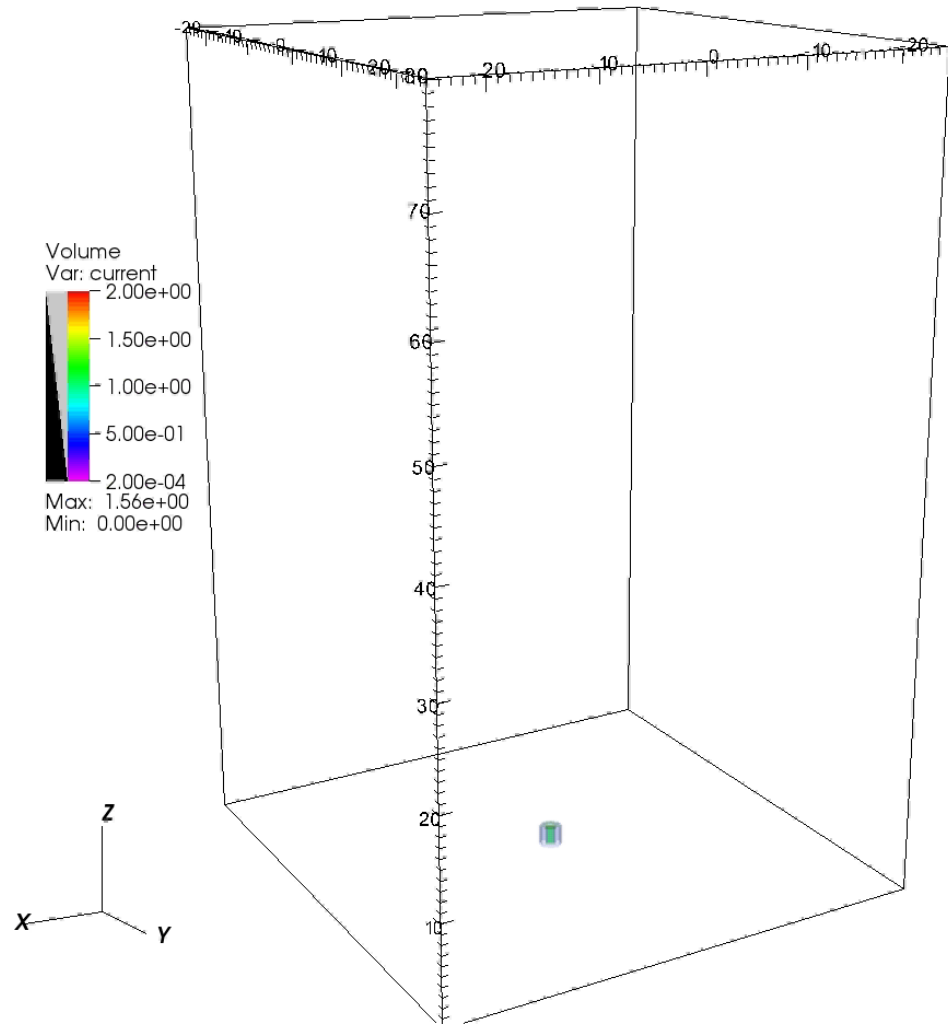
# Jet Structure

## ➤ Front-end regions:

- rapid variability
- strong interaction with the ambient

➤ For strong magnetization  
→ formation of twisted helical structures.

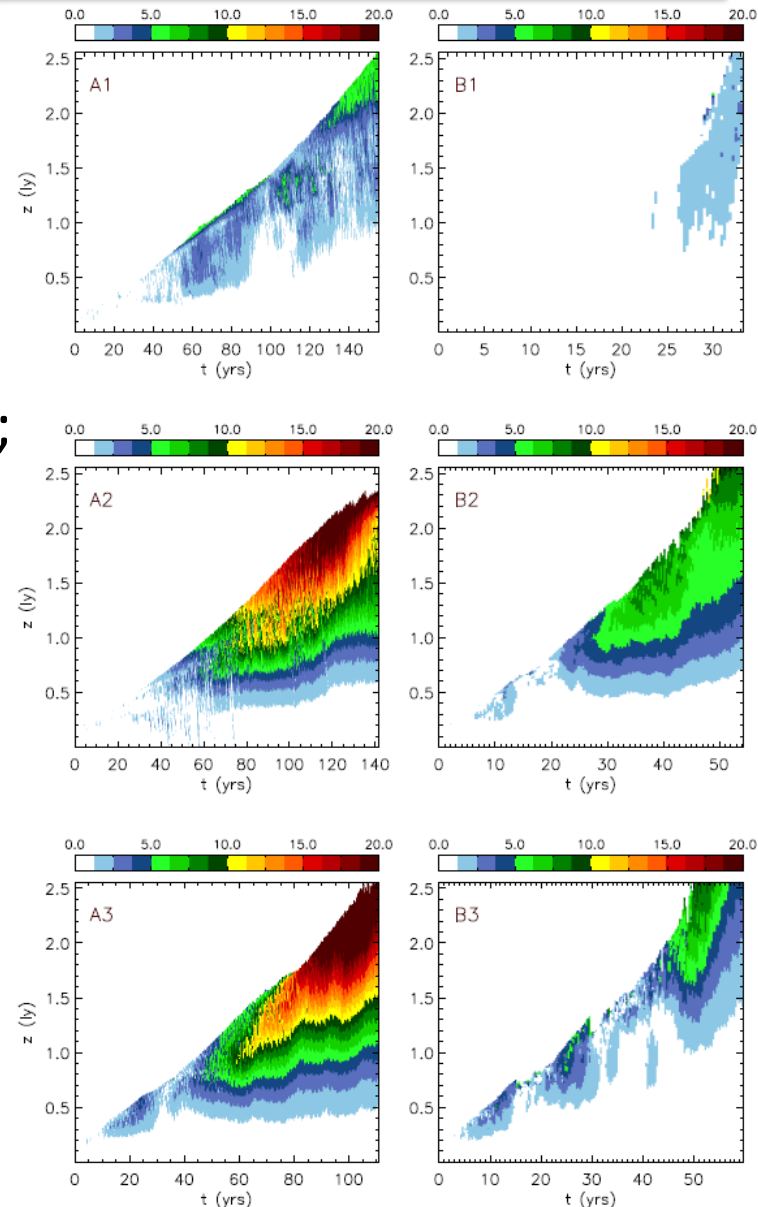
Case A3, t=0.00 (yrs)



$$J = \nabla \times B$$

# Jet Deflections

- Center of mass  $\rightarrow$  amount of deflection;
- Low-speed ( $\gamma \approx 2$ ), magnetized ( $\sigma \gtrsim 1$ ) jets show the largest bending ( $\gtrsim 20 R_j$ );
- Larger Lorentz factors ( $\gamma \approx 4$ ) have a stabilizing effect<sup>1</sup>;
- Weakly magnetized jets less affected by the growth of instability;

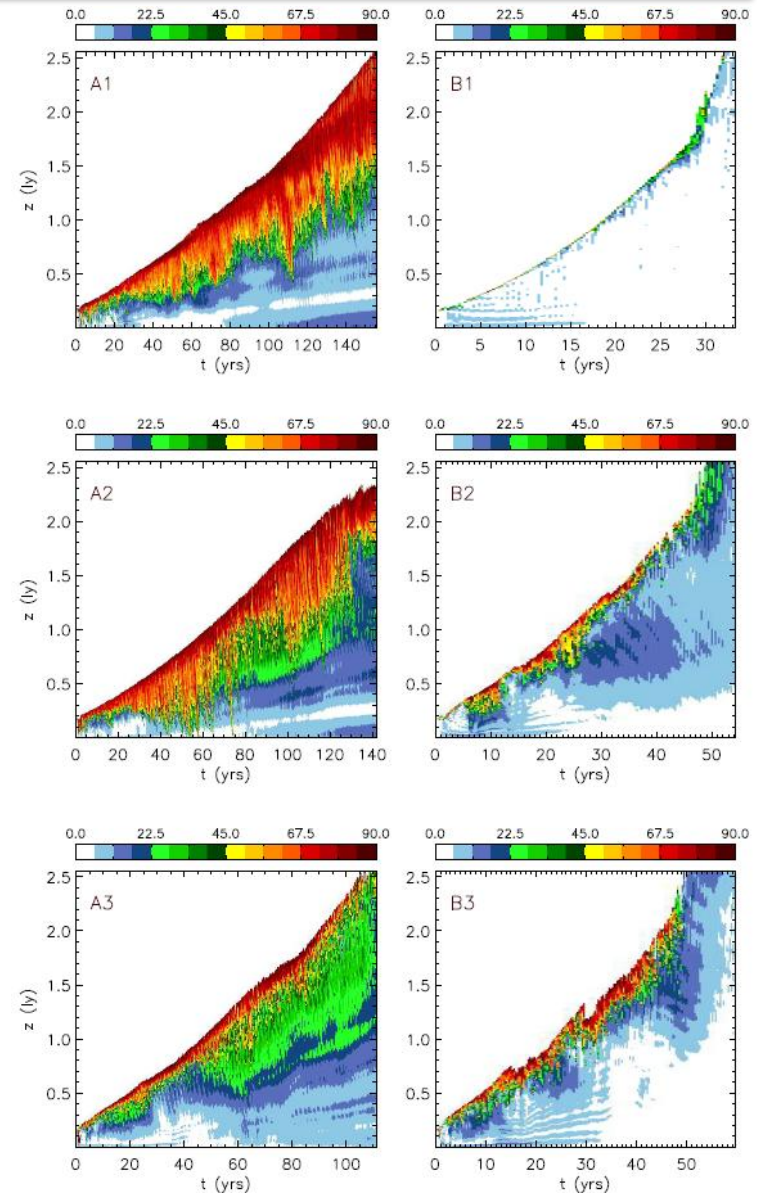


<sup>1</sup>Bodo et al. *MNRAS* (2013, accepted)

# Flow Direction

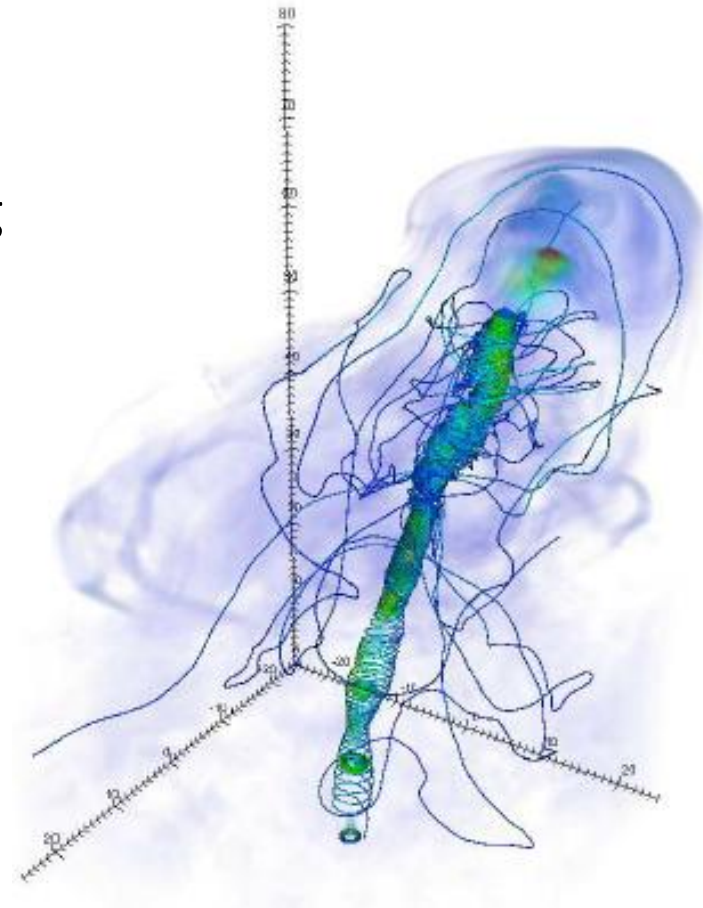
$$\bar{\theta}_{\pm} = \text{acos} \left\langle \frac{v_{z,\pm}}{|\mathbf{v}|}, \chi_j \right\rangle$$

- Change in trajectory → variation of the average propagation velocity.
- Low-speed jets → large-scale curved structure with  $\theta$  gradually changing from  $0^\circ$  (base) to  $90^\circ$  (head);
- High-speed jets stabilized by the larger inertia, build large kicks at the head.



# Magnetic Fields

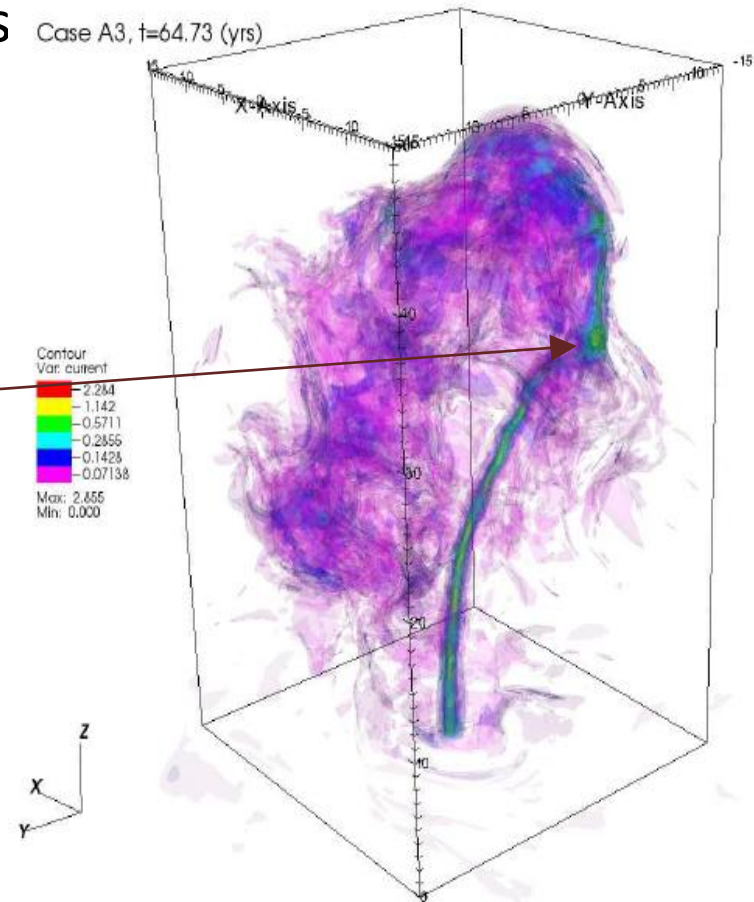
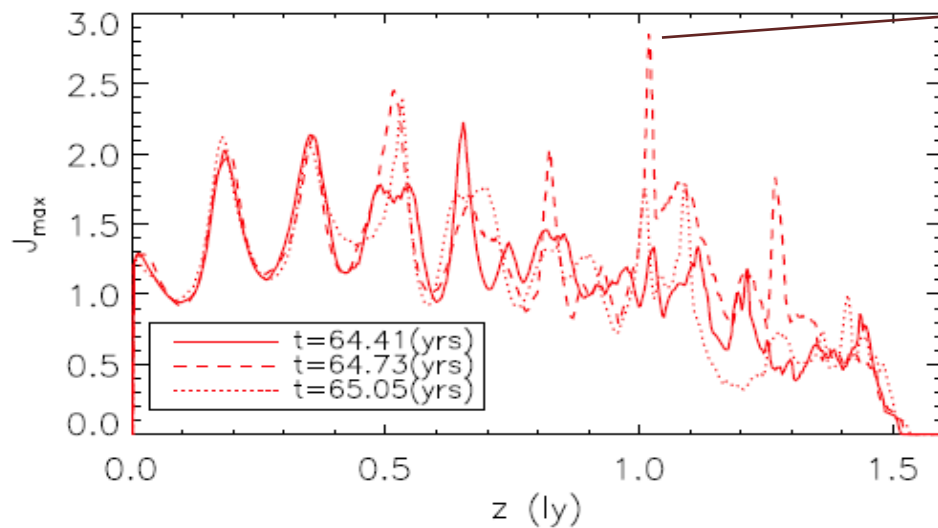
- Magnetic field remains mainly toroidal or helical during the propagation;
- Azimuthal field “shields” the core preventing interaction with the surrounding<sup>1</sup>.
- Poynting flux efficiently diverted at the termination shock and scattered via the backflow to feed the cocoon.
- Magnetic field dissipates and becomes turbulent in the cocoon (→ randomization<sup>2</sup>)



<sup>1</sup>Mignone et al, *MNRAS* (2010) 402, 7; <sup>2</sup>Porth et al., *MNRAS* (2013)

# Current Sheets

- Current sheets localized in two regions:
  - at conical pinch shocks
    - quasi-steady, periodic
  - at jet “kinks” → short-lived episodes
- Magnetic reconnection
  - particle acceleration regions ?



# Summary

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- 3D models of azimuthally confined relativistic jets are very different from 2D axisymmetric models:
  - Kink-unstable non-axisymmetric structures with large time-variability;
  - Large  $\sigma$  ( $\gtrsim 1$ ) leads to considerable jet deflections, one-sided propagation;
  - Jet wiggling progressively more pronounced towards the jet head
  - Larger Lorentz factors  $\rightarrow$  stabilizing effect;
  - Multiple shocks observed at pinching regions and deflection sites where flow changes direction;
- Low-speed ( $\gamma \lesssim 2$ ), moderately/highly magnetized jets ( $\sigma \simeq 1-10$ ) are promising candidates for explaining the morphology of the Crab jet.
- Future models will consider the jet-torus connection in 3D

*Thank you*