

# News from AFTER@LHC

**Jean-Philippe Lansberg**  
IPN Orsay, Université Paris-Sud

SPhN, IRFU, CEA-Saclay, France



## **Part 1:** Why a new fixed-target experiment for HEP now ?

# Outline

**Part 1:** Why a new fixed-target experiment for HEP now ?

**Part 2:** A Fixed-Target Experiment using LHC beams: AFTER@LHC

# Outline

**Part 1:** Why a new fixed-target experiment for HEP now ?

**Part 2:** A Fixed-Target Experiment using LHC beams: AFTER@LHC

**Part 3:** Flagship studies and news

# Outline

**Part 1:** Why a new fixed-target experiment for HEP now ?

**Part 2:** A Fixed-Target Experiment using LHC beams: AFTER@LHC

**Part 3:** Flagship studies and news

**Part 4:** Back to the future

# Outline

**Part 1:** Why a new fixed-target experiment for HEP now ?

**Part 2:** A Fixed-Target Experiment using LHC beams: AFTER@LHC

**Part 3:** Flagship studies and news

**Part 4:** Back to the future

**Conclusions and Outlooks**

# Part I

Why a new fixed-target experiment for  
HEP now ?

# Decisive advantages of Fixed-target experiments

- They brought essential contributions to particle & nuclear physics

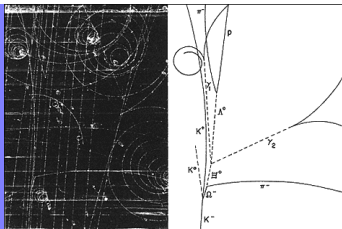


# Decisive advantages of Fixed-target experiments

- They brought essential contributions to particle & nuclear physics
  - particle discoveries ( $\Omega^-$  (sss),  $J/\psi$ ,  $\Upsilon$ , ...)

# Decisive advantages of Fixed-target experiments

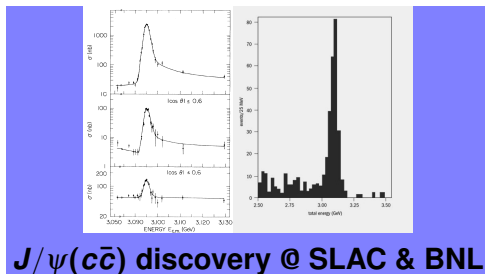
- They brought essential contributions to particle & nuclear physics
  - particle discoveries ( $\Omega^-(sss)$ ,  $J/\psi$ ,  $\Upsilon$ , ...)



**$\Omega(sss)$  discovery @ BNL**

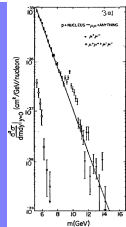
# Decisive advantages of Fixed-target experiments

- They brought essential contributions to particle & nuclear physics
  - particle discoveries ( $\Omega^-$  (sss),  $J/\psi$ ,  $\Upsilon$ , ...)



# Decisive advantages of Fixed-target experiments

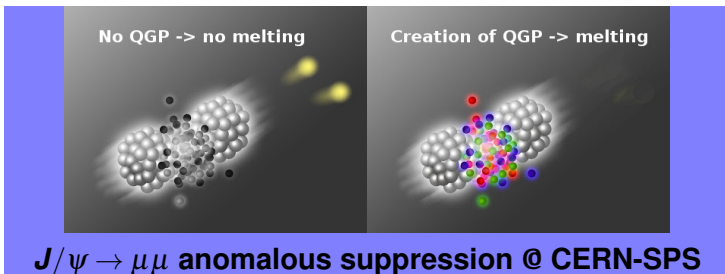
- They brought essential contributions to particle & nuclear physics
  - particle discoveries ( $\Omega^-$  (sss),  $J/\psi$ ,  $\Upsilon$ , ...)



$\Upsilon(b\bar{b})$  discovery @ Fermilab

# Decisive advantages of Fixed-target experiments

- They brought essential contributions to particle & nuclear physics
  - particle discoveries ( $\Omega^-(sss)$ ,  $J/\psi$ ,  $\Upsilon, \dots$ )
  - evidence for the novel dynamics of quarks and gluons in HIC (QGP)



# Decisive advantages of Fixed-target experiments

- They brought essential contributions to particle & nuclear physics
  - particle discoveries ( $\Omega^- (sss)$ ,  $J/\psi$ ,  $\Upsilon$ , ...)
  - evidence for the novel dynamics of quarks and gluons in HIC (QGP)
  - observation of surprising QCD phenomena
    - breakdown of the Lam-Tung relation,
    - colour transparency,
    - higher-twist effects in forward meson production ,
    - anomalously large Single & Double Spin Asymmetries,
    - factorisation breakdown in forward  $J/\psi$  production in  $pA$

# Decisive advantages of Fixed-target experiments

- They brought essential contributions to particle & nuclear physics
  - particle discoveries ( $\Omega^-(sss)$ ,  $J/\psi$ ,  $\Upsilon$ ,...)
  - evidence for the novel dynamics of quarks and gluons in HIC (QGP)
  - observation of surprising QCD phenomena
    - breakdown of the Lam-Tung relation,
    - colour transparency,
    - higher-twist effects in forward meson production ,
    - anomalously large Single & Double Spin Asymmetries,
    - factorisation breakdown in forward  $J/\psi$  production in  $pA$
- Fixed-target experiments offer specific **advantages** that are still nowadays **difficult to challenge by collider experiments**

# Decisive advantages of Fixed-target experiments

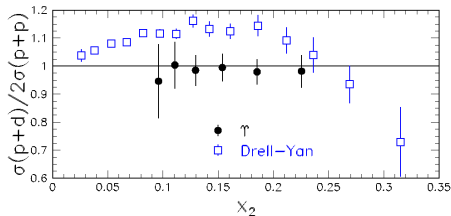
- They brought essential contributions to particle & nuclear physics
  - particle discoveries ( $\Omega^-(sss)$ ,  $J/\psi$ ,  $\Upsilon, \dots$ )
  - evidence for the novel dynamics of quarks and gluons in HIC (QGP)
  - observation of surprising QCD phenomena
    - breakdown of the Lam-Tung relation,
    - colour transparency,
    - higher-twist effects in forward meson production ,
    - anomalously large Single & Double Spin Asymmetries,
    - factorisation breakdown in forward  $J/\psi$  production in  $pA$
- Fixed-target experiments offer specific **advantages** that are still nowadays **difficult to challenge by collider experiments**
- They exhibit 4 decisive features,
  - accessing the **high** Feynman  $x_F$  domain ( $x_F \equiv p_z/p_{z\max}$ )
  - achieving **high luminosities** with dense targets,
  - **varying** the atomic mass of the **target** almost at will,
  - **polarising** the target.



## E866 at Fermilab with the Tevatron beam

– Precision  $\Upsilon$  studies in  $pp$  and  $pd$  collisions

E866 PRL 100 (2008) 062301

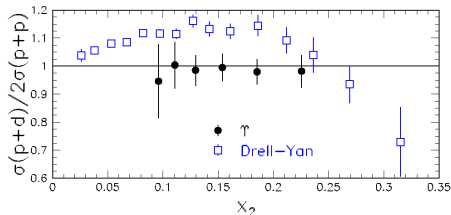


Precision: necessary to show a different behaviour from DY

## E866 at Fermilab with the Tevatron beam

– Precision  $\Upsilon$  studies in  $pp$  and  $pd$  collisions

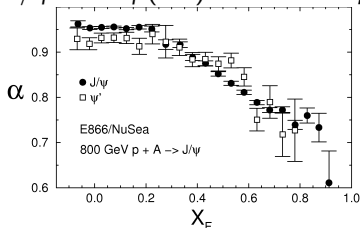
E866 PRL 100 (2008) 062301



Precision: necessary to show a different behaviour from DY

– Precision  $J/\psi$  and  $\psi(2S)$  studies in  $pA$  collisions

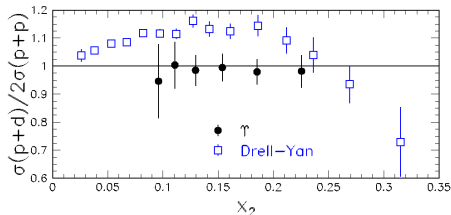
E866 PRL 84 (2000) 3256

Precision: necessary to show a different behaviour of  $\psi(2S)$  vs.  $J/\psi$

## E866 at Fermilab with the Tevatron beam

– Precision  $\Upsilon$  studies in  $pp$  and  $pd$  collisions

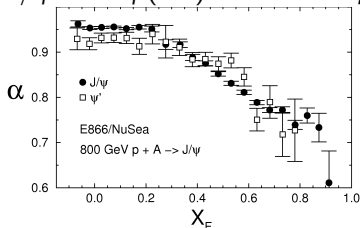
E866 PRL 100 (2008) 062301



Precision: necessary to show a different behaviour from DY

– Precision  $J/\psi$  and  $\psi(2S)$  studies in  $pA$  collisions

E866 PRL 84 (2000) 3256

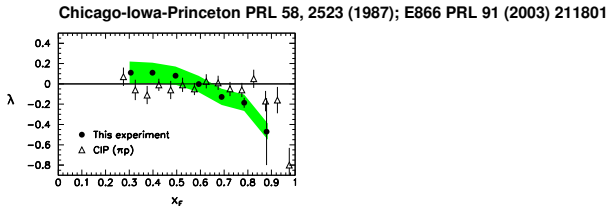


vs. 1 single preliminary  
 $\psi(2S)$  point at RHIC in  
 $dAu$  collisions

Precision: necessary to show a different behaviour of  $\psi(2S)$  vs.  $J/\psi$

## E866 and before

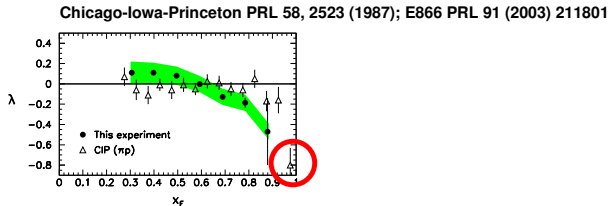
- **Precision  $J/\psi$  polarisation** (in the CS frame) **studies at large  $x_F$**



Precision and reach in  $x_F$ : necessary to show the change of pol. pattern

## E866 and before

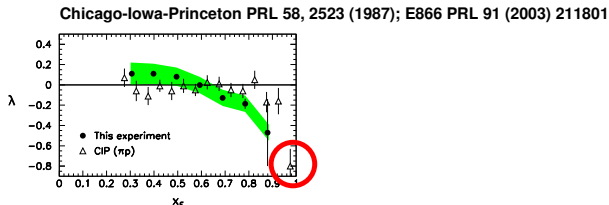
- **Precision**  $J/\psi$  polarisation (in the CS frame) studies at large  $x_F$



Precision and reach in  $x_F$ : necessary to show the change of pol. pattern

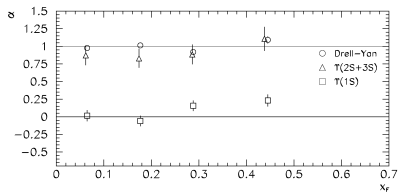
## E866 and before

- **Precision  $J/\psi$  polarisation** (in the CS frame) **studies at large  $x_F$**



Precision and reach in  $x_F$ : necessary to show the change of pol. pattern

- **Precision  $\Upsilon(nS)$  polarisation** (in the CS frame) **studies**

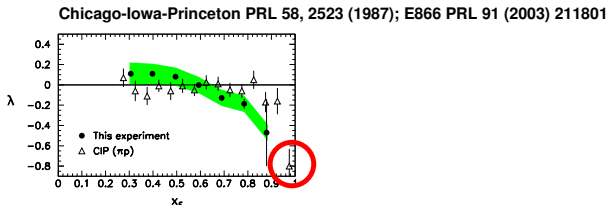


E866 PRL 86 2529 (2001); CMS PRL 110, 081802 (2013)

Precision: necessary to show the different polarisation pattern between 1S and 2S+3S

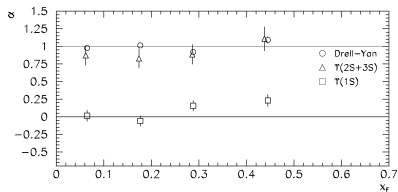
## E866 and before

- **Precision  $J/\psi$  polarisation** (in the CS frame) **studies at large  $x_F$**

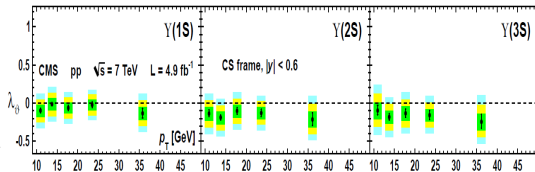


Precision and reach in  $x_F$ : necessary to show the change of pol. pattern

- **Precision  $\Upsilon(nS)$  polarisation** (in the CS frame) **studies**



E866 PRL 86 2529 (2001); CMS PRL 110, 081802 (2013)



Precision: necessary to show the different polarisation pattern between 1S and 2S+3S

# The European strategy for particle physics

Approved by the CERN council at the special Session held in Lisbon on 14 July, 2006



# The European strategy for particle physics

Approved by the CERN council at the special Session held in Lisbon on 14 July, 2006

9. A variety of important research lines are at the interface between particle and nuclear physics requiring dedicated experiments; *Council will seek to work with NuPECC in areas of mutual interest, and maintain the capability to perform fixed target experiments at CERN.*

pg. 37 of the Strategy Brochure

# The European strategy for particle physics

Approved by the CERN council at the special Session held in Lisbon on 14 July, 2006

9. A variety of important research lines are at the interface between particle and nuclear physics requiring dedicated experiments; *Council will seek to work with NuPECC in areas of mutual interest, and maintain the capability to perform **fixed target experiments at CERN.***

pg. 37 of the Strategy Brochure

# The European strategy for particle physics

Approved by the CERN council at the special Session held in Lisbon on 14 July, 2006

9. A variety of important research lines are at the interface between particle and nuclear physics requiring dedicated experiments; *Council will seek to work with NuPECC in areas of mutual interest, and maintain the capability to perform **fixed target experiments at CERN.***

pg. 37 of the Strategy Brochure

Using the LHC beams, for the first time,  
**the 100-GeV frontier can be broken at a fixed target experiment,**

# The European strategy for particle physics

Approved by the CERN council at the special Session held in Lisbon on 14 July, 2006

9. A variety of important research lines are at the interface between particle and nuclear physics requiring dedicated experiments; *Council will seek to work with NuPECC in areas of mutual interest, and maintain the capability to perform **fixed target experiments at CERN.***

pg. 37 of the Strategy Brochure

Using the LHC beams, for the first time,  
**the 100-GeV frontier can be broken at a fixed target experiment,**

- without affecting the LHC performance
- with an extracted beam line using a bent crystal

# The European strategy for particle physics

Approved by the CERN council at the special Session held in Lisbon on 14 July, 2006

9. A variety of important research lines are at the interface between particle and nuclear physics requiring dedicated experiments; *Council will seek to work with NuPECC in areas of mutual interest, and maintain the capability to perform **fixed target experiments at CERN.***

pg. 37 of the Strategy Brochure

Using the LHC beams, for the first time,  
**the 100-GeV frontier can be broken at a fixed target experiment,**

- without affecting the LHC performance
- with an extracted beam line using a bent crystal
- with the possibility of polarising the target
- without target-species limitation

# The European strategy for particle physics

Approved by the CERN council at the special Session held in Lisbon on 14 July, 2006

9. A variety of important research lines are at the interface between particle and nuclear physics requiring dedicated experiments; *Council will seek to work with NuPECC in areas of mutual interest, and maintain the capability to perform **fixed target experiments at CERN.***

pg. 37 of the Strategy Brochure

Using the LHC beams, for the first time,  
**the 100-GeV frontier can be broken at a fixed target experiment,**

- without affecting the LHC performance
- with an extracted beam line using a bent crystal
- with the possibility of polarising the target
- without target-species limitation
- with an outstanding luminosity
- with virtually no limit on particle-species studies (except top quark)

# The European strategy for particle physics

Approved by the CERN council at the special Session held in Lisbon on 14 July, 2006

9. A variety of important research lines are at the interface between particle and nuclear physics requiring dedicated experiments; *Council will seek to work with NuPECC in areas of mutual interest, and maintain the capability to perform **fixed target experiments at CERN.***

pg. 37 of the Strategy Brochure

Using the LHC beams, for the first time,  
**the 100-GeV frontier can be broken at a fixed target experiment,**

- without affecting the LHC performance
- with an extracted beam line using a bent crystal
- with the possibility of polarising the target
- without target-species limitation
- with an outstanding luminosity
- with virtually no limit on particle-species studies (except top quark)
- with modern detection techniques

## Part II

A fixed-target experiment using the LHC  
beam(s): AFTER@LHC



# Generalities

- $pp$  or  $pA$  collisions with a 7 TeV  $p^+$  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

# Generalities

- $pp$  or  $pA$  collisions with a 7 TeV  $p^+$  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode,  $\sqrt{s} = 2E_p$ , *i.e.* much larger

# Generalities

- $pp$  or  $pA$  collisions with a 7 TeV  $p^+$  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode,  $\sqrt{s} = 2E_p$ , *i.e.* much larger
- Benefit of the fixed target mode : **boost**:  $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$

# Generalities

- $pp$  or  $pA$  collisions with a 7 TeV  $p^+$  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode,  $\sqrt{s} = 2E_p$ , *i.e.* much larger
- Benefit of the fixed target mode : **boost**:  $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$ 
  - Consider a **photon emitted at  $90^\circ$**  w.r.t. the z-axis (beam) in the CM:  
( $p_{z,CM} = 0$ ,  $E_{CM}^\gamma = p_T$ )

# Generalities

- $pp$  or  $pA$  collisions with a 7 TeV  $p^+$  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode,  $\sqrt{s} = 2E_p$ , *i.e.* much larger

- Benefit of the fixed target mode : **boost**:  $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$

- Consider a **photon emitted at  $90^\circ$**  w.r.t. the z-axis (beam) in the CM:

$$\begin{pmatrix} E_{Lab} \\ p_{z,Lab} \end{pmatrix} = \begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix} \begin{pmatrix} p_T \\ 0 \end{pmatrix} \quad (p_{z,CM} = 0, E_{CM}^\gamma = p_T)$$

# Generalities

- $pp$  or  $pA$  collisions with a **7 TeV  $p^+$**  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode,  $\sqrt{s} = 2E_p$ , *i.e.* much larger
- Benefit of the fixed target mode : **boost**:  $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$ 
  - Consider a **photon emitted at  $90^\circ$**  w.r.t. the z-axis (beam) in the CM:
    - $$\begin{pmatrix} E_{Lab} \\ p_{z,Lab} \end{pmatrix} = \begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix} \begin{pmatrix} p_T \\ 0 \end{pmatrix} \quad (p_{z,CM} = 0, E_{CM}^\gamma = p_T)$$
    - $p_{z,Lab} \simeq 60p_T$  ! [A 67 MeV  $\gamma$  from a  $\pi^0$  at rest in the CM can easily be detected.]

# Generalities

- $pp$  or  $pA$  collisions with a **7 TeV  $p^+$**  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode,  $\sqrt{s} = 2E_p$ , *i.e.* much larger
- Benefit of the fixed target mode : **boost**:  $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$ 
  - Consider a **photon emitted at  $90^\circ$**  w.r.t. the z-axis (beam) in the CM:
    - $\begin{pmatrix} E_{Lab} \\ p_{z,Lab} \end{pmatrix} = \begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix} \begin{pmatrix} p_T \\ 0 \end{pmatrix}$  ( $p_{z,CM} = 0$ ,  $E_{CM}^\gamma = p_T$ )
    - $p_{z,Lab} \simeq 60p_T$  ! [A 67 MeV  $\gamma$  from a  $\pi^0$  at rest in the CM can easily be detected.]
- Angle in the Lab. frame:  $\tan \theta = \frac{p_T}{p_{z,Lab}} = \frac{1}{\gamma\beta} \Rightarrow \theta \simeq 1^\circ$ .  
[Rapidity shift:  $\Delta y = \tanh^{-1} \beta \simeq 4.8$ ]

# Generalities

- $pp$  or  $pA$  collisions with a **7 TeV  $p^+$**  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode,  $\sqrt{s} = 2E_p$ , *i.e.* much larger
- Benefit of the fixed target mode : **boost**:  $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$ 
  - Consider a **photon emitted at  $90^\circ$**  w.r.t. the z-axis (beam) in the CM:
    - $\begin{pmatrix} E_{Lab} \\ p_{z,Lab} \end{pmatrix} = \begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix} \begin{pmatrix} p_T \\ 0 \end{pmatrix}$  ( $p_{z,CM} = 0$ ,  $E_{CM}^\gamma = p_T$ )
    - $p_{z,Lab} \simeq 60p_T$  ! [A 67 MeV  $\gamma$  from a  $\pi^0$  at rest in the CM can easily be detected.]
- Angle in the Lab. frame:  $\tan \theta = \frac{p_T}{p_{z,Lab}} = \frac{1}{\gamma\beta} \Rightarrow \theta \simeq 1^\circ$ .  
[Rapidity shift:  $\Delta y = \tanh^{-1} \beta \simeq 4.8$ ]
- The entire forward CM hemisphere ( $y_{CM} > 0$ ) within  $0^\circ \leq \theta_{Lab} \leq 1^\circ$   
[ $y_{CM} = 0 \Rightarrow y_{Lab} \simeq 4.8$ ]



# Generalities

- $pp$  or  $pA$  collisions with a **7 TeV  $p^+$**  on a fixed target occur at a CM energy

$$\sqrt{s} = \sqrt{2m_N E_p} \simeq 115 \text{ GeV}$$

- In a symmetric collider mode,  $\sqrt{s} = 2E_p$ , *i.e.* much larger

- Benefit of the fixed target mode : **boost**:  $\gamma_{CM}^{Lab} = \frac{\sqrt{s}}{2m_p} \simeq 60$

- Consider a **photon emitted at  $90^\circ$**  w.r.t. the z-axis (beam) in the CM:

$$\begin{pmatrix} E_{Lab} \\ p_{z,Lab} \end{pmatrix} = \begin{pmatrix} \gamma & \gamma\beta \\ \gamma\beta & \gamma \end{pmatrix} \begin{pmatrix} p_T \\ 0 \end{pmatrix} \quad (p_{z,CM} = 0, E_{CM}^\gamma = p_T)$$

- $p_{z,Lab} \simeq 60p_T$  ! [A 67 MeV  $\gamma$  from a  $\pi^0$  at rest in the CM can easily be detected.]

- Angle in the Lab. frame:  $\tan \theta = \frac{p_T}{p_{z,Lab}} = \frac{1}{\gamma\beta} \Rightarrow \theta \simeq 1^\circ$ .

[Rapidity shift:  $\Delta y = \tanh^{-1} \beta \simeq 4.8$ ]

- The entire forward CM hemisphere ( $y_{CM} > 0$ ) within  $0^\circ \leq \theta_{Lab} \leq 1^\circ$

[ $y_{CM} = 0 \Rightarrow y_{Lab} \simeq 4.8$ ]

- **Good thing**: small forward detector  $\equiv$  large acceptance

- **Bad thing**: high multiplicity  $\Rightarrow$  absorber  $\Rightarrow$  physics limitation

# Backward physics ?

- Let's adopt a **novel strategy** and look at larger angles
  - particles with sufficient  $p_T$  to be detected
  - heavy particles whose decay product have enough  $p_T$  to be detected  
[not very heavy in fact:  $J/\psi \rightarrow \mu\mu$  or  $D \rightarrow K\pi$  are fine for current detectors]

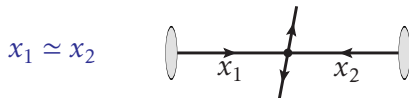
# Backward physics ?

- Let's adopt a **novel strategy** and look at larger angles
  - particles with sufficient  $p_T$  to be detected
  - heavy particles whose decay product have enough  $p_T$  to be detected  
 [not very heavy in fact:  $J/\psi \rightarrow \mu\mu$  or  $D \rightarrow K\pi$  are fine for current detectors]
- Advantages:
  - reduced multiplicities at large(r) angles
  - **access to partons with momentum fraction  $x \rightarrow 1$  in the target**
  - last, but not least, **no geometrical constrain** (e.g. beam pipe) **at  $\theta_{CM} \simeq 180^\circ$**

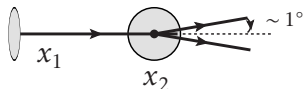
# Backward physics ?

- Let's adopt a **novel strategy** and look at larger angles
  - particles with sufficient  $p_T$  to be detected
  - heavy particles whose decay product have enough  $p_T$  to be detected  
[not very heavy in fact:  $J/\psi \rightarrow \mu\mu$  or  $D \rightarrow K\pi$  are fine for current detectors]
- Advantages:
  - reduced multiplicities at large(r) angles
  - access to partons with momentum fraction  $x \rightarrow 1$  in the target**
  - last, but not least, **no geometrical constrain** (e.g. beam pipe) **at  $\theta_{CM} \simeq 180^\circ$**

Hadron center-of-mass system



Target rest frame

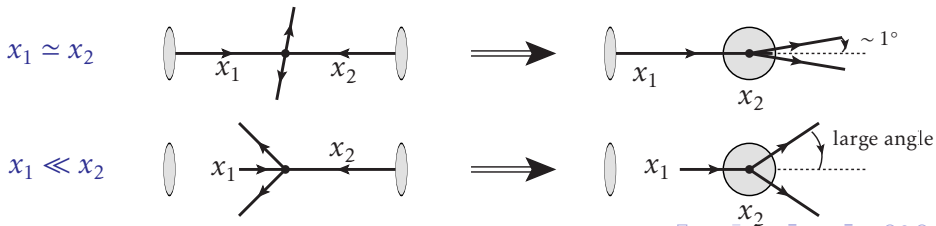


# Backward physics ?

- Let's adopt a **novel strategy** and look at larger angles
  - particles with sufficient  $p_T$  to be detected
  - heavy particles whose decay product have enough  $p_T$  to be detected  
[not very heavy in fact:  $J/\psi \rightarrow \mu\mu$  or  $D \rightarrow K\pi$  are fine for current detectors]
- Advantages:
  - reduced multiplicities at large(r) angles
  - access to partons with momentum fraction  $x \rightarrow 1$  in the target**
  - last, but not least, **no geometrical constrain** (e.g. beam pipe) **at  $\theta_{CM} \simeq 180^\circ$**

Hadron center-of-mass system

Target rest frame

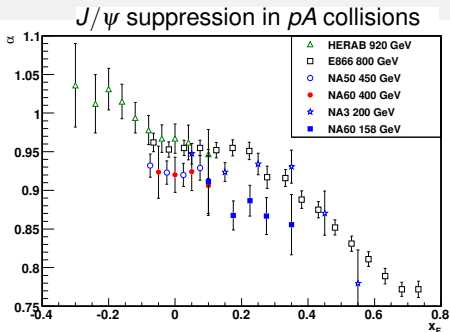


# First systematic access to the target-rapidity region

( $x_F \rightarrow -1$ )

# First systematic access to the target-rapidity region

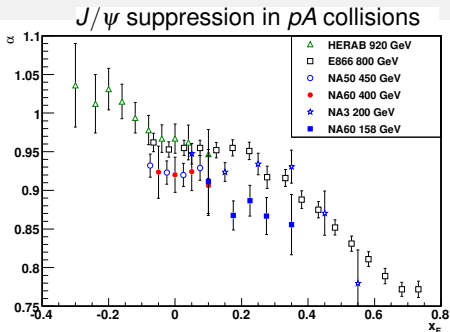
( $x_F \rightarrow -1$ )



- $x_F$  systematically studied at fixed target experiments up to +1

# First systematic access to the target-rapidity region

( $x_F \rightarrow -1$ )



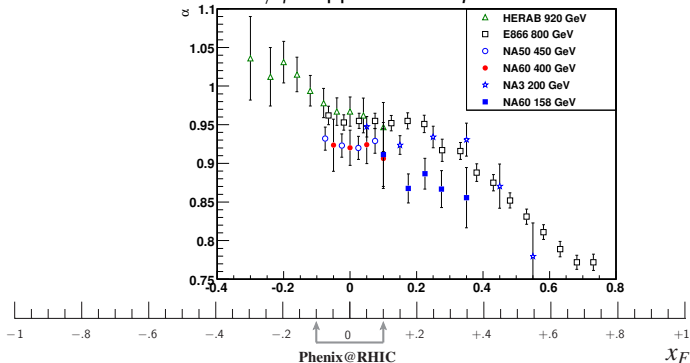
- $x_F$  systematically studied at fixed target experiments up to +1
- Hera-B was the only one to really explore  $x_F < 0$ , up to -0.3



# First systematic access to the target-rapidity region

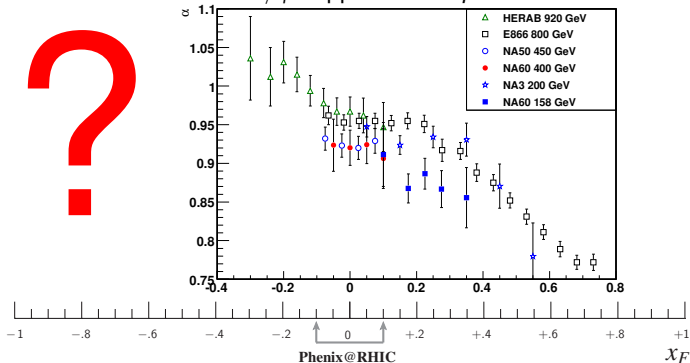
( $x_F \rightarrow -1$ )

## $J/\psi$ suppression in $pA$ collisions



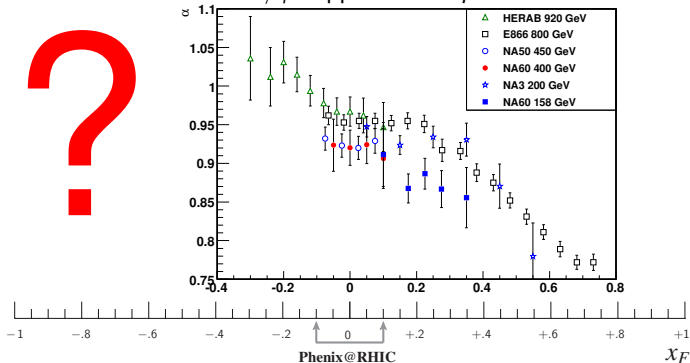
- $x_F$  systematically studied at fixed target experiments **up to +1**
- Hera-B was the only one to really explore  $x_F < 0$ , up to -0.3
- PHENIX @ RHIC:  $-0.1 < x_F < 0.1$  [could be wider with  $\Upsilon$ , but low stat.]
- CMS/ATLAS:  $|x_F| < 5 \cdot 10^{-3}$ ; LHCb:  $5 \cdot 10^{-3} < x_F < 4 \cdot 10^{-2}$

## First systematic access to the target-rapidity region

 $(x_F \rightarrow -1)$  $J/\psi$  suppression in  $pA$  collisions

- $x_F$  systematically studied at fixed target experiments **up to +1**
- Hera-B was the only one to really explore  $x_F < 0$ , up to -0.3
- PHENIX @ RHIC:  $-0.1 < x_F < 0.1$  [could be wider with  $\Upsilon$ , but low stat.]
- CMS/ATLAS:  $|x_F| < 5 \cdot 10^{-3}$ ; LHCb:  $5 \cdot 10^{-3} < x_F < 4 \cdot 10^{-2}$

## First systematic access to the target-rapidity region

 $(x_F \rightarrow -1)$  $J/\psi$  suppression in  $pA$  collisions

- $x_F$  systematically studied at fixed target experiments **up to +1**
- Hera-B was the only one to really explore  $x_F < 0$ , up to -0.3
- PHENIX @ RHIC:  $-0.1 < x_F < 0.1$  [could be wider with  $\Upsilon$ , but low stat.]
- CMS/ATLAS:  $|x_F| < 5 \cdot 10^{-3}$ ; LHCb:  $5 \cdot 10^{-3} < x_F < 4 \cdot 10^{-2}$
- If we measure  $\Upsilon(b\bar{b})$  at  $y_{\text{cms}} \simeq -2.5 \Rightarrow x_F \simeq \frac{2m_\Upsilon}{\sqrt{s}} \sinh(y_{\text{cms}}) \simeq -1$

# The lead-ion beam

- Design LHC lead-beam energy: **2.76 TeV** per nucleon

# The lead-ion beam

- Design LHC lead-beam energy: **2.76 TeV** per nucleon
- In the fixed target mode, PbA collisions at  $\sqrt{s_{NN}} \simeq$  **72 GeV**

# The lead-ion beam

- Design LHC lead-beam energy: **2.76 TeV** per nucleon
- In the fixed target mode, PbA collisions at  $\sqrt{s_{NN}} \simeq 72 \text{ GeV}$
- Half way **between BNL-RHIC** (AuAu, CuCu @ **200 GeV**) and **CERN-SPS** (PbPb @ **17.2 GeV**)

# The lead-ion beam

- Design LHC lead-beam energy: **2.76 TeV** per nucleon
- In the fixed target mode, PbA collisions at  $\sqrt{s_{NN}} \simeq 72 \text{ GeV}$
- Half way **between BNL-RHIC** (AuAu, CuCu @ 200 GeV) and **CERN-SPS** (PbPb @ 17.2 GeV)
- Example of motivations:

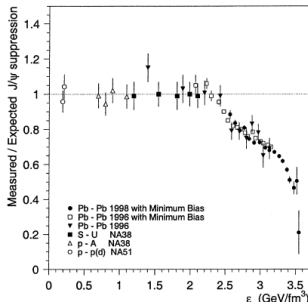


Fig. 7. Measured  $J/\psi$  production yields, normalised to the yields expected assuming that the only source of suppression is the ordinary absorption by the nuclear medium. The data is shown as a function of the energy density reached in the several collision systems.

# The lead-ion beam

- Design LHC lead-beam energy: **2.76 TeV** per nucleon
- In the fixed target mode, PbA collisions at  $\sqrt{s_{NN}} \simeq 72 \text{ GeV}$
- Half way **between BNL-RHIC** (AuAu, CuCu @ **200 GeV**) and **CERN-SPS** (PbPb @ **17.2 GeV**)
- Example of motivations:

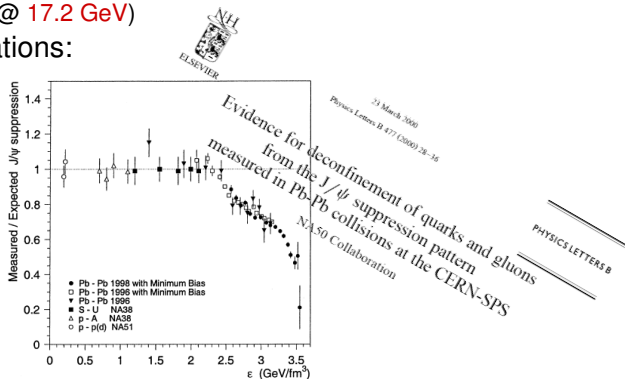


Fig. 7. Measured  $J/\psi$  production yields, normalised to the yields expected assuming that the only source of suppression is the ordinary absorption by the nuclear medium. The data is shown as a function of the energy density reached in the several collision systems.



# The beam extraction

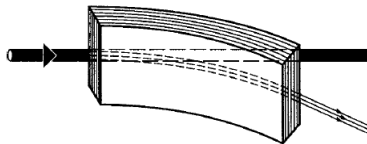
- ★ The LHC beam may be extracted using “Strong crystalline field”  
**without any decrease in performance of the LHC !**

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131

# The beam extraction

- ★ The LHC beam may be extracted using “Strong crystalline field”  
**without any decrease in performance of the LHC !**

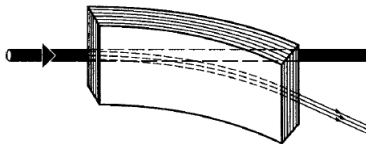
E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131



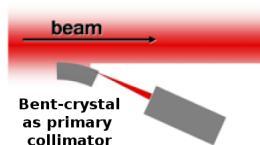
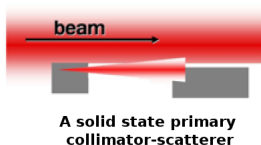
# The beam extraction

- ★ The LHC beam may be extracted using “Strong crystalline field”  
**without any decrease in performance of the LHC !**

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131



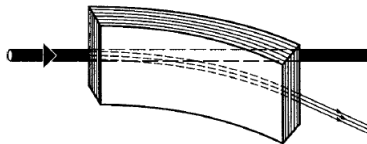
- ★ **Illustration for collimation**



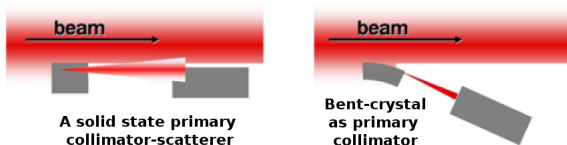
# The beam extraction

- ★ The LHC beam may be extracted using “Strong crystalline field”  
**without any decrease in performance of the LHC !**

E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31, Rev. Mod. Phys. 77 (2005) 1131



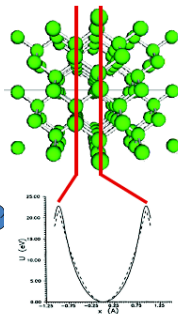
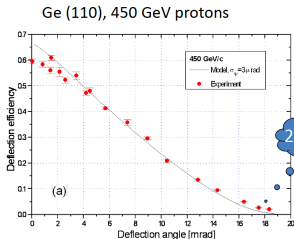
- ★ **Illustration for collimation**



- ★ **Tests** will be performed on the **LHC beam**:  
LUA9 proposal approved by the LHCC

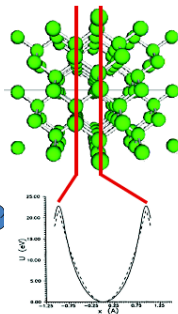
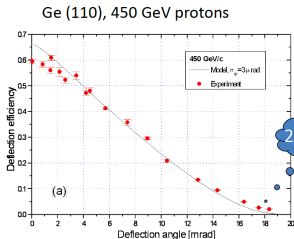
# The beam extraction

- Inter-crystalline fields are huge



# The beam extraction

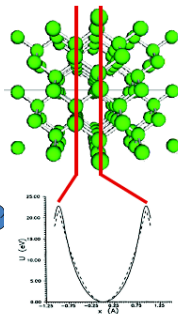
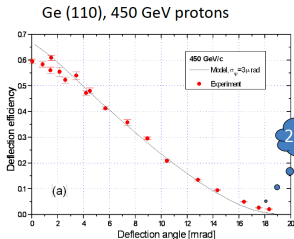
- Inter-crystalline fields are huge



- The **channeling efficiency** is high for a deflection of a few mrad

# The beam extraction

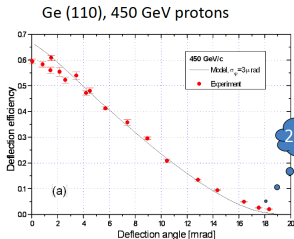
- Inter-crystalline fields are huge



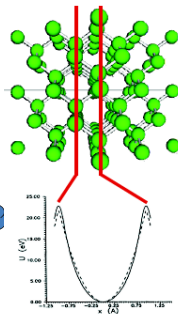
- The **channeling efficiency** is high for a deflection of a few mrad
- One can **extract** a significant part of the **beam loss** ( $10^9 p^+ s^{-1}$ )

# The beam extraction

- Inter-crystalline fields are huge



2000 T !



- The **channeling efficiency** is high for a deflection of a few mrad
- One can **extract** a significant part of the **beam loss** ( $10^9 p^+ s^{-1}$ )
- Simple and robust way to extract the most energetic beam ever:





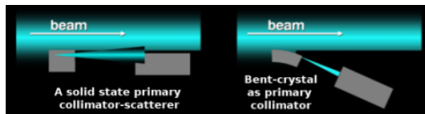
# The beam extraction: news

[S. Montesano, *Physics at AFTER using LHC beams, ECT\* Trento, Feb. 2013*]

Goal : assess the possibility to use bent crystals as primary collimators in hadronic accelerators and colliders



UA9 installation in the SPS



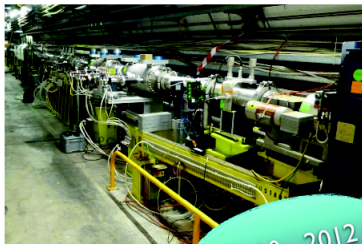
Prototype crystal collimation system at SPS :

- local beam loss reduction (5÷20x reduction for proton beam)
- beam loss map show average loss reduction in the entire SPS ring
- halo extraction efficiency  
70÷80% for protons (50÷70% for Pb)

# The beam extraction: news

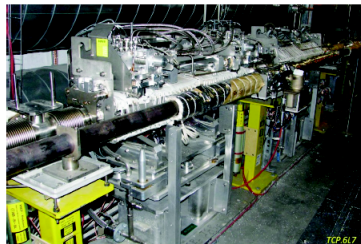
[S. Montesano, *Physics at AFTER using LHC beams, ECT\* Trento, Feb. 2013*]

Goal : assess the possibility to **use bent crystals as primary collimators** in hadronic accelerators and colliders



UA9 installation in the SPS

2010 - 2012



LUA9 future installation in LHC

Prototype crystal collimation system at SPS :

- local **beam loss reduction** (5+20x reduction for proton beam)
- beam loss map show average loss reduction in the entire SPS ring
- **halo extraction efficiency**  
70+80% for protons (50+70% for Pb)

# The beam extraction: news

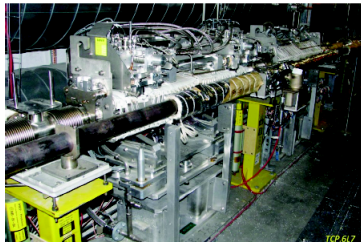
[S. Montesano, *Physics at AFTER using LHC beams, ECT\* Trento, Feb. 2013*]

Goal : assess the possibility to **use bent crystals as primary collimators** in hadronic accelerators and colliders



UA9 installation in the SPS

2010 - 2012



LUA9 future installation in LHC

Prototype crystal collimation system at SPS :

- local **beam loss reduction** (5+20x reduction for proton beam)
- beam loss map show average loss reduction in the entire SPS ring
- **halo extraction efficiency**  
70+80% for protons (50+70% for Pb)

Towards an installation in the LHC : propose and **install during LSI** a min. number of devices

- 2 crystals

Long term plan is ambitious : **propose a collimation system based on bent crystals** for the upgrade of the current LHC collimation system

# Luminosities

- Expected proton flux  $\Phi_{beam} = 5 \times 10^8 p^+ s^{-1}$

# Luminosities

- Expected **proton flux**  $\Phi_{beam} = 5 \times 10^8 \text{ p}^+ \text{ s}^{-1}$
- Instantaneous **Luminosity**:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

[  $\ell$ : target thickness (for instance 1cm)]

# Luminosities

- Expected **proton flux**  $\Phi_{beam} = 5 \times 10^8 p^+ s^{-1}$
- Instantaneous **Luminosity**:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

[  $\ell$ : target thickness (for instance 1cm)]

- Integrated luminosity:  $\int dt \mathcal{L}$  over  $10^7$  s for  $p^+$  and  $10^6$  for Pb

[the so-called LHC years]

# Luminosities

- Expected **proton flux**  $\Phi_{beam} = 5 \times 10^8 p^+ s^{-1}$
- Instantaneous **Luminosity**:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times \ell \times \mathcal{N}_A) / A$$

[  $\ell$ : target thickness (for instance 1cm)]

- Integrated luminosity:  $\int dt \mathcal{L}$  over  $10^7$  s for  $p^+$  and  $10^6$  for Pb

[the so-called LHC years]

Target	$\rho$ (g.cm <sup>-3</sup> )	A	$\mathcal{L}$ ( $\mu\text{b}^{-1}.\text{s}^{-1}$ )	$\int \mathcal{L}$ (pb <sup>-1</sup> .yr <sup>-1</sup> )
Sol. H <sub>2</sub>	0.09	1	<b>26</b>	<b>260</b>
Liq. H <sub>2</sub>	0.07	1	<b>20</b>	<b>200</b>
Liq. D <sub>2</sub>	0.16	2	<b>24</b>	<b>240</b>
Be	1.85	9	<b>62</b>	<b>620</b>
Cu	8.96	64	<b>42</b>	<b>420</b>
W	19.1	185	<b>31</b>	<b>310</b>
Pb	11.35	207	<b>16</b>	<b>160</b>

# Luminosities

- 1 meter-long liquid  $H_2$  &  $D_2$  targets can be used (see NA51, ...)



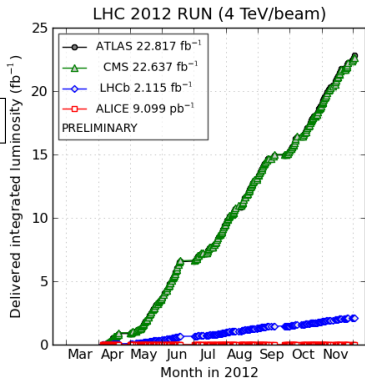
# Luminosities

- 1 meter-long liquid  $H_2$  &  $D_2$  targets can be used (see NA51, ...)
- This gives:  $\mathcal{L}_{H_2/D_2} \simeq 20 \text{ fb}^{-1} \text{ y}^{-1}$

# Luminosities

- 1 meter-long liquid  $H_2$  &  $D_2$  targets can be used (see NA51, ...)
- This gives:  $\mathcal{L}_{H_2/D_2} \simeq 20 \text{ fb}^{-1} \text{ y}^{-1}$
- Recycling the LHC beam loss, one gets

a luminosity comparable to the LHC itself !



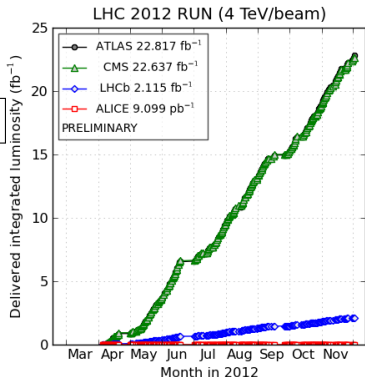
(generated 2012-12-02 18:23 including fill 3360)

# Luminosities

- 1 meter-long liquid  $H_2$  &  $D_2$  targets can be used (see NA51, ...)
- This gives:  $\mathcal{L}_{H_2/D_2} \simeq 20 \text{ fb}^{-1} \text{ y}^{-1}$
- Recycling the LHC beam loss, one gets

**a luminosity comparable to the LHC itself !**

- PHENIX lumi in their decadal plan
  - Run14pp  $12 \text{ pb}^{-1}$  @  $\sqrt{s_{NN}} = 200 \text{ GeV}$
  - Run14dAu  $0.15 \text{ pb}^{-1}$  @  $\sqrt{s_{NN}} = 200 \text{ GeV}$



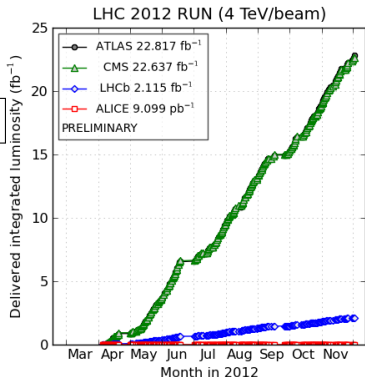
(generated 2012-12-02 18:23 including fill 3360)

# Luminosities

- 1 meter-long liquid  $H_2$  &  $D_2$  targets can be used (see NA51, ...)
- This gives:  $\mathcal{L}_{H_2/D_2} \simeq 20 \text{ fb}^{-1} \text{ y}^{-1}$
- Recycling the LHC beam loss, one gets

**a luminosity comparable to the LHC itself !**

- PHENIX lumi in their decadal plan
  - Run14pp  $12 \text{ pb}^{-1}$  @  $\sqrt{s_{NN}} = 200 \text{ GeV}$
  - Run14dAu  $0.15 \text{ pb}^{-1}$  @  $\sqrt{s_{NN}} = 200 \text{ GeV}$
- AFTER vs PHENIX@RHIC:  
3 orders of magnitude larger



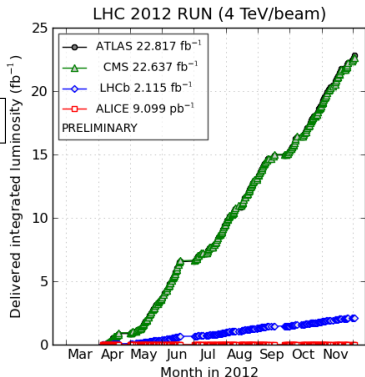
(generated 2012-12-02 18:23 including fill 3360)

# Luminosities

- 1 meter-long liquid  $H_2$  &  $D_2$  targets can be used (see NA51, ...)
- This gives:  $\mathcal{L}_{H_2/D_2} \simeq 20 \text{ fb}^{-1} \text{ y}^{-1}$
- Recycling the LHC beam loss, one gets

**a luminosity comparable to the LHC itself !**

- PHENIX lumi in their decadal plan
  - Run14pp  $12 \text{ pb}^{-1}$  @  $\sqrt{s_{NN}} = 200 \text{ GeV}$
  - Run14dAu  $0.15 \text{ pb}^{-1}$  @  $\sqrt{s_{NN}} = 200 \text{ GeV}$
- AFTER vs PHENIX@RHIC:  
3 orders of magnitude larger
- Lumi for Pb runs in the backup slides  
(roughly 10 times that planned for the LHC)



# Luminosities

- Instantaneous Luminosity:

$$\mathcal{L} = \Phi_{beam} \times N_{target} = N_{beam} \times (\rho \times l \times \mathcal{N}_A) / A$$

$$\Phi_{beam} = 2 \times 10^5 \text{ Pb s}^{-1}, \quad l = 1 \text{ cm (target thickness)}$$

- Integrated luminosity  $\int dt \mathcal{L} = \mathcal{L} \times 10^6 \text{ s}$  for Pb
- Expected luminosities with  $2 \times 10^5 \text{ Pb s}^{-1}$  extracted (1cm-long target)

Target	$\rho$ (g.cm <sup>-3</sup> )	A	$\mathcal{L}$ (mb <sup>-1</sup> .s <sup>-1</sup> ) = $\int \mathcal{L}$ (nb <sup>-1</sup> .yr <sup>-1</sup> )
Sol. H <sub>2</sub>	0.09	1	<b>11</b>
Liq. H <sub>2</sub>	0.07	1	<b>8</b>
Liq. D <sub>2</sub>	0.16	2	<b>10</b>
Be	1.85	9	<b>25</b>
Cu	8.96	64	<b>17</b>
W	19.1	185	<b>13</b>
Pb	11.35	207	<b>7</b>

- Planned lumi for PHENIX Run15AuAu 2.8 nb<sup>-1</sup> (0.13 nb<sup>-1</sup> at 62 GeV)
- Nominal LHC lumi for PbPb 0.5 nb<sup>-1</sup>

## A few figures on the (extracted) proton beam

- Beam loss:  $10^9 p^+s^{-1}$
- Extracted intensity:  $5 \times 10^8 p^+s^{-1}$  (1/2 the beam loss) E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31

## A few figures on the (extracted) proton beam

- Beam loss:  $10^9 p^+ s^{-1}$
- Extracted intensity:  $5 \times 10^8 p^+ s^{-1}$  (1/2 the beam loss) E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31
- Number of  $p^+$ : 2808 bunches of  $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$



## A few figures on the (extracted) proton beam

- Beam loss:  $10^9 p^+ s^{-1}$
- Extracted intensity:  $5 \times 10^8 p^+ s^{-1}$  (1/2 the beam loss) E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31
- Number of  $p^+$ : 2808 bunches of  $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$
- Revolution frequency: Each bunch passes the extraction point at a rate of  $3 \cdot 10^5 \text{ km} \cdot \text{s}^{-1} / 27 \text{ km} \simeq 11 \text{ kHz}$

## A few figures on the (extracted) proton beam

- Beam loss:  $10^9 p^+ s^{-1}$
- Extracted intensity:  $5 \times 10^8 p^+ s^{-1}$  (1/2 the beam loss) E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31
- Number of  $p^+$ : 2808 bunches of  $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$
- Revolution frequency: Each bunch passes the extraction point at a rate of  $3.10^5 \text{ km.s}^{-1} / 27 \text{ km} \simeq 11 \text{ kHz}$
- Extracted “mini” bunches:
  - the crystal sees  $2808 \times 11000 \text{ s}^{-1} \simeq 3.10^7 \text{ bunches s}^{-1}$
  - one extracts  $5.10^8 / 3.10^7 \simeq 15 p^+$  from each bunch at each pass
  - Provided that the probability of interaction with the target is below 5%,  
no pile-up...

## A few figures on the (extracted) proton beam

- Beam loss:  $10^9 p^+ s^{-1}$
- Extracted intensity:  $5 \times 10^8 p^+ s^{-1}$  (1/2 the beam loss) E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31
- Number of  $p^+$ : 2808 bunches of  $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$
- Revolution frequency: Each bunch passes the extraction point at a rate of  $3.10^5 \text{ km.s}^{-1} / 27 \text{ km} \simeq 11 \text{ kHz}$
- Extracted “mini” bunches:
  - the crystal sees  $2808 \times 11000 \text{ s}^{-1} \simeq 3.10^7 \text{ bunches s}^{-1}$
  - one extracts  $5.10^8 / 3.10^7 \simeq 15 p^+$  from each bunch at each pass
  - Provided that the probability of interaction with the target is below 5%,  
no pile-up...
- Extraction over a 10h fill:
  - $5 \times 10^8 p^+ \times 3600 \text{ s h}^{-1} \times 10 \text{ h} = 1.8 \times 10^{13} p^+ \text{ fill}^{-1}$
  - This means  $1.8 \times 10^{13} / 3.2 \times 10^{14} \simeq 5.6\%$  of the  $p^+$  in the beam  
*These protons are lost anyway !*

## A few figures on the (extracted) proton beam

- Beam loss:  $10^9 p^+ s^{-1}$
- Extracted intensity:  $5 \times 10^8 p^+ s^{-1}$  (1/2 the beam loss) E. Uggerhøj, U.I Uggerhøj, NIM B 234 (2005) 31
- Number of  $p^+$ : 2808 bunches of  $1.15 \times 10^{11} p^+ = 3.2 \times 10^{14} p^+$
- Revolution frequency: Each bunch passes the extraction point at a rate of  $3.10^5 \text{ km.s}^{-1} / 27 \text{ km} \simeq 11 \text{ kHz}$
- Extracted “mini” bunches:
  - the crystal sees  $2808 \times 11000 \text{ s}^{-1} \simeq 3.10^7 \text{ bunches s}^{-1}$
  - one extracts  $5.10^8 / 3.10^7 \simeq 15 p^+$  from each bunch at each pass
  - Provided that the probability of interaction with the target is below 5%,  
no pile-up...
- Extraction over a 10h fill:
  - $5 \times 10^8 p^+ \times 3600 \text{ s h}^{-1} \times 10 \text{ h} = 1.8 \times 10^{13} p^+ \text{ fill}^{-1}$
  - This means  $1.8 \times 10^{13} / 3.2 \times 10^{14} \simeq 5.6\%$  of the  $p^+$  in the beam  
*These protons are lost anyway !*
- similar figures for the Pb-beam extraction

# Part III

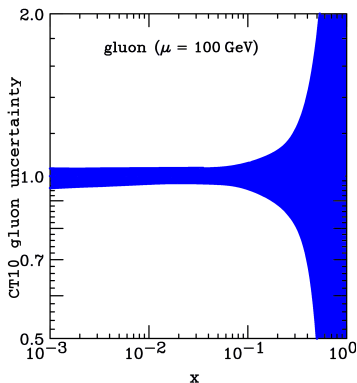
## AFTER: flagships measurements

# Key studies: gluons in the proton

- **Gluon distribution** at mid, high and ultra-high  $x_B$  in the proton

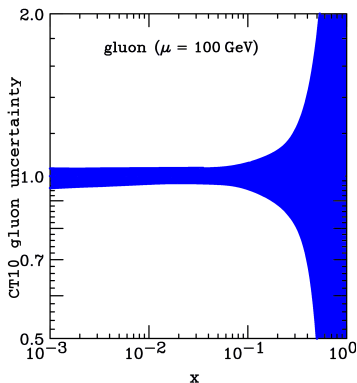
# Key studies: gluons in the proton

- **Gluon distribution** at mid, high and ultra-high  $x_B$  in the proton
  - Not easily accessible in DIS



# Key studies: gluons in the proton

- **Gluon distribution** at mid, high and ultra-high  $x_B$  in the proton
  - Not easily accessible in DIS
  - Very large uncertainties

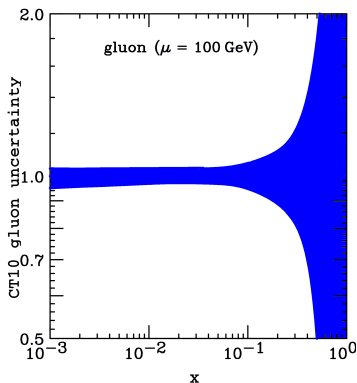




# Key studies: gluons in the proton

- **Gluon distribution** at mid, high and ultra-high  $x_B$  in the proton
  - Not easily accessible in DIS
  - Very large uncertainties

Accessible thanks gluon sensitive probes,

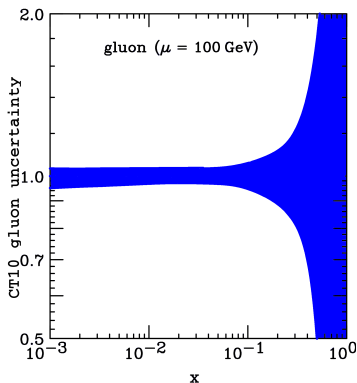


# Key studies: gluons in the proton

- **Gluon distribution** at mid, high and ultra-high  $x_B$  in the proton
  - Not easily accessible in DIS
  - Very large uncertainties

Accessible thanks gluon sensitive probes,

- **quarkonia**

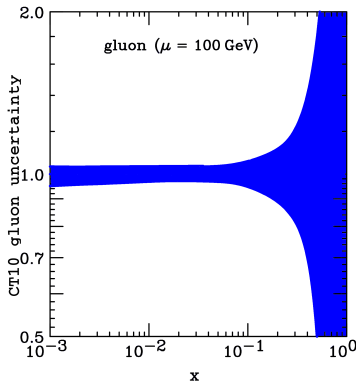


# Key studies: gluons in the proton

- **Gluon distribution** at mid, high and ultra-high  $x_B$  in the proton
  - Not easily accessible in DIS
  - Very large uncertainties

Accessible thanks gluon sensitive probes,

- **quarkonia**
- Isolated **photon**

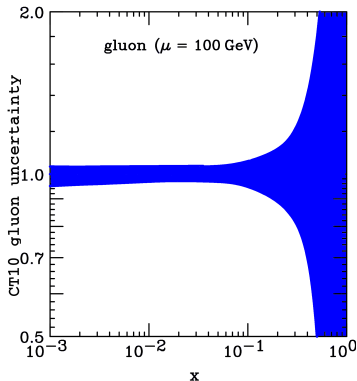


# Key studies: gluons in the proton

- **Gluon distribution** at mid, high and ultra-high  $x_B$  in the proton
  - Not easily accessible in DIS
  - Very large uncertainties

Accessible thanks gluon sensitive probes,

- **quarkonia**
- Isolated **photon**
- **jets** (  $P_T \in [20, 40]$  GeV)



# Isolated- $\gamma$ in p(7 TeV)-p(rest): $\sqrt{s} \sim 115$ GeV

- p-p photon kinematics at fixed-target LHC (central rapidities):  
To access  $x > 0.3$  one needs isolated- $\gamma$  at:  $p_T = x_T \sqrt{s}/2 > 20$  GeV/c
- JETPHOX NLO  
pQCD calculations:

p-p at  $\sqrt{s}=115$  GeV

$|y| < 0.5$ ,  $p_T > 20$  GeV/c

Isolation:  $R=0.4$ ,  $E_T^{\text{had}} < 5$  GeV

$\mathcal{L}$  (10 cm  $H_2$ -target)  $\sim 2 \cdot 10^3$  pb $^{-1}$ /year

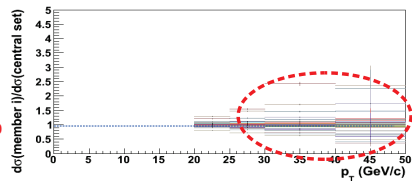
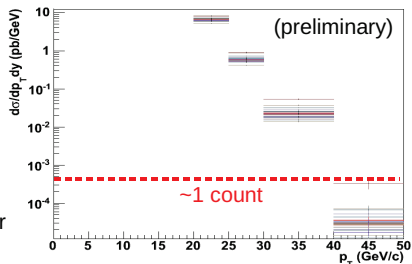
PDF: CT10 52 eigenval. (90% CL)

Scales:  $\mu_i = p_T$

FF = BFG-II

x-section uncertainties<sup>(\*)</sup> of  $\pm 150\%$

<sup>(\*)</sup> (68%CL)/(90% CL)  $\sim 1.65$



# Isolated- $\gamma$ in p(7 TeV)-p(rest): $\sqrt{s} \sim 115$ GeV

- p-p photon kinematics at fixed-target LHC (backwards rapidities):  
To access  $x > 0.3$  one needs isolated- $\gamma$  at:  $p_T = x_T \sqrt{s} / 2e^y > 10$  GeV/c
- JETPHOX NLO  
pQCD calculations:

p-p at  $\sqrt{s}=115$  GeV

$0 < y < -3.$ ,  $p_T > 20$  GeV/c

Isolation:  $R=0.4$ ,  $E_T^{\text{had}} < 5$  GeV

$\mathcal{L}$  (10 cm  $H_2$ -target)  $\sim 2 \cdot 10^3$  pb $^{-1}$ /year

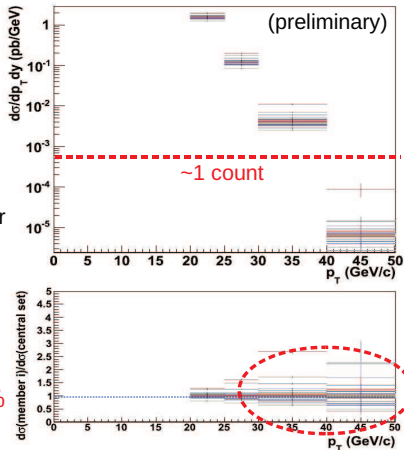
PDF: CT10 52 eigenval. (90% CL)

Scales:  $\mu_i = p_T$

FF = BFG-II

x-section uncertainties<sup>(\*)</sup> of  $\pm 170\%$

<sup>(\*)</sup> (68%CL)/(90% CL)  $\sim 1.65$



# Accessing the large $x$ glue with quarkonia

PYTHIA simulation  
 $\sigma(y) / \sigma(y=0.4)$   
 statistics for one month  
 5% acceptance considered

Statistical relative uncertainty  
 Large statistics allow to access  
 very backward region

Gluon uncertainty from  
 MSTWPDF  
 - only for the gluon content of  
 the target  
 - assuming

$$x_g = M_{J/\psi} / \sqrt{s} e^{-y_{CM}}$$

$J/\psi$

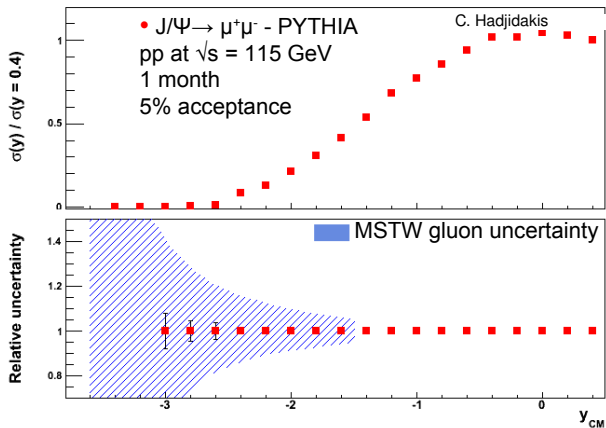
$$y_{CM} \sim 0 \rightarrow x_g = 0.03$$

$$y_{CM} \sim -3.6 \rightarrow x_g = 1$$

$Y$ : larger  $x_g$  for same  $y_{CM}$

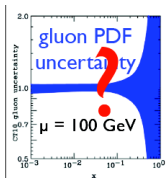
$$y_{CM} \sim 0 \rightarrow x_g = 0.08$$

$$y_{CM} \sim -2.4 \rightarrow x_g = 1$$



⇒ Backward measurements allow to access large  $x$  gluon pdf

# Key studies: gluons in the neutron



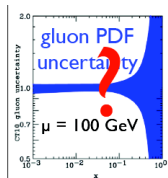
gluon PDF unknown for neutron

exp. probes :

- ▶ heavy quarkonia
- ▶ isolated photons
- ▶ high  $p_T$  jets



## Key studies: gluons in the neutron

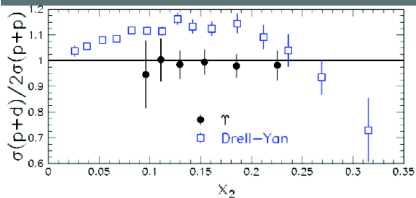


gluon PDF unknown for neutron

exp. probes :

- ▶ heavy quarkonia
- ▶ isolated photons
- ▶ high  $p_T$  jets

[ E866, PRL 100 (2008) 062301 ]



Pioneering measurement by E866 @ Fermilab :

- ▶ using  $\Upsilon$
- ▶ at  $Q^2 \sim 100 \text{ GeV}^2$  similar gluon distribution in proton and neutron

could be extended using  $J/\psi$  :

- ▶ to ( $\sim 10x$ ) lower  $x$
- ▶ to lower  $Q^2$

Need high luminosity.

[ Lansberg et al., FBS 53 (2012) 11 ]

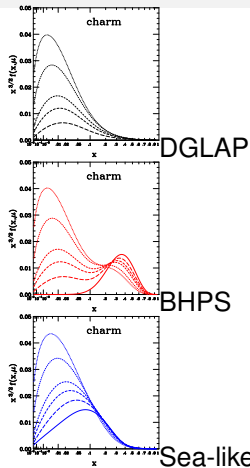
target	yearly lumi( $\text{fb}^{-1}$ )	$B_{ll} \frac{dN_{J/\psi}}{dy} \Big _{y=0}$	$B_{ll} \frac{dN_{\Upsilon}}{dy} \Big _{y=0}$
l m Liq. H <sub>2</sub>	20	$4.0 \cdot 10^8$	$8.0 \cdot 10^5$
l m Lid. D <sub>2</sub>	24	$9.6 \cdot 10^8$	$1.9 \cdot 10^6$

# Key studies

- Heavy-quark distributions (at high  $x_B$ )

# Key studies

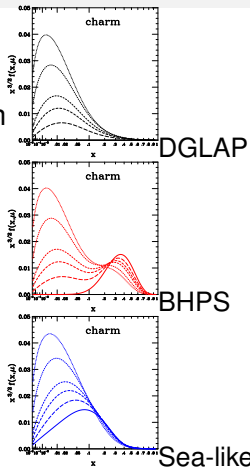
- Heavy-quark distributions (at high  $x_B$ )
  - Pin down **intrinsic charm**, ... at last



All 3 compatible  
with DIS data  
(Pumplin *et al.*)

# Key studies

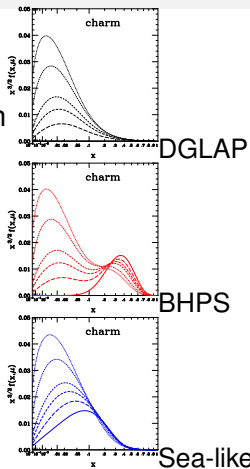
- Heavy-quark distributions (at high  $x_B$ )
  - Pin down **intrinsic charm**, ... at last
  - **Total open charm and beauty** cross section (aim: down to  $P_T \rightarrow 0$ )



All 3 compatible  
with DIS data  
(Pumplin *et al.*)

# Key studies

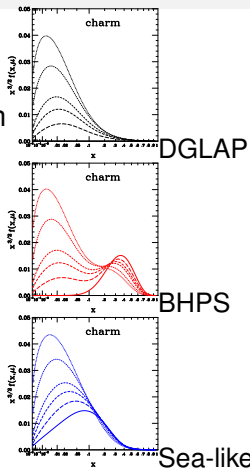
- Heavy-quark distributions (at high  $x_B$ )
    - Pin down **intrinsic charm**, ... at last
    - **Total open charm and beauty** cross section (aim: down to  $P_T \rightarrow 0$ )
- requires



All 3 compatible  
with DIS data  
(Pumplin *et al.*)

# Key studies

- **Heavy-quark** distributions (at high  $x_B$ )
    - Pin down **intrinsic charm**, ... at last
    - **Total open charm and beauty** cross section (aim: down to  $P_T \rightarrow 0$ )
- requires
- several **complementary** measurements



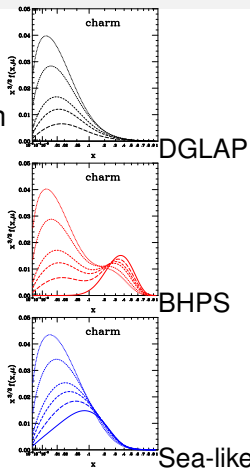
All 3 compatible  
with DIS data  
(Pumplin *et al.*)

# Key studies

- **Heavy-quark** distributions (at high  $x_B$ )
  - Pin down **intrinsic charm**, ... at last
  - **Total open charm and beauty** cross section (aim: down to  $P_T \rightarrow 0$ )

requires

- several **complementary** measurements
- good coverage in the **target-rapidity region**



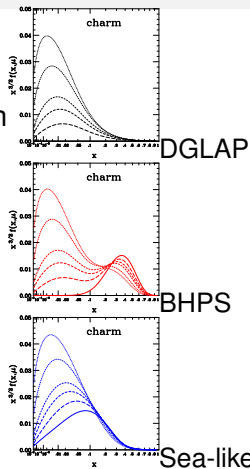
All 3 compatible  
with DIS data  
(Pumplin *et al.*)

# Key studies

- **Heavy-quark** distributions (at high  $x_B$ )
  - Pin down **intrinsic charm**, ... at last
  - **Total open charm and beauty** cross section (aim: down to  $P_T \rightarrow 0$ )

requires

- several **complementary** measurements
- good coverage in the **target-rapidity region**
- high **luminosity** to reach **large  $x_B$**



All 3 compatible  
with DIS data  
(Pumplin *et al.*)

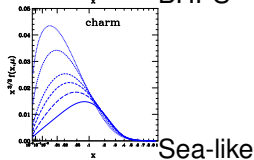
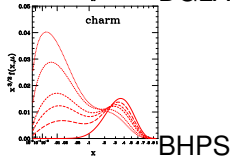
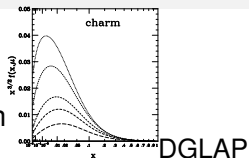
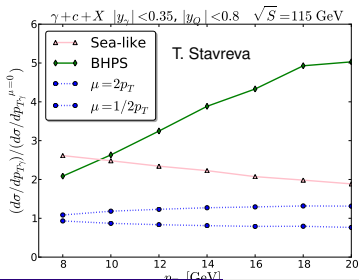


# Key studies

- **Heavy-quark** distributions (at high  $x_B$ )
  - Pin down **intrinsic charm**, ... at last
  - **Total open charm and beauty** cross section (aim: down to  $P_T \rightarrow 0$ )

requires

- several **complementary** measurements
- good coverage in the **target-rapidity** region
- high **luminosity** to reach **large  $x_B$**



All 3 compatible  
with DIS data  
(Pumplin *et al.*)

# Key studies: gluon contribution to the proton spin

- **Gluon Sivers effect: correlation between the gluon transverse momentum & the proton spin**

# Key studies: gluon contribution to the proton spin



- **Gluon Sivers effect: correlation** between the **gluon transverse momentum** & the **proton spin**
  - **Transverse single spin asymmetries**  
using **gluon sensitive probes**

# Key studies: gluon contribution to the proton spin



- **Gluon Sivers effect: correlation** between the **gluon transverse momentum** & the **proton spin**
  - Transverse **single spin asymmetries** using **gluon sensitive probes**
  - quarkonia ( $J/\psi$ ,  $\Upsilon$ ,  $\chi_C$ , ...)

F. Yuan, PRD 78 (2008) 014024

# Key studies: gluon contribution to the proton spin



- **Gluon Sivers effect: correlation** between the **gluon transverse momentum** & the **proton spin**
  - Transverse **single spin asymmetries** using **gluon sensitive probes**
  - quarkonia ( $J/\psi$ ,  $\Upsilon$ ,  $\chi_C$ , ...)
  - $B$  &  $D$  meson production

F. Yuan, PRD 78 (2008) 014024

# Key studies: gluon contribution to the proton spin



- **Gluon Sivers effect: correlation** between the **gluon transverse momentum** & the **proton spin**
  - Transverse **single spin asymmetries** using **gluon sensitive probes**
- quarkonia ( $J/\psi$ ,  $\Upsilon$ ,  $\chi_C$ , ...)
- $B$  &  $D$  meson production
- $\gamma$ ,  $\gamma$ -jet,  $\gamma - \gamma$

F. Yuan, PRD 78 (2008) 014024

A. Bacchetta, *et al.*, PRL 99 (2007) 212002  
 J.W. Qiu, *et al.*, PRL 107 (2011) 062001

# Key studies: gluon contribution to the proton spin



- **Gluon Sivers effect: correlation** between the **gluon transverse momentum** & the **proton spin**
  - Transverse **single spin asymmetries** using **gluon sensitive probes**
  - quarkonia ( $J/\psi$ ,  $\Upsilon$ ,  $\chi_C$ , ... ) F. Yuan, PRD 78 (2008) 014024
  - $B$  &  $D$  meson production
  - $\gamma$ ,  $\gamma$ -jet,  $\gamma - \gamma$  A. Bacchetta, *et al.*, PRL 99 (2007) 212002  
J.W. Qiu, *et al.*, PRL 107 (2011) 062001
- the target-rapidity region corresponds to **high  $x^\uparrow$**  where the  **$k_T$ -spin correlation is the largest**

# Key studies: gluon contribution to the proton spin



- **Gluon Sivers effect: correlation** between the **gluon transverse momentum** & the **proton spin**
  - Transverse **single spin asymmetries** using **gluon sensitive probes**
  - quarkonia ( $J/\psi$ ,  $\Upsilon$ ,  $\chi_C$ , ... ) F. Yuan, PRD 78 (2008) 014024
  - $B$  &  $D$  meson production
  - $\gamma$ ,  $\gamma$ -jet,  $\gamma - \gamma$  A. Bacchetta, *et al.*, PRL 99 (2007) 212002  
J.W. Qiu, *et al.*, PRL 107 (2011) 062001
- the target-rapidity region corresponds to **high  $x^\uparrow$**  where the  **$k_T$ -spin correlation is the largest**
- In general, one can carry out an extensive spin-physics program



# Key studies: gluon contribution to the proton spin



- **Gluon Sivers effect: correlation** between the **gluon transverse momentum** & the **proton spin**

- Trans

expected Sivers asymmetry in  
D-Y@AFTER, sign change,  
no TMD evolution

ve probes

- quark

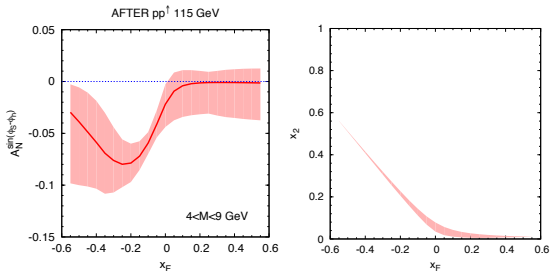
78 (2008) 014024

- B & L

- $\gamma, \gamma^*j$

212002  
1001

- the target



ne largest

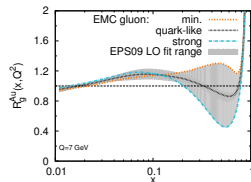
- In general **M. Anselmino, Trento, February 2013**

rogram

# Key studies: large- $x$ gluon content of the nucleus

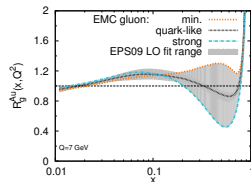
# Key studies: large-x gluon content of the nucleus

- Large-x gluon nPDF: unknown
- Gluon EMC effect ?



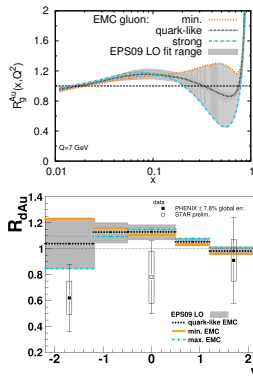
# Key studies: large-x gluon content of the nucleus

- Large-x gluon nPDF: unknown
- Gluon EMC effect ?
- Hint from  $\Upsilon$  data at RHIC



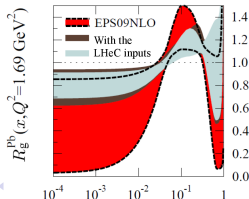
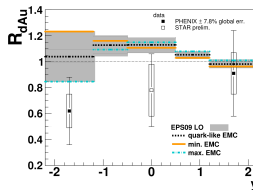
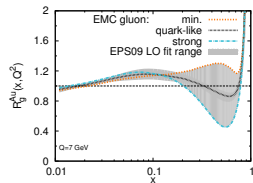
# Key studies: large-x gluon content of the nucleus

- Large-x gluon nPDF: unknown
- Gluon EMC effect ?
- Hint from  $\Upsilon$  data at RHIC
- Strongly limited in terms of statistics  
after 10 years of RHIC:



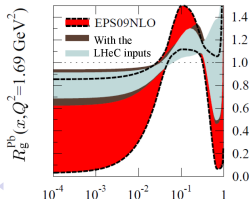
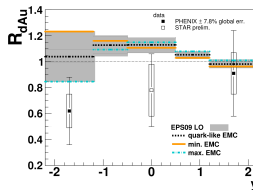
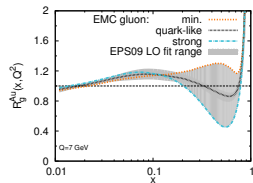
# Key studies: large-x gluon content of the nucleus

- Large-x gluon nPDF: unknown
- Gluon EMC effect ?
- Hint from  $\Upsilon$  data at RHIC
- Strongly limited in terms of statistics  
after 10 years of RHIC:
- DIS contribution expected for low x mainly  
projected contribution of LHeC:



# Key studies: large-x gluon content of the nucleus

- Large-x gluon nPDF: unknown
- Gluon EMC effect ?
- Hint from  $\Upsilon$  data at RHIC
- Strongly limited in terms of statistics  
after 10 years of RHIC:
- DIS contribution expected for low x mainly  
projected contribution of LHeC:
- AFTER allows for extensive studies of  
gluon sensitive probes in  $pA$
- Unique potential for gluons at  $x > 0.1$



# Key studies: W/Z production at threshold

- For the first time, one would study **W/Z production** in their **threshold region** ( $m_{W/Z}/\sqrt{s_{AFTER}} \sim 1$ )



# Key studies: W/Z production at threshold

- For the first time, one would study **W/Z production**  
in their **threshold region** ( $m_{W/Z}/\sqrt{s_{AFTER}} \sim 1$ )
  - Unique opportunity to measure QCD/threshold effects  
on W/Z production

## Key studies: W/Z production at threshold

- For the first time, one would study **W/Z production**  
in their **threshold region** ( $m_{W/Z}/\sqrt{s_{AFTER}} \sim 1$ )
  - Unique opportunity to measure QCD/threshold effects  
on W/Z production
  - If  $W'/Z'$  exist, their production may share similar threshold corrections to that of W/Z, but at LHC energies  
( $m_{W'/Z'}/\sqrt{s_{LHC}} \sim 1$  ?)

## Key studies: W/Z production at threshold

- For the first time, one would study **W/Z production** in their **threshold region** ( $m_{W/Z}/\sqrt{s_{AFTER}} \sim 1$ )
  - Unique opportunity to measure QCD/threshold effects on W/Z production
  - If  $W'/Z'$  exist, their production may share similar threshold corrections to that of W/Z, but at LHC energies ( $m_{W'/Z'}/\sqrt{s_{LHC}} \sim 1$  ?)
  - Reconstructed rate are most likely between **a few dozen to a few thousand / year**

## Further key studies ?

- **Multiply heavy baryons**: discovery potential ? ( $\Omega^{++}(ccc)$ , ...)

## Further key studies ?

- **Multiply heavy baryons**: discovery potential ? ( $\Omega^{++}(ccc)$ , ...)
- Very forward (backward) physics

## Further key studies ?

- **Multiply heavy baryons**: discovery potential ? ( $\Omega^{++}(ccc)$ , ...)
- Very forward (backward) physics
  - it is of course conceivable to cover the most forward region for cosmic-ray studies

## Further key studies ?

- **Multiply heavy baryons**: discovery potential ? ( $\Omega^{++}(ccc)$ , ...)
- Very forward (backward) physics
  - it is of course conceivable to cover the most forward region for cosmic-ray studies
  - need for a long detector  $\rightarrow$  need for a strong physics case

## Further key studies ?

- **Multiply heavy baryons**: discovery potential ? ( $\Omega^{++}(ccc)$ , ...)
- Very forward (backward) physics
  - it is of course conceivable to cover the most forward region for cosmic-ray studies
  - need for a long detector  $\rightarrow$  need for a strong physics case
- Semi-diffractive events

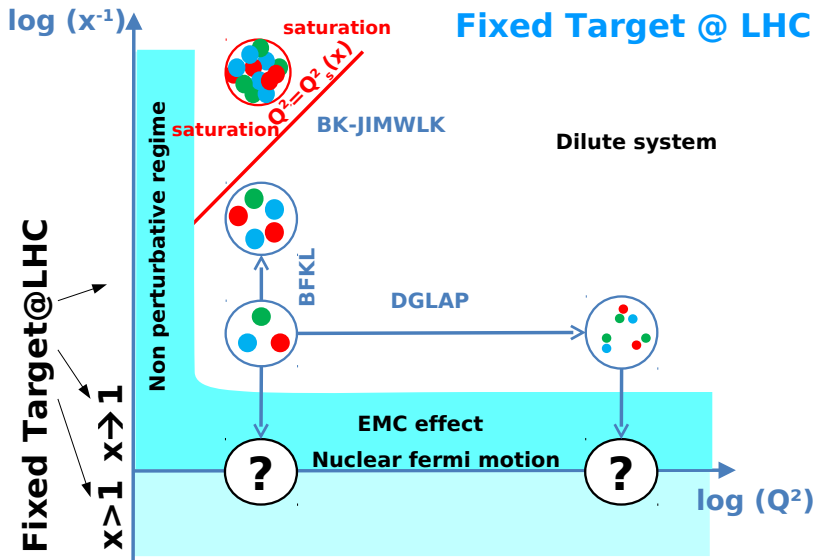


## Further key studies ?

- **Multiply heavy baryons**: discovery potential ? ( $\Omega^{++}(ccc)$ , ...)
- Very forward (backward) physics
  - it is of course conceivable to cover the most forward region for cosmic-ray studies
  - need for a long detector  $\rightarrow$  need for a strong physics case
- Semi-diffractive events
- Ultra-peripheral collisions via  $\gamma p$  interaction
  - $\gamma_{\text{lab}}^{\text{beam}} \simeq 7000$
  - $E_{\gamma,\text{lab}}^{\text{max}} \simeq \gamma_{\text{lab}}^{\text{beam}} \times 30 \text{ MeV}$
  - $\sqrt{s_{\gamma p}} = \sqrt{2m_p E_\gamma}$  up to 20 GeV

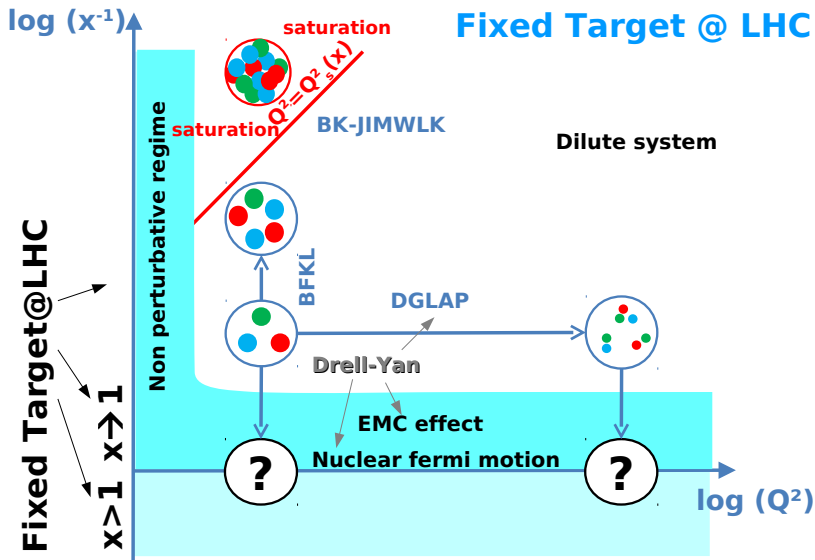
## Overall

## Fixed Target @ LHC



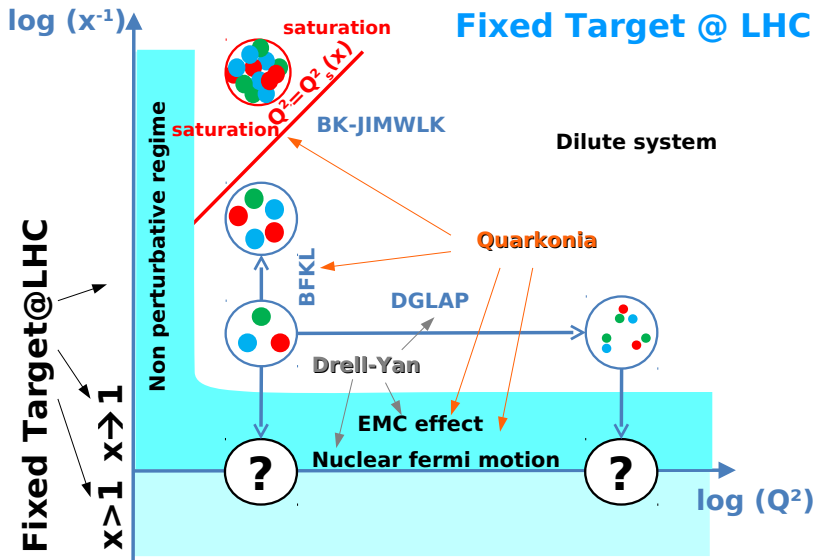
## Overall

## Fixed Target @ LHC



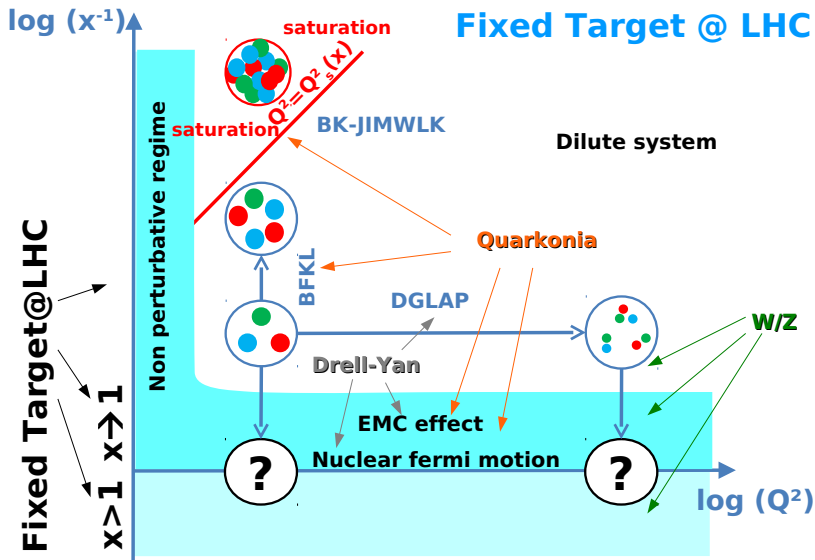
## Overall

## Fixed Target @ LHC



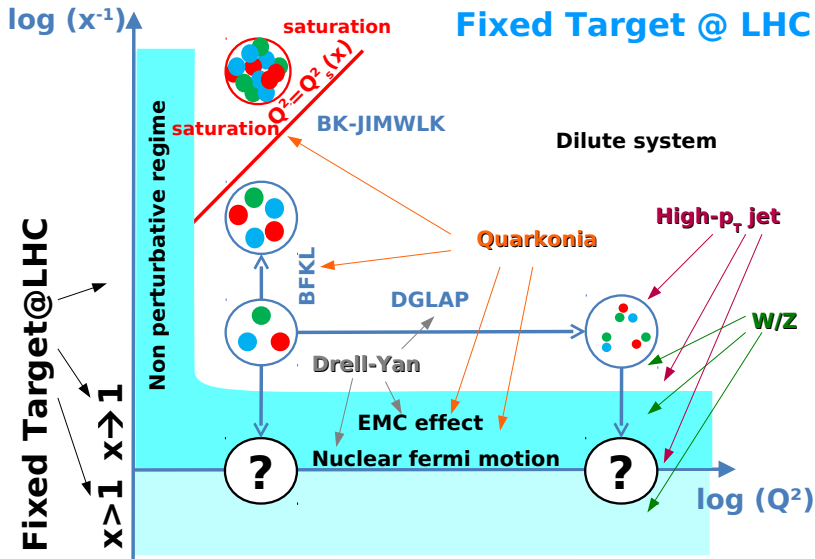
## Overall

## Fixed Target @ LHC



## Overall

## Fixed Target @ LHC



# More details in

Physics Reports 522 (2013) 239–255



Contents lists available at SciVerse ScienceDirect

## Physics Reports

journal homepage: [www.elsevier.com/locate/physrep](http://www.elsevier.com/locate/physrep)



## Physics opportunities of a fixed-target experiment using LHC beams

S.J. Brodsky<sup>a</sup>, F. Fleuret<sup>b</sup>, C. Hadjidakis<sup>c</sup>, J.P. Lansberg<sup>c,\*</sup>

<sup>a</sup> SLAC National Accelerator Laboratory, Stanford University, Menlo Park, CA 94025, USA

<sup>b</sup> Laboratoire Leprince Ringuet, Ecole polytechnique, CNRS/IN2P3, 91128 Palaiseau, France

<sup>c</sup> IPNO, Université Paris-Sud, CNRS/IN2P3, 91406 Orsay, France

### Contents

1. Introduction.....	6. Deconfinement in heavy-ion collisions .....
2. Key numbers and features .....	6.1. Quarkonium studies .....
3. Nucleon partonic structure .....	6.2. Jet quenching .....
3.1. Drell–Yan.....	6.3. Direct photon .....
3.2. Gluons in the proton at large $x$ .....	6.4. Deconfinement and the target rest frame .....
3.2.1. Quarkonia .....	6.5. Nuclear-matter baseline.....
3.2.2. Jets .....	7. $W$ and $Z$ boson production in $pp$ , $pd$ and $pA$ collisions.....
3.2.3. Direct/isolated photons.....	7.1. First measurements in $pA$ .....
3.3. Gluons in the deuteron and in the neutron.....	7.2. $W/Z$ production in $pp$ and $pd$ .....
3.4. Charm and bottom in the proton.....	8. Exclusive, semi-exclusive and backward reactions .....
3.4.1. Open-charm production.....	8.1. Ultra-peripheral collisions .....
3.4.2. $J/\psi + D$ meson production .....	8.2. Hard diffractive reactions .....
3.4.3. Heavy-quark plus photon production .....	8.3. Heavy-hadron (diffractive) production at $x_F \rightarrow -1$ .....
4. Spin physics.....	8.4. Very backward physics.....
4.1. Transverse SSA and DY .....	8.5. Direct hadron production.....
4.2. Quarkonium and heavy-quark transverse SSA .....	9. Further potentialities of a high-energy fixed-target set-up.....
4.3. Transverse SSA and photon .....	9.1. $D$ and $B$ physics .....
4.4. Spin asymmetries with a final state polarization .....	9.2. Secondary beams .....
5. Nuclear matter .....	9.3. Forward studies in relation with cosmic shower .....
5.1. Quark nPDF: Drell–Yan in $pA$ and $PbPb$ .....	10. Conclusions.....
5.2. Gluon nPDF.....	Acknowledgments .....
5.2.1. Isolated photons and photon–jet correlations.....	References.....
5.2.2. Precision quarkonium and heavy-flavour studies.....	
5.3. Color filtering, energy loss, Sudakov suppression and hadron break-up in the nucleus .....	

# Part IV

Back to the future ...



Nuclear Instruments and Methods in Physics Research A 333 (1993) 125–135  
North-Holland

## LHB, a fixed target experiment at LHC to measure CP violation in B mesons

Flavio Costantini

*University of Pisa and INFN, Italy*

A fixed target experiment at LHC to measure CP violation in B mesons is presented. A description of the proposed apparatus is given together with its sensitivity on the CP violation asymmetry measurement for the two benchmark decay channels  $B^0 \rightarrow J/\psi + K_s^0$ ,  $B^0 \rightarrow \pi^+ \pi^-$ . The possibility of obtaining an extracted LHC beam hinges on channeling in a bent silicon crystal. Recent results on beam extraction efficiencies measured at CERN SPS based on this technique are presented.

# LHB

Our idea is not completely new

## 1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about  $10^8$  protons/s allowing the production of as many as  $10^{10}$   $B\bar{B}$  pairs per year, i.e. about two orders of magnitude more than what could be produced by an  $e^+e^-$  asymmetric B factory with  $10^{34}$   $\text{cm}^{-2}\text{s}^{-1}$  luminosity [5].



# LHB

Our idea is not completely new

## 1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about  $10^8$  protons/s allowing the production of as many as  $10^{10}$   $B\bar{B}$  pairs per year, i.e. about two orders of magnitude more than what could be produced by an  $e^+e^-$  asymmetric B factory with  $10^{34}$   $\text{cm}^{-2}\text{s}^{-1}$  luminosity [5].

$10^{10}$   $B\bar{B}$  pairs per year

- B-factories:  $1 \text{ ab}^{-1}$  means  $10^9 B\bar{B}$  pairs



# LHB

Our idea is not completely new

## 1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about  $10^8$  protons/s allowing the production of as many as  $10^{10}$   $B\bar{B}$  pairs per year, i.e. about two orders of magnitude more than what could be produced by an  $e^+e^-$  asymmetric B factory with  $10^{34}$   $\text{cm}^{-2}\text{s}^{-1}$  luminosity [5].

$10^{10}$   $B\bar{B}$  pairs per year



- *B*-factories:  $1 \text{ ab}^{-1}$  means  $10^9 B\bar{B}$  pairs
- For LHCb, typically  $1 \text{ fb}^{-1}$  means  $\simeq 2 \times 10^{11} B\bar{B}$  pairs at 14 TeV

# LHB

Our idea is not completely new

## 1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about  $10^8$  protons/s allowing the production of as many as  $10^{10}$   $B\bar{B}$  pairs per year, i.e. about two orders of magnitude more than what could be produced by an  $e^+e^-$  asymmetric B factory with  $10^{34}$   $\text{cm}^{-2}\text{s}^{-1}$  luminosity [5].

$10^{10}$   $B\bar{B}$  pairs per year



- B-factories:  $1 \text{ ab}^{-1}$  means  $10^9 B\bar{B}$  pairs
- For LHCb, typically  $1 \text{ fb}^{-1}$  means  $\simeq 2 \times 10^{11} B\bar{B}$  pairs at 14 TeV
- LHB turned down in favour of LHCb mainly because of the **fear of a premature degradation of the bent crystal** due to radiation damages.

## 1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about  $10^8$  protons/s allowing the production of as many as  $10^{10}$   $B\bar{B}$  pairs per year, i.e. about two orders of magnitude more than what could be produced by an  $e^+e^-$  asymmetric B factory with  $10^{34}$   $\text{cm}^{-2}\text{s}^{-1}$  luminosity [5].

$10^{10} B\bar{B}$  pairs per year



- B-factories:  $1 \text{ ab}^{-1}$  means  $10^9 B\bar{B}$  pairs
- For LHCb, typically  $1 \text{ fb}^{-1}$  means  $\simeq 2 \times 10^{11} B\bar{B}$  pairs at 14 TeV
- LHB turned down in favour of LHCb mainly because of the **fear of a premature degradation of the bent crystal** due to radiation damages.
- Nowadays, degradation is known to be  $\simeq 6\%$  per  $10^{20}$  particles/ $\text{cm}^2$
- $10^{20}$  particles/ $\text{cm}^2$  : one year of operation for realistic conditions

### 1. Introduction

...

This paper presents a fixed target experiment to measure CP violation in the B system based on the possibility of extracting the 8 TeV LHC proton beam using a bent silicon crystal [4]. A 10% extraction efficiency of the LHC beam halo will give an extracted beam intensity of about  $10^8$  protons/s allowing the production of as many as  $10^{10}$   $B\bar{B}$  pairs per year, i.e. about two orders of magnitude more than what could be produced by an  $e^+e^-$  asymmetric B factory with  $10^{34}$   $\text{cm}^{-2}\text{s}^{-1}$  luminosity [5].

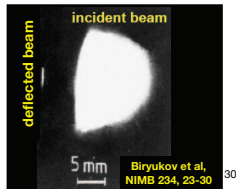
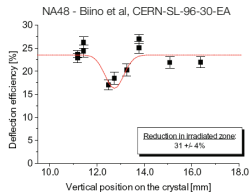
$10^{10}$   $B\bar{B}$  pairs per year



- B-factories:  $1 \text{ ab}^{-1}$  means  $10^9 B\bar{B}$  pairs
- For LHCb, typically  $1 \text{ fb}^{-1}$  means  $\simeq 2 \times 10^{11} B\bar{B}$  pairs at 14 TeV
- LHB turned down in favour of LHCb mainly because of the **fear of a premature degradation of the bent crystal** due to radiation damages.
- Nowadays, degradation is known to be  $\simeq 6\%$  per  $10^{20}$  particles/ $\text{cm}^2$
- $10^{20}$  particles/ $\text{cm}^2$  : one year of operation for realistic conditions
- After a year, one simply moves the crystal by less than one mm ...

# Crystal resistance to irradiation

- **IHEP U-70** (Biryukov et al, NIMB 234, 23-30):
  - 70 GeV protons, 50 ms spills of  **$10^{14}$  protons every 9.6 s**, several minutes irradiation
  - equivalent to 2 nominal LHC bunches for 500 turns every 10 s
  - 5 mm silicon crystal, **channeling efficiency unchanged**
- **SPS North Area - NA48** (Biino et al, CERN-SL-96-30-EA):
  - 450 GeV protons, 2.4 s spill of  $5 \times 10^{12}$  protons every 14.4 s, one year irradiation,  **$2.4 \times 10^{20}$  protons/cm<sup>2</sup>** in total,
  - equivalent to several year of operation for a primary collimator in LHC
  - $10 \times 50 \times 0.9$  mm<sup>3</sup> silicon crystal,  $0.8 \times 0.3$  mm<sup>2</sup> area irradiated, **channeling efficiency reduced by 30%**.
- **HRMT16-UA9CRY** (HiRadMat facility, November 2012):
  - 440 GeV protons, up to 288 bunches **in 7.2  $\mu$ s**,  $1.1 \times 10^{11}$  protons per bunch ( **$3 \times 10^{13}$  protons** in total)
  - energy deposition comparable to an asynchronous beam dump in LHC
  - 3 mm long silicon crystal, **no damage to the crystal after accurate visual inspection**, more tests planned to assess possible crystal lattice damage
    - **accurate FLUKA simulation of energy deposition** and residual dose





# Part V

## Conclusion and outlooks

# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments

# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments
- Extracting a few per cent of the beam  $\rightarrow 5 \times 10^8$  protons per sec

# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments
- Extracting a few per cent of the beam  $\rightarrow 5 \times 10^8$  protons per sec
- This allows for high luminosity  $pp$ ,  $pA$  and  $PbA$  collisions at  $\sqrt{s} = 115$  GeV and  $\sqrt{s_{NN}} = 72$  GeV

# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments
- Extracting a few per cent of the beam  $\rightarrow 5 \times 10^8$  protons per sec
- This allows for high luminosity  $pp$ ,  $pA$  and  $PbA$  collisions at  $\sqrt{s} = 115$  GeV and  $\sqrt{s_{NN}} = 72$  GeV
- **Example:** precision quarkonium studies taking advantage of

# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments
- Extracting a few per cent of the beam  $\rightarrow 5 \times 10^8$  protons per sec
- This allows for high luminosity  $pp$ ,  $pA$  and  $PbA$  collisions at  $\sqrt{s} = 115$  GeV and  $\sqrt{s_{NN}} = 72$  GeV
- **Example:** precision quarkonium studies taking advantage of
  - high luminosity (reach in  $y$ ,  $P_T$ , small BR channels)
  - target versatility (nuclear effects, strongly limited at colliders)
  - modern detection techniques (e.g.  $\gamma$  detection with high multiplicity)

# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments
- Extracting a few per cent of the beam  $\rightarrow 5 \times 10^8$  protons per sec
- This allows for high luminosity  $pp$ ,  $pA$  and  $PbA$  collisions at  $\sqrt{s} = 115$  GeV and  $\sqrt{s_{NN}} = 72$  GeV
- **Example:** precision quarkonium studies taking advantage of
  - high luminosity (reach in  $y$ ,  $P_T$ , small BR channels)
  - target versatility (nuclear effects, strongly limited at colliders)
  - modern detection techniques (e.g.  $\gamma$  detection with high multiplicity)
- This would likely prepare the ground for  $g(x, Q^2)$  extraction

# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments
- Extracting **a few per cent of the beam**  $\rightarrow 5 \times 10^8$  **protons per sec**
- This allows for **high luminosity**  $pp$ ,  $pA$  and  $PbA$  collisions at  $\sqrt{s} = 115$  GeV and  $\sqrt{s_{NN}} = 72$  GeV
- **Example: precision quarkonium studies** taking advantage of
  - high luminosity (reach in  $y$ ,  $P_T$ , small BR channels)
  - target versatility (nuclear effects, strongly limited at colliders)
  - modern detection techniques (e.g.  $\gamma$  detection with high multiplicity)
- This would likely prepare the ground for  **$g(x, Q^2)$  extraction**
- A wealth of possible measurements:  
DY, Open  $b/c$ , jet correlation, UPC... (not mentioning secondary beams)



# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments
- Extracting **a few per cent of the beam**  $\rightarrow 5 \times 10^8$  **protons per sec**
- This allows for **high luminosity**  $pp$ ,  $pA$  and  $PbA$  collisions at  $\sqrt{s} = 115$  GeV and  $\sqrt{s_{NN}} = 72$  GeV
- **Example: precision quarkonium studies** taking advantage of
  - high luminosity (reach in  $y$ ,  $P_T$ , small BR channels)
  - target versatility (nuclear effects, strongly limited at colliders)
  - modern detection techniques (e.g.  $\gamma$  detection with high multiplicity)
- This would likely prepare the ground for  **$g(x, Q^2)$  extraction**
- A wealth of possible measurements:  
DY, Open  $b/c$ , jet correlation, UPC... (not mentioning secondary beams)
- Planned LHC long shutdown ( $< 2020$  ?) could be used to install the extraction system

# Conclusion

- Both  $p$  and  $Pb$  LHC beams can be extracted without disturbing the other experiments
- Extracting **a few per cent of the beam**  $\rightarrow 5 \times 10^8$  **protons per sec**
- This allows for **high luminosity**  $pp$ ,  $pA$  and  $PbA$  collisions at  $\sqrt{s} = 115$  GeV and  $\sqrt{s_{NN}} = 72$  GeV
- **Example: precision quarkonium studies** taking advantage of
  - high luminosity (reach in  $y$ ,  $P_T$ , small BR channels)
  - target versatility (nuclear effects, strongly limited at colliders)
  - modern detection techniques (e.g.  $\gamma$  detection with high multiplicity)
- This would likely prepare the ground for  **$g(x, Q^2)$  extraction**
- A wealth of possible measurements:  
DY, Open  $b/c$ , jet correlation, UPC... (not mentioning secondary beams)
- Planned LHC long shutdown ( $< 2020$  ?) could be used to install the extraction system
- Very good **complementarity** with electron-ion programs

# Outlooks

- First physics paper [Physics Reports 522 \(2013\) 239](#)

# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>

# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>
- NEW: Workshop in **Les Houches on 12-17 January 2014**  
and we plan for a 3 day workshop at CERN  
in the 2nd part of October within the GDR PH-QCD

# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>
- NEW: Workshop in **Les Houches on 12-17 January 2014**  
and we plan for a 3 day workshop at CERN  
in the 2nd part of October within the GDR PH-QCD
- We are looking for **more partners** to

# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>
- NEW: Workshop in **Les Houches on 12-17 January 2014**  
and we plan for a 3 day workshop at CERN  
in the 2nd part of October within the GDR PH-QCD
- We are looking for **more partners** to
  - do first **simulations** (we are getting ready for fast simulations)

# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>
- NEW: Workshop in **Les Houches on 12-17 January 2014**  
and we plan for a 3 day workshop at CERN  
in the 2nd part of October within the GDR PH-QCD
- We are looking for **more partners** to
  - do first **simulations** (we are getting ready for fast simulations)
  - think about **possible designs**



# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>
- NEW: Workshop in **Les Houches on 12-17 January 2014**  
and we plan for a 3 day workshop at CERN  
in the 2nd part of October within the GDR PH-QCD
- We are looking for **more partners** to
  - do first **simulations** (we are getting ready for fast simulations)
  - think about **possible designs**
  - think about the optimal **detector technologies**
  - enlarge the physics case (cosmic rays, flavour physics, ...)

# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>
- NEW: Workshop in **Les Houches on 12-17 January 2014**  
and we plan for a 3 day workshop at CERN  
in the 2nd part of October within the GDR PH-QCD
- We are looking for **more partners** to
  - do first **simulations** (we are getting ready for fast simulations)
  - think about **possible designs**
  - think about the optimal **detector technologies**
  - enlarge the physics case (cosmic rays, flavour physics, ...)
- **Theorist colleagues** are encouraged to think about **additional ideas of physics**

already 2 papers on the physics at AFTER:  
T. Liu, B.Q. Ma, EPJC (2012) 72:2037  
D. Boer, C. Pisano, Phys.Rev. D86 (2012) 094007

# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>
- NEW: Workshop in **Les Houches on 12-17 January 2014**  
and we plan for a 3 day workshop at CERN  
in the 2nd part of October within the GDR PH-QCD
- We are looking for **more partners** to
  - do first **simulations** (we are getting ready for fast simulations)
  - think about **possible designs**
  - think about the optimal **detector technologies**
  - enlarge the physics case (cosmic rays, flavour physics, ...)
- **Theorist colleagues** are encouraged to think about  
**additional ideas of physics**  
already 2 papers on the physics at AFTER:  
T. Liu, B.Q. Ma, EPJC (2012) 72:2037  
D. Boer, C. Pisano, Phys.Rev. D86 (2012) 094007
- Do not hesitate to contact us

# Outlooks

- First physics paper **Physics Reports 522 (2013) 239**
- A 10-day exploratory workshop at ECT\* Trento, February 4-13, 2013  
slides at <http://indico.in2p3.fr/event/AFTER@ECTstar>
- NEW: Workshop in **Les Houches on 12-17 January 2014**  
and we plan for a 3 day workshop at CERN  
in the 2nd part of October within the GDR PH-QCD
- We are looking for **more partners** to
  - do first **simulations** (we are getting ready for fast simulations)
  - think about **possible designs**
  - think about the optimal **detector technologies**
  - enlarge the physics case (cosmic rays, flavour physics, ...)
- **Theorist colleagues** are encouraged to think about  
**additional ideas of physics**
- Do not hesitate to contact us
- Webpage: <http://after.in2p3.fr>

already 2 papers on the physics at AFTER:  
T. Liu, B.Q. Ma, EPJC (2012) 72:2037  
D. Boer, C. Pisano, Phys.Rev. D86 (2012) 094007



# Part VI

## Backup slides

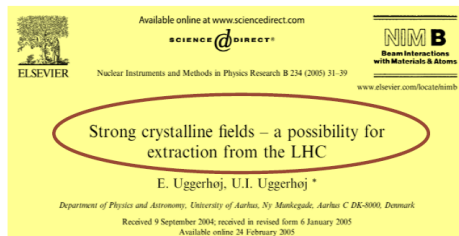
# Beam extraction

## Beam extraction @ LHC

... there are extremely promising possibilities to extract 7 TeV protons from the circulating beam by means of a bent crystal.

... The idea is to put a bent, single crystal of either Si or Ge (W would perform slightly better but needs substantial improvements in crystal quality) at a distance of  $\simeq 7\sigma$  to the beam where it can intercept and deflect part of the beam halo by an angle similar to the one the foreseen dump kicking system will apply to the circulating beam.

... ions with the same momentum per charge as protons are deflected in a crystal with similar efficiencies



If the crystal is positioned at the kicking section, the whole dump system can be used for slow extraction of parts of the beam halo, the particles that are anyway lost subsequently at collimators.

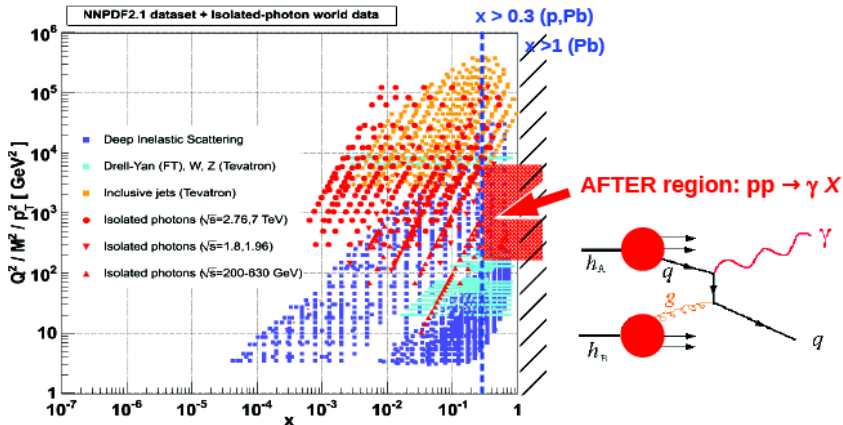
NEW!

# $(x, Q^2)$ map of AFTER isolated- $\gamma$

[D. d'E &amp; J. Rojo, NPB 860 (2012) 311]

## ■ p-p kinematics at fixed-target LHC:

To access  $x > 0.3$  one needs isolated- $\gamma$  with:  $p_T = x_T \sqrt{s}/2 > 10\text{-}20 \text{ GeV}/c$

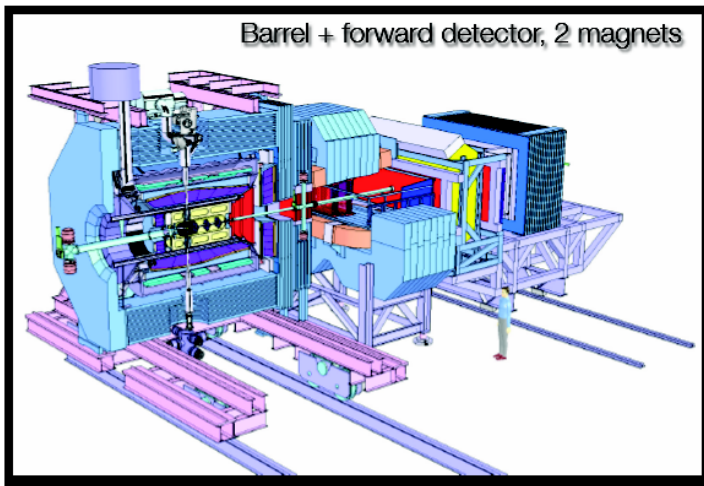


[D. D'Enterria, Physics at AFTER using LHC beams, ICT\* Trento, Feb 2013]



AFTER@LHC

# Detector : could be inspired by PANDA



EmCal could be based on ultragranular CALICE, developed for ILC



AFTER, among other things, a quarkonium observatory in  $pp$ 

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. $\text{H}_2$	20	$4.0 \cdot 10^8$	$8.0 \cdot 10^5$
1 m Liq. $\text{D}_2$	24	$9.6 \cdot 10^8$	$1.9 \cdot 10^6$
LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	$3.6 \cdot 10^7$ $1.4 \cdot 10^9$	$1.8 \cdot 10^5$ $7.2 \cdot 10^6$
RHIC pp 200GeV	$1.2 \cdot 10^{-2}$	$4.8 \cdot 10^5$	$1.2 \cdot 10^3$

AFTER, among other things, a quarkonium observatory in  $pp$ 

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(\text{J}/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
<b>1 m Liq. H<sub>2</sub></b>	<b>20</b>	<b>4.0 10<sup>8</sup></b>	<b>8.0 10<sup>5</sup></b>
<b>1 m Liq. D<sub>2</sub></b>	<b>24</b>	<b>9.6 10<sup>8</sup></b>	<b>1.9 10<sup>6</sup></b>
<b>LHC pp 14 Tev (low pT)</b>	<b>0.05 (ALICE) 2 LHCb</b>	<b>3.6 10<sup>7</sup> 1.4 10<sup>9</sup></b>	<b>1.8 10<sup>5</sup> 7.2 10<sup>6</sup></b>
<b>RHIC pp 200GeV</b>	<b>1.2 10<sup>-2</sup></b>	<b>4.8 10<sup>5</sup></b>	<b>1.2 10<sup>3</sup></b>

- 1000 times higher than at RHIC; comparable to ALICE/LHCb at the LHC

AFTER, among other things, a quarkonium observatory in  $pp$ 

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
<b>1 m Liq. H<sub>2</sub></b>	<b>20</b>	<b>4.0 10<sup>8</sup></b>	<b>8.0 10<sup>5</sup></b>
<b>1 m Liq. D<sub>2</sub></b>	<b>24</b>	<b>9.6 10<sup>8</sup></b>	<b>1.9 10<sup>6</sup></b>
<b>LHC pp 14 Tev (low pT)</b>	<b>0.05 (ALICE) 2 LHCb</b>	<b>3.6 10<sup>7</sup> 1.4 10<sup>9</sup></b>	<b>1.8 10<sup>5</sup> 7.2 10<sup>6</sup></b>
<b>RHIC pp 200GeV</b>	<b>1.2 10<sup>-2</sup></b>	<b>4.8 10<sup>5</sup></b>	<b>1.2 10<sup>3</sup></b>

- 1000 times higher than at RHIC; comparable to ALICE/LHCb at the LHC
- Numbers are for only one unit of rapidity about 0

AFTER, among other things, a quarkonium observatory in  $pp$ 

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
<b>1 m Liq. H<sub>2</sub></b>	<b>20</b>	<b>4.0 10<sup>8</sup></b>	<b>8.0 10<sup>5</sup></b>
<b>1 m Liq. D<sub>2</sub></b>	<b>24</b>	<b>9.6 10<sup>8</sup></b>	<b>1.9 10<sup>6</sup></b>
<b>LHC pp 14 Tev (low pT)</b>	<b>0.05 (ALICE) 2 LHCb</b>	<b>3.6 10<sup>7</sup> 1.4 10<sup>9</sup></b>	<b>1.8 10<sup>5</sup> 7.2 10<sup>6</sup></b>
<b>RHIC pp 200GeV</b>	<b>1.2 10<sup>-2</sup></b>	<b>4.8 10<sup>5</sup></b>	<b>1.2 10<sup>3</sup></b>

- 1000 times higher than at RHIC; comparable to ALICE/LHCb at the LHC
- Numbers are for only one unit of rapidity about 0
- Unique access in the backward region

AFTER, among other things, a quarkonium observatory in  $pp$ 

- Interpolating the world data set:

Target	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}B\sigma_{\Upsilon}$
1 m Liq. $H_2$	20	4.0 $10^8$	8.0 $10^5$
1 m Liq. $D_2$	24	9.6 $10^8$	1.9 $10^6$
LHC pp 14 Tev (low pT)	0.05 (ALICE) 2 LHCb	3.6 $10^7$ 1.4 $10^9$	1.8 $10^5$ 7.2 $10^6$
RHIC pp 200GeV	1.2 $10^{-2}$	4.8 $10^5$	1.2 $10^3$

- 1000 times higher than at RHIC; comparable to ALICE/LHCb at the LHC
- Numbers are for only one unit of rapidity about 0
- Unique access in the backward region
- Probe of the (very) large  $x$  in the target

# Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**

# Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**
  - in photo/lepto production (DIS)
  - but also  $pp$  collisions in  $gg$ -fusion process
  - mainly because of the presence of a natural “hard” scale:  $m_Q$
  - and the good detectability of a dimuon pair

# Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**
  - in photo/lepto production (DIS)
  - but also  $pp$  collisions in  $gg$ -fusion process
  - mainly because of the presence of a natural “hard” scale:  $m_Q$
  - and the good detectability of a dimuon pair

PHYSICAL REVIEW D

VOLUME 37, NUMBER 5

1 MARCH 1988

## Structure-function analysis and $\psi$ , jet, $W$ , and $Z$ production: Determining the gluon distribution

A. D. Martin

*Department of Physics, University of Durham, Durham, England*

R. G. Roberts

*Rutherford Appleton Laboratory, Didcot, Oxon, England*

W. J. Stirling

*Department of Physics, University of Durham, Durham, England*

(Received 27 July 1987)

We perform a next-to-leading-order structure-function analysis of deep-inelastic  $\mu N$  and  $\nu N$  scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) “soft,” (2) “hard,” and (3) which behave as  $xG(x) \sim 1/\sqrt{x}$  at small  $x$ .  $J/\psi$  and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the “soft”-gluon distribution, is favored.  $W$ ,  $Z$ , and jet production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for  $\sigma_W$  and  $\sigma_Z$  allow the collider measurements to yield information on the number of light neutrinos and the mass of the top quark. Finally we discuss how the gluon distribution at very small  $x$  may be directly measured at DESY HERA.



# Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**
  - in photo/lepto production (DIS)
  - but also  $pp$  collisions in  $gg$ -fusion process
  - mainly because of the presence of a natural “hard” scale:  $m_Q$
  - and the good detectability of a dimuon pair

PHYSICAL REVIEW D

VOLUME 37, NUMBER 5

1 MARCH 1988

## Structure-function analysis and $\psi$ , jet, $W$ , and $Z$ production: Determining the gluon distribution

A. D. Martin

*Department of Physics, University of Durham, Durham, England*

R. G. Roberts

*Rutherford Appleton Laboratory, Didcot, Oxon, England*

W. J. Stirling

*Department of Physics, University of Durham, Durham, England*

(Received 27 July 1987)

We perform a next-to-leading-order structure-function analysis of deep-inelastic  $\mu N$  and  $\nu N$  scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) “soft,” (2) “hard,” and (3) which behave as  $xG(x) \sim 1/\sqrt{x}$  at small  $x$ .  $J/\psi$  and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the “soft”-gluon distribution, is favored.  $W$ ,  $Z$ , and jet production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for  $\sigma_W$  and  $\sigma_Z$  allow the collider measurements to yield information on the number of light neutrinos and the mass of the top quark. Finally we discuss how the gluon distribution at very small  $x$  may be directly measured at DESY HERA.

- Production **puzzle** → quarkonium not used anymore in global fits

# Need for a quarkonium observatory

- Many **hopes** were put in **quarkonium studies** to extract **gluon PDF**
  - in photo/lepto production (DIS)
  - but also  $pp$  collisions in  $gg$ -fusion process
  - mainly because of the presence of a natural “hard” scale:  $m_Q$
  - and the good detectability of a dimuon pair

PHYSICAL REVIEW D

VOLUME 37, NUMBER 5

1 MARCH 1988

## Structure-function analysis and $\psi$ , jet, $W$ , and $Z$ production: Determining the gluon distribution

A. D. Martin

*Department of Physics, University of Durham, Durham, England*

R. G. Roberts

*Rutherford Appleton Laboratory, Didcot, Oxon, England*

W. J. Stirling

*Department of Physics, University of Durham, Durham, England*

(Received 27 July 1987)

We perform a next-to-leading-order structure-function analysis of deep-inelastic  $\mu N$  and  $\nu N$  scattering data and find acceptable fits for a range of input gluon distributions. We show three equally acceptable sets of parton distributions which correspond to gluon distributions which are (1) “soft,” (2) “hard,” and (3) which behave as  $xG(x) \sim 1/\sqrt{x}$  at small  $x$ .  $J/\psi$  and prompt photon hadroproduction data are used to discriminate between the three sets. Set 1, with the “soft”-gluon distribution, is favored.  $W$ ,  $Z$ , and jet production data from the CERN collider are well described but do not distinguish between the sets of structure functions. The precision of the predictions for  $\sigma_W$  and  $\sigma_Z$  allow the collider measurements to yield information on the number of light neutrinos and the mass of the top quark. Finally we discuss how the gluon distribution at very small  $x$  may be directly measured at DESY HERA.

- Production **puzzle**  $\rightarrow$  quarkonium not used anymore in global fits
- With systematic studies, one would **restore its status as gluon probe**

AFTER: also a quarkonium observatory in  $pA$ 

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Upsilon}$
<b>1cm Be</b>	9	<b>0.62</b>	<b>1.1 10<sup>8</sup></b>	<b>2.2 10<sup>5</sup></b>
<b>1cm Cu</b>	64	<b>0.42</b>	<b>5.3 10<sup>8</sup></b>	<b>1.1 10<sup>6</sup></b>
<b>1cm W</b>	185	<b>0.31</b>	<b>1.1 10<sup>9</sup></b>	<b>2.3 10<sup>6</sup></b>
<b>1cm Pb</b>	207	<b>0.16</b>	<b>6.7 10<sup>8</sup></b>	<b>1.3 10<sup>6</sup></b>
<b>LHC pPb 8.8 TeV</b>	207	<b>10<sup>-4</sup></b>	<b>1.0 10<sup>7</sup></b>	<b>7.5 10<sup>4</sup></b>
<b>RHIC dAu 200GeV</b>	198	<b>1.5 10<sup>-4</sup></b>	<b>2.4 10<sup>6</sup></b>	<b>5.9 10<sup>3</sup></b>
<b>RHIC dAu 62GeV</b>	198	<b>3.8 10<sup>-6</sup></b>	<b>1.2 10<sup>4</sup></b>	<b>18</b>

- In principle, one can get **300 times more  $J/\psi$**  –not counting the likely wider  $y$  coverage– than at RHIC, allowing for

AFTER: also a quarkonium observatory in  $pA$ 

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Upsilon}$
1cm Be	9	0.62	1.1 $10^8$	2.2 $10^5$
1cm Cu	64	0.42	5.3 $10^8$	1.1 $10^6$
1cm W	185	0.31	1.1 $10^9$	2.3 $10^6$
1cm Pb	207	0.16	6.7 $10^8$	1.3 $10^6$
LHC pPb 8.8 TeV	207	$10^{-4}$	1.0 $10^7$	7.5 $10^4$
RHIC dAu 200GeV	198	$1.5 \cdot 10^{-4}$	2.4 $10^6$	5.9 $10^3$
RHIC dAu 62GeV	198	$3.8 \cdot 10^{-6}$	1.2 $10^4$	18

- In principle, one can get **300 times more  $J/\psi$**  –not counting the likely wider  $y$  coverage– than at RHIC, allowing for
  - $\chi_c$  measurement in  $pA$  via  $J/\psi + \gamma$  (extending Hera-B studies)

AFTER: also a quarkonium observatory in  $pA$ 

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Upsilon}$
<b>1cm Be</b>	9	<b>0.62</b>	<b>1.1 10<sup>8</sup></b>	<b>2.2 10<sup>5</sup></b>
<b>1cm Cu</b>	64	<b>0.42</b>	<b>5.3 10<sup>8</sup></b>	<b>1.1 10<sup>6</sup></b>
<b>1cm W</b>	185	<b>0.31</b>	<b>1.1 10<sup>9</sup></b>	<b>2.3 10<sup>6</sup></b>
<b>1cm Pb</b>	207	<b>0.16</b>	<b>6.7 10<sup>8</sup></b>	<b>1.3 10<sup>6</sup></b>
<b>LHC pPb 8.8 TeV</b>	207	<b>10<sup>-4</sup></b>	<b>1.0 10<sup>7</sup></b>	<b>7.5 10<sup>4</sup></b>
<b>RHIC dAu 200GeV</b>	198	<b>1.5 10<sup>-4</sup></b>	<b>2.4 10<sup>6</sup></b>	<b>5.9 10<sup>3</sup></b>
<b>RHIC dAu 62GeV</b>	198	<b>3.8 10<sup>-6</sup></b>	<b>1.2 10<sup>4</sup></b>	<b>18</b>

- In principle, one can get **300 times more  $J/\psi$**  –not counting the likely wider  $y$  coverage– than at RHIC, allowing for
  - $\chi_c$  measurement in  $pA$  via  $J/\psi + \gamma$  (extending Hera-B studies)
  - **Polarisation** measurement as **the centrality,  $y$  or  $P_T$**

AFTER: also a quarkonium observatory in  $pA$ 

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1})$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Upsilon}$
<b>1cm Be</b>	9	<b>0.62</b>	<b>1.1 <math>10^8</math></b>	<b>2.2 <math>10^5</math></b>
<b>1cm Cu</b>	64	<b>0.42</b>	<b>5.3 <math>10^8</math></b>	<b>1.1 <math>10^6</math></b>
<b>1cm W</b>	185	<b>0.31</b>	<b>1.1 <math>10^9</math></b>	<b>2.3 <math>10^6</math></b>
<b>1cm Pb</b>	207	<b>0.16</b>	<b>6.7 <math>10^8</math></b>	<b>1.3 <math>10^6</math></b>
<b>LHC pPb 8.8 TeV</b>	207	<b><math>10^{-4}</math></b>	<b>1.0 <math>10^7</math></b>	<b>7.5 <math>10^4</math></b>
<b>RHIC dAu 200GeV</b>	198	<b><math>1.5 \cdot 10^{-4}</math></b>	<b>2.4 <math>10^6</math></b>	<b>5.9 <math>10^3</math></b>
<b>RHIC dAu 62GeV</b>	198	<b><math>3.8 \cdot 10^{-6}</math></b>	<b>1.2 <math>10^4</math></b>	<b>18</b>

- In principle, one can get **300 times more  $J/\psi$**  –not counting the likely wider  $y$  coverage– than at RHIC, allowing for
  - $\chi_c$  measurement in  $pA$  via  $J/\psi + \gamma$  (extending Hera-B studies)
  - **Polarisation** measurement as **the centrality,  $y$  or  $P_T$**
  - Ratio  $\psi'$  over **direct  $J/\psi$**  measurement in  $pA$

AFTER: also a quarkonium observatory in  $pA$ 

Target	A	$\int \mathcal{L} \text{ (fb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= A\mathcal{L}\mathcal{B}\sigma_{\Upsilon}$
<b>1cm Be</b>	9	<b>0.62</b>	<b>1.1 10<sup>8</sup></b>	<b>2.2 10<sup>5</sup></b>
<b>1cm Cu</b>	64	<b>0.42</b>	<b>5.3 10<sup>8</sup></b>	<b>1.1 10<sup>6</sup></b>
<b>1cm W</b>	185	<b>0.31</b>	<b>1.1 10<sup>9</sup></b>	<b>2.3 10<sup>6</sup></b>
<b>1cm Pb</b>	207	<b>0.16</b>	<b>6.7 10<sup>8</sup></b>	<b>1.3 10<sup>6</sup></b>
<b>LHC pPb 8.8 TeV</b>	207	<b>10<sup>-4</sup></b>	<b>1.0 10<sup>7</sup></b>	<b>7.5 10<sup>4</sup></b>
<b>RHIC dAu 200GeV</b>	198	<b>1.5 10<sup>-4</sup></b>	<b>2.4 10<sup>6</sup></b>	<b>5.9 10<sup>3</sup></b>
<b>RHIC dAu 62GeV</b>	198	<b>3.8 10<sup>-6</sup></b>	<b>1.2 10<sup>4</sup></b>	<b>18</b>

- In principle, one can get **300 times more  $J/\psi$**  –not counting the likely wider  $y$  coverage– than at RHIC, allowing for
  - $\chi_c$  measurement in  $pA$  via  $J/\psi + \gamma$  (extending Hera-B studies)
  - **Polarisation** measurement as **the centrality,  $y$  or  $P_T$**
  - Ratio  $\psi'$  over **direct  $J/\psi$**  measurement in  $pA$
  - not to mention ratio with **open charm, Drell-Yan**, etc ...

# What for ?

- The **target versatility** of a fixed-target experiment is undisputable



# What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
  - a precise analysis of **gluon nuclear PDF**:  $y, p_T \leftrightarrow x_2$
  - a handle on **formation time effects**

# What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
  - a precise analysis of **gluon nuclear PDF**:  $y, p_T \leftrightarrow x_2$
  - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements

# What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
  - a precise analysis of **gluon nuclear PDF**:  $y, p_T \leftrightarrow x_2$
  - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements
- The **backward kinematics** is very useful for large- $x_{target}$  studies

# What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
  - a precise analysis of **gluon nuclear PDF**:  $y, p_T \leftrightarrow x_2$
  - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements
- The **backward kinematics** is very useful for large- $x_{target}$  studies
  - What is the amount of Intrinsic charm ? Is it color filtered ?

# What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
  - a precise analysis of **gluon nuclear PDF**:  $y, p_T \leftrightarrow x_2$
  - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements
- The **backward kinematics** is very useful for large- $x_{target}$  studies
  - What is the amount of Intrinsic charm ? Is it color filtered ?
  - **Is there an EMC effect for gluon ?** (reminder: EMC region  $0.3 < x < 0.7$ )

# What for ?

- The **target versatility** of a fixed-target experiment is undisputable
- A **wide rapidity coverage** is needed for:
  - a precise analysis of **gluon nuclear PDF**:  $y, p_T \leftrightarrow x_2$
  - a handle on **formation time effects**
- Strong need for **cross checks** from **various** measurements
- The **backward kinematics** is very useful for large- $x_{target}$  studies
  - What is the amount of Intrinsic charm ? Is it color filtered ?
  - **Is there an EMC effect for gluon ?** (reminder: EMC region  $0.3 < x < 0.7$ )
- One should be careful with factorization breaking effects:
 

This calls for **multiple measurements** to (in)validate factorization

# AFTER: also an heavy-flavour observatory in $PbA$

- Luminosities and yields with the extracted 2.76 TeV Pb beam  
( $\sqrt{s_{NN}} = 72$  GeV)

Target	A.B	$\int \mathcal{L} \text{ (nb}^{-1}\cdot\text{yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ $= AB\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= AB\mathcal{L}B\sigma_{\Upsilon}$
<b>1 m Liq. H<sub>2</sub></b>	207.1	<b>800</b>	<b>3.4 10<sup>6</sup></b>	<b>6.9 10<sup>3</sup></b>
<b>1cm Be</b>	207.9	<b>25</b>	<b>9.1 10<sup>5</sup></b>	<b>1.9 10<sup>3</sup></b>
<b>1cm Cu</b>	207.64	<b>17</b>	<b>4.3 10<sup>6</sup></b>	<b>0.9 10<sup>3</sup></b>
<b>1cm W</b>	207.185	<b>13</b>	<b>9.7 10<sup>6</sup></b>	<b>1.9 10<sup>4</sup></b>
<b>1cm Pb</b>	207.207	<b>7</b>	<b>5.7 10<sup>6</sup></b>	<b>1.1 10<sup>4</sup></b>
<b>LHC PbPb 5.5 TeV</b>	207.207	<b>0.5</b>	<b>7.3 10<sup>6</sup></b>	<b>3.6 10<sup>4</sup></b>
<b>RHIC AuAu 200GeV</b>	198.198	<b>2.8</b>	<b>4.4 10<sup>6</sup></b>	<b>1.1 10<sup>4</sup></b>
<b>RHIC AuAu 62GeV</b>	198.198	<b>0.13</b>	<b>4.0 10<sup>4</sup></b>	<b>61</b>

# AFTER: also an heavy-flavour observatory in $PbA$

- Luminosities and yields with the extracted 2.76 TeV Pb beam  
( $\sqrt{s_{NN}} = 72$  GeV)

Target	A.B	$\int \mathcal{L} \text{ (nb}^{-1}\text{.yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Upsilon}$
<b>1 m Liq. H<sub>2</sub></b>	207.1	<b>800</b>	<b>3.4 10<sup>6</sup></b>	<b>6.9 10<sup>3</sup></b>
<b>1cm Be</b>	207.9	<b>25</b>	<b>9.1 10<sup>5</sup></b>	<b>1.9 10<sup>3</sup></b>
<b>1cm Cu</b>	207.64	<b>17</b>	<b>4.3 10<sup>6</sup></b>	<b>0.9 10<sup>3</sup></b>
<b>1cm W</b>	207.185	<b>13</b>	<b>9.7 10<sup>6</sup></b>	<b>1.9 10<sup>4</sup></b>
<b>1cm Pb</b>	207.207	<b>7</b>	<b>5.7 10<sup>6</sup></b>	<b>1.1 10<sup>4</sup></b>
<b>LHC PbPb 5.5 TeV</b>	207.207	<b>0.5</b>	<b>7.3 10<sup>6</sup></b>	<b>3.6 10<sup>4</sup></b>
<b>RHIC AuAu 200GeV</b>	198.198	<b>2.8</b>	<b>4.4 10<sup>6</sup></b>	<b>1.1 10<sup>4</sup></b>
<b>RHIC AuAu 62GeV</b>	198.198	<b>0.13</b>	<b>4.0 10<sup>4</sup></b>	<b>61</b>

- Yields **similar** to those of RHIC at 200 GeV,  
**100 times** those of RHIC at 62 GeV



# AFTER: also an heavy-flavour observatory in $PbA$

- Luminosities and yields with the extracted 2.76 TeV Pb beam  
( $\sqrt{s_{NN}} = 72$  GeV)

Target	A.B	$\int \mathcal{L} \text{ (nb}^{-1}\text{.yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ $= AB\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ $= AB\mathcal{L}B\sigma_{\Upsilon}$
<b>1 m Liq. H<sub>2</sub></b>	207.1	<b>800</b>	<b>3.4 10<sup>6</sup></b>	<b>6.9 10<sup>3</sup></b>
<b>1cm Be</b>	207.9	<b>25</b>	<b>9.1 10<sup>5</sup></b>	<b>1.9 10<sup>3</sup></b>
<b>1cm Cu</b>	207.64	<b>17</b>	<b>4.3 10<sup>6</sup></b>	<b>0.9 10<sup>3</sup></b>
<b>1cm W</b>	207.185	<b>13</b>	<b>9.7 10<sup>6</sup></b>	<b>1.9 10<sup>4</sup></b>
<b>1cm Pb</b>	207.207	<b>7</b>	<b>5.7 10<sup>6</sup></b>	<b>1.1 10<sup>4</sup></b>
<b>LHC PbPb 5.5 TeV</b>	207.207	<b>0.5</b>	<b>7.3 10<sup>6</sup></b>	<b>3.6 10<sup>4</sup></b>
<b>RHIC AuAu 200GeV</b>	198.198	<b>2.8</b>	<b>4.4 10<sup>6</sup></b>	<b>1.1 10<sup>4</sup></b>
<b>RHIC AuAu 62GeV</b>	198.198	<b>0.13</b>	<b>4.0 10<sup>4</sup></b>	<b>61</b>

- Yields **similar** to those of RHIC at 200 GeV,  
**100 times** those of RHIC at 62 GeV
- Also **very competitive** compared to the **LHC**.

# AFTER: also an heavy-flavour observatory in $PbA$

- Luminosities and yields with the extracted 2.76 TeV Pb beam  
( $\sqrt{s_{NN}} = 72$  GeV)

Target	A.B	$\int \mathcal{L} \text{ (nb}^{-1}\text{.yr}^{-1}\text{)}$	$N(J/\Psi) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Psi}$	$N(\Upsilon) \text{ yr}^{-1}$ = $AB\mathcal{L}B\sigma_{\Upsilon}$
<b>1 m Liq. H<sub>2</sub></b>	207.1	<b>800</b>	<b>3.4 10<sup>6</sup></b>	<b>6.9 10<sup>3</sup></b>
<b>1cm Be</b>	207.9	<b>25</b>	<b>9.1 10<sup>5</sup></b>	<b>1.9 10<sup>3</sup></b>
<b>1cm Cu</b>	207.64	<b>17</b>	<b>4.3 10<sup>6</sup></b>	<b>0.9 10<sup>3</sup></b>
<b>1cm W</b>	207.185	<b>13</b>	<b>9.7 10<sup>6</sup></b>	<b>1.9 10<sup>4</sup></b>
<b>1cm Pb</b>	207.207	<b>7</b>	<b>5.7 10<sup>6</sup></b>	<b>1.1 10<sup>4</sup></b>
<b>LHC PbPb 5.5 TeV</b>	207.207	<b>0.5</b>	<b>7.3 10<sup>6</sup></b>	<b>3.6 10<sup>4</sup></b>
<b>RHIC AuAu 200GeV</b>	198.198	<b>2.8</b>	<b>4.4 10<sup>6</sup></b>	<b>1.1 10<sup>4</sup></b>
<b>RHIC AuAu 62GeV</b>	198.198	<b>0.13</b>	<b>4.0 10<sup>4</sup></b>	<b>61</b>

- Yields **similar** to those of RHIC at 200 GeV,  
**100 times** those of RHIC at 62 GeV
- Also **very competitive** compared to the **LHC**.

The same picture also holds for **open heavy flavour**

# What for ?

Observation of  $J/\psi$  sequential suppression **seems to be hindered** by

- the **Cold Nuclear Matter effects**: non trivial and  
... not well understood

# What for ?

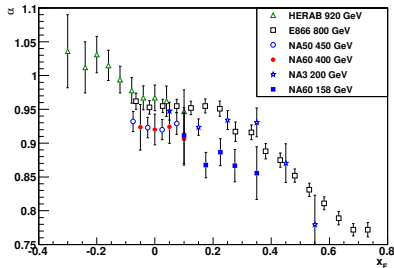
- Observation of  $J/\psi$  sequential suppression **seems to be hindered** by
- the **Cold Nuclear Matter effects**: non trivial and ... not well understood
  - the difficulty to observe directly the **excited states** which would melt before the ground states
    - $\chi_c$  **never studied in AA** collisions
    - $\psi(2S)$  **not yet studied in AA collisions at RHIC**

# What for ?

- Observation of  $J/\psi$  sequential suppression **seems to be hindered** by
- the **Cold Nuclear Matter effects**: non trivial and ... not well understood
  - the difficulty to observe directly the **excited states** which would melt before the ground states
    - $\chi_c$  **never studied in AA** collisions
    - $\psi(2S)$  **not yet studied in AA collisions at RHIC**
  - the possibilities for  **$c\bar{c}$  recombination**
    - **Open charm** studies are **difficult** where recombination matters most i.e. at **low  $P_T$**
    - Only indirect indications –from the  $y$  and  $P_T$  dependence of  $R_{AA}$ – that recombination may be at work
    - CNM effects may show a non-trivial  $y$  and  $P_T$  dependence ...

# SPS and Hera-B

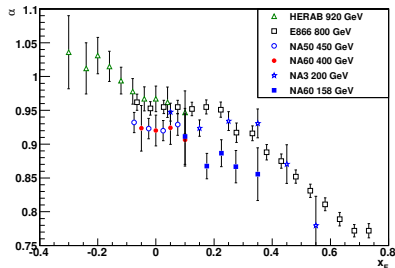
## – $J/\psi$ data in $pA$ collisions



NA60 Phys.Lett. B 706 (2012) 263  
 NA 50 Eur.Phys.J. C48 (2006) 329  
 NA 3 Z.Phys. C20 (1983)  
 HERA-B Eur.Phys.J. C60 (2009) 525

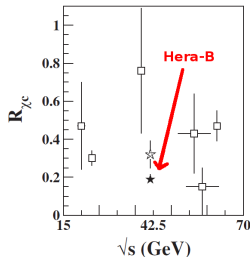
# SPS and Hera-B

## – $J/\psi$ data in $pA$ collisions



NA60 Phys.Lett. B 706 (2012) 263  
 NA 50 Eur.Phys.J. C48 (2006) 329  
 NA 3 Z.Phys. C20 (1983)  
 HERA-B Eur.Phys.J. C60 (2009) 525

## – $\chi_c$ data in $pA$ collisions



HERA-B PRD 79 (2009) 012001, and ref. therein