

From the EDF method to Ab initio approaches and back again – A DRF-DAM success story –

J.-P. Ebran

CEA,DAM,DIF

Séminaire DPhN

21/04/2023

T. Duguet, V. Somà



J.-P. Ebran

M. Frosini

Outline

- 1. General context**
- 2. Recent work on empirical EDFs**
- 3. EDF-inspired ab initio methods**
- 4. Towards a first-principle formulation of EDFs**

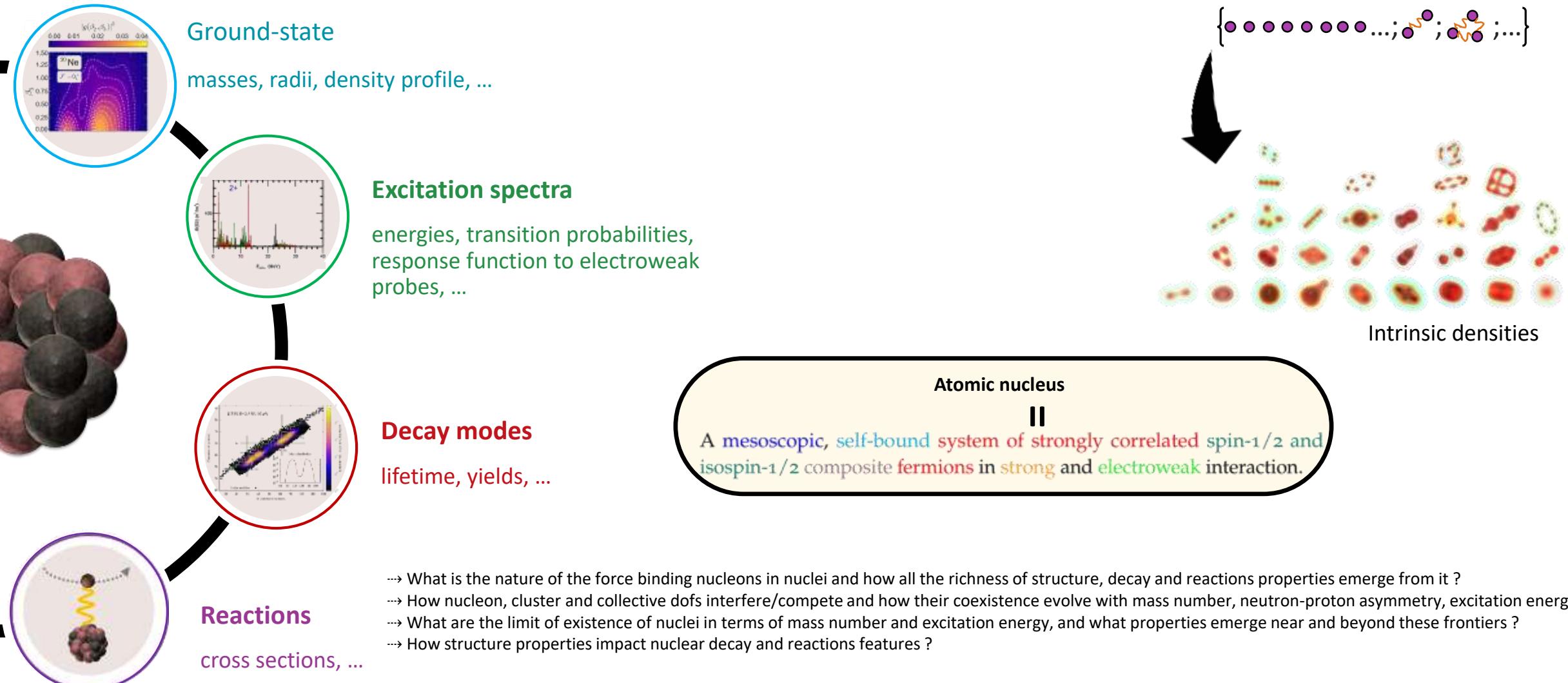


1 ■ General context



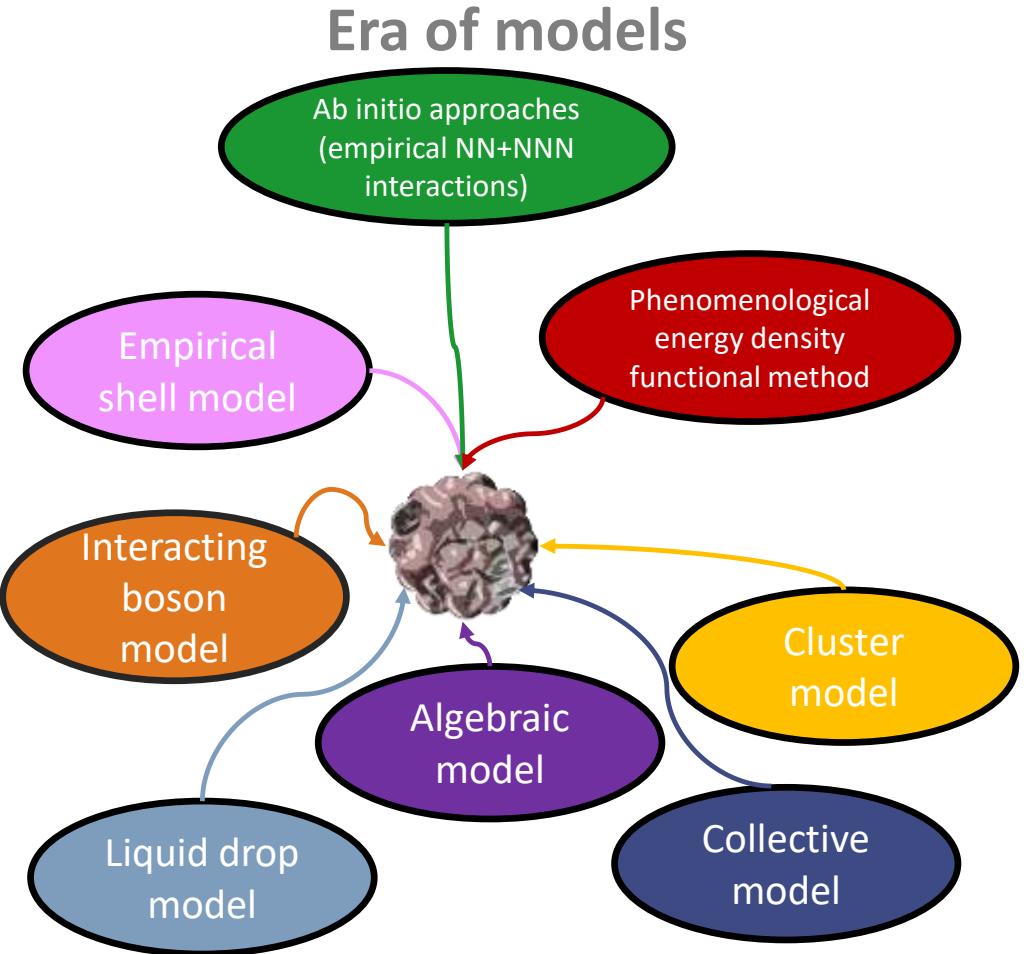
Context : General goal of nuclear structure theory

- Starting from the hadronic level of organization (nucleons + interactions), what novel structures emerge and how they evolve with E_{ex} , N , Z , ...





Strategies



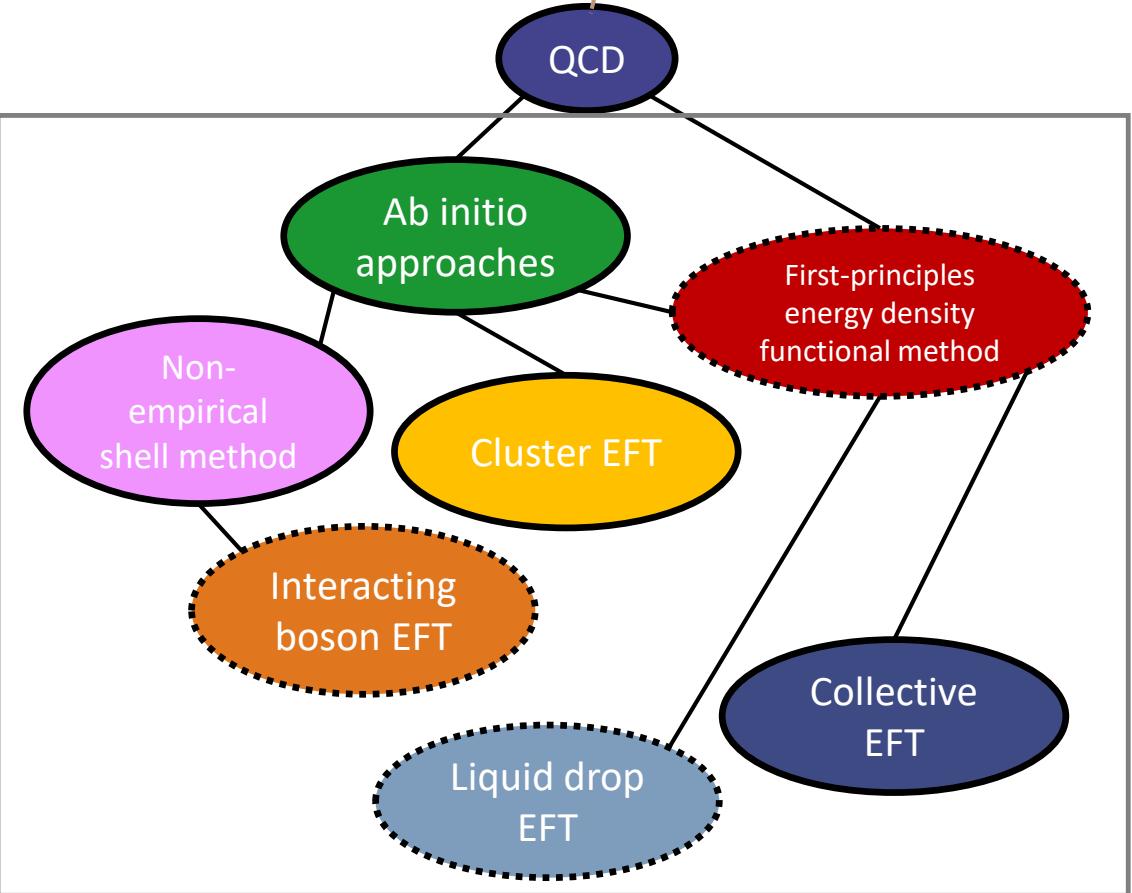
- Gives insight about relevant scales/dofs
- Ready to be used
- Lack of control
⇒ double counting issues, error compensation, no error assessment

- Achieve a

accurate
predictive
computationally affordable

description ?

Era of effective (field) theories



- Full control ⇒ systematically improvable, no error compensation, no double counting, possibility of error estimation, ...
- Force you to step back and rethink





Microscopic viewpoint

- 1) Nucleus: A interacting, structure-less nucleons
- 2) Structure & dynamic encoded in Hamiltonian, Functional, ...
- 3) Solve A-nucleon Schrödinger/Dirac equation to desired accuracy

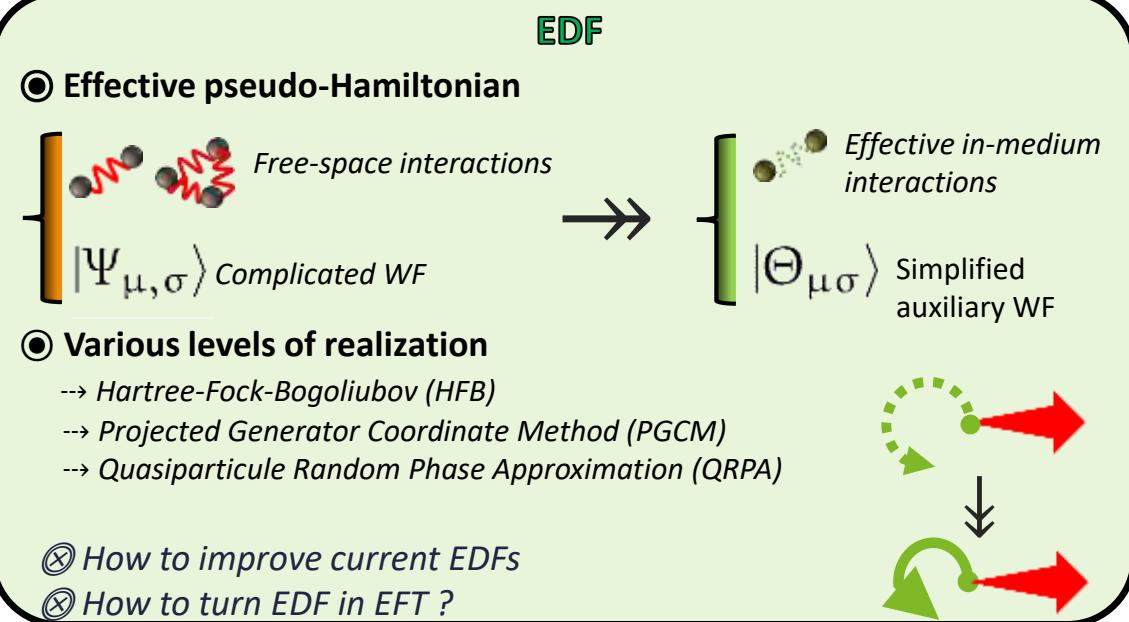
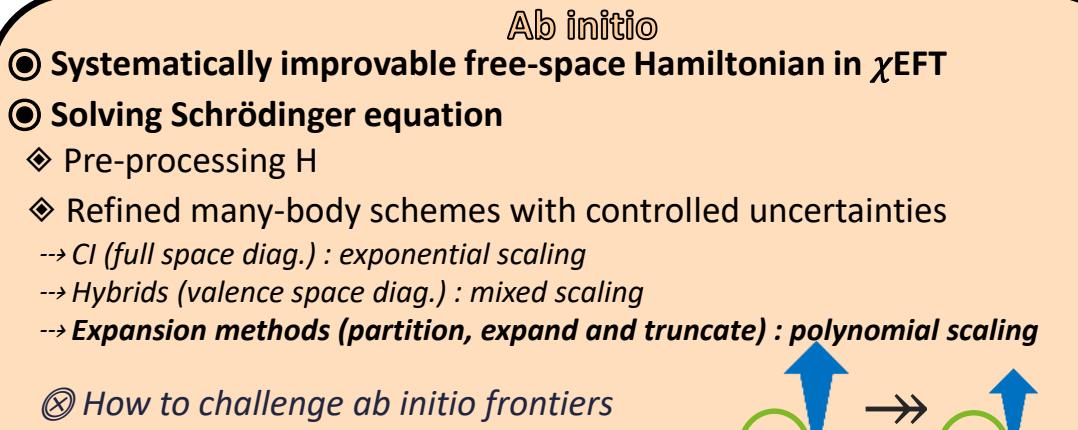
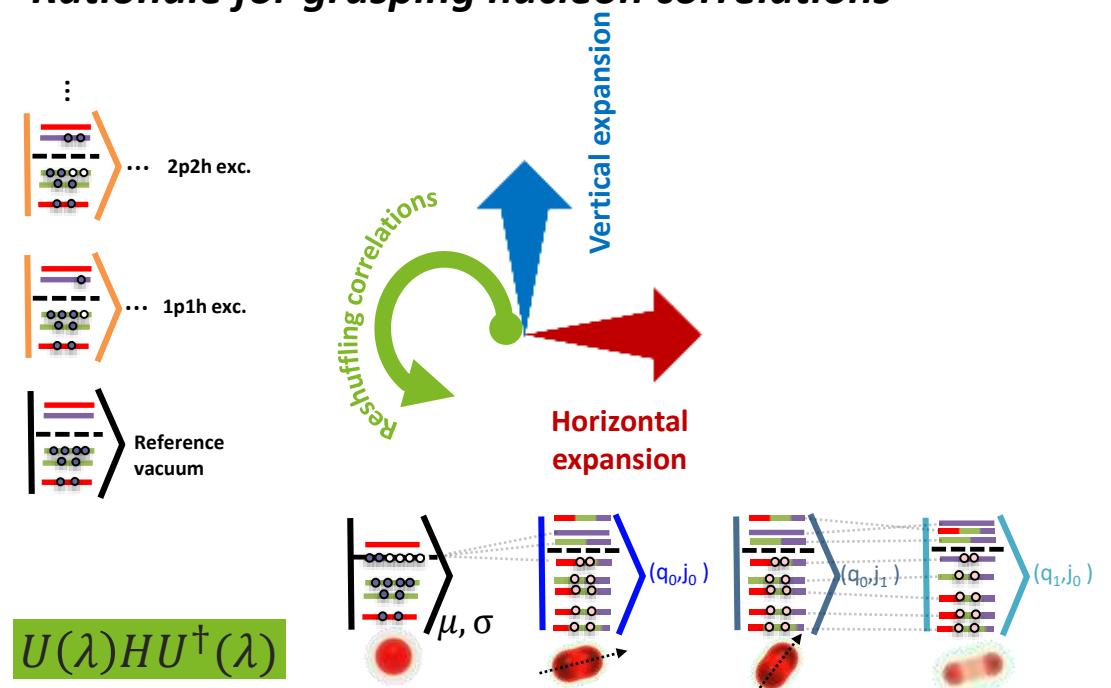
$$H(\text{Nucleus}, \dots) |\Psi_{\mu, \sigma}\rangle = E_{\mu, \sigma} |\Psi_{\mu, \sigma}\rangle$$

$N_{\text{FCI}} \propto \binom{A}{A}$

Strongly correlated WF

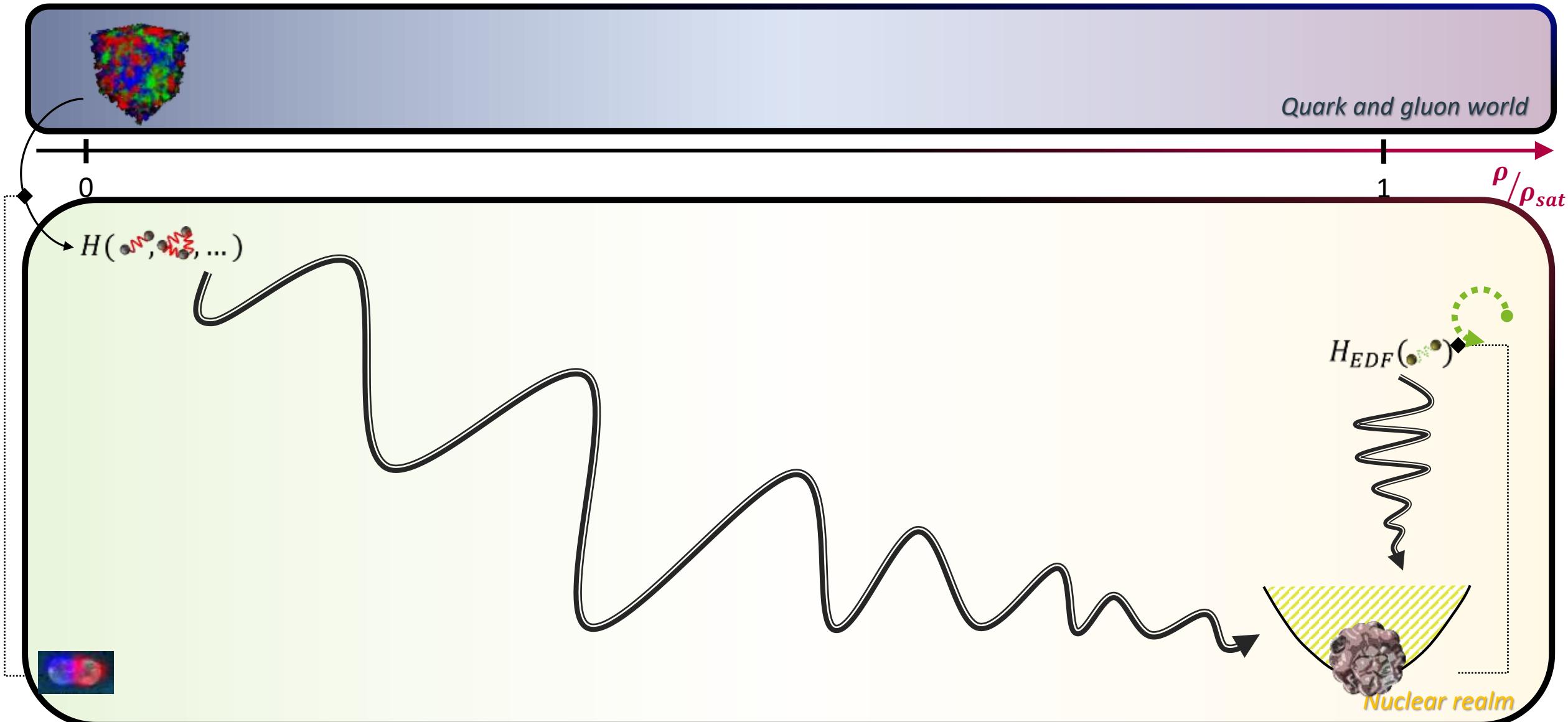
$$|\Psi_{\text{gs}}\rangle = \sum_{i_1 < \dots < i_A}^L C_{i_1 \dots i_A} |\phi_{i_1} \dots \phi_{i_A}\rangle \equiv \sum_I C_I |\Phi_I\rangle$$

Rationale for grasping nucleon correlations





Microscopic viewpoint



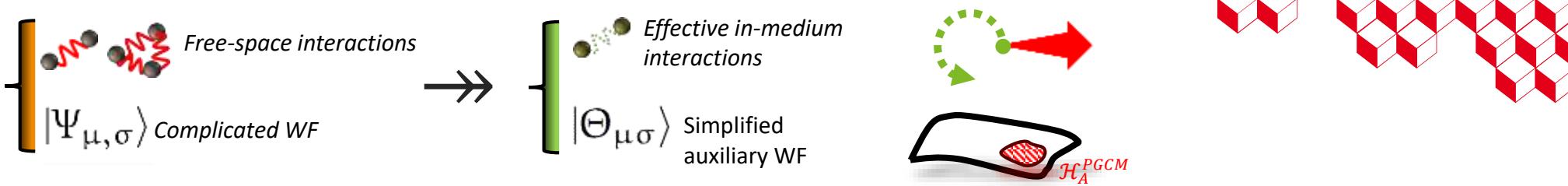
Outline

- 1. General context**
- 2. Recent work on empirical EDFs**
- 3. EDF-inspired ab initio methods**
- 4. Towards a first-principle formulation of EDFs**

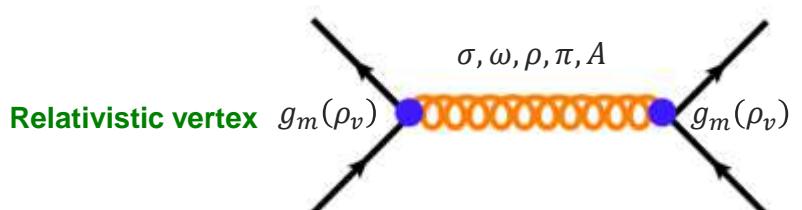
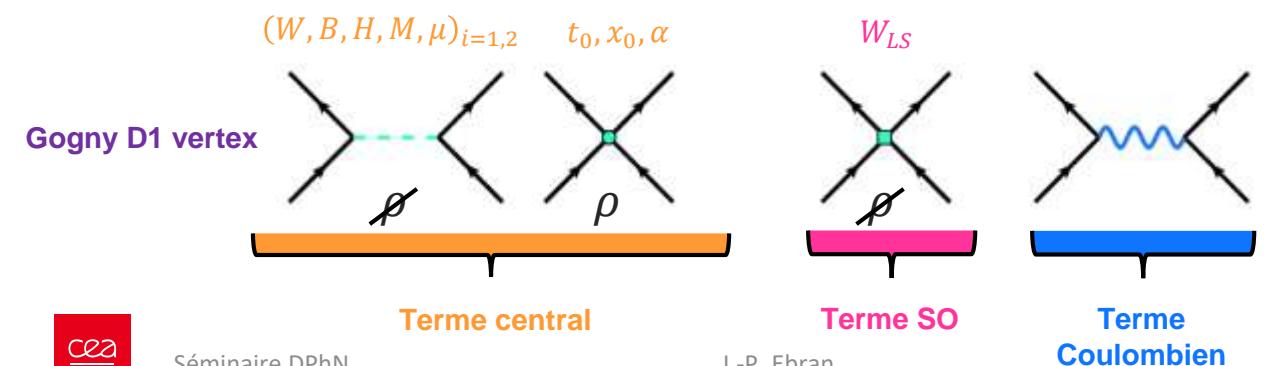
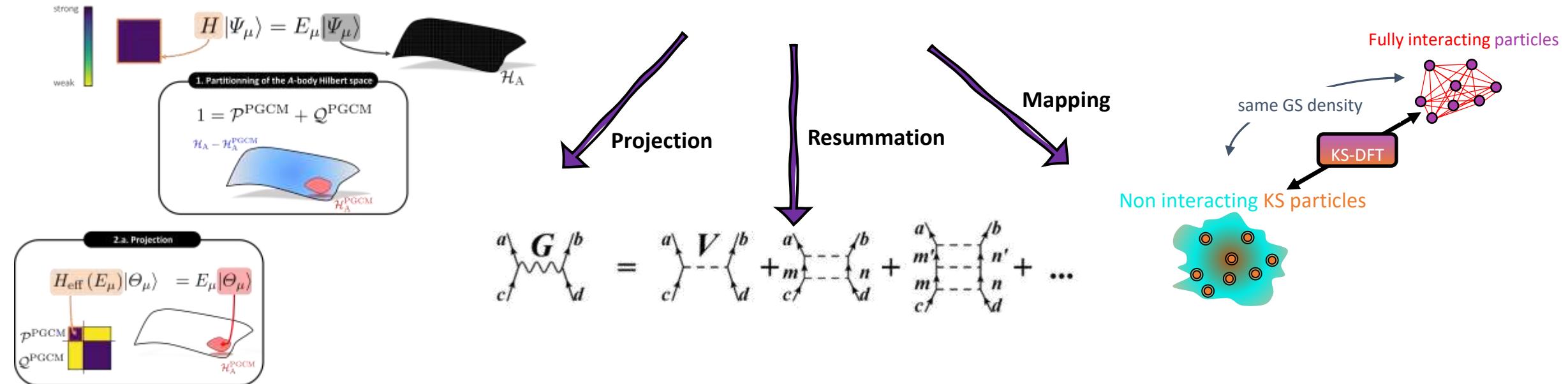


2 ■ Recent work on empirical EDFs

Empirical EDFs



Ground- and low-lying excited states computed from $H_{EDF}(\bullet)$ and $|\Theta_{\mu\sigma}\rangle \Leftrightarrow$ the ones computed from $H(\bullet)$ and $|\Psi_{\mu,\sigma}\rangle$

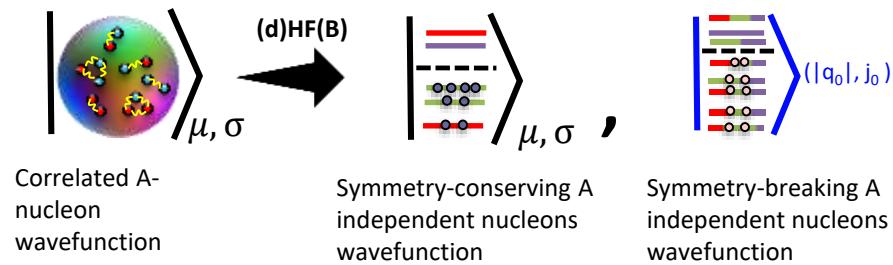




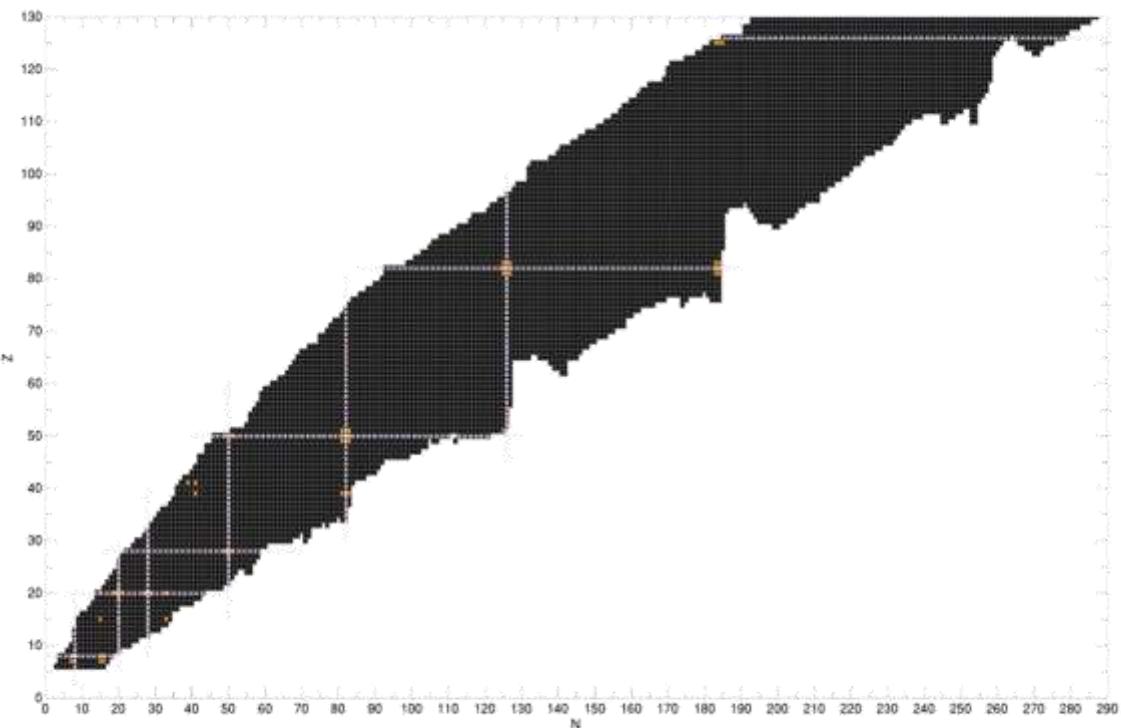
EDF : HFB realization

● HFB treatment

--> A -nucleon problem $\rightarrow A$ 1-nucleon problems



--> SSB : Efficient way for capturing so-called static correlations



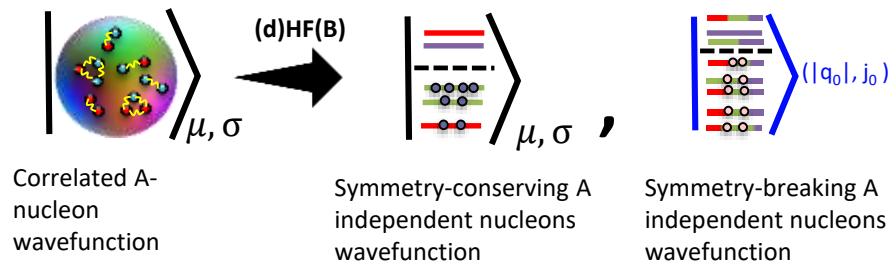
No SSBs allowed : First level of description for ~ 30 nuclei



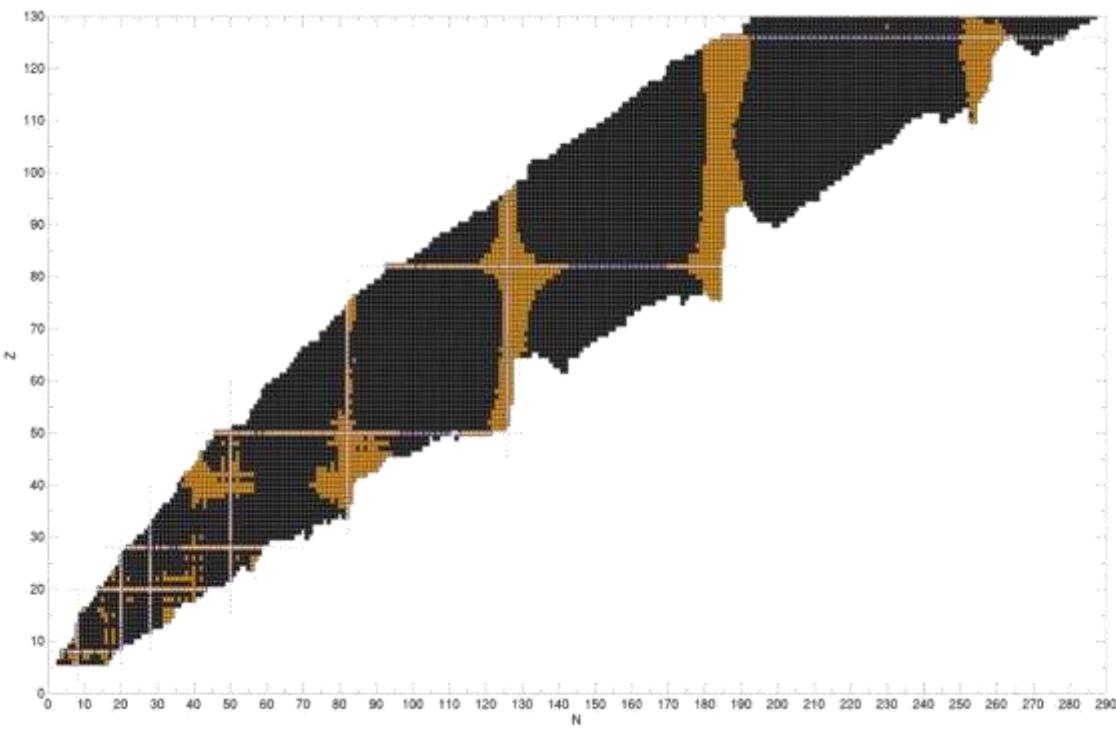
EDF : HFB realization

● HFB treatment

--> A -nucleon problem $\rightarrow A$ 1-nucleon problems



--> SSB : Efficient way for capturing so-called static correlations

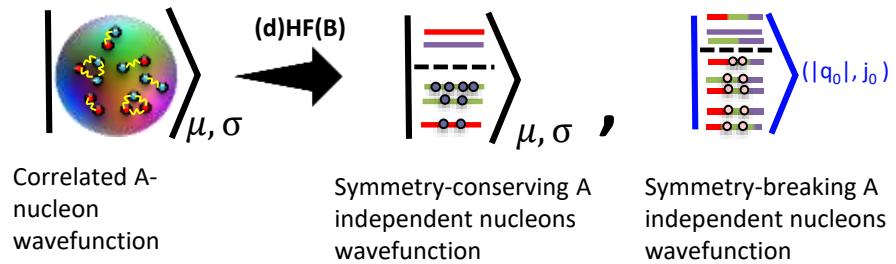


U(1) SSB allowed : First level of description for ~ 300 nuclei

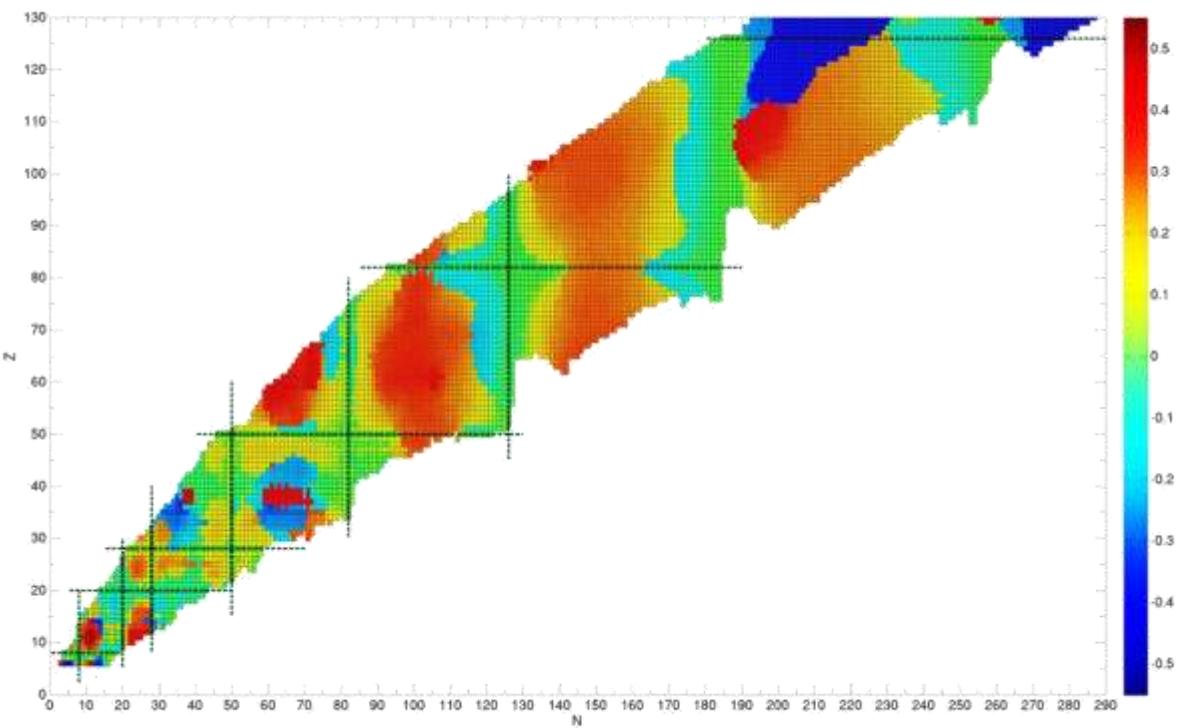
EDF : HFB realization

- HFB treatment

--> A -nucleon problem $\rightarrow A$ 1-nucleon problems



--> SSB : Efficient way for capturing so-called static correlations



SU(2) & U(1) SSB allowed : First level of description for all nuclei

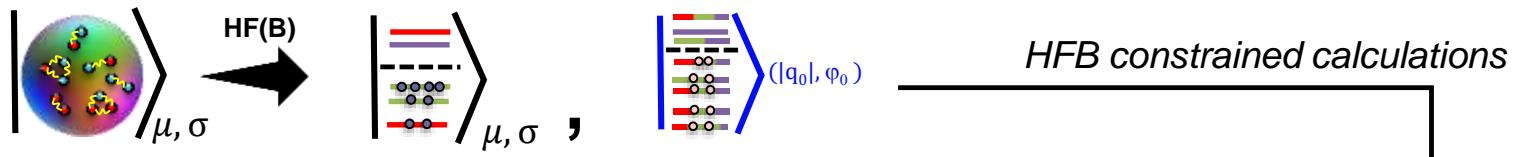
- Typical physical quantities computed at the HFB level :
 - ◆ Mass, radius, intrinsic density, ... of the GS (and some isomeric states)
 - ◆ Single-particle energies and wfs
 - ◆ Barrier, inertia tensor, ...
- Refined description of properties already accessible at the HFB level, or access to new types of properties (essentially spectroscopic ones) require going beyond the HFB realization



Horizontal expansion

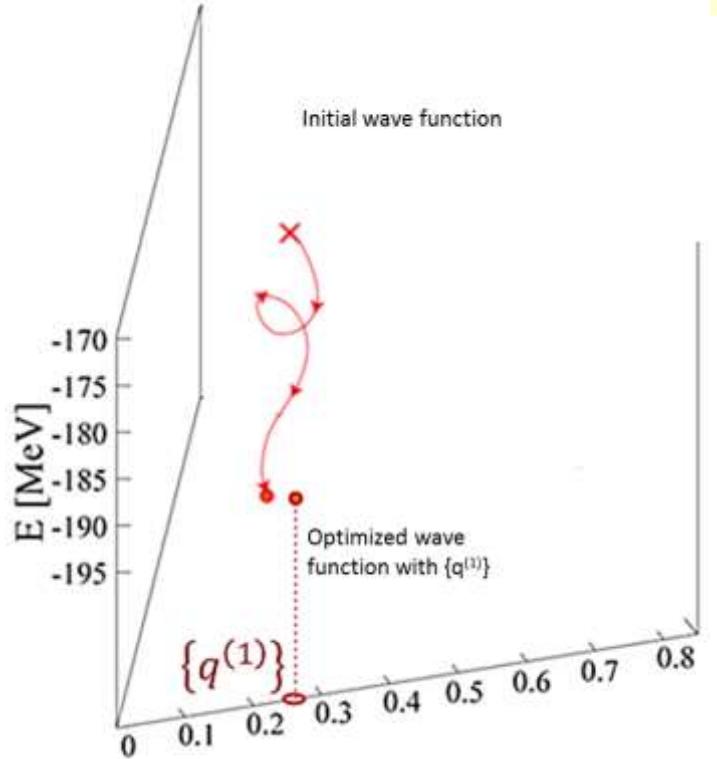
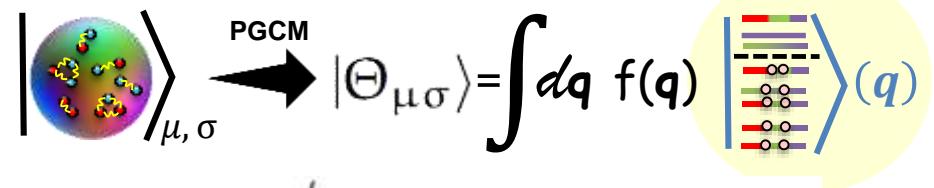
- HFB treatment

→ A -nucleon problem → A 1-nucleon problems



- Post-HFB treatment : PGCM

→ Symmetry-conserving (non orthogonal) mixture of symmetry-breaking HFB vacua

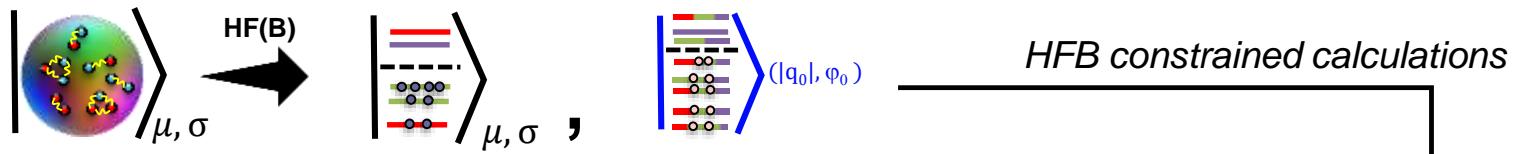




Horizontal expansion

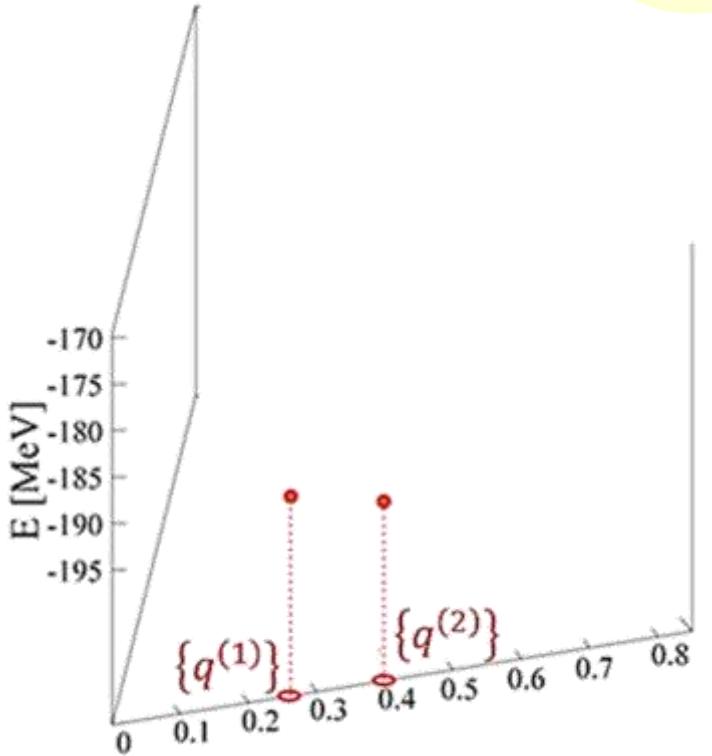
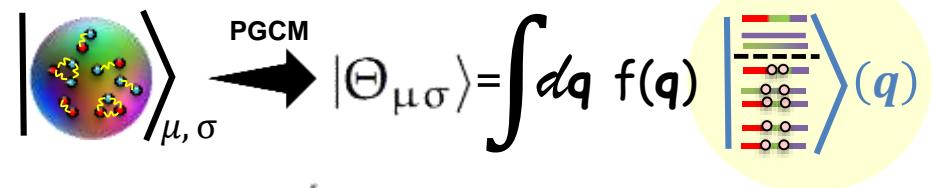
- HFB treatment

→ A -nucleon problem → A 1-nucleon problems



- Post-HFB treatment : PGCM

→ Symmetry-conserving (non orthogonal) mixture of symmetry-breaking HFB vacua

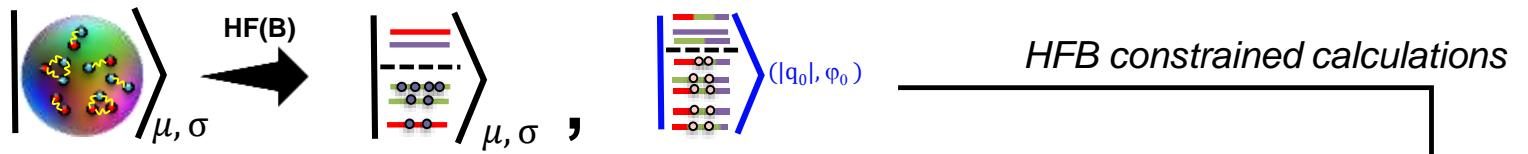




Horizontal expansion

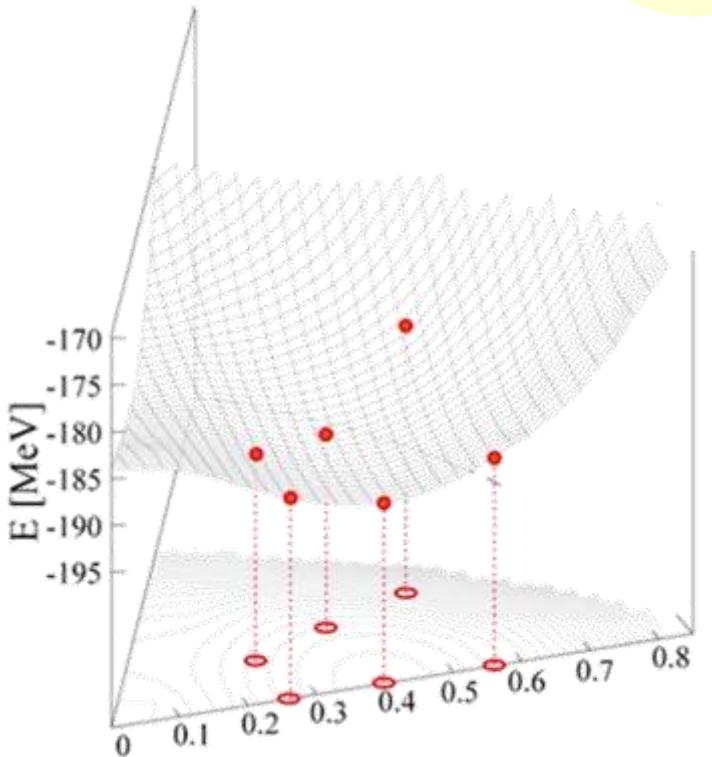
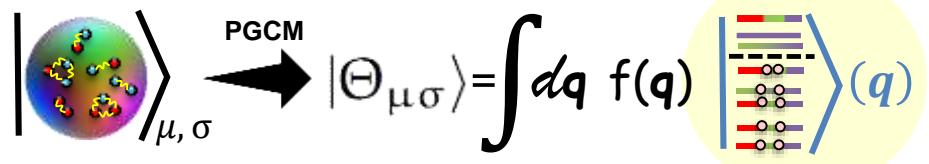
- HFB treatment

→ A -nucleon problem → A 1-nucleon problems



- Post-HFB treatment : PGCM

→ Symmetry-conserving (non orthogonal) mixture of symmetry-breaking HFB vacua

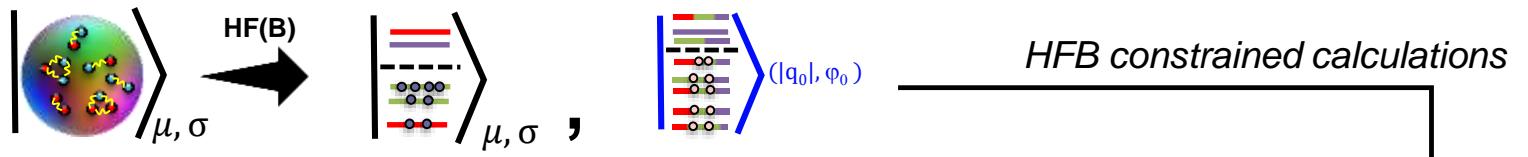




Horizontal expansion

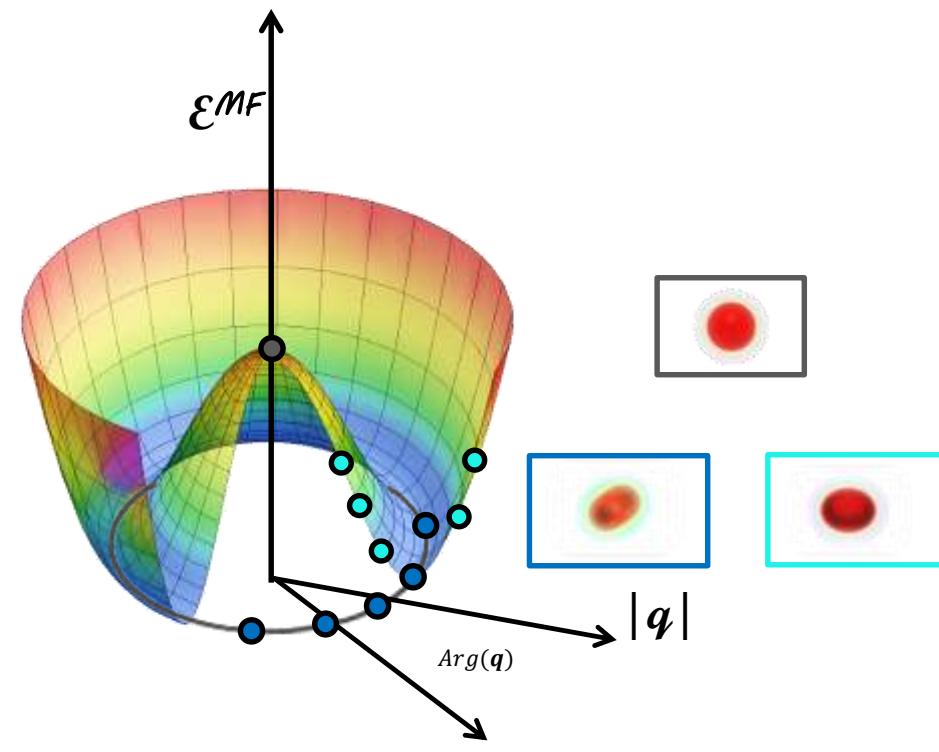
- HFB treatment

→ A -nucleon problem → A 1-nucleon problems

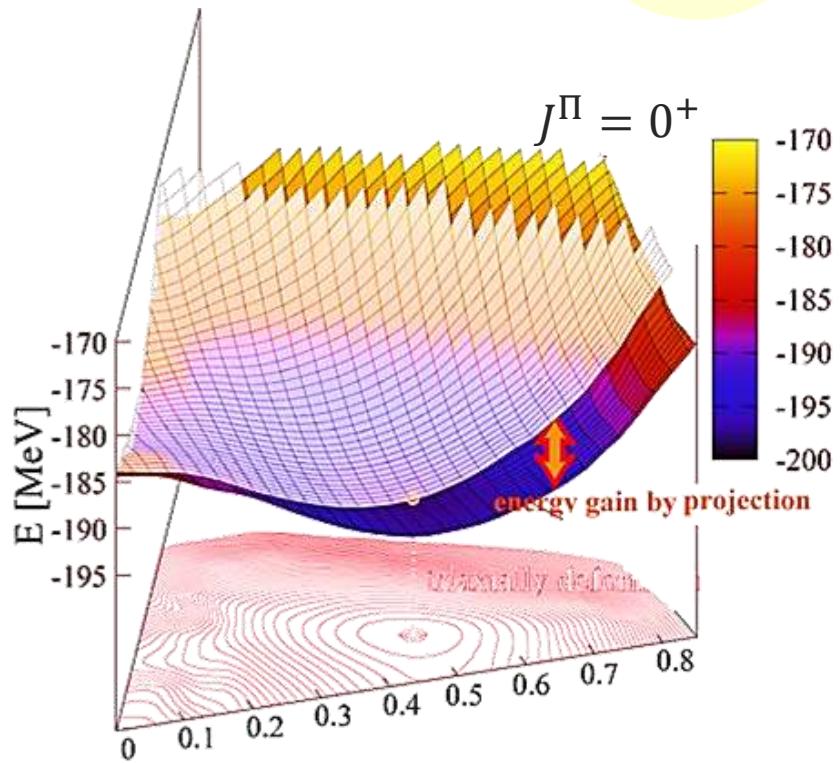


- Post-HFB treatment : PGCM

→ Symmetry-conserving (non orthogonal) mixture of symmetry-breaking HFB vacua



$$|\Theta_{\mu\sigma}\rangle = \int d\mathbf{q} f(\mathbf{q}) |\langle \mathbf{q} | (q)\rangle_{\mu\sigma}$$

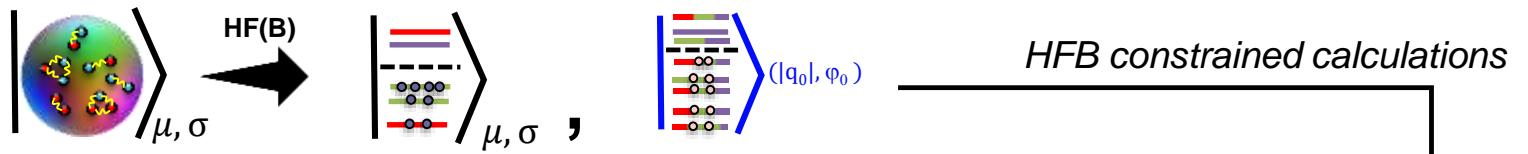




Horizontal expansion

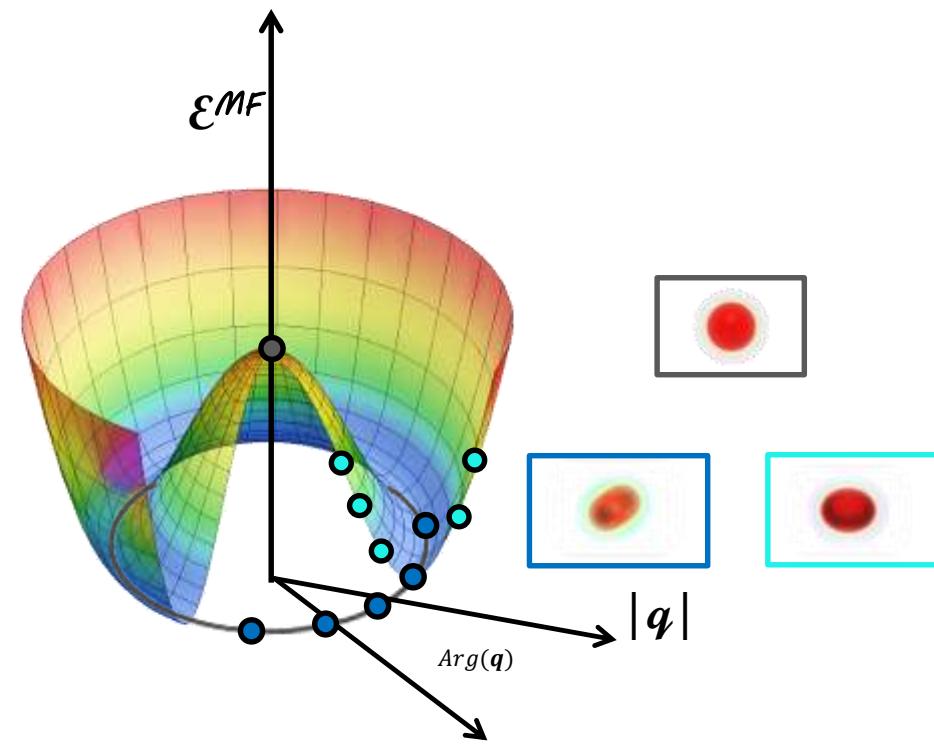
- HFB treatment

→ A -nucleon problem → A 1-nucleon problems

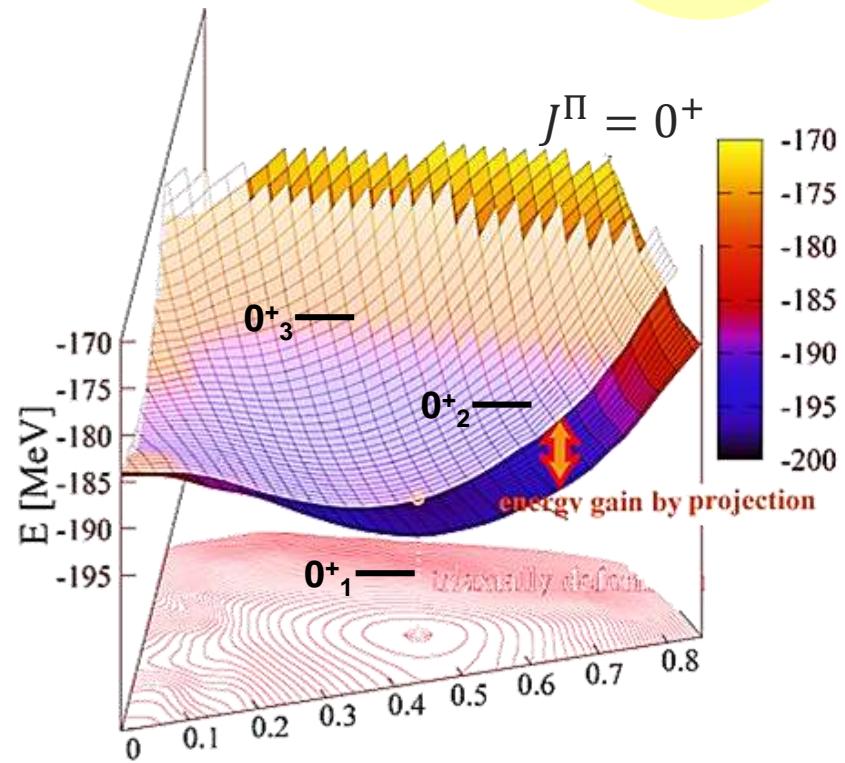


- Post-HFB treatment : PGCM

→ Symmetry-conserving (non orthogonal) mixture of symmetry-breaking HFB vacua



$$|\Theta_{\mu\sigma}\rangle = \int d\mathbf{q} f(\mathbf{q}) |\langle \mathbf{q} | \rangle$$

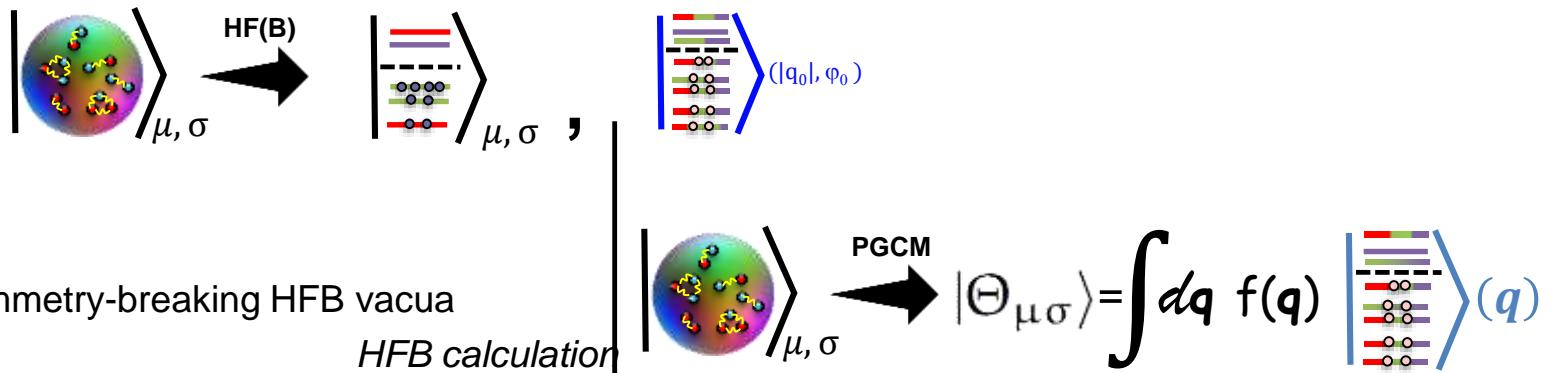




Horizontal expansion

- HFB treatment

→ A -nucleon problem → A 1-nucleon problems



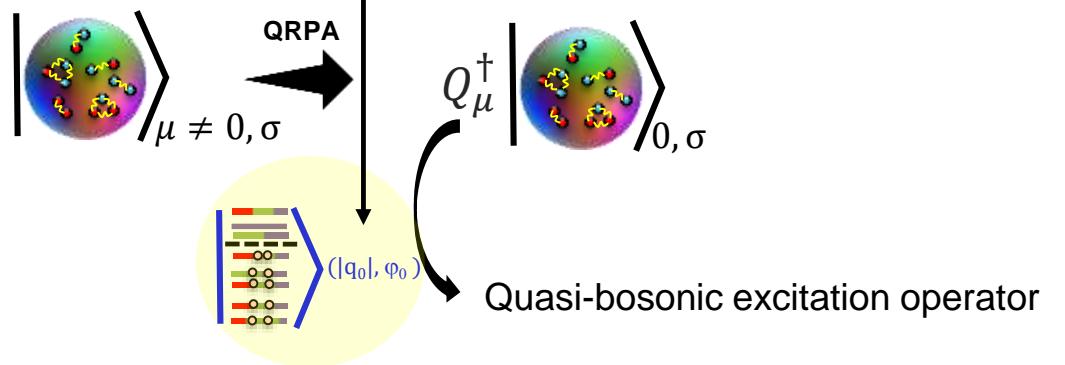
- Post-HFB treatment : PGCM

→ Symmetry-conserving (non orthogonal) mixture of symmetry-breaking HFB vacua

- Post-HFB : QRPA

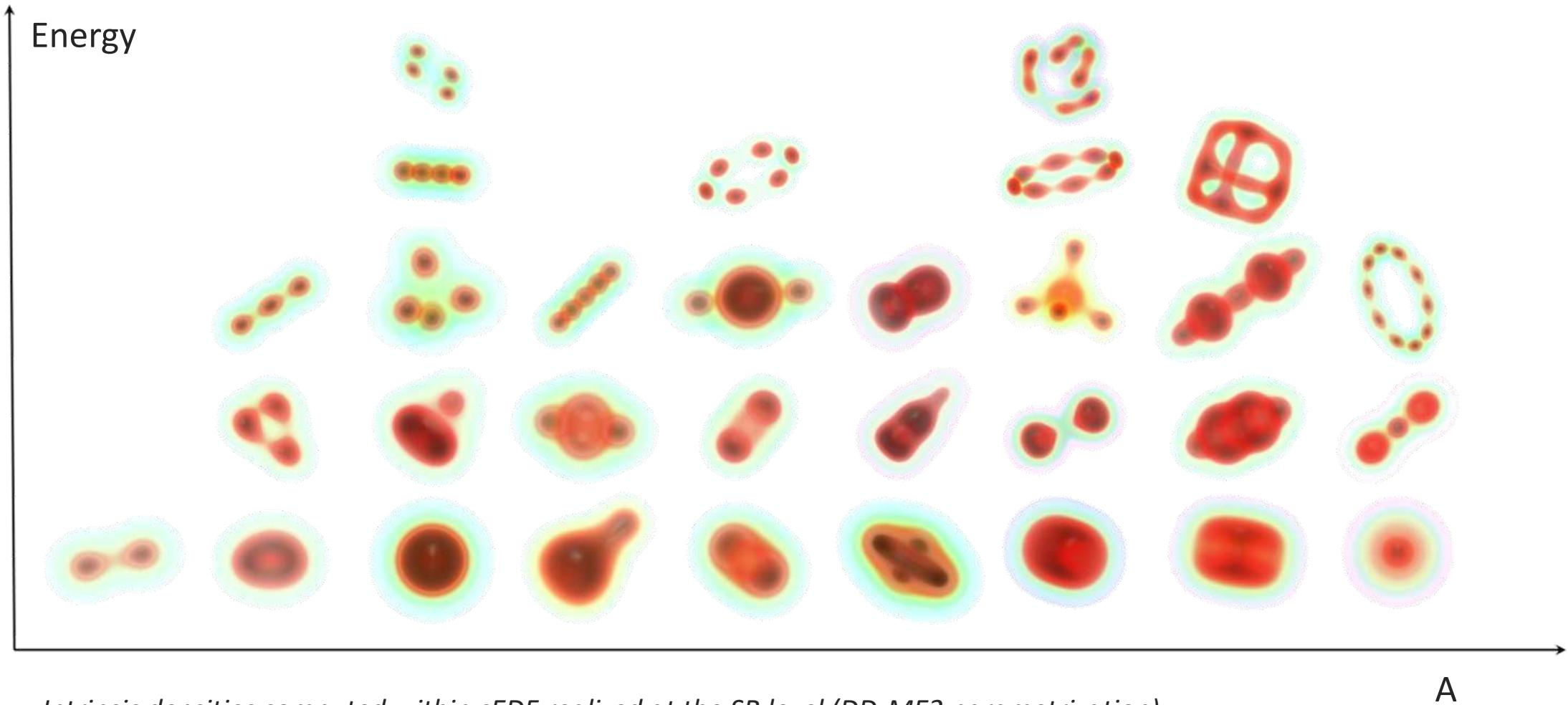
→ Excitations = coherent mixture of 2-qp excitations

→ Harmonic limit of the GCM



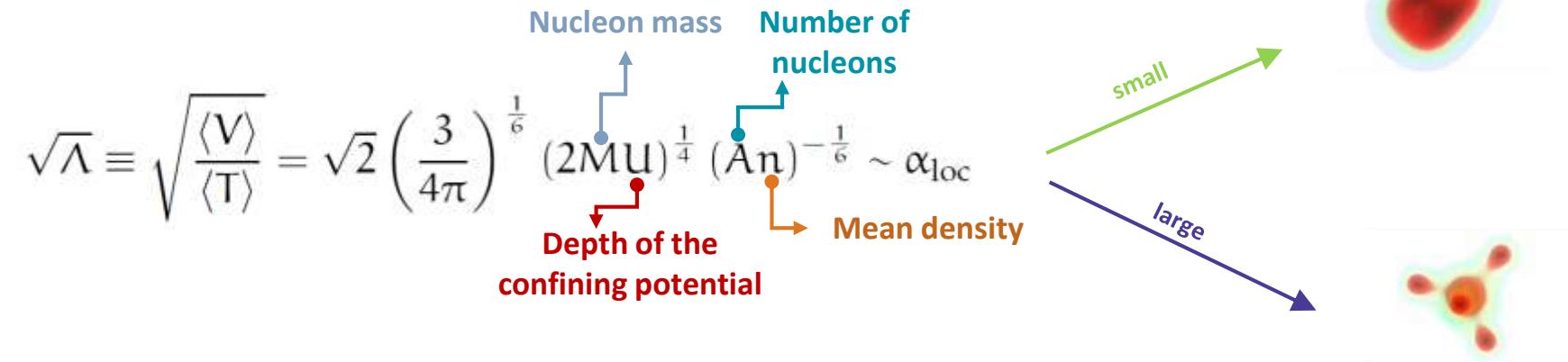
Example of applications : Nuclear clustering

- Clustering = nucleons clumping together into sub-groups within the nucleus



Strength of correlations

- Strength of correlations measured by dimensionless ratios



Clustering favored

- For deep confining potential
- For light nuclei
- In regions at low-density

- Formation/dissolution of clusters : Mott parameter

$$\frac{R_X}{d_{Mott}^X} \sim 1 \Rightarrow n_{Mott}^X \sim \frac{\rho_{sat}}{A_X}$$

Size of the nucleus X
inter-nucleon average distance

$$n_{Mott}^\alpha \sim 0.25\rho_{sat}$$

$$\sim \frac{\rho_{sat}}{3}$$

Size of an α in free-space

0.9 size of an α in free-space

Ebran, Girod, Khan, Lasseri, Schuck, PRC 2020

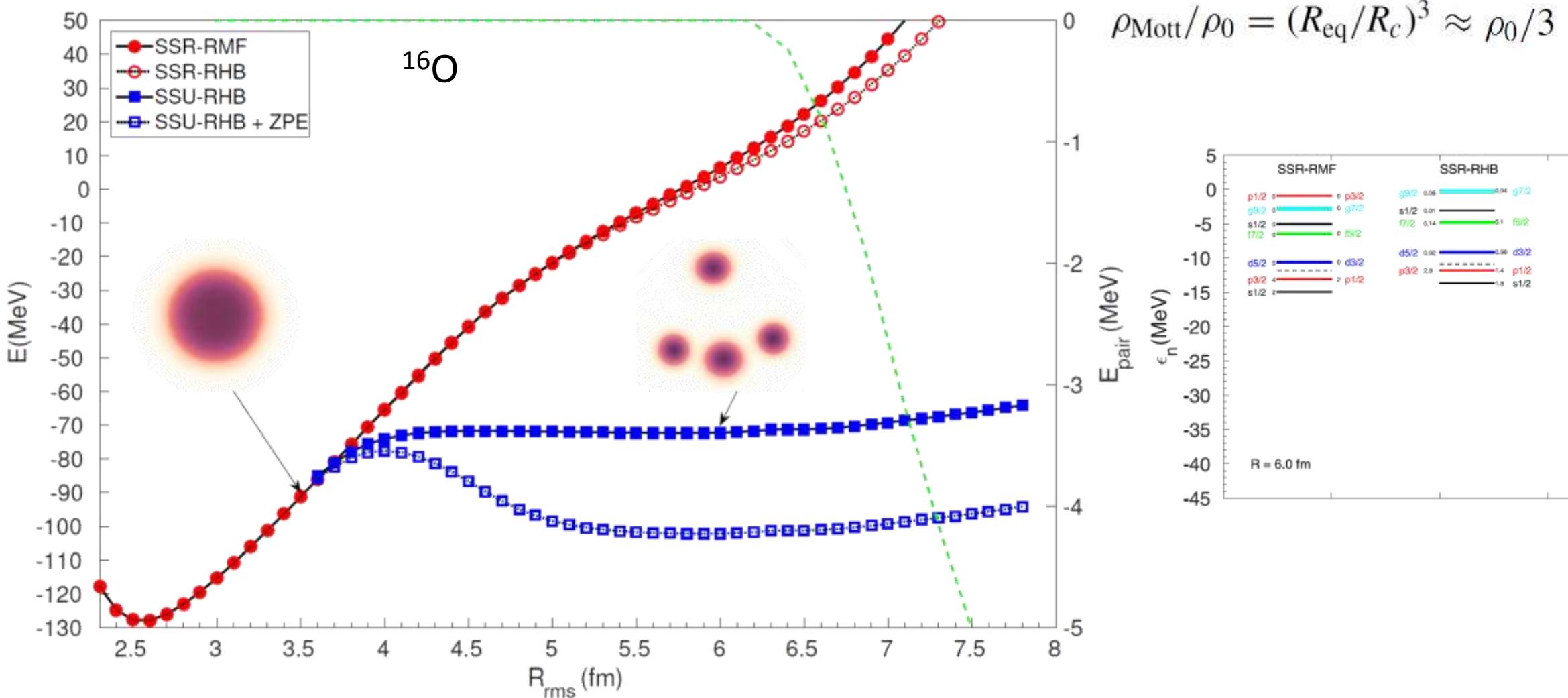
Ebran, Khan, Niksic, Vretenar, PRC 2014

Ebran, Khan, Niksic & Vretenar PRC 2013

Ebran, Khan, Niksic & Vretenar Nature 2012

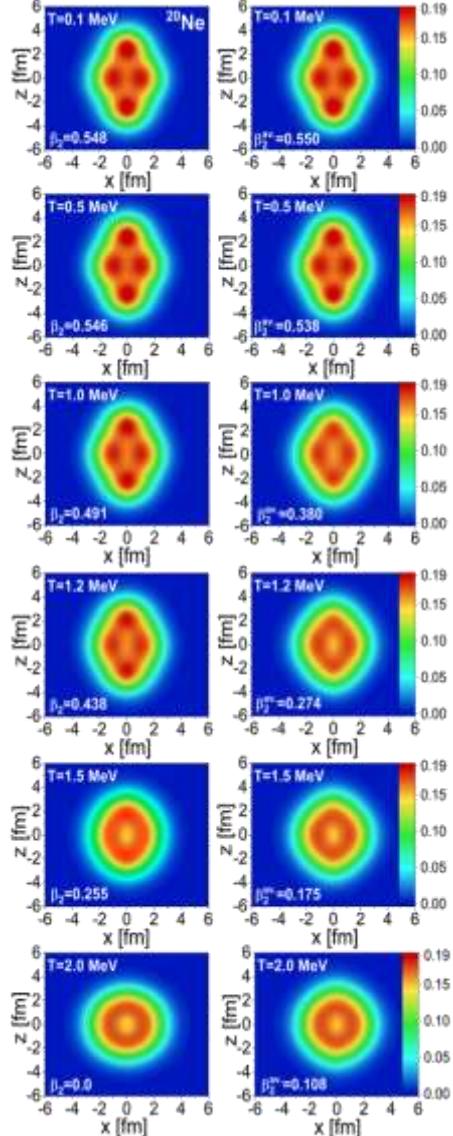
Quantum Mott-like phase transition

- Isotropically inflate ^{16}O by constraining its r.m.s. radius while imposing a global quadrupole moment to be zero

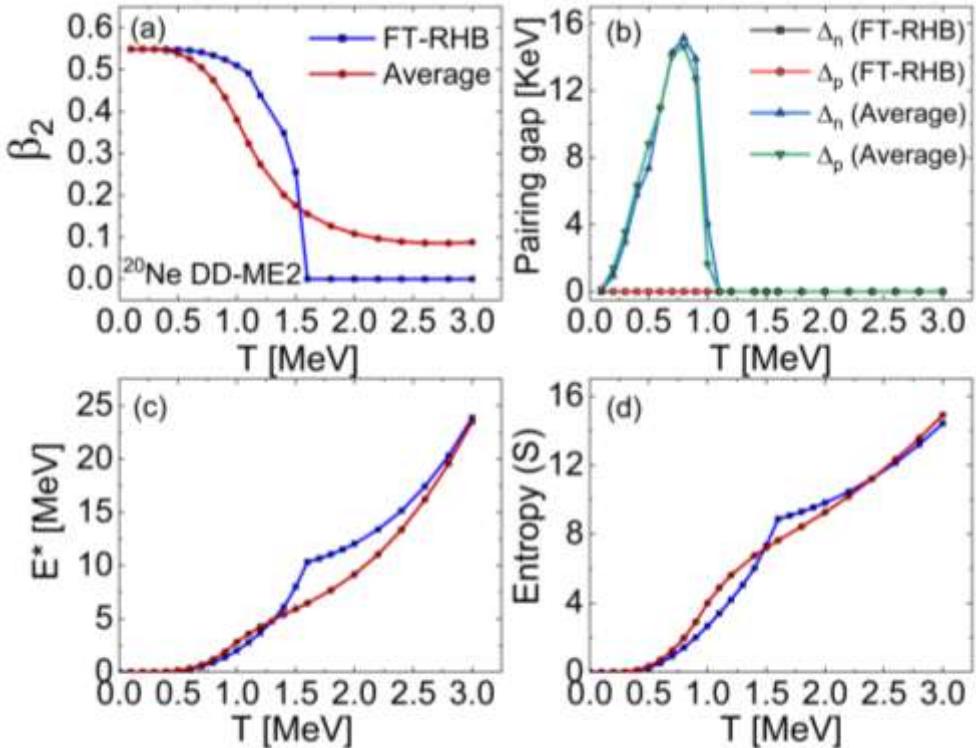


Thermal phase transition

- Isotropically inflate ^{16}O by constraining its r.m.s. radius while imposing a global quadrupole moment to be zero



$$\overline{O} = \frac{\int d\beta_2 O(\beta_2, T) \exp(-\Delta F(\beta_2, T)/T)}{\int d\beta_2 \exp(-\Delta F(\beta_2, T)/T)}.$$



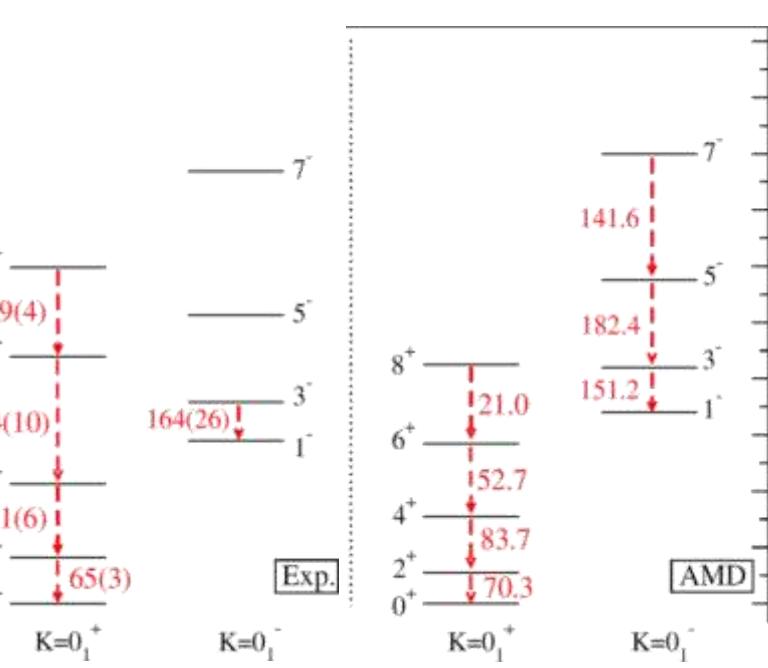
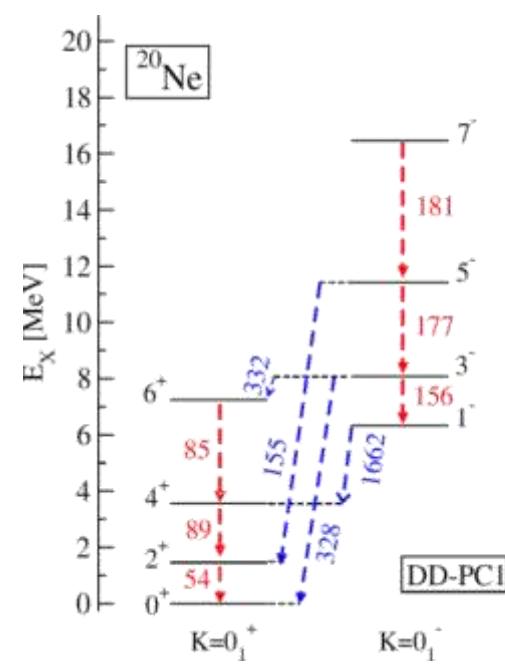
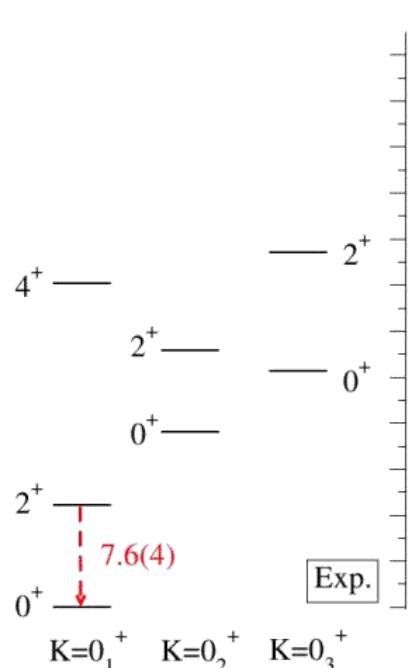
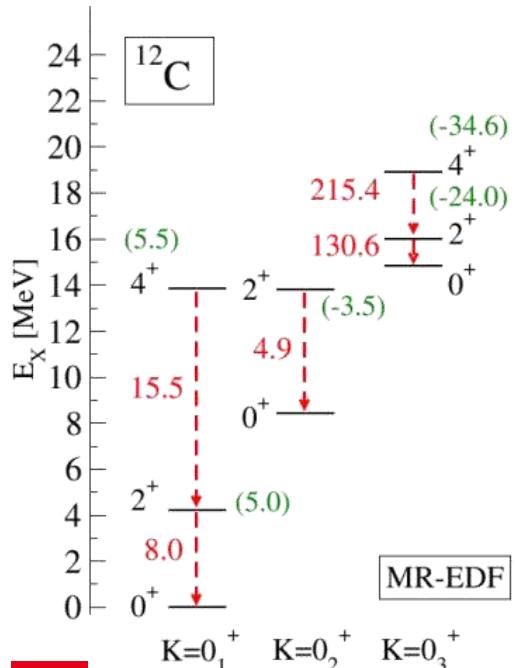
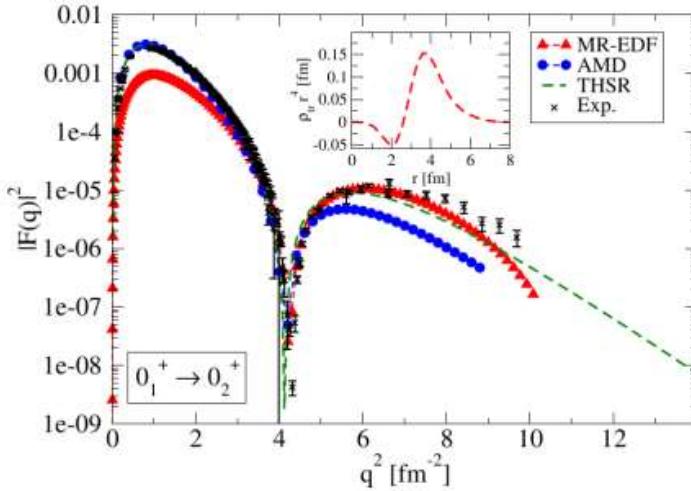
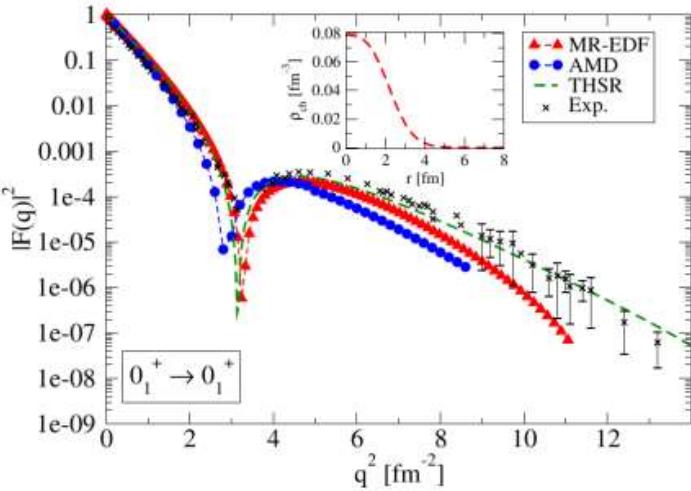
Yüksel, Mercier, Ebran, Khan PRC 2022

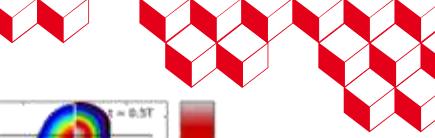


Nuclear clustering & PGCM

● Spectroscopy

Marević, Ebran, Khan, Nikšić, and Vretenar, 2019

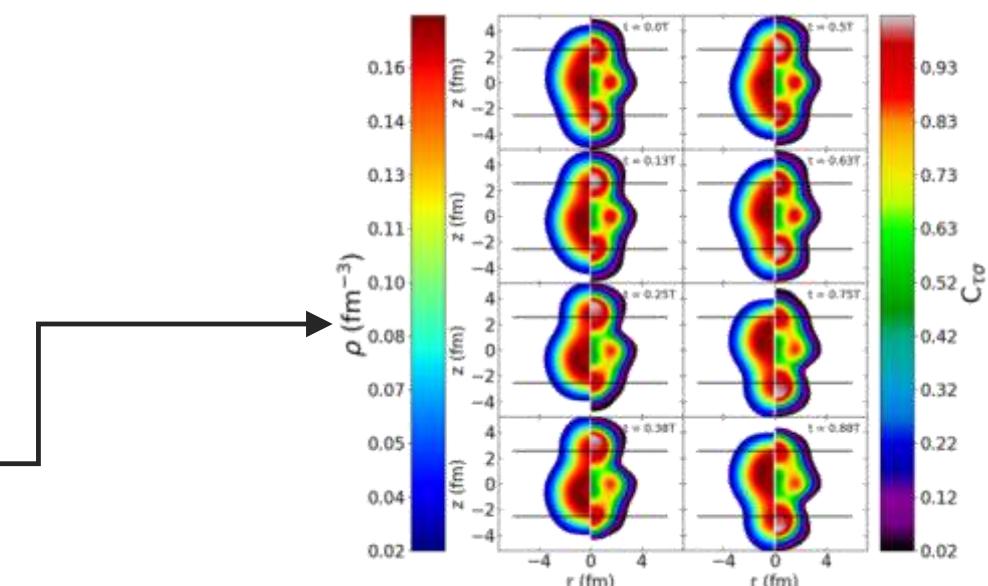
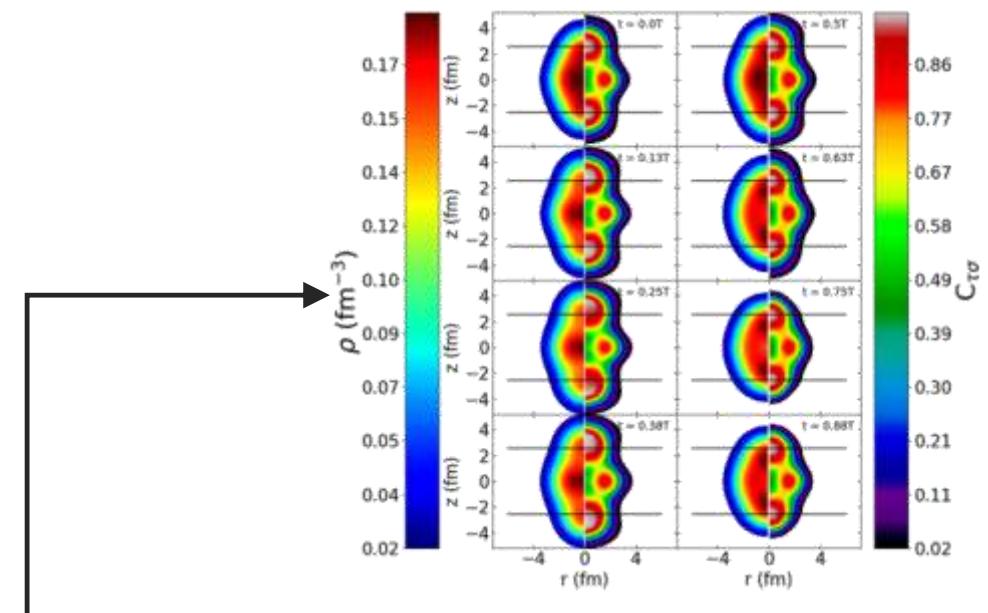
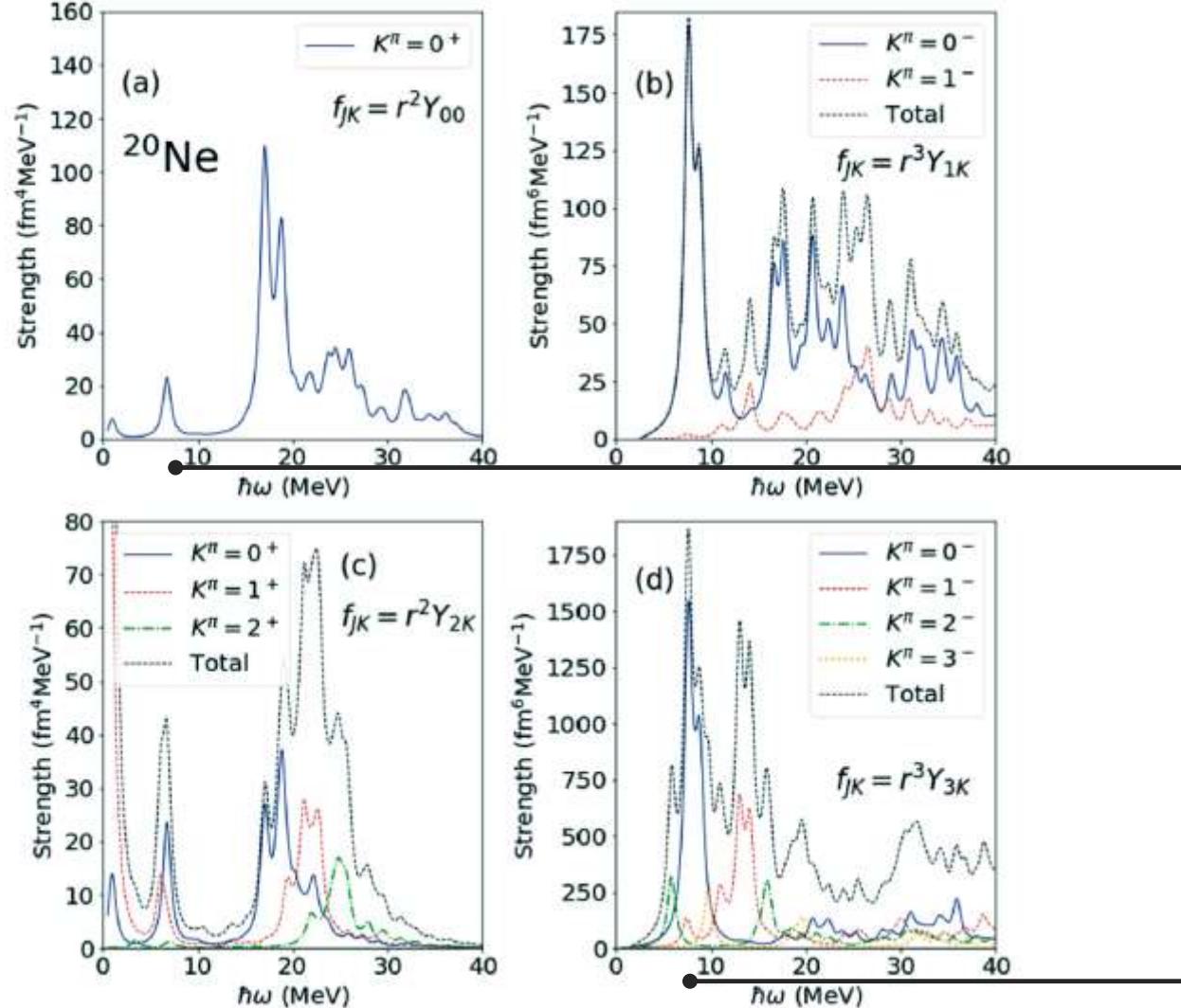




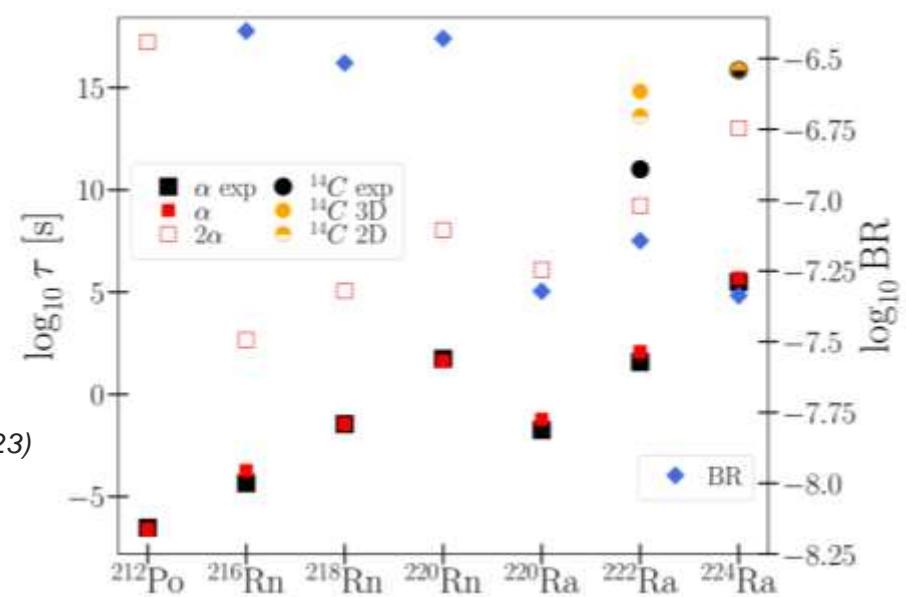
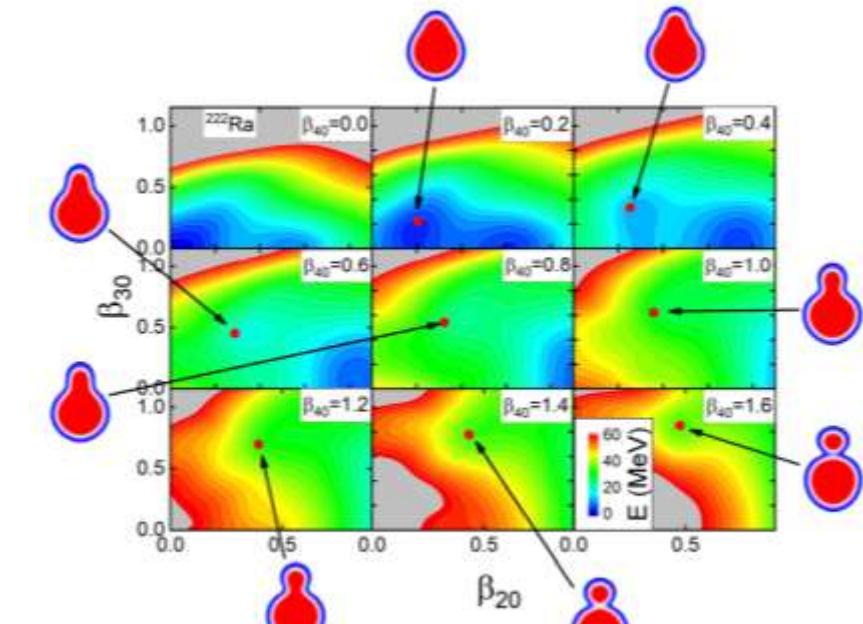
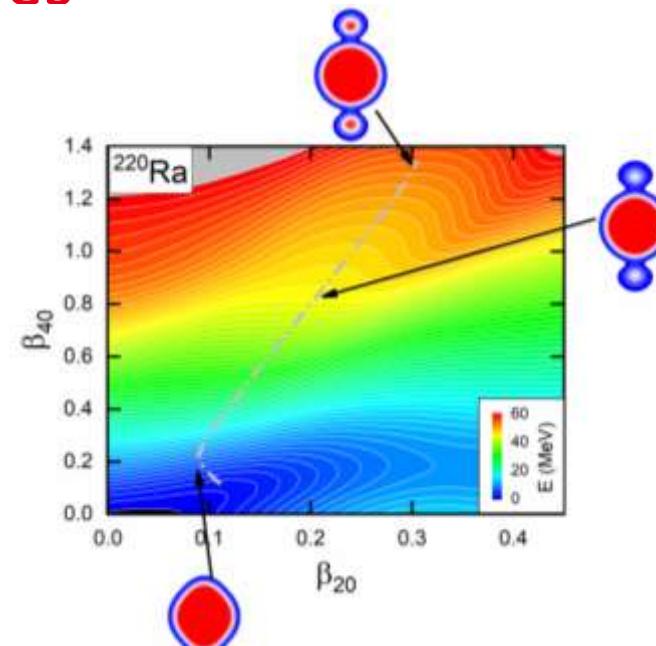
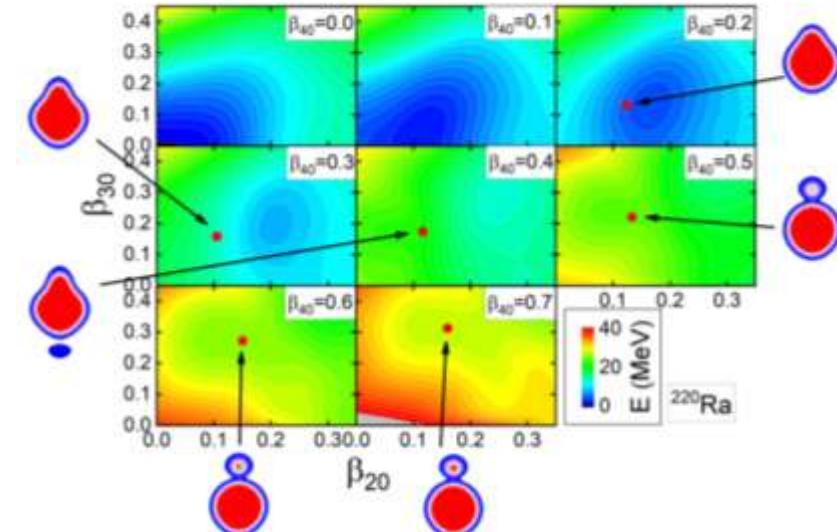
Nuclear clustering & QRPA

● Cluster vibration

Mercier, Bjelčić, Nikšić, Ebran, Khan, Vretenar 2021
Mercier, Ebran, Khan 2022



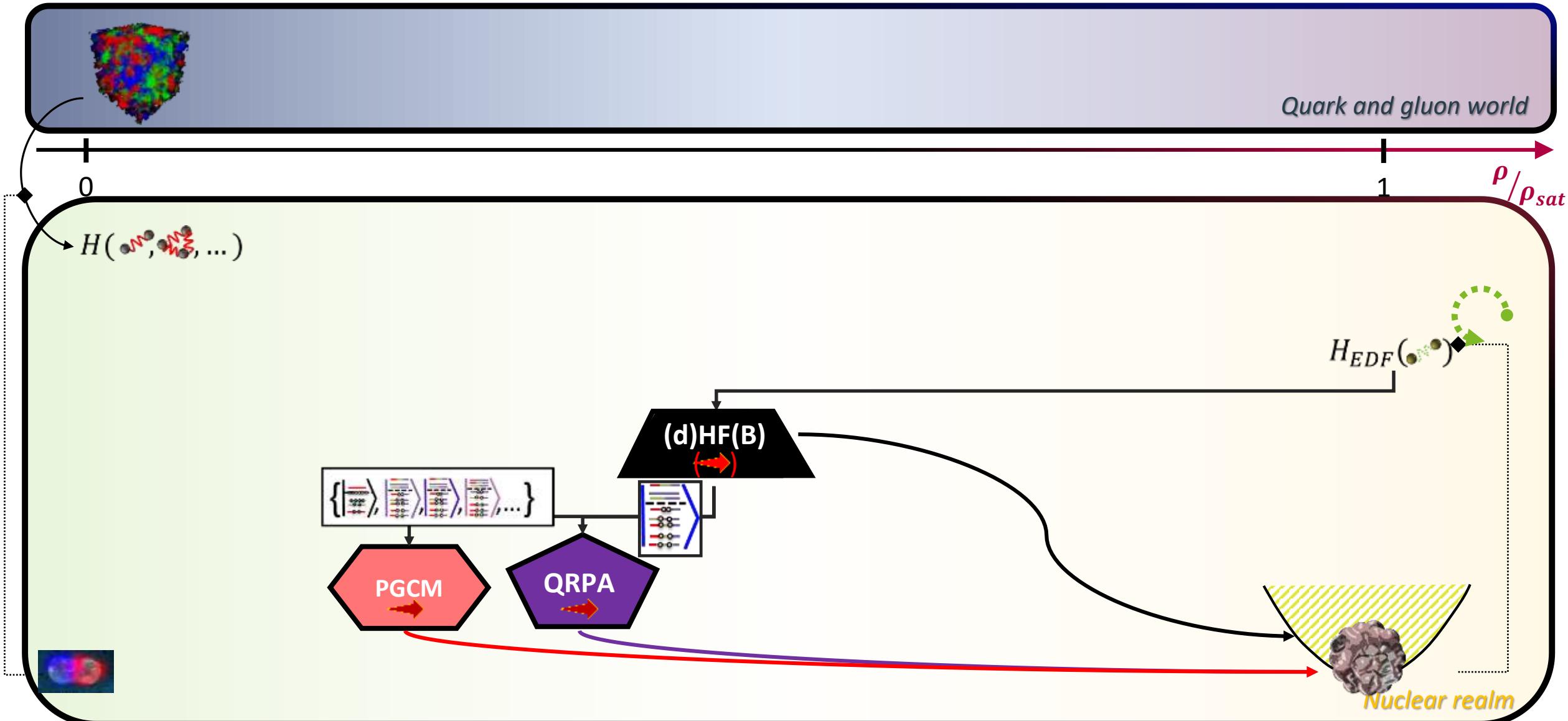
Cluster, α and 2α radioactivities



Zhao, Ebran, Heitz , Khan , Mercier, Nikšić, Vretenar PRC (2023)
 Mercier, Zhao, Ebran, Khan, Nikšić, Vretenar PRL (2021)



EDF workflow



Outline

- 1. General context**
- 2. Recent work on empirical EDFs**
- 3. EDF-inspired ab initio methods**
- 4. Towards a first-principle formulation of EDFs**



3 ■ EDF-inspired ab initio methods



Ab initio strategy

● Solve in a controlled way, to some desired accuracy $H(\bullet\text{-}\bullet, \bullet\text{-}\bullet, \dots) |\Psi_{\mu,\sigma}\rangle = E_{\mu\tilde{\sigma}} |\Psi_{\mu,\sigma}\rangle$

Which part of correlations should be treated here ?

--> Ab initio WFT : Expansion many-body methods

$$H = H_0 + H_1$$

$$|\Theta\rangle$$

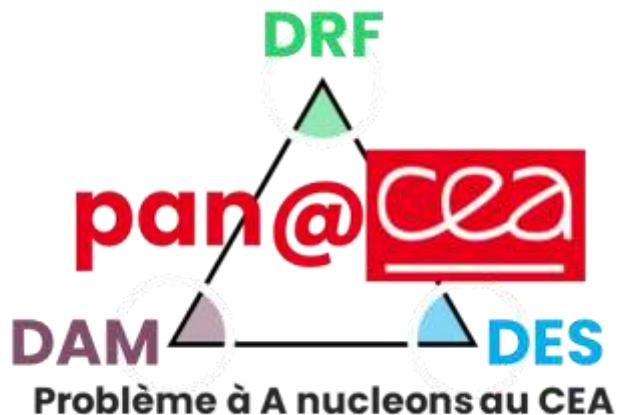
$$|\Psi\rangle = \Omega|\Theta\rangle$$

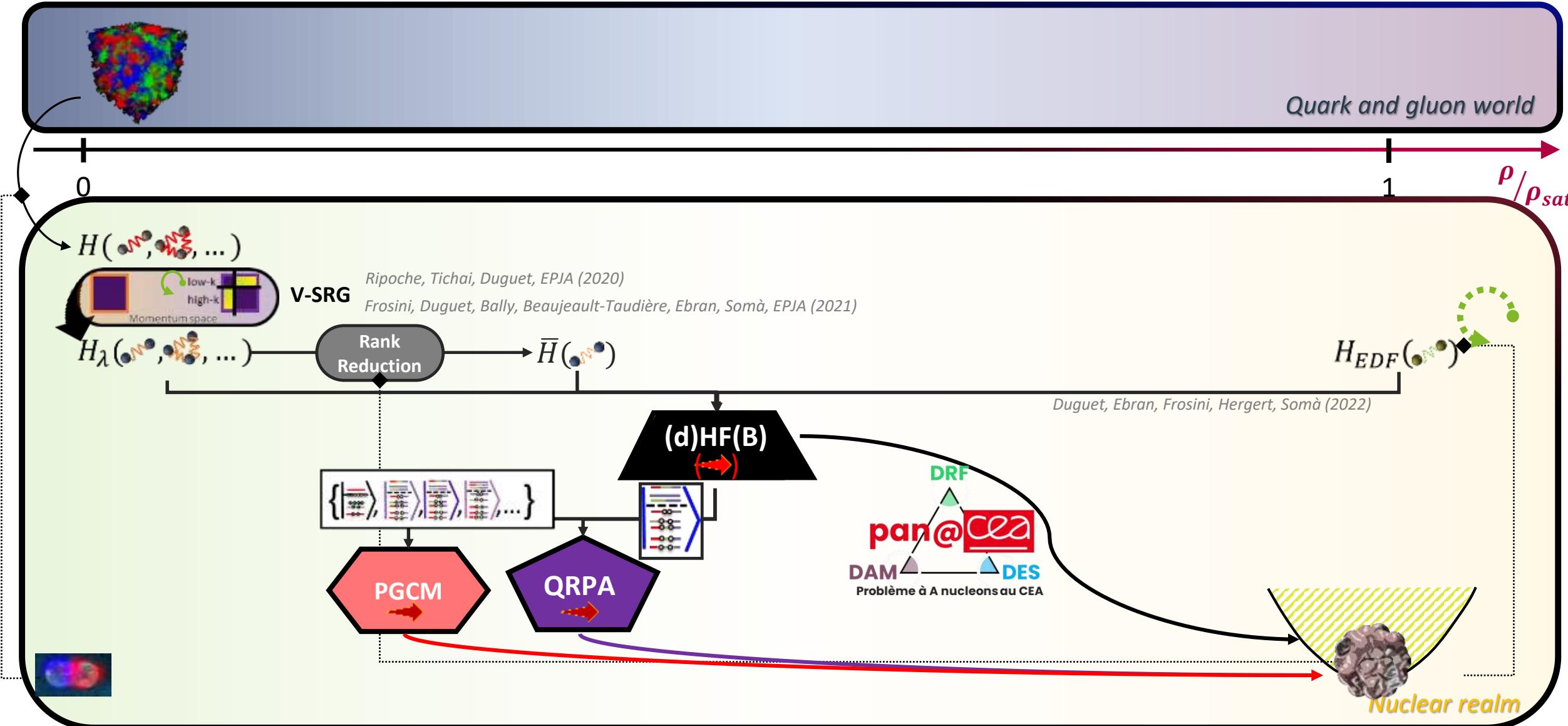
--> Get inspiration from EDFs to design new expansion methods working for both closed- and open-shell nuclei

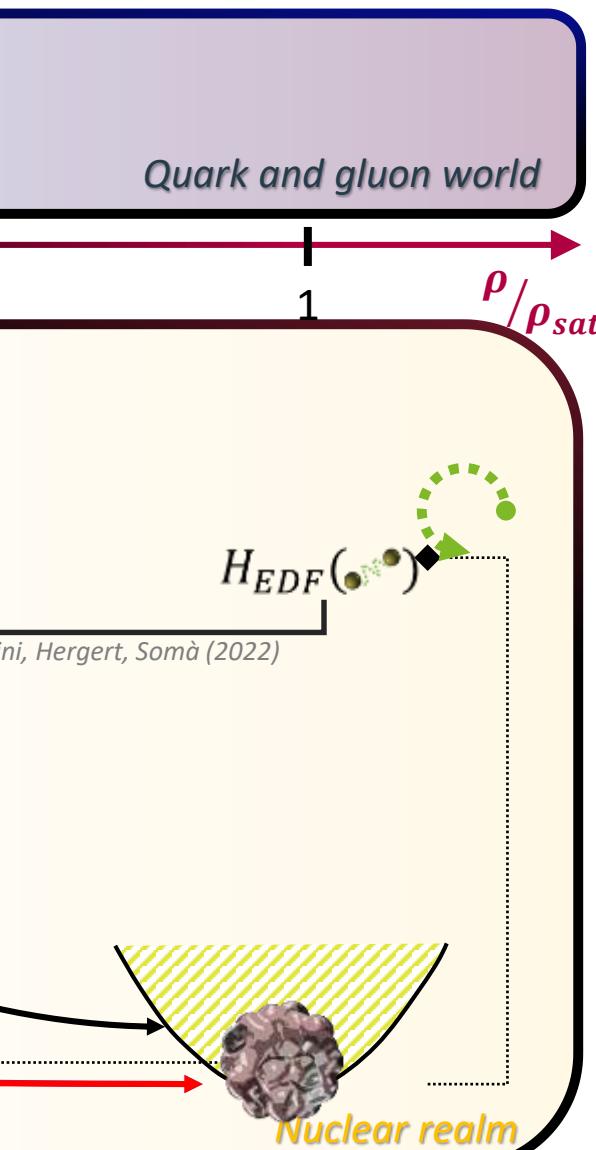
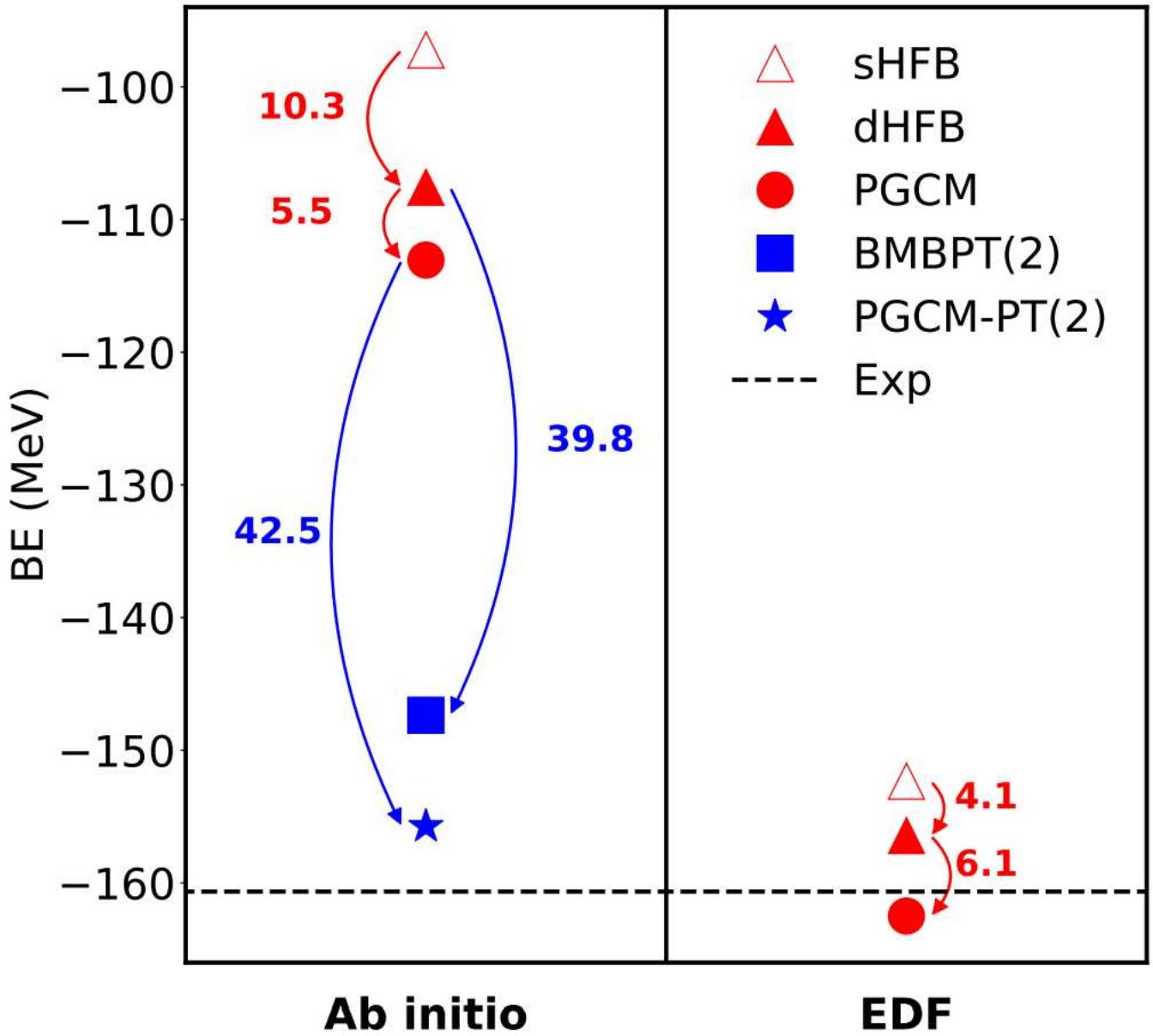
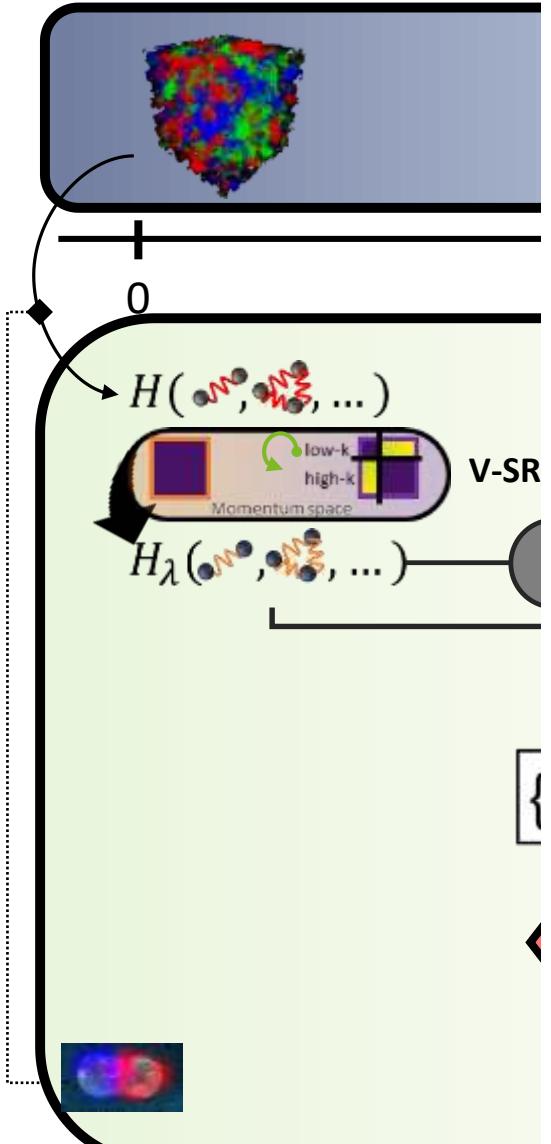
Pioneering work by T. Duguet

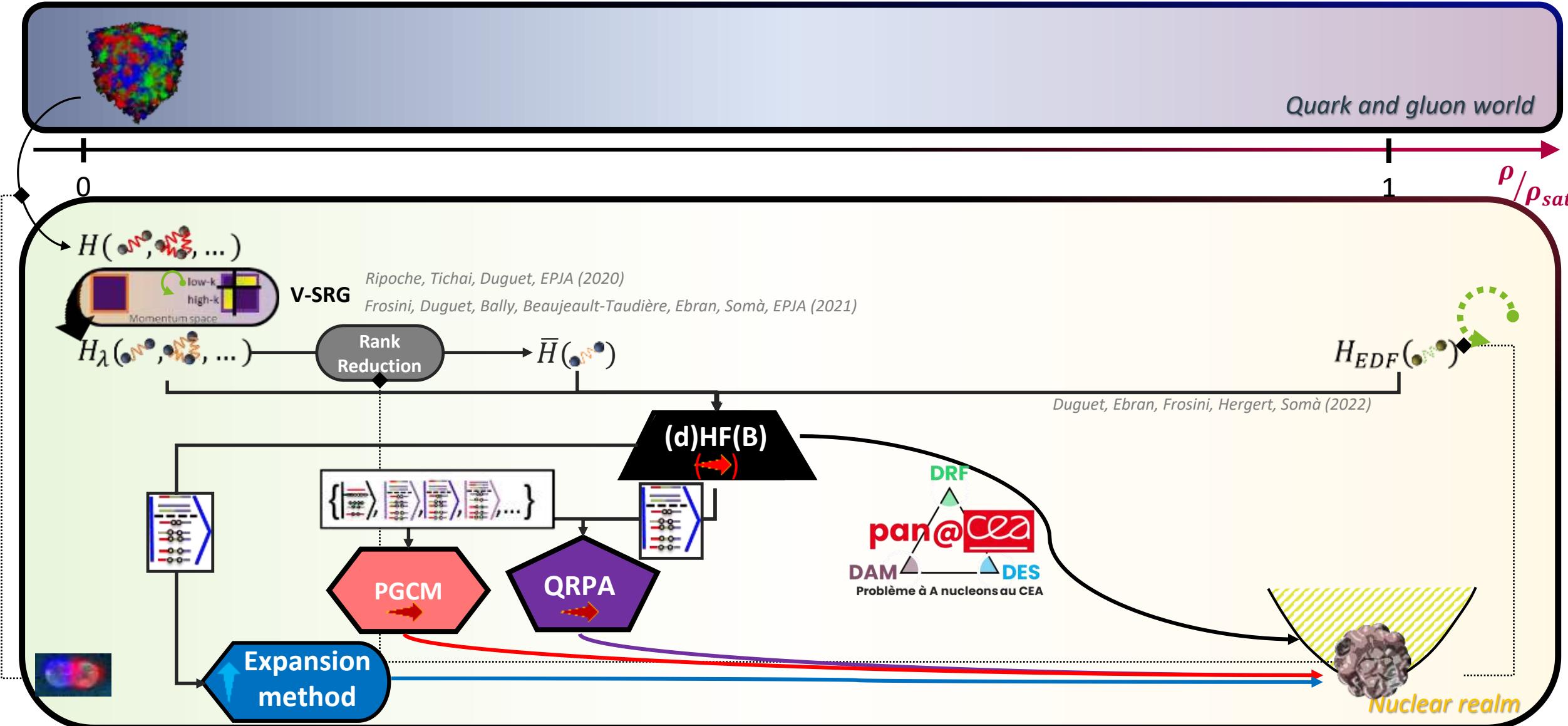


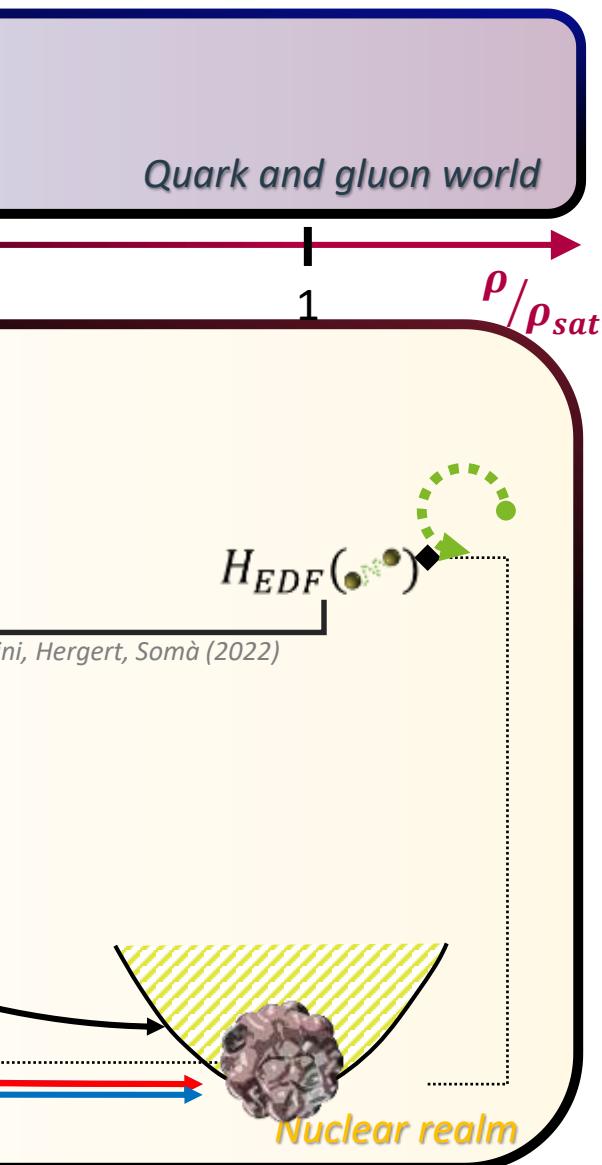
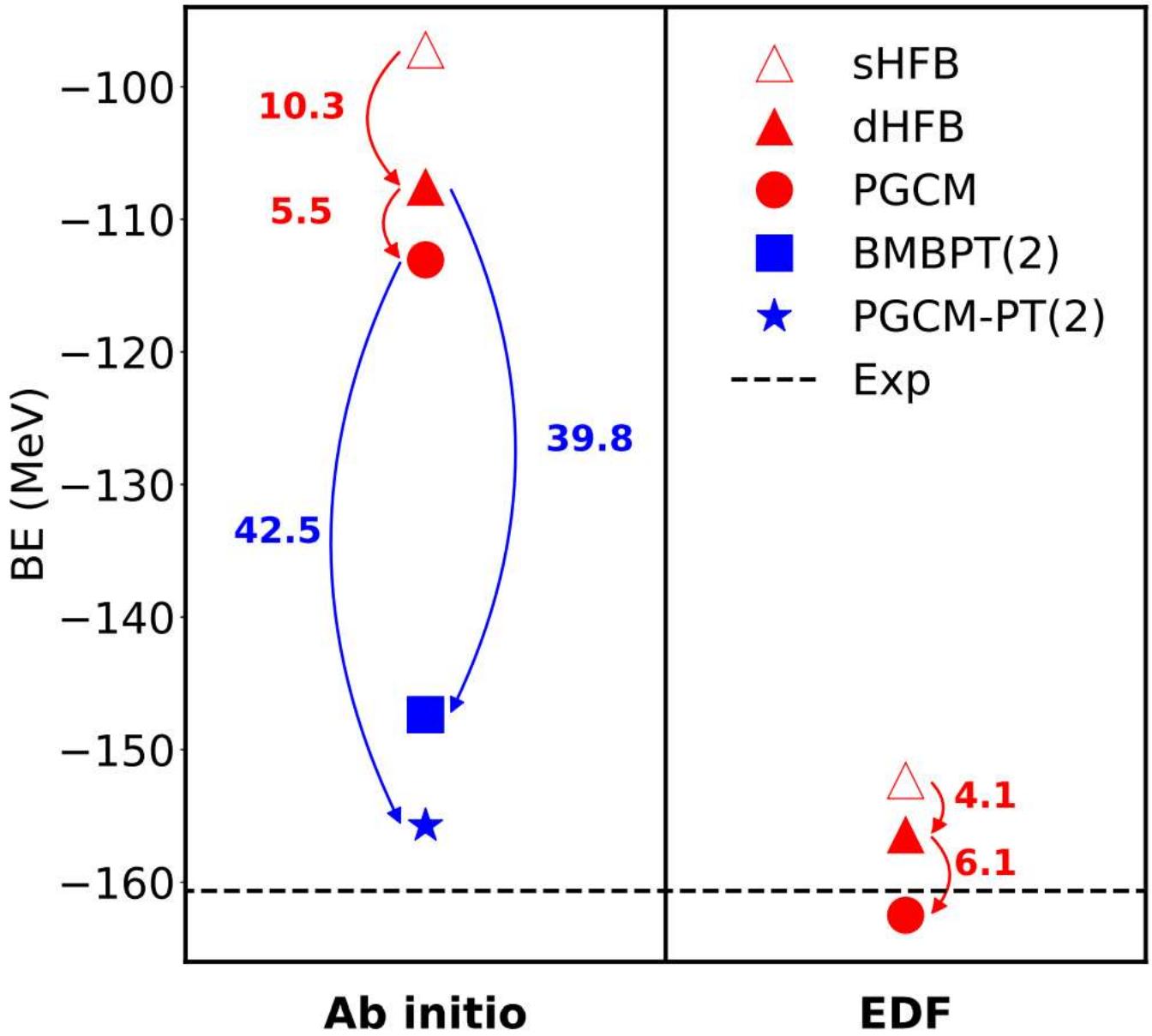
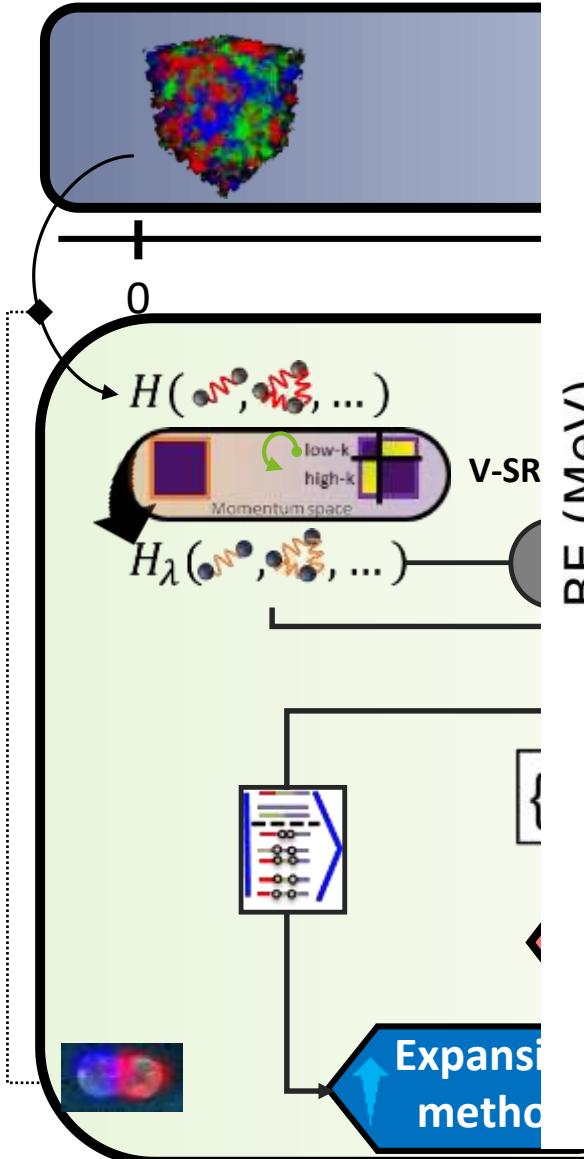
Somà, Barbieri, Duguet, PRC (2014)
 Duguet, JPG:NPP (2014)
 Signoracci, Duguet, Hagen, Jansen PRC (2015)
 Duguet, Signoracci JPG (2017)
 Tichai, Arthuis, Duguet, Hergert, Somà, Roth PLB (2018)
 Arthuis, Duguet, Tichai, Lasseri, Ebran CPC (2019)

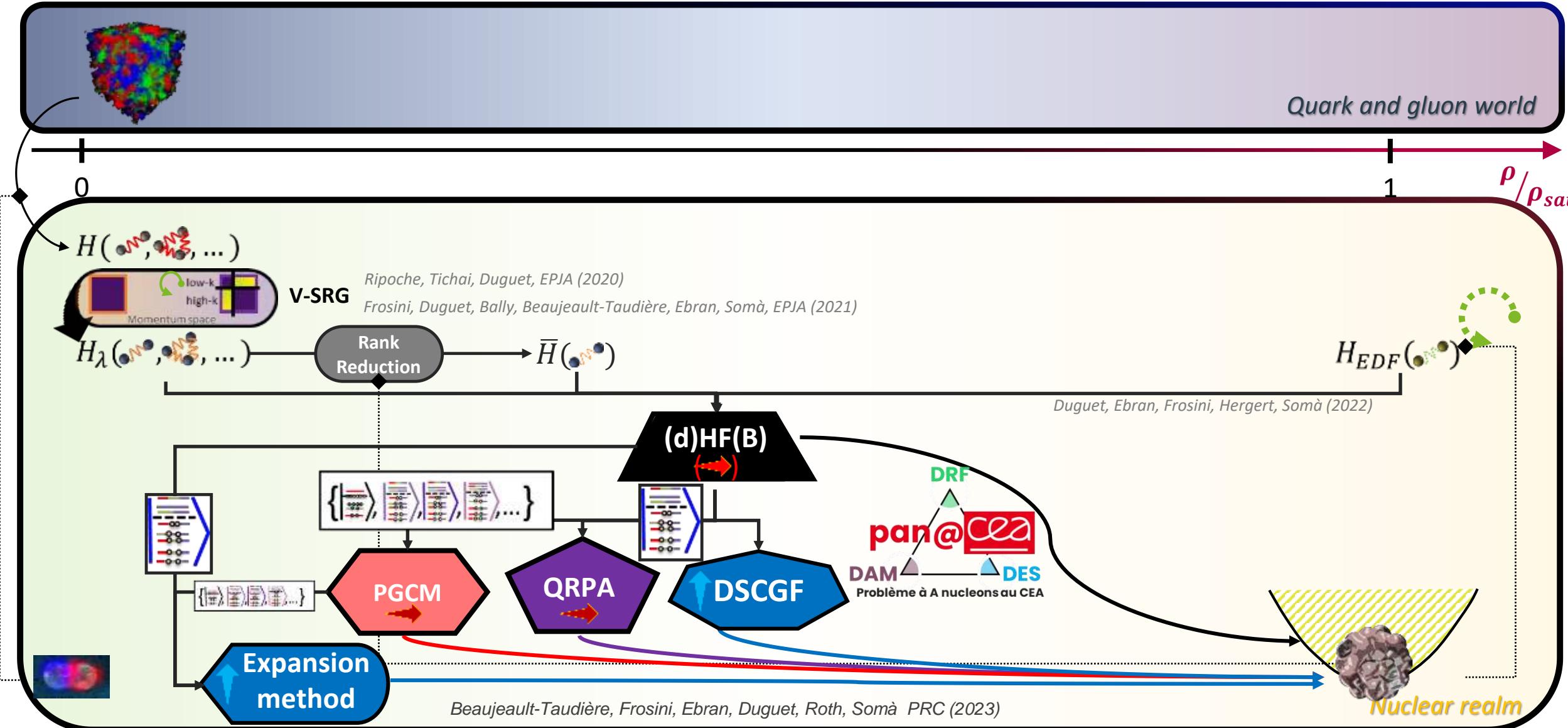








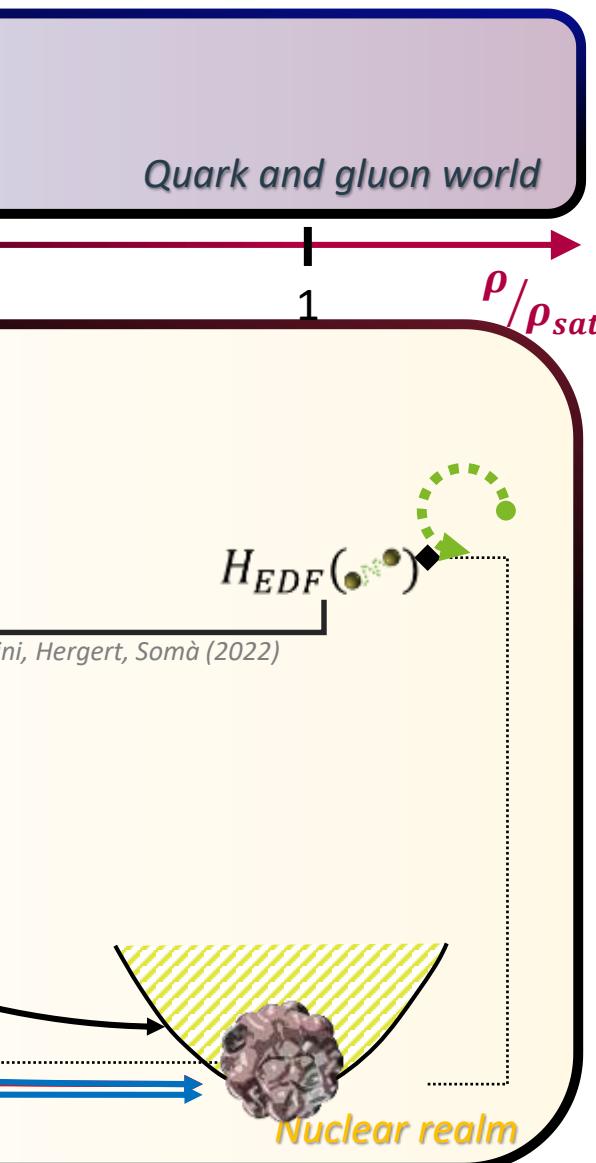
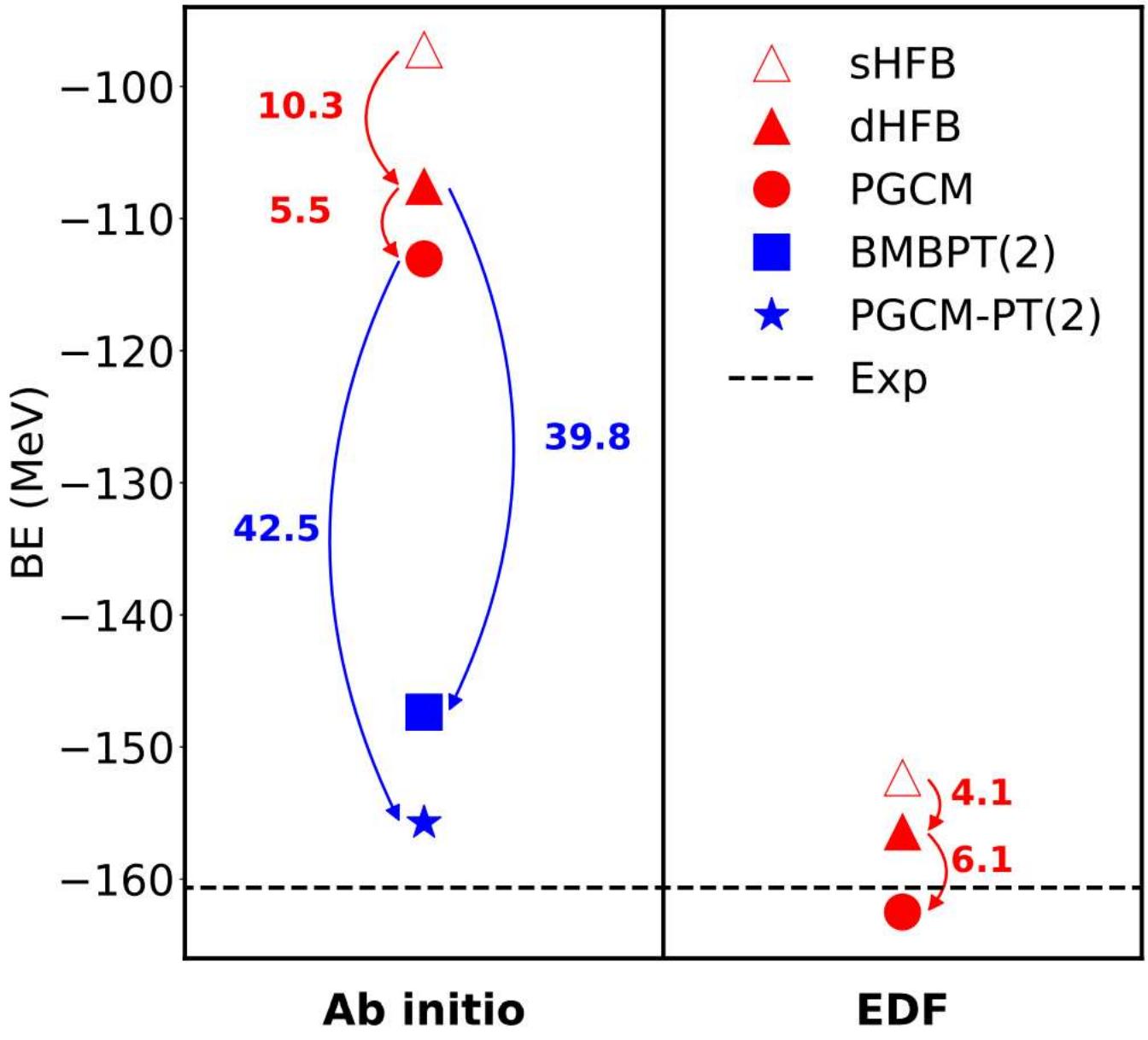
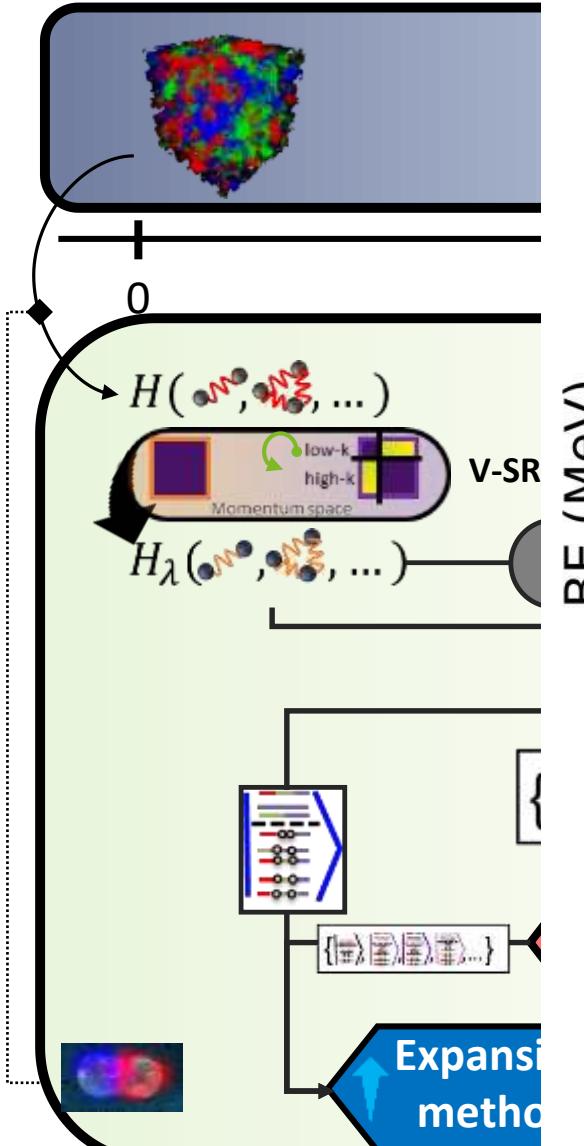




Frosini, Duguet, Ebran, Somà, EPJA (2022)

Frosini, Duguet, Ebran, Bally, Mongelli, Rodriguez, Roth, Somà, EPJA (2022)

Frosini, Duguet, Ebran, Bally, Hergert, Rodriguez, Roth, Yao, Somà, EPJA (2022)

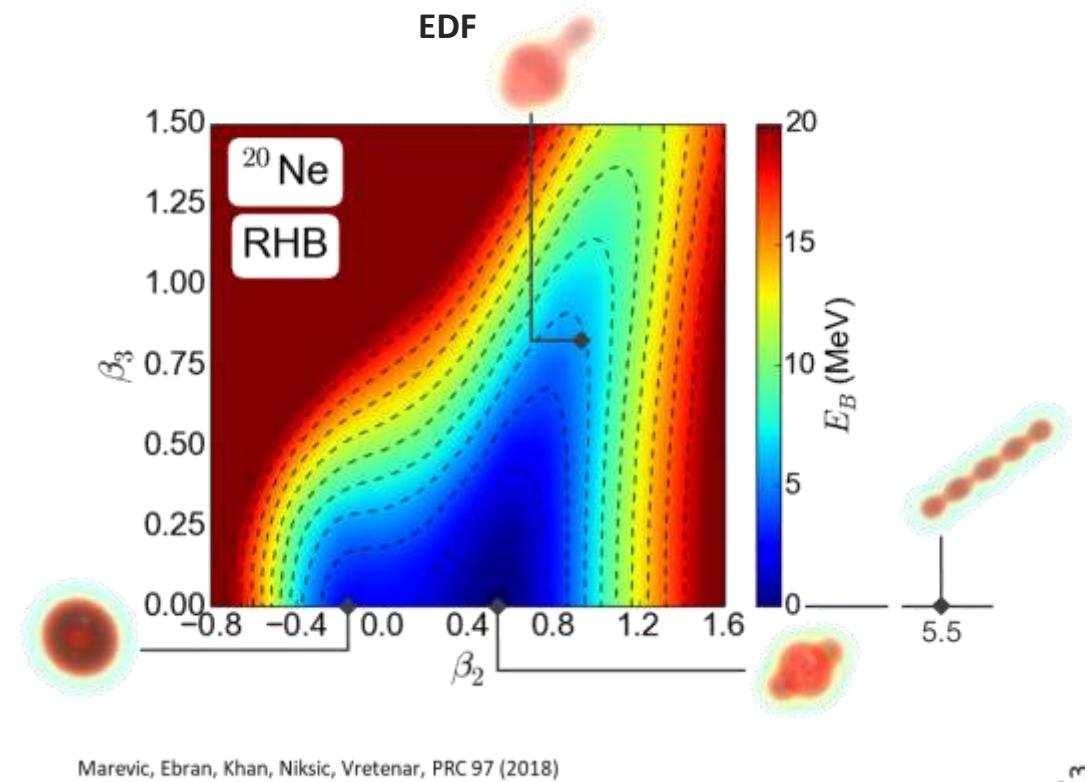


Frosini, Duguet, Ebran, Somà, EPJA (2022)

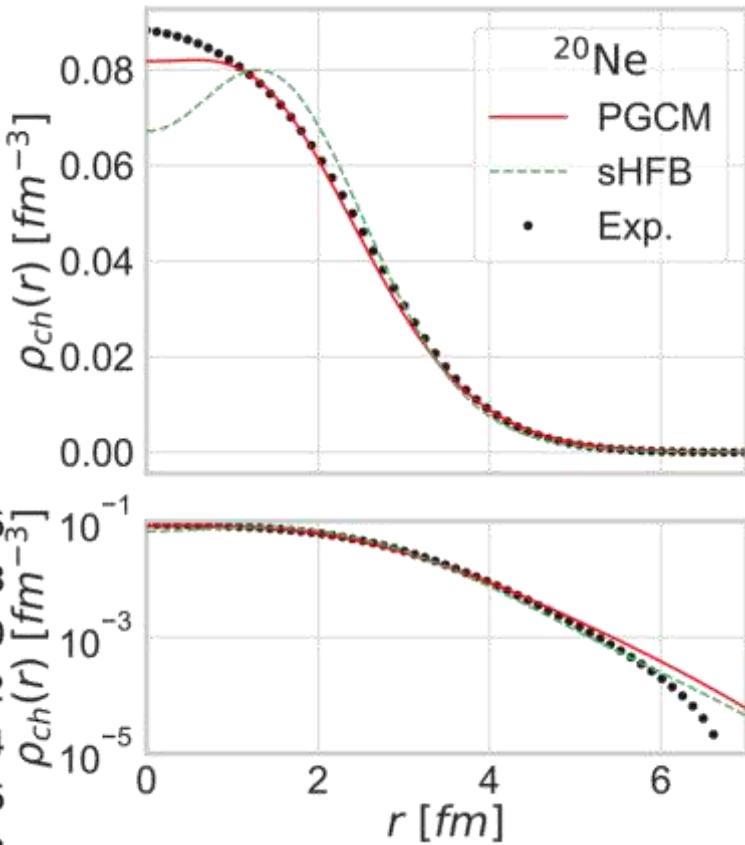
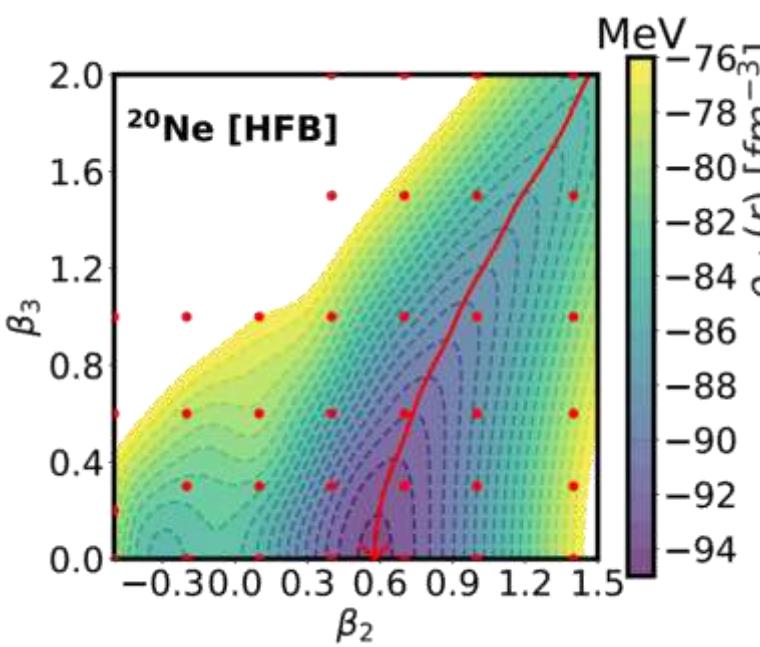
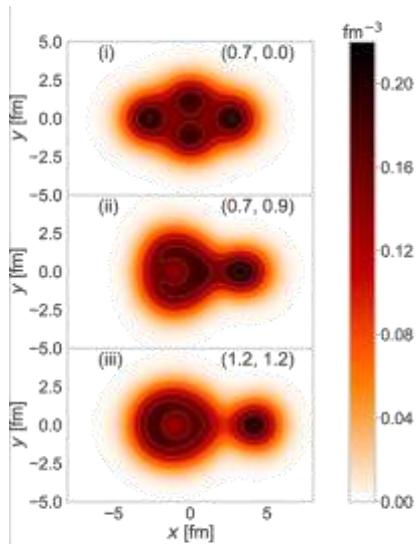
Frosini, Duguet, Ebran, Bally, Mongelli, Rodriguez, Roth, Somà, EPJA (2022)

Frosini, Duguet, Ebran, Bally, Hergert, Rodriguez, Roth, Yao, Somà, EPJA (2022)

● Correlated GS

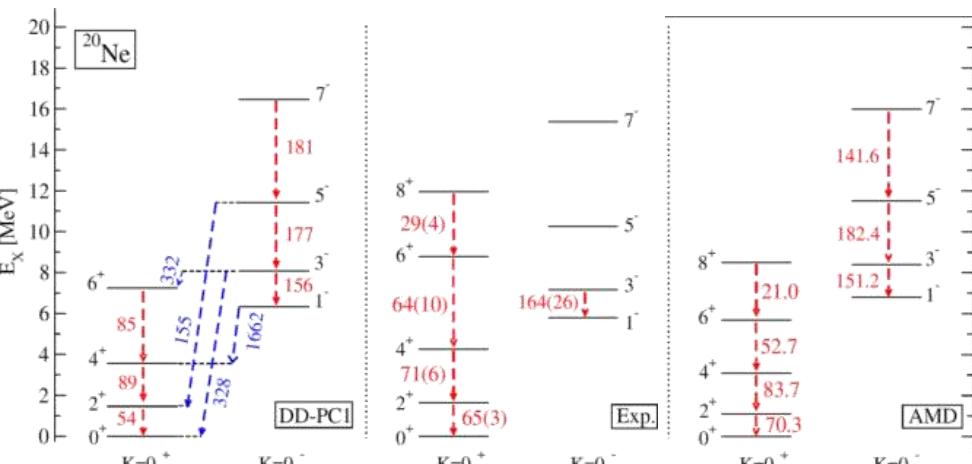


Ab initio

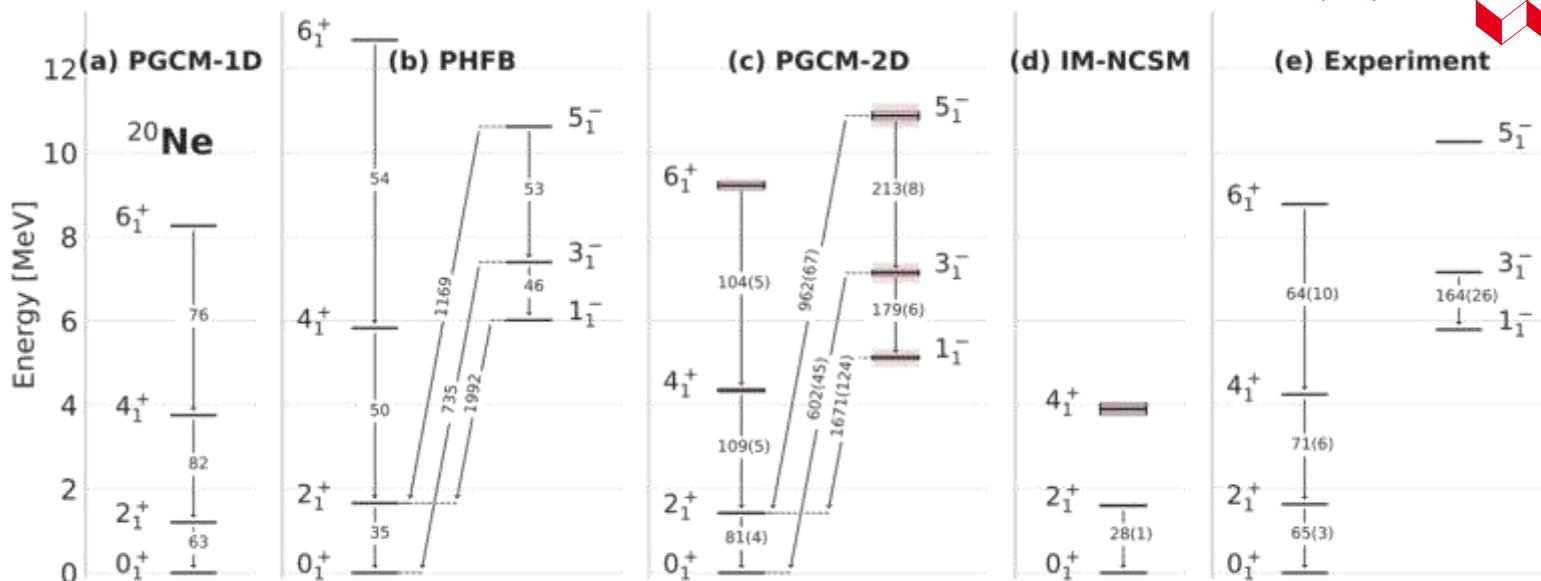


Frosini, Duguet, Ebran, Bally, Mongelli, Rodriguez, Roth, Somà, EPJA 2022

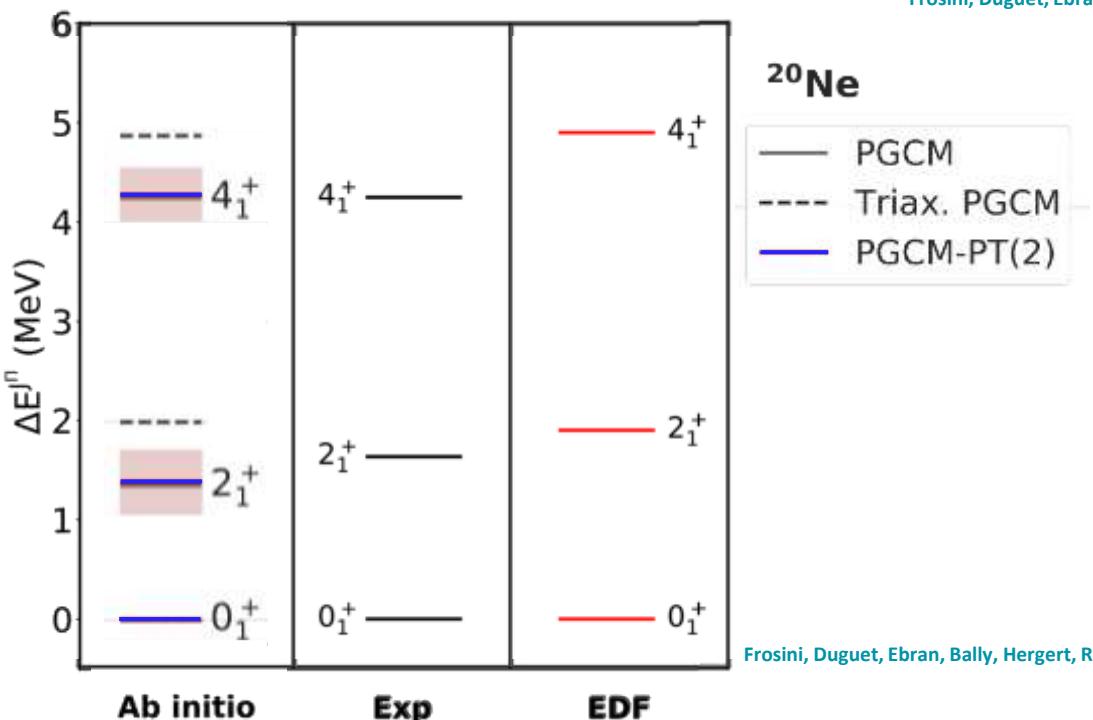
EDF



Marevic, Ebran, Khan, Niksic, Vretenar, 2018



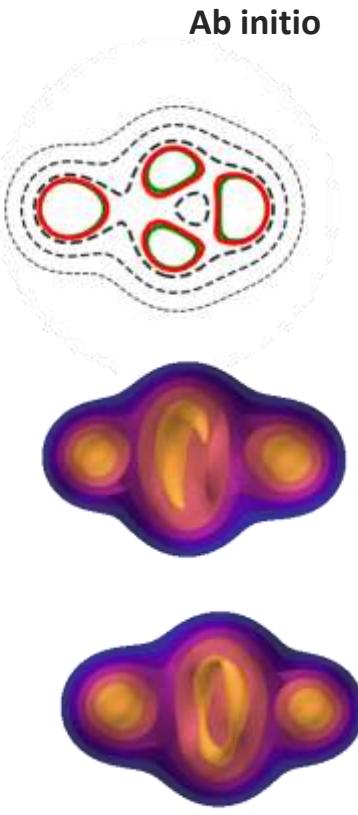
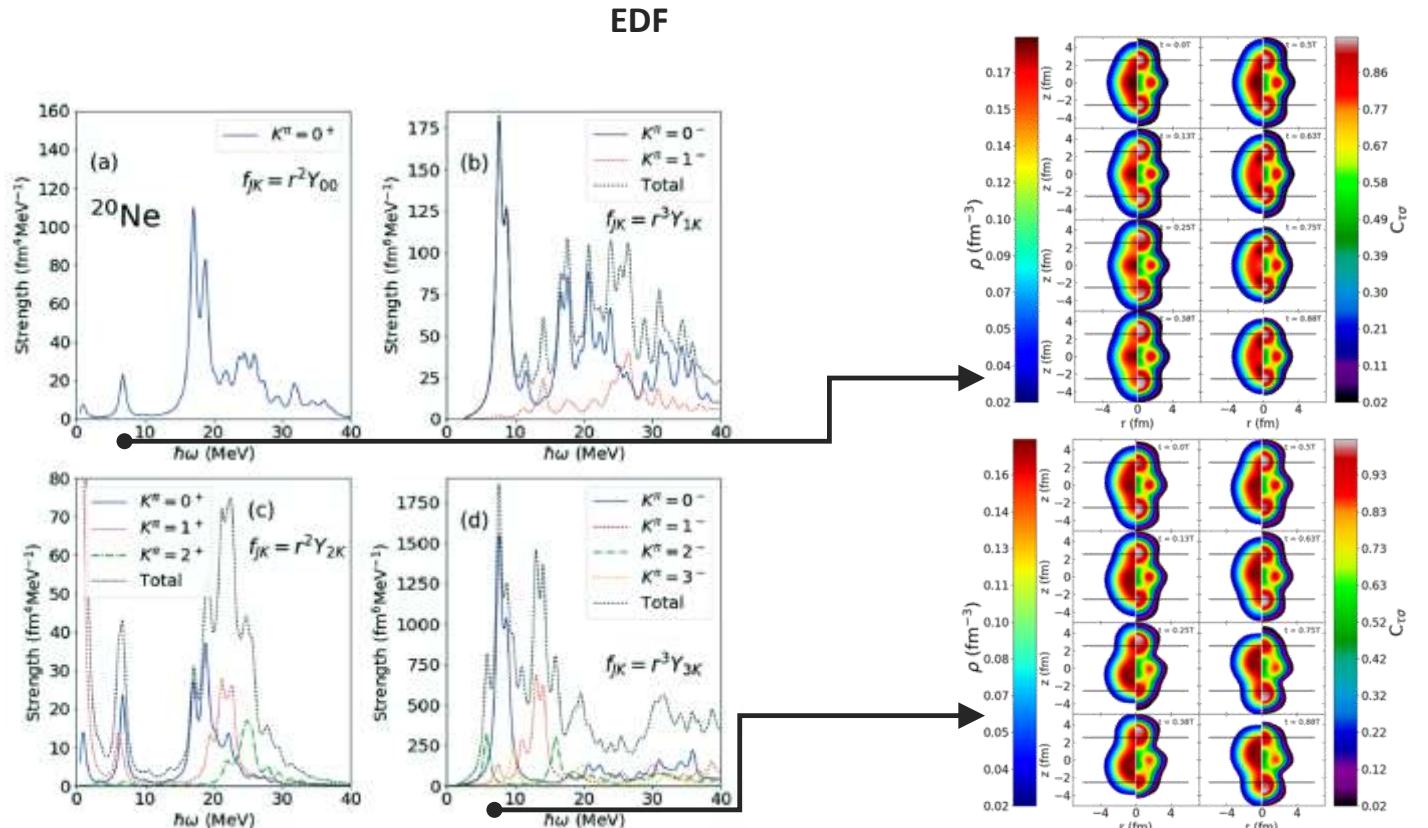
Frosini, Duguet, Ebran, Bally, Mongelli, Rodriguez, Roth, Somà, EPJA 2022



Frosini, Duguet, Ebran, Bally, Hergert, Rodriguez, Roth, Yao, Somà, EPJA 2022

QRPA (FAM)

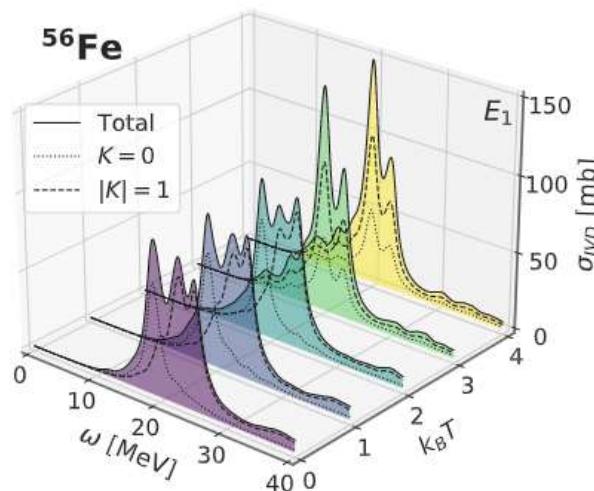
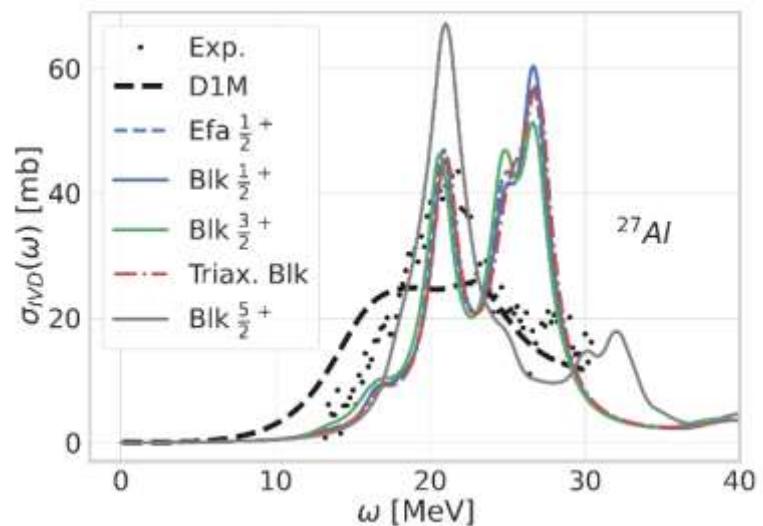
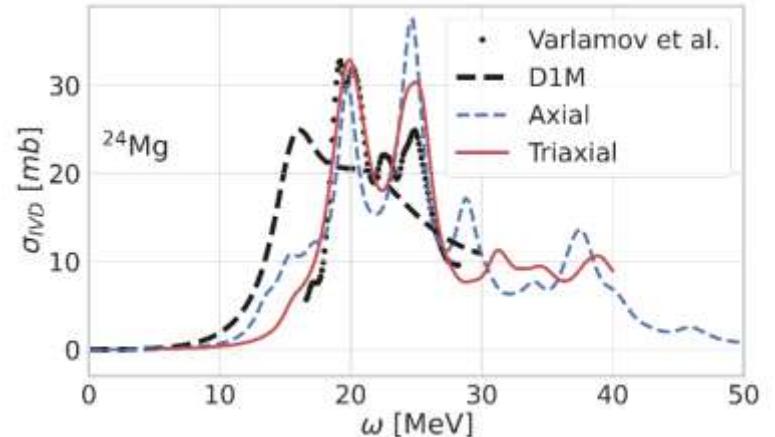
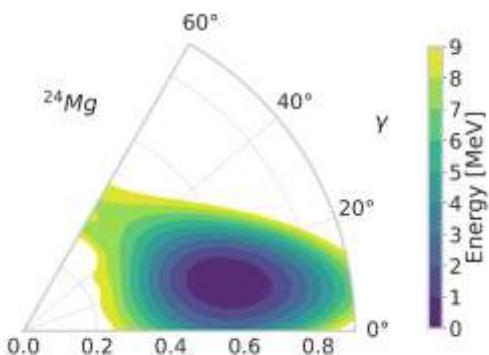
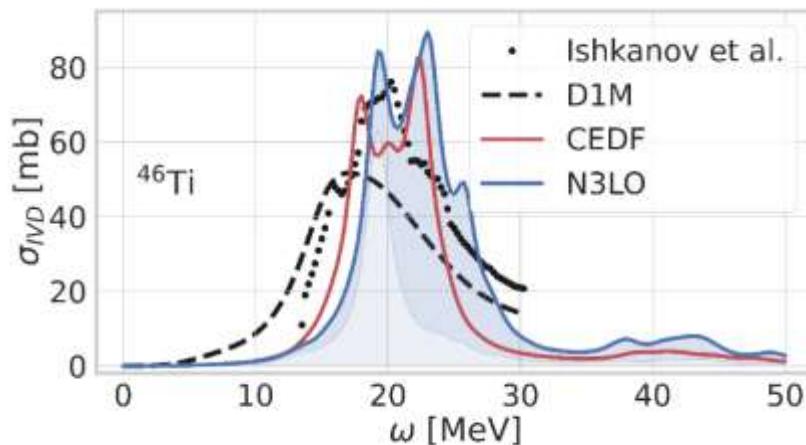
● Cluster vibration



Ab initio QFAM time-dependent intrinsic density
Frosini, Ebran, Duguet, Somà, unpublished

Mercier, Bjelčić, Nikšić, Ebran, Khan, Vretenar 2021
Mercier, Ebran, Khan 2022

QRPA (FAM)



Beaujeault-Taudière, Frosini, Ebran, Duguet, Roth, Somà PRC (2023)
Frosini, Duguet, Ebran, Somà, unpublished

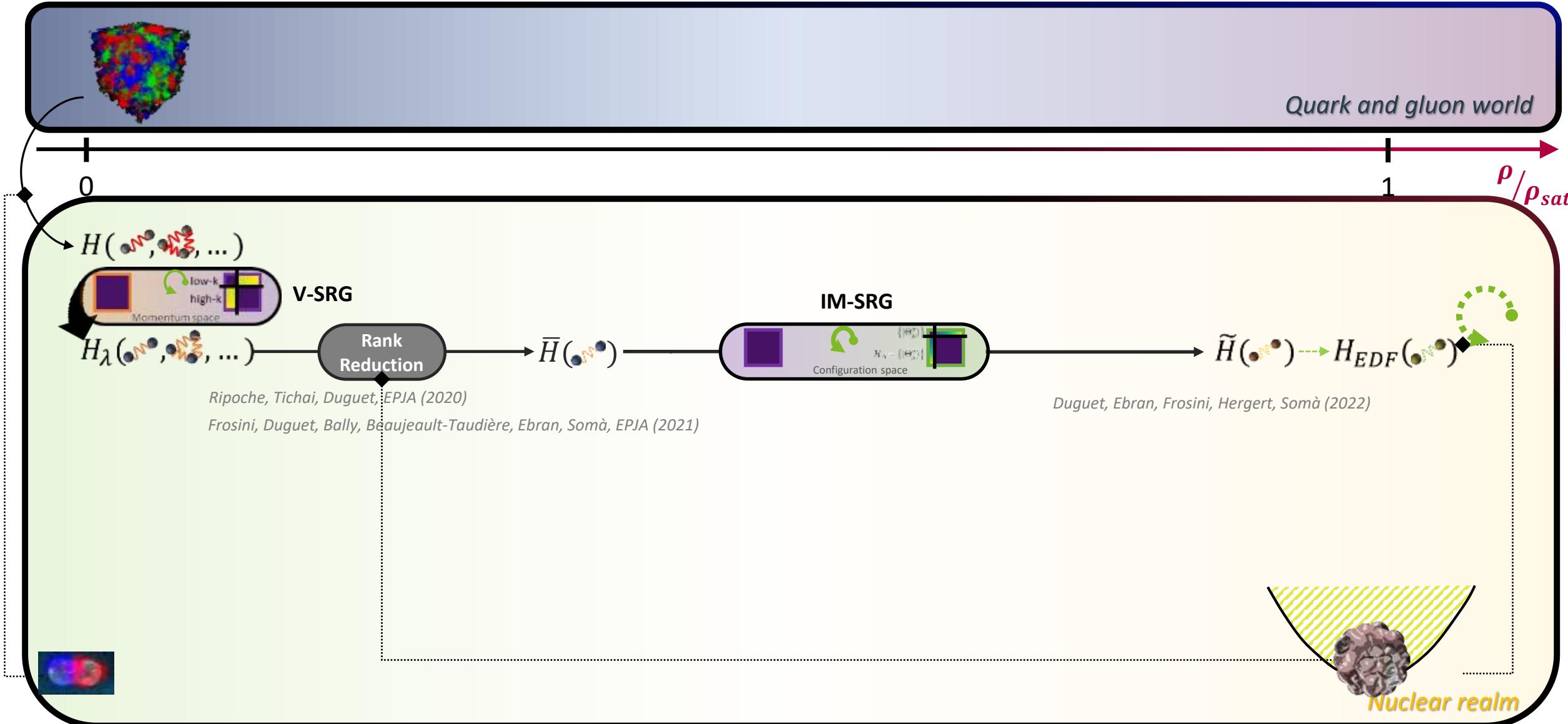
Outline

- 1. General context**
- 2. Recent work on empirical EDFs**
- 3. EDF-inspired ab initio methods**
- 4. Towards a first-principle formulation of EDFs**

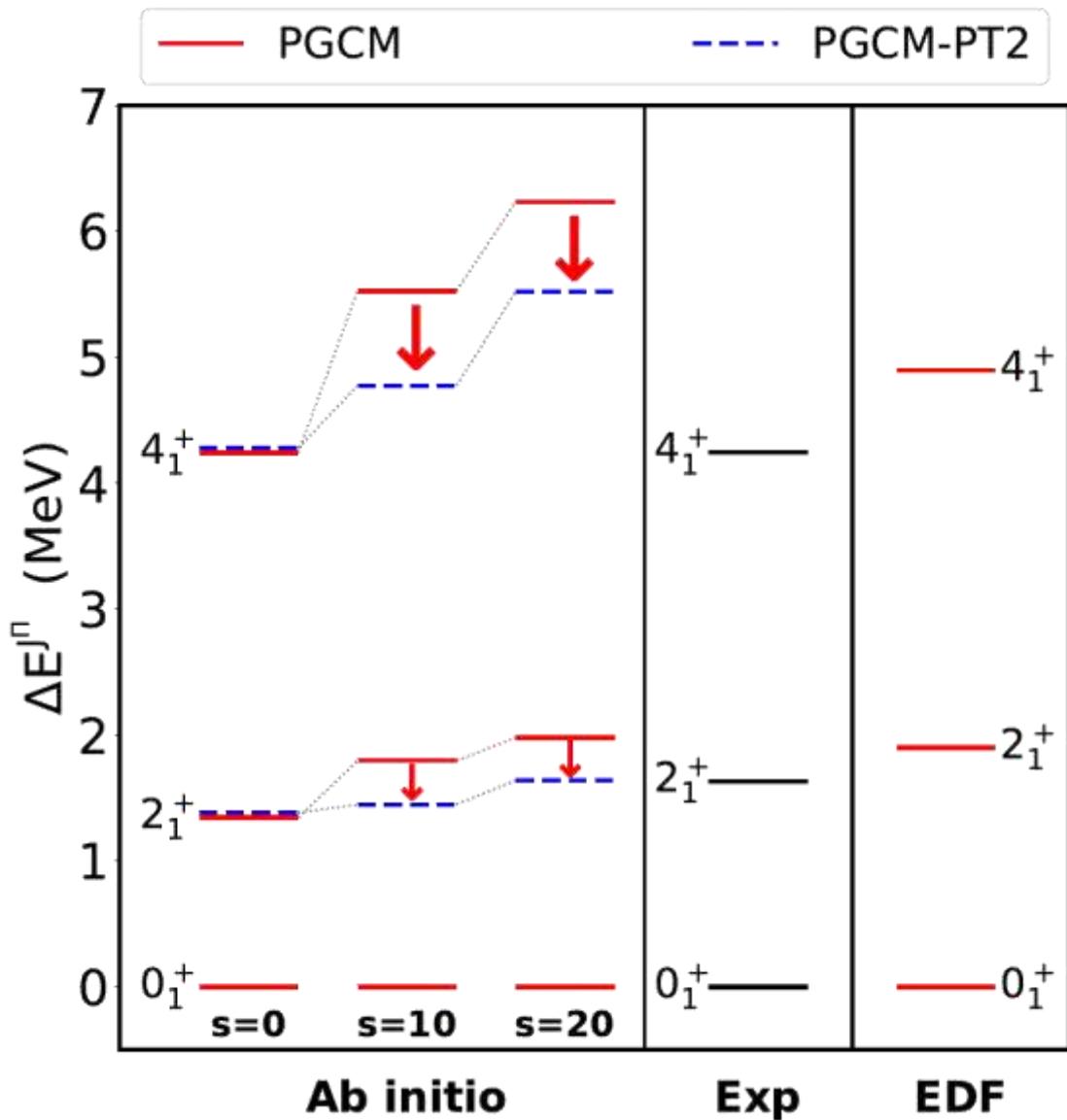
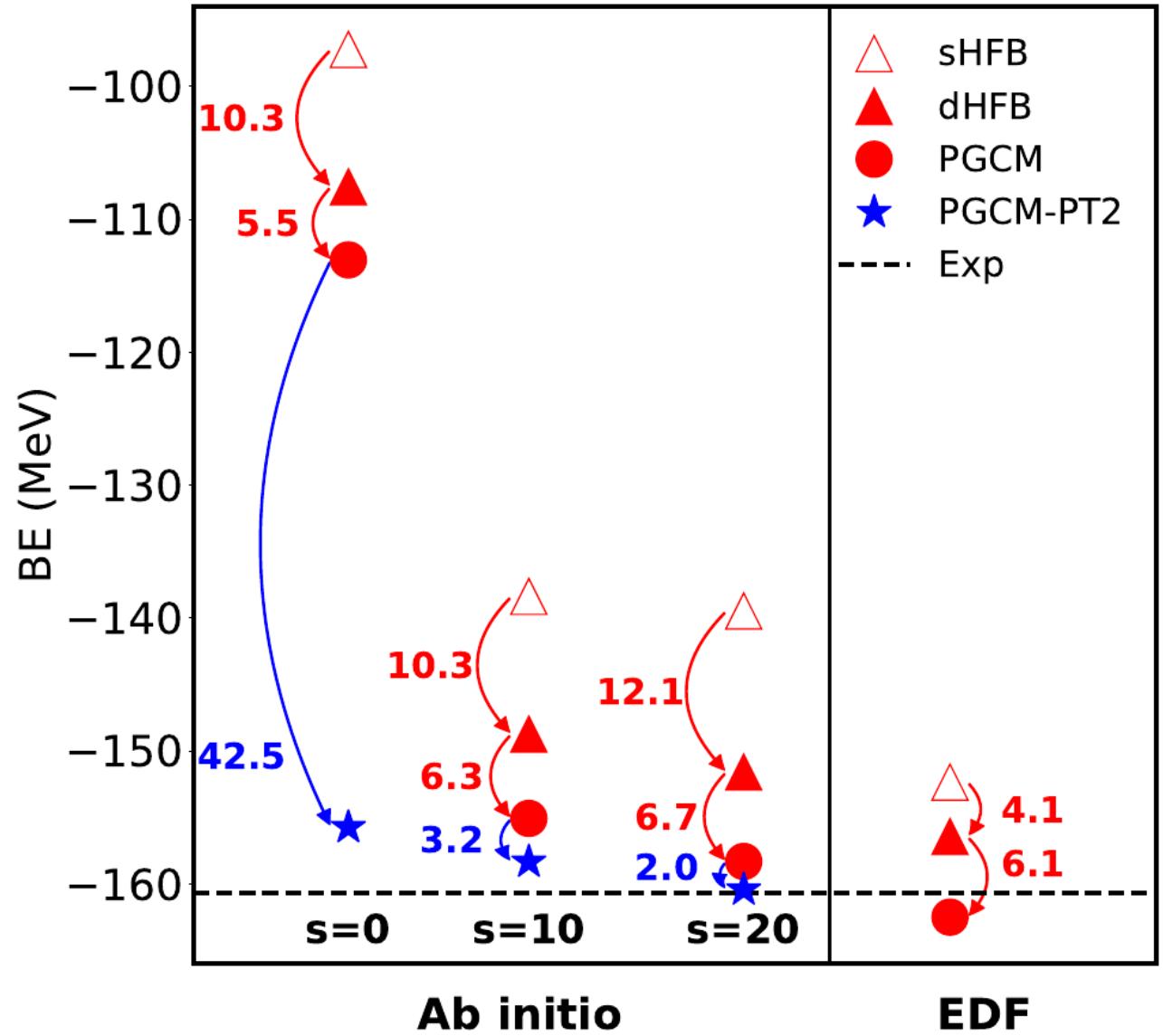


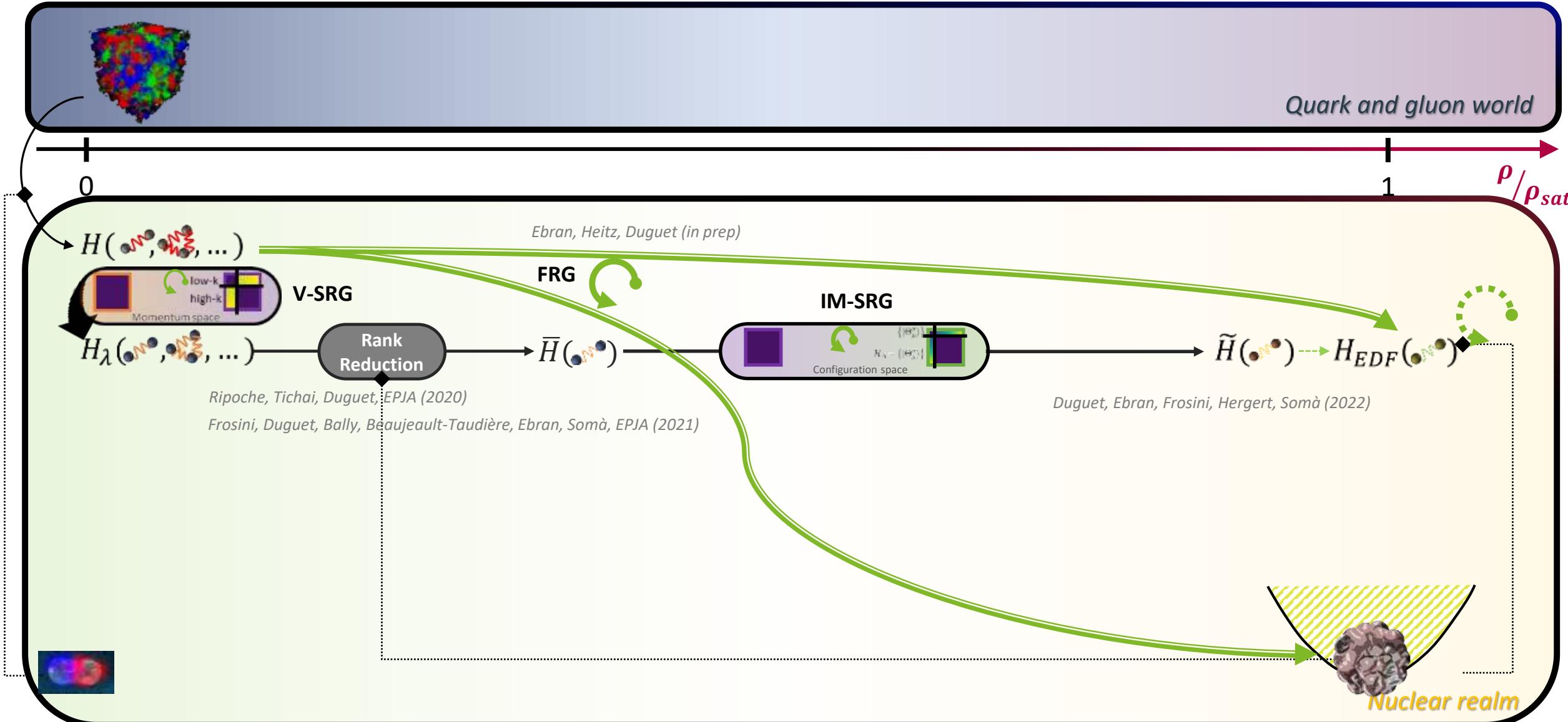
4. Towards a first-principle formulation of EDFs

(EDF-inspired ab initio method)-inspired EDFs



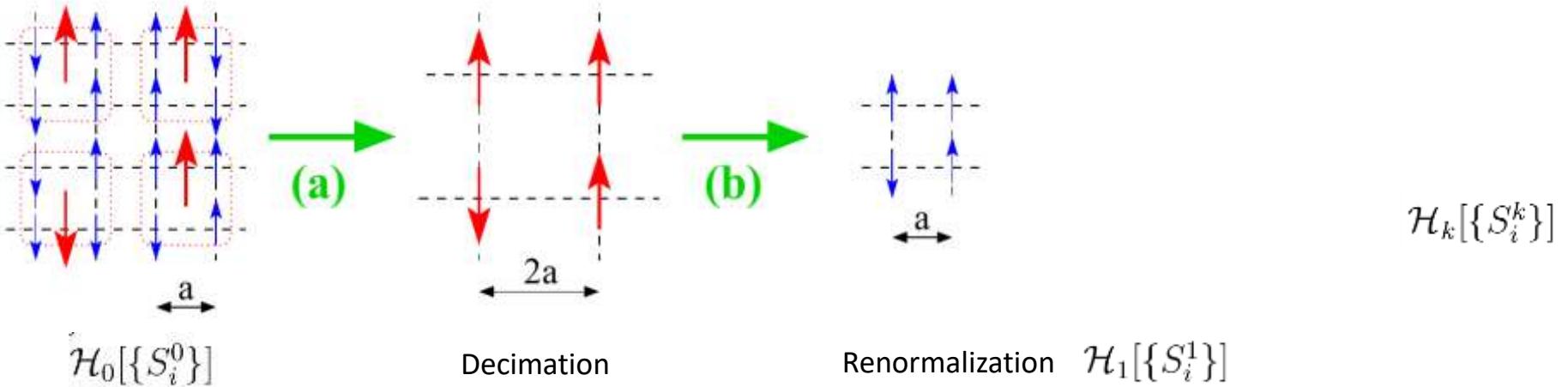
Non empirical EDFs via IM-SRG





Non empirical EDFs via FRG

● Renormalization group transformation : Wilson-Kadanoff procedure



● Renormalization group transformation : FRG

--> Central object of FRG : scale-dependent (or average) effective action Γ_k interpolating between the S and Γ

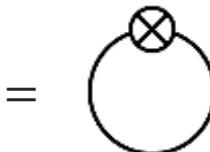
--> Mass term $S + \Delta S_k$

$$\Delta S_k = \frac{1}{2} \int \psi^\dagger(q) R_k(q) \psi(-q)$$

--> Exact RG (or Wetterich) equation

$$\partial_k \Gamma_k[\varphi] = \frac{1}{2} \text{tr} \left\{ \partial_k R_k \left(\Gamma_k^{(2)}[\varphi] + R_k \right)^{-1} \right\}$$

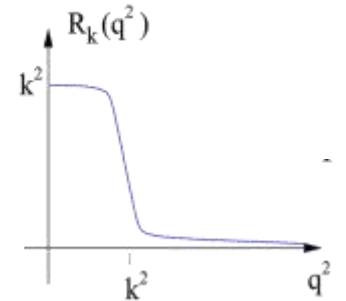
$$\partial_k \Gamma_k[\bar{\psi}, \psi] = - \text{tr} \left\{ \partial_k R_k \left(\Gamma_k^{(1,1)}[\bar{\psi}, \psi] + R_k \right)^{-1} \right\}$$



$$\begin{array}{c} \Gamma_{k=\Lambda}[\varphi_\Lambda] = S[\psi] \\ \downarrow \\ \Gamma_k[\varphi_k] \\ \downarrow \\ \Gamma_{k=0}[\varphi_0] = \Gamma[\varphi] \end{array}$$

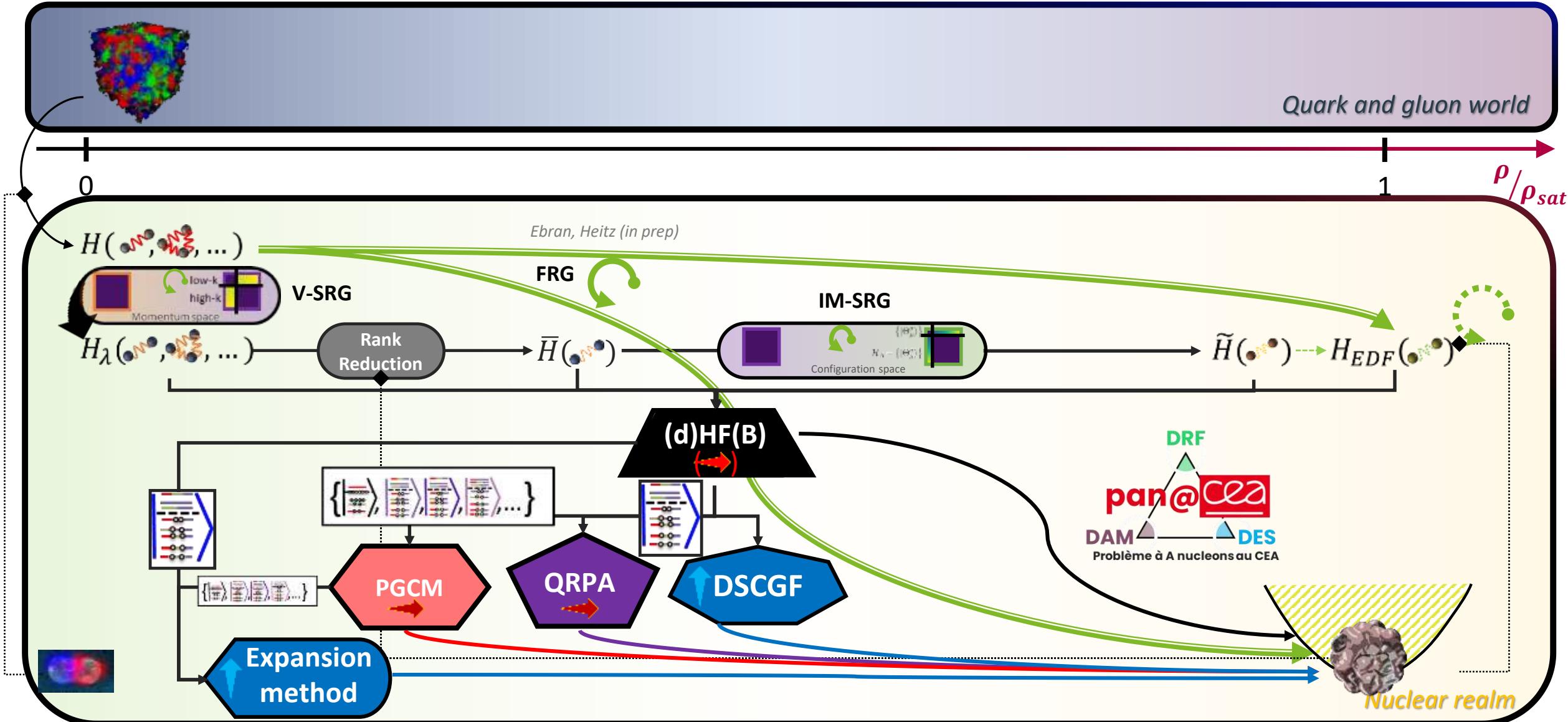
$k = \Lambda$
 k
 $k = 0$
 Incorporation
of fluctuations
of modes with
 $q > k$

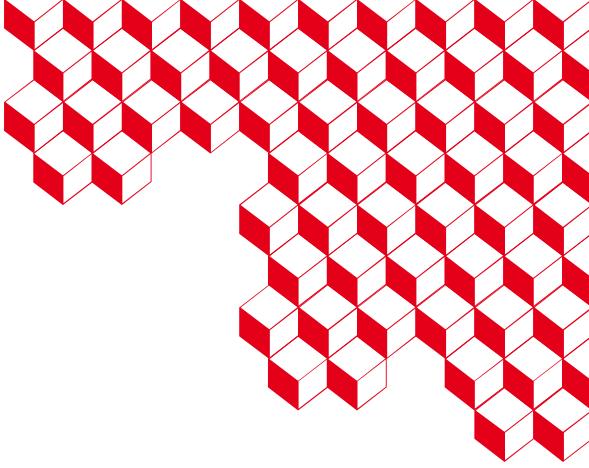
$$\begin{aligned} \Gamma_k^{(a)}[\varphi] &= \frac{\delta^a}{\delta \varphi^a} \Gamma_k[\varphi] \\ \Gamma_k^{(a,b)}[\psi, \bar{\psi}] &= \frac{\overrightarrow{\delta^a}}{\delta \bar{\psi}^a} \Gamma_k[\psi, \bar{\psi}] \frac{\overleftarrow{\delta^b}}{\delta \psi^b} \end{aligned}$$



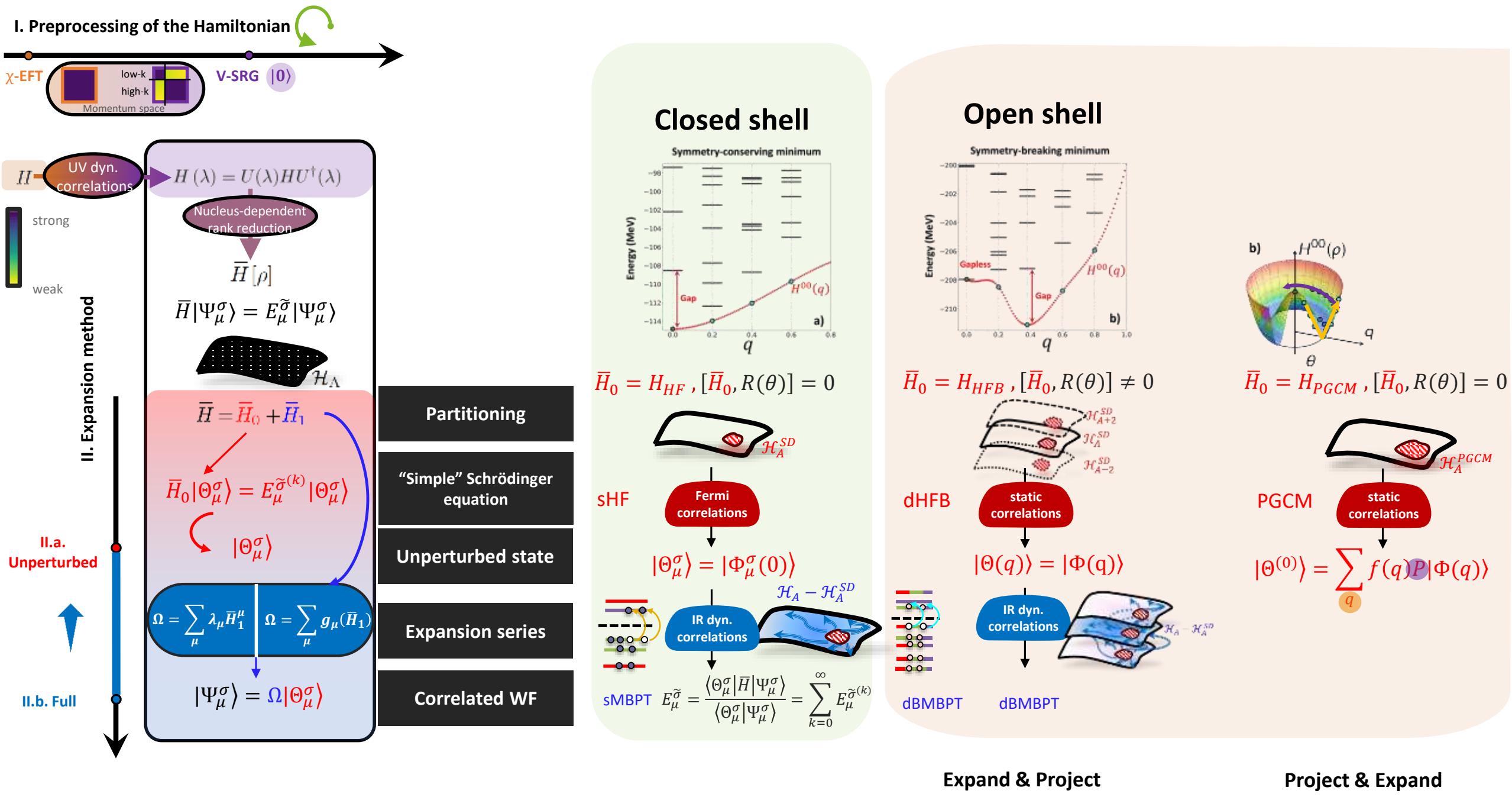


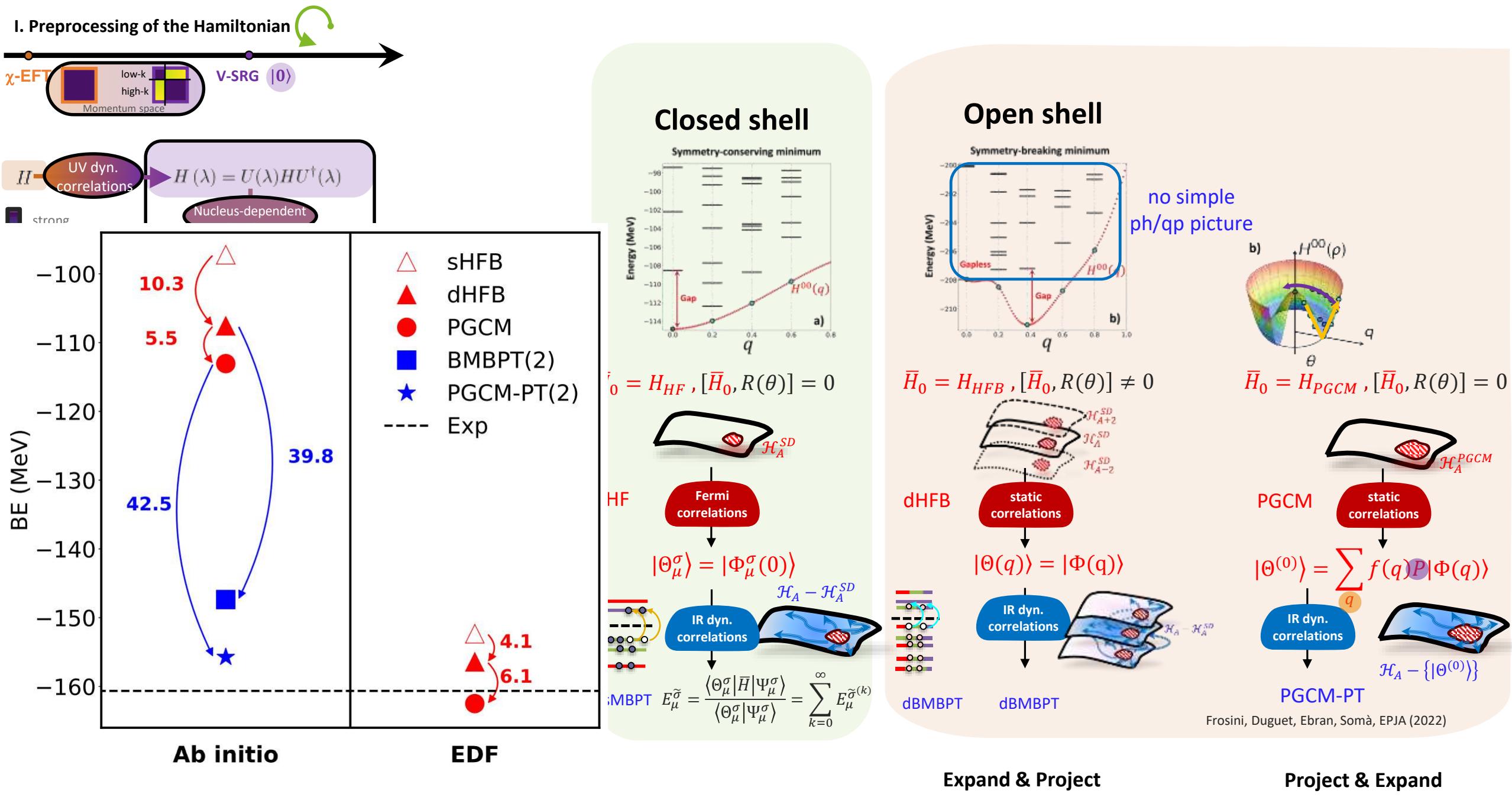
Conclusion





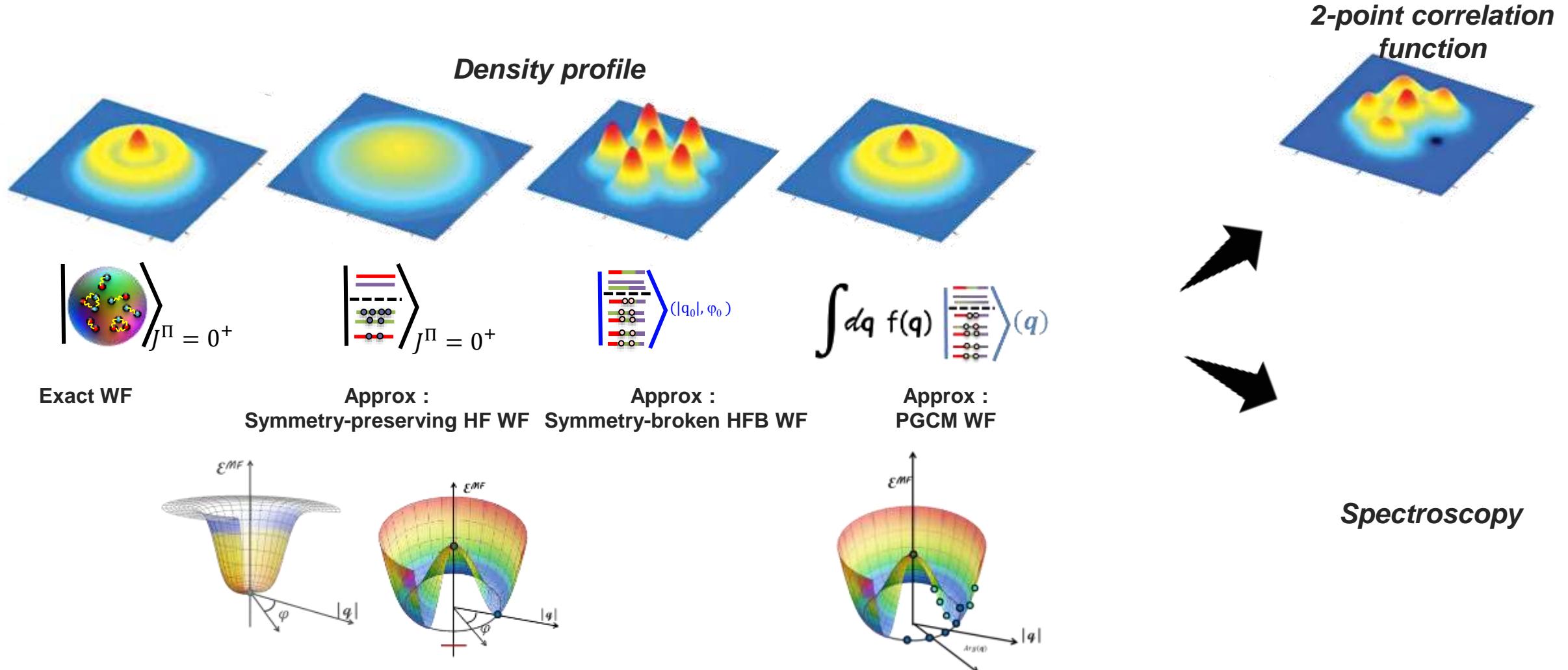
Thank you for your attention







Nuclear clustering & PGCM



Yannouleas & Landman, 2017

Nuclear clustering & PGCM

