

Adam Falkowski

LPT ORSAY

Supersymmetry

After the First LHC Results

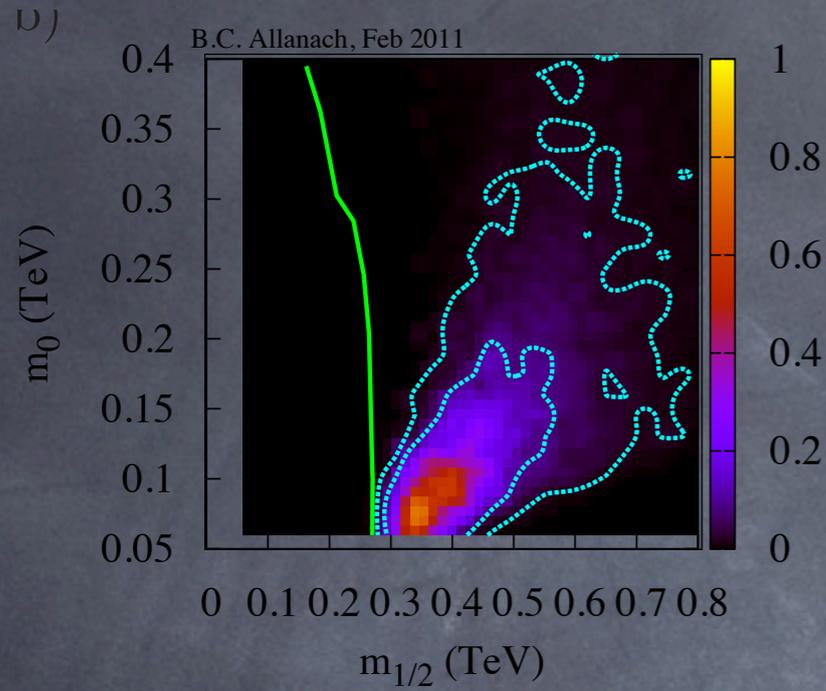


SUSY's portrait borrowed from C.Bernet's talk at Moriond'11

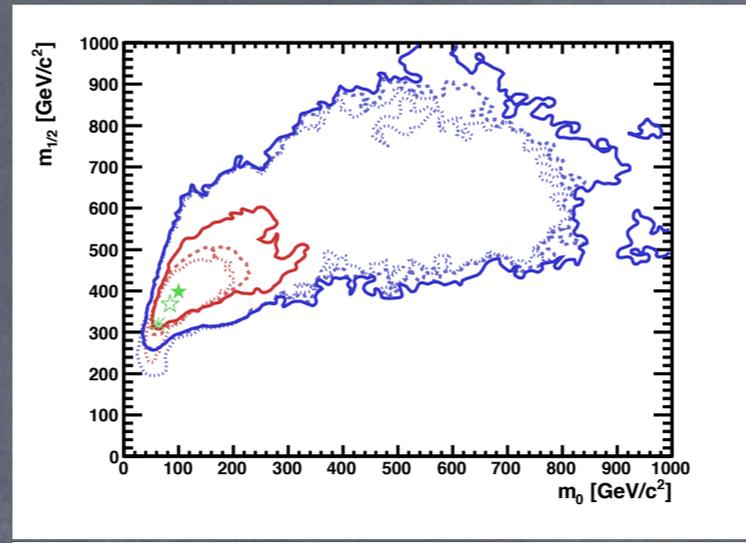
What is the impact of
the first LHC limits on
viability of
SUSY



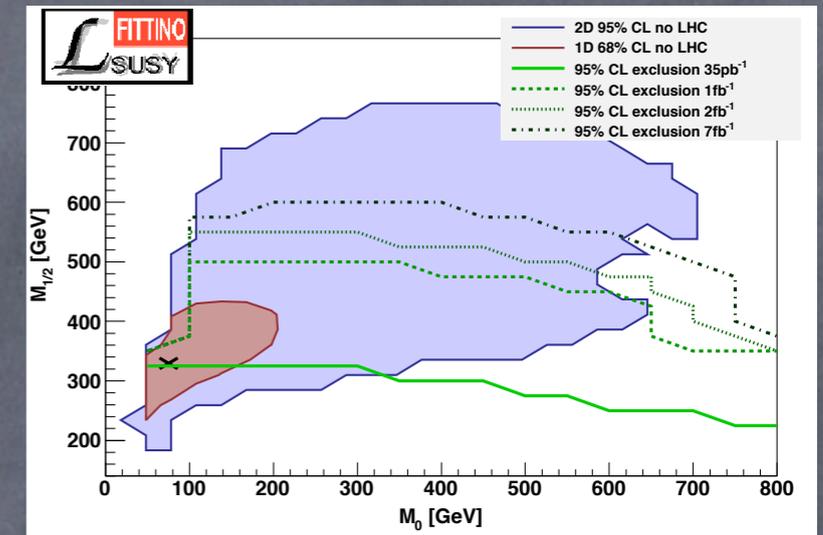
What Do Experts Say



Allanach, 1102.3149



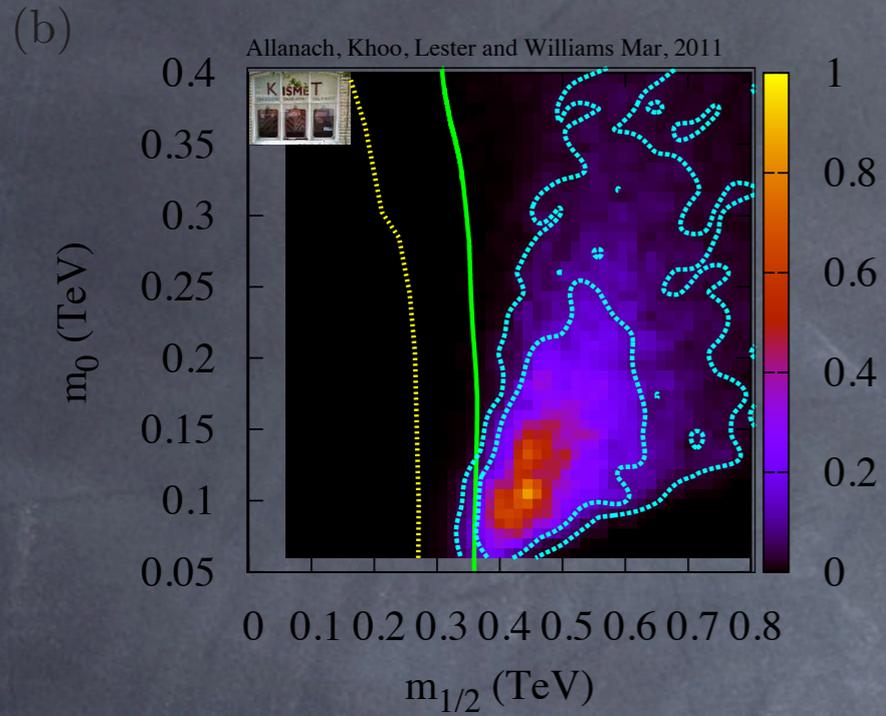
Buchmuller et al, 1102.4585



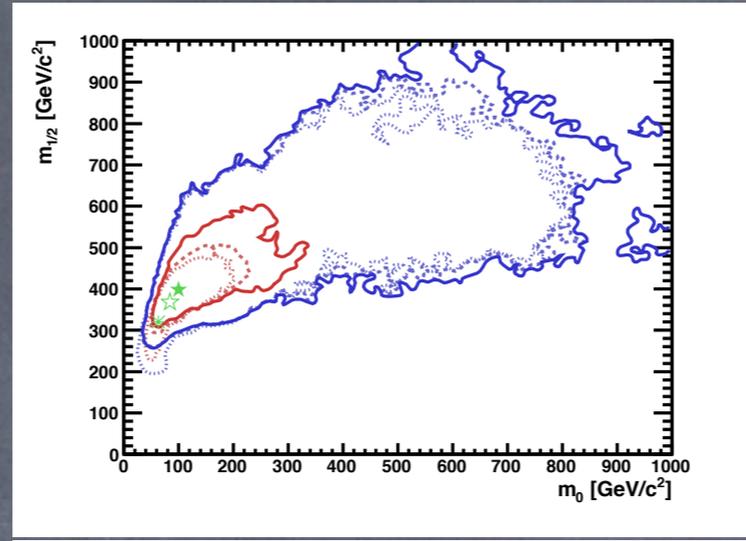
Bechtel et al, 1102.4693

Supersymmetry is just behind the corner!!!

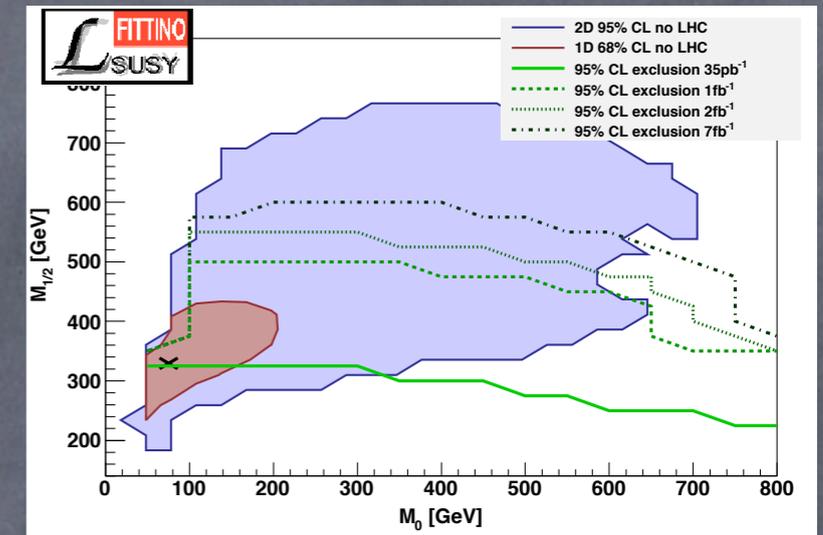
What Do Experts Say



Allanach et al, 1103.0969



Buchmuller et al, 1102.4585

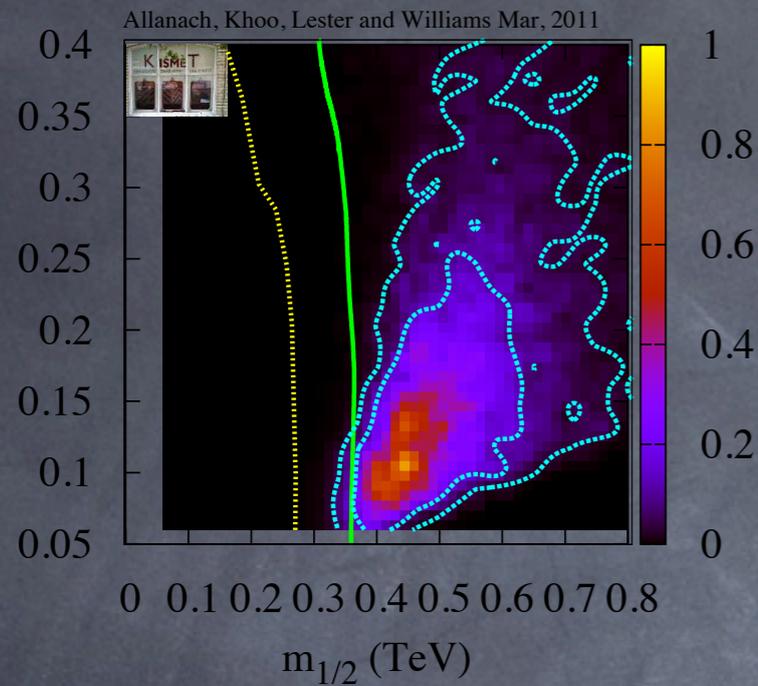


Bechtel et al, 1102.4693

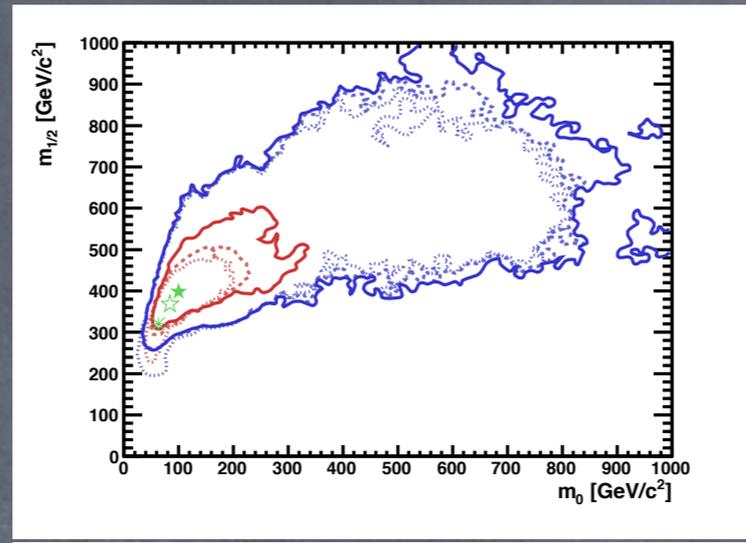
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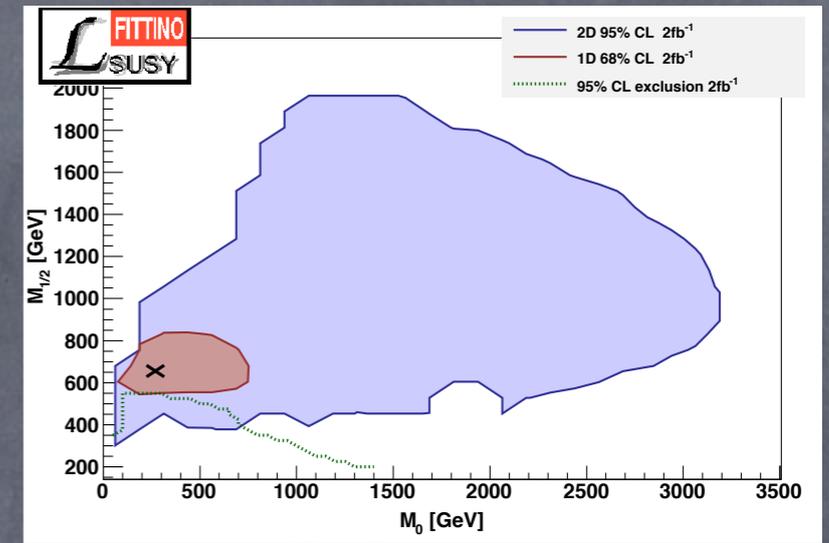
(b)



Allanach et al, 1103.0969



Buchmuller et al, 1102.4585



Bechtle et al, 1102.4693

Supersymmetry is just behind the corner!!!

(although the corner may move around)

"...the more he looked
inside the more Piglet
wasn't there..."



Plan

- Who ordered SUSY?
- How do we search for SUSY?
- What have we learnt?

EW symmetry breaking chart

Strongly Coupled

Weakly Coupled

Standard Model

5D Higgsless

Technicolor

Composite Higgs

Little Higgs

Supersymmetry

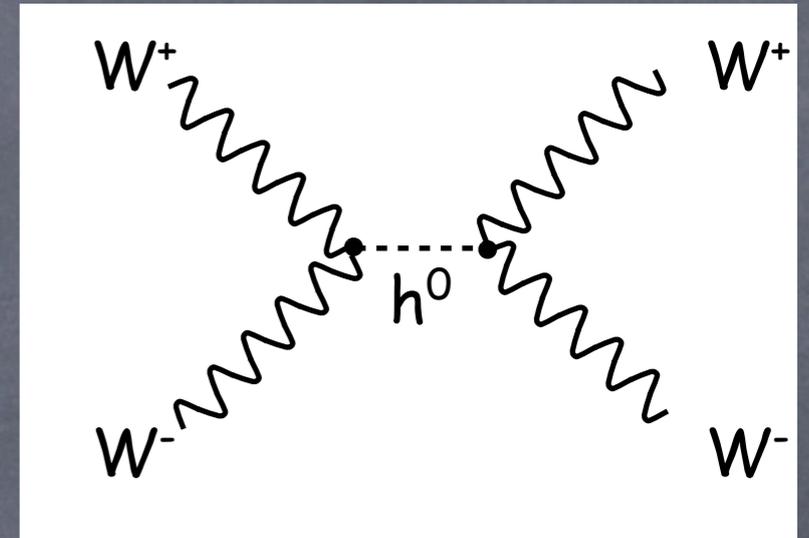
No Higgs

Higgs

In the following, weakly coupled only...

Weakly Coupled Models

- Contain a narrow scalar particle(s) coupled to W and Z bosons whose contribution unitarizes VV scattering
- Simplest example is the Standard Model

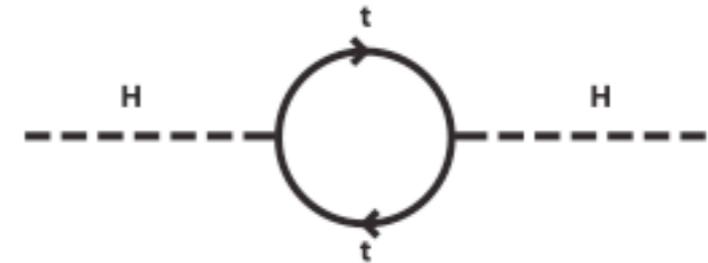


Standard Model

One Higgs field that acquires vacuum expectation value

$$V(H) = m_0^2 H^\dagger H + \lambda (H^\dagger H)^2$$
$$\langle H \rangle = \begin{pmatrix} 0 \\ v/\sqrt{2} \end{pmatrix} \quad v = \sqrt{\frac{-m_0^2}{\lambda}}$$

Problem: the mass term, and therefore the vev, receive large quantum corrections

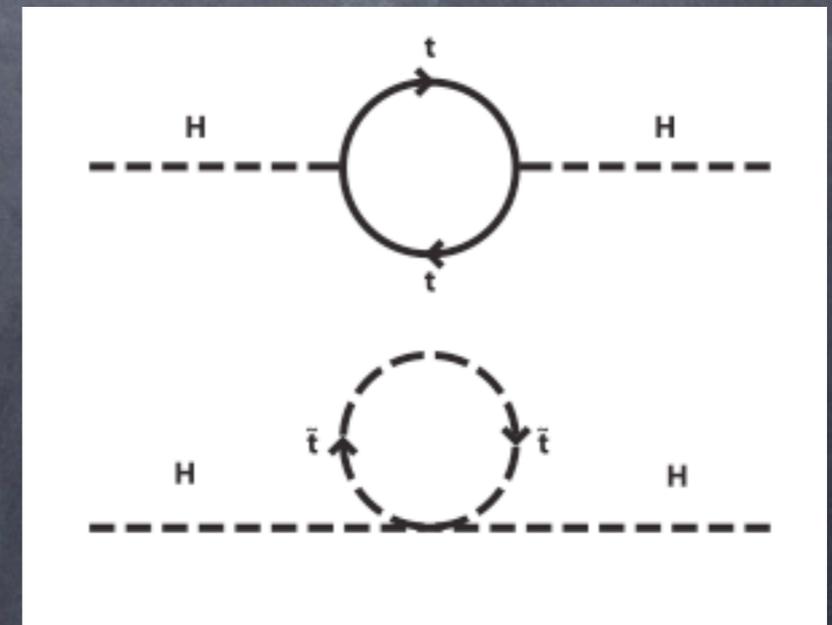
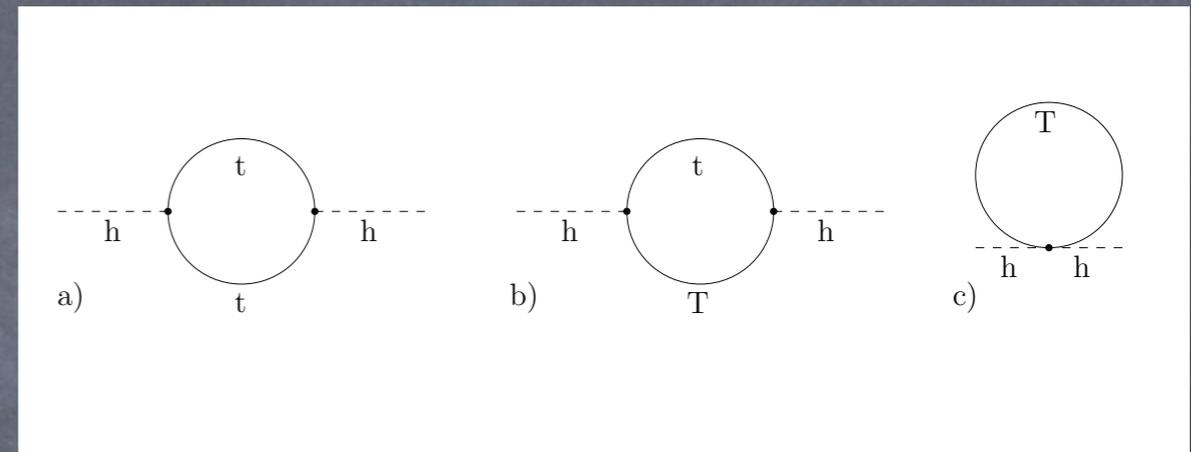


$$\Delta m_0^2 \approx \frac{m_{top}^2}{4\pi^2 v^2} \Lambda_{UV}^2 + \dots$$

If the cutoff $\gg 1$ TeV, we either need find tuning or new structure/particles that soften the quantum corrections

Who can stabilize Higgs potential ?

- Global symmetries:
quantum corrections
cancel between particles
of the same statistics
- Supersymmetry:
quantum corrections
cancel between particles
of opposite statistics



Supersymmetry: the good and the bad and the ugly

- Stabilizes Higgs
- (together with R-parity) Provides dark matter candidates
- Allows for better gauge coupling unification
- Flavor problem
- CP problem
- μ problem
- Doublet-triplet splitting problem

Problems in the Nutshell

- SUSY says: there is ~ 100 new degrees of freedom at the weak scale
- Experiment says: no new degrees of freedom show up in kaon mixing, D-mixing, B-mixing, $B_s \rightarrow \mu\mu$, neutron EDMs, $\mu \rightarrow e\gamma$, electroweak precision tests, proton decay, etc
- Generically they could be indirectly seen in numerous low energy experiments
- Approximate symmetries of the SM unexpectedly well respected by new physics

But at least SUSY solves the hierarchy problem... does she?

Higgs potential in the MSSM depends on Higgs soft masses and μ -term:

$$V = (|\mu|^2 + m_{H_u}^2)|H_u^0|^2 + (|\mu|^2 + m_{H_d}^2)|H_d^0|^2 - (b H_u^0 H_d^0 + \text{c.c.}) + \frac{1}{8}(g^2 + g'^2)(|H_u^0|^2 - |H_d^0|^2)^2.$$

The Higgs vev (equivalently Z-boson mass) can be expressed by these parameters:

$$m_Z^2 = \frac{|m_{H_d}^2 - m_{H_u}^2|}{\sqrt{1 - \sin^2(2\beta)}} - m_{H_u}^2 - m_{H_d}^2 - 2|\mu|^2.$$

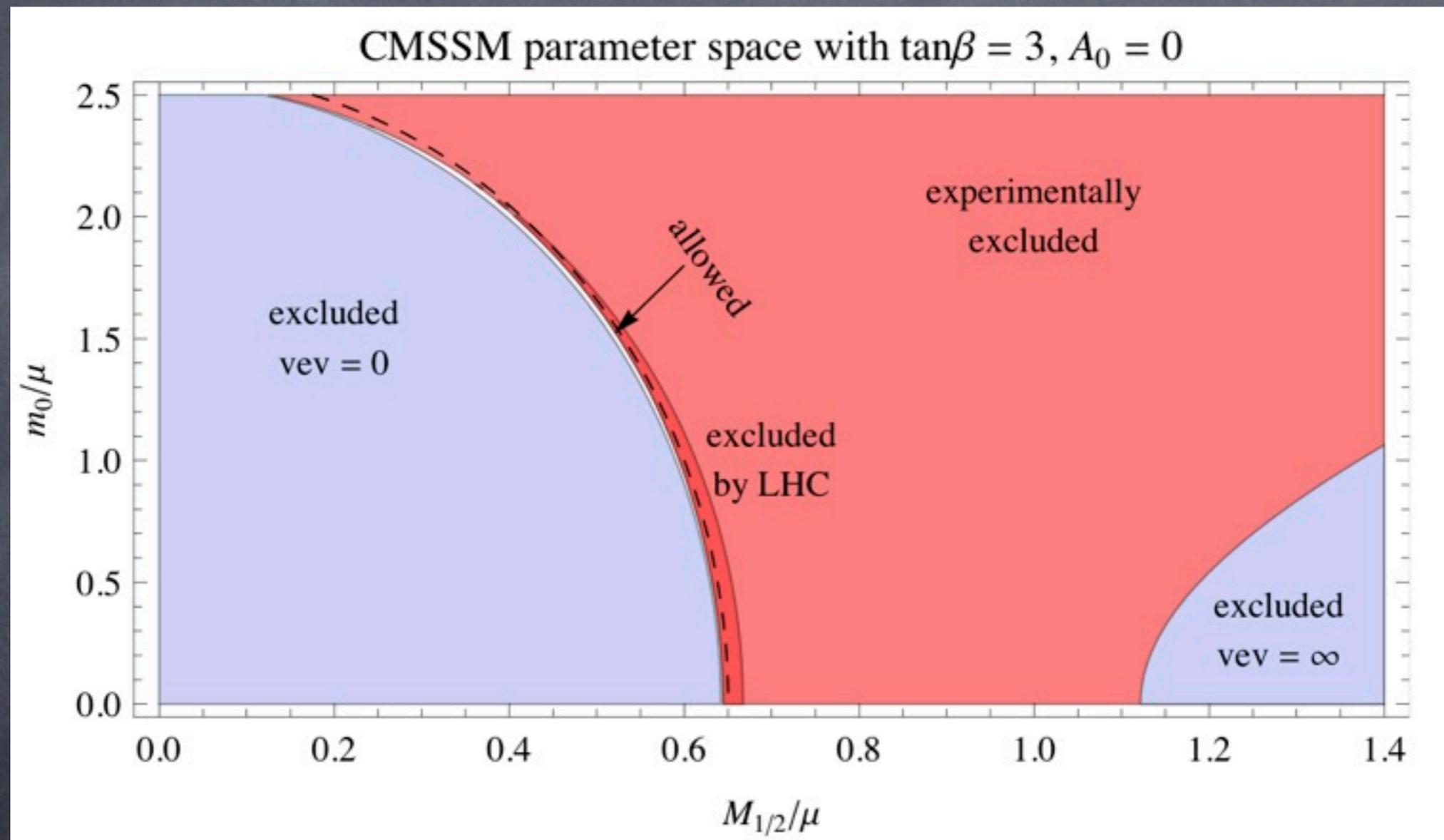
But due to loop corrections:

$$m_{H_u}^2 \sim \frac{\log \Lambda}{4\pi^2} M_{Susy}^2 \sim M_{Susy}^2 \Rightarrow M_{Susy} \sim m_Z$$

Susy particles should be at the Z-boson mass, otherwise fine-tuning!

Fine-tuning in CMSSM, Strumia 1101.2195

$$M_Z^2 \approx 0.2m_0^2 + 0.7M_3^2 - 2\mu^2 = (91 \text{ GeV})^2 \times 50 \left(\frac{M_3}{780 \text{ GeV}} \right)^2 + \dots$$



Already LEP constraints on SUSY and Higgs seriously constrain the parameter space. Only a small strip on the boundary between EW breaking and no EW breaking remains viable

Frog Metaphor

If a frog is placed in boiling water, it will jump out, but if it is placed in cold water that is slowly heated, it will never jump out.



Measuring S and T parameters at LEP1 was like hot water for Technicolor Frogs, and like cold water for SUSY Frogs. SUSY and Higgs searches at LEP2 and Tevatron are like heating the water.

Enters LHC

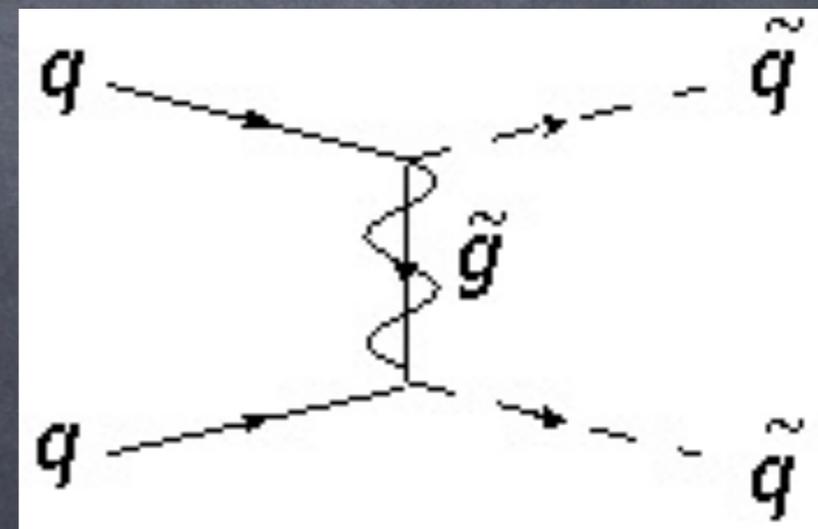
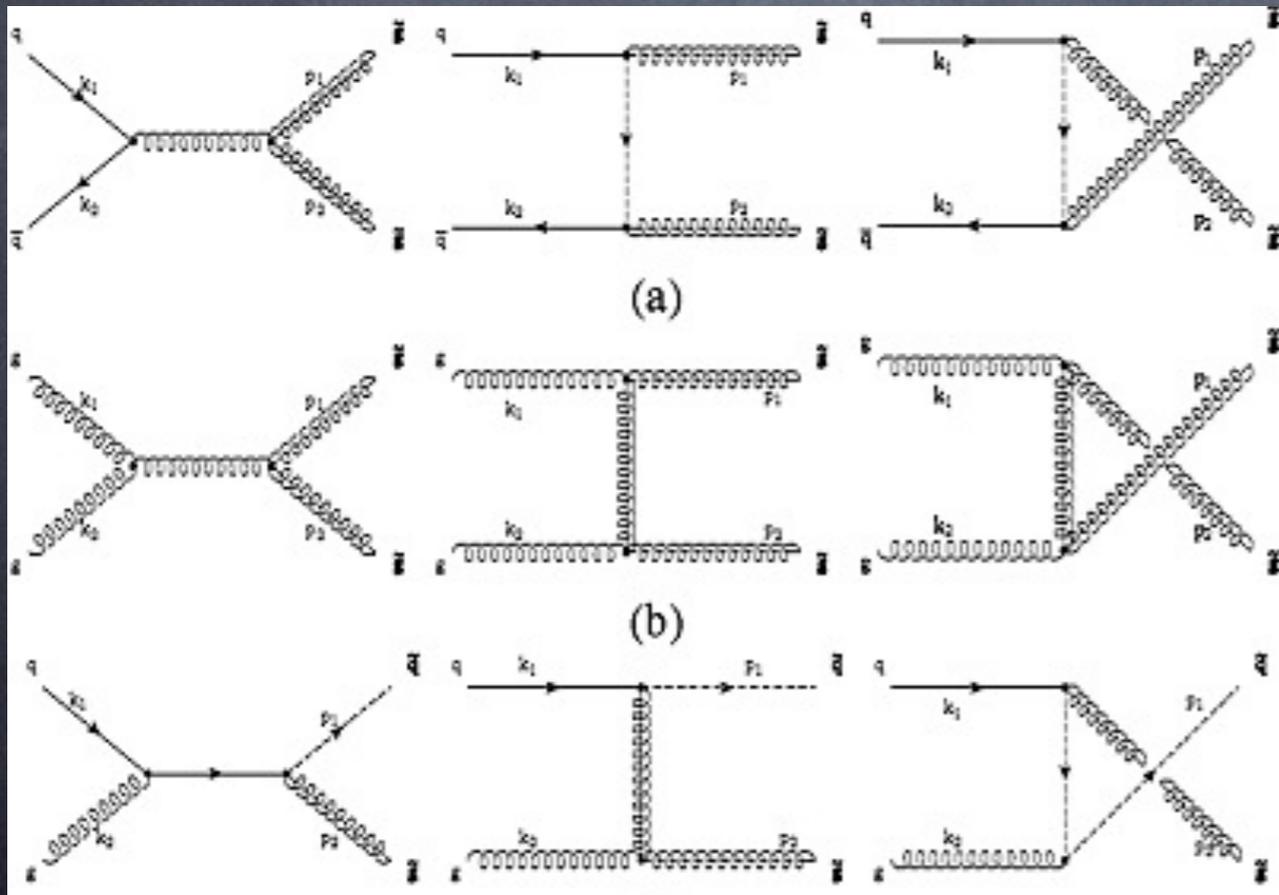
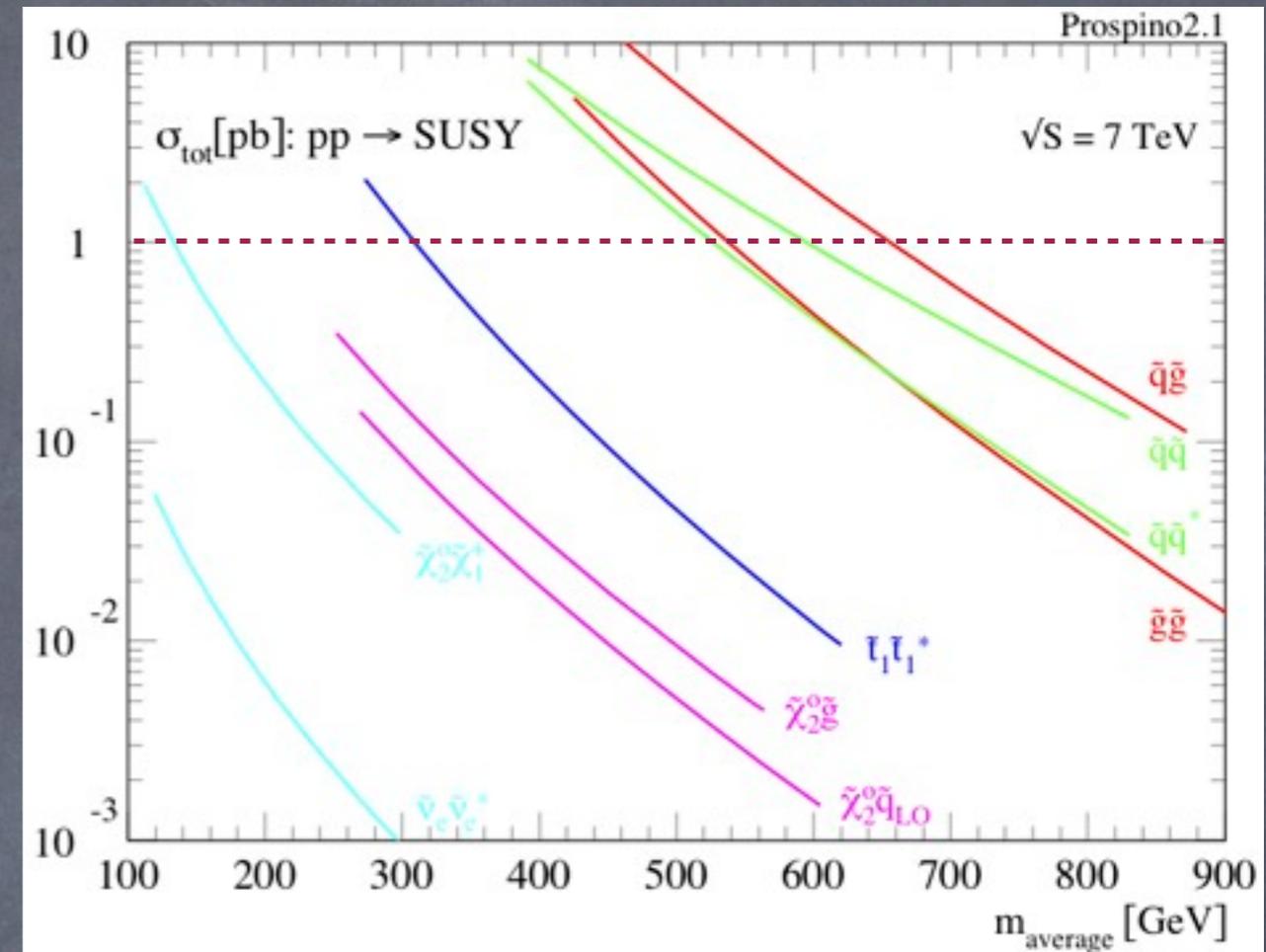
In the rest of this talk:

- How LHC searches for SUSY
- What have we learnt so far

I'm borrowing heavily from results and plots presented at Moriond by CMS (M.Chiorboli, C. Bernet) and ATLAS (N. Barlow, S. Caron)

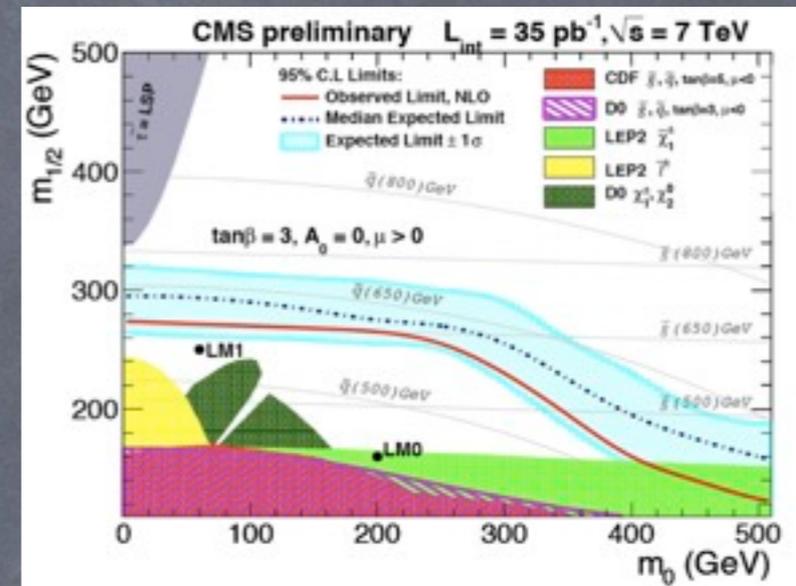
SUSY @ LHC

- In the early phase LHC can only produce colored superpartners
- In 2010, sensitivity to ~ 500 GeV 1st generation squark and gluinos (~ 700 GeV if both are present), and to ~ 300 GeV stops.



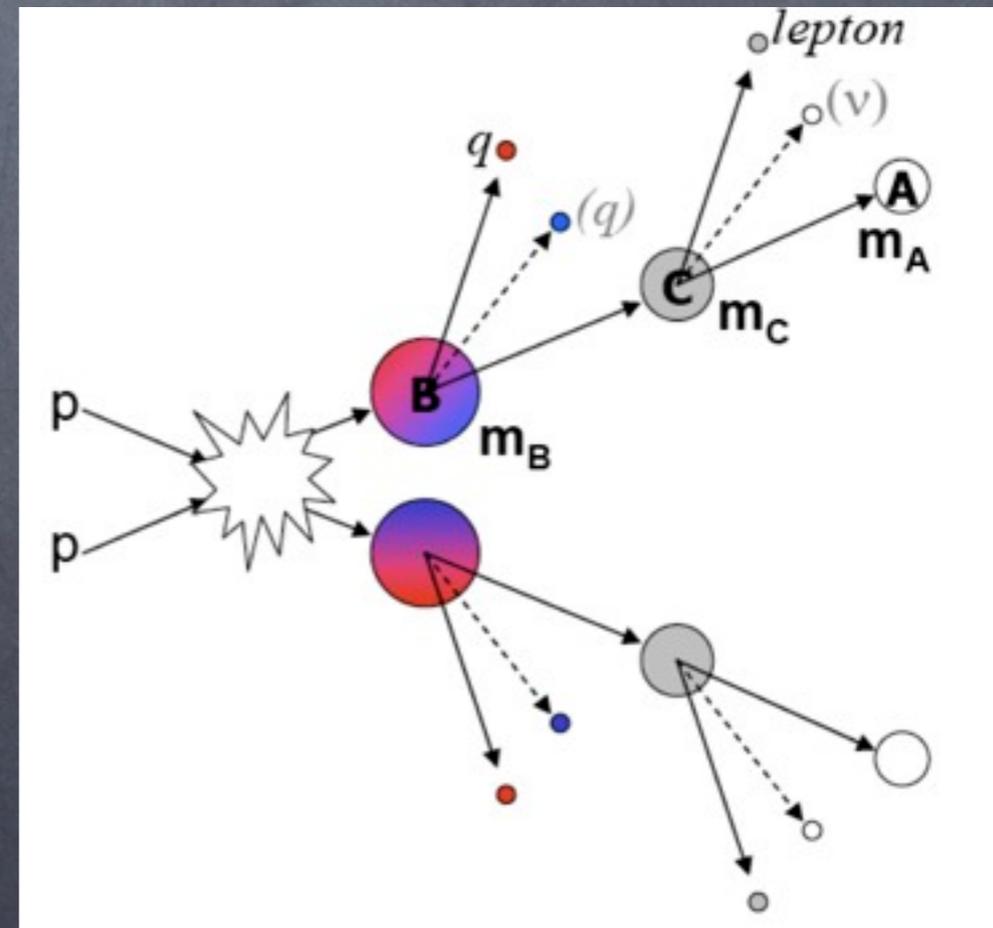
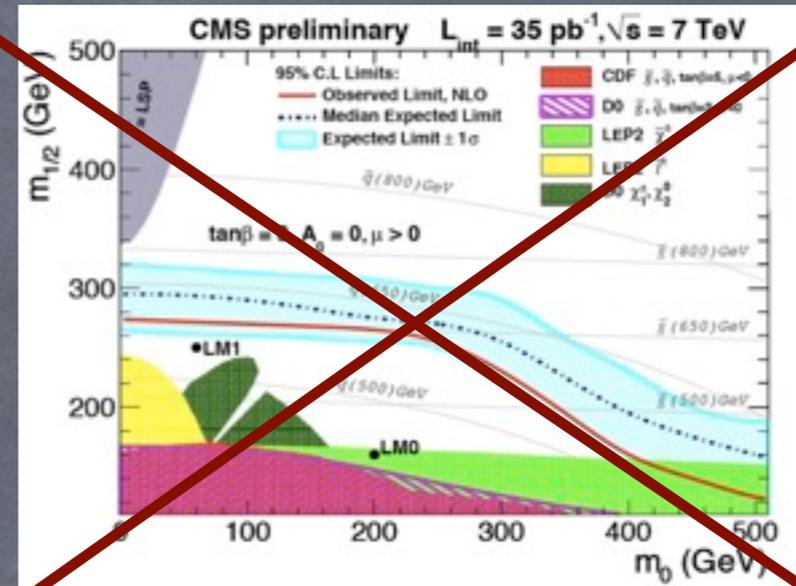
SUSY Limits: Topological Approach

- Results of SUSY searches usually presented as limits on mSUGRA parameters



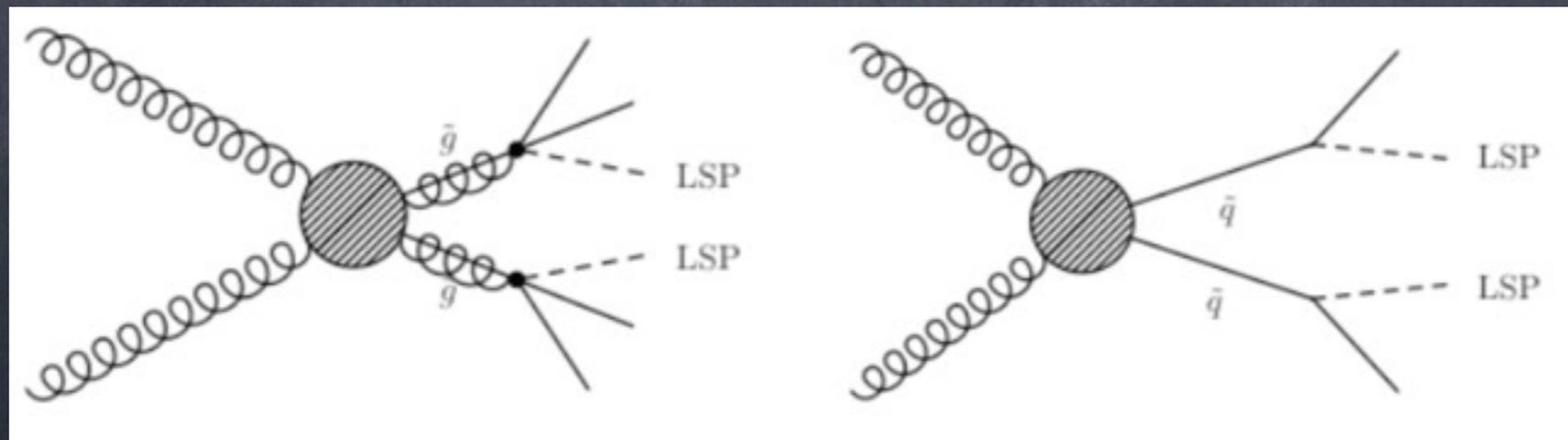
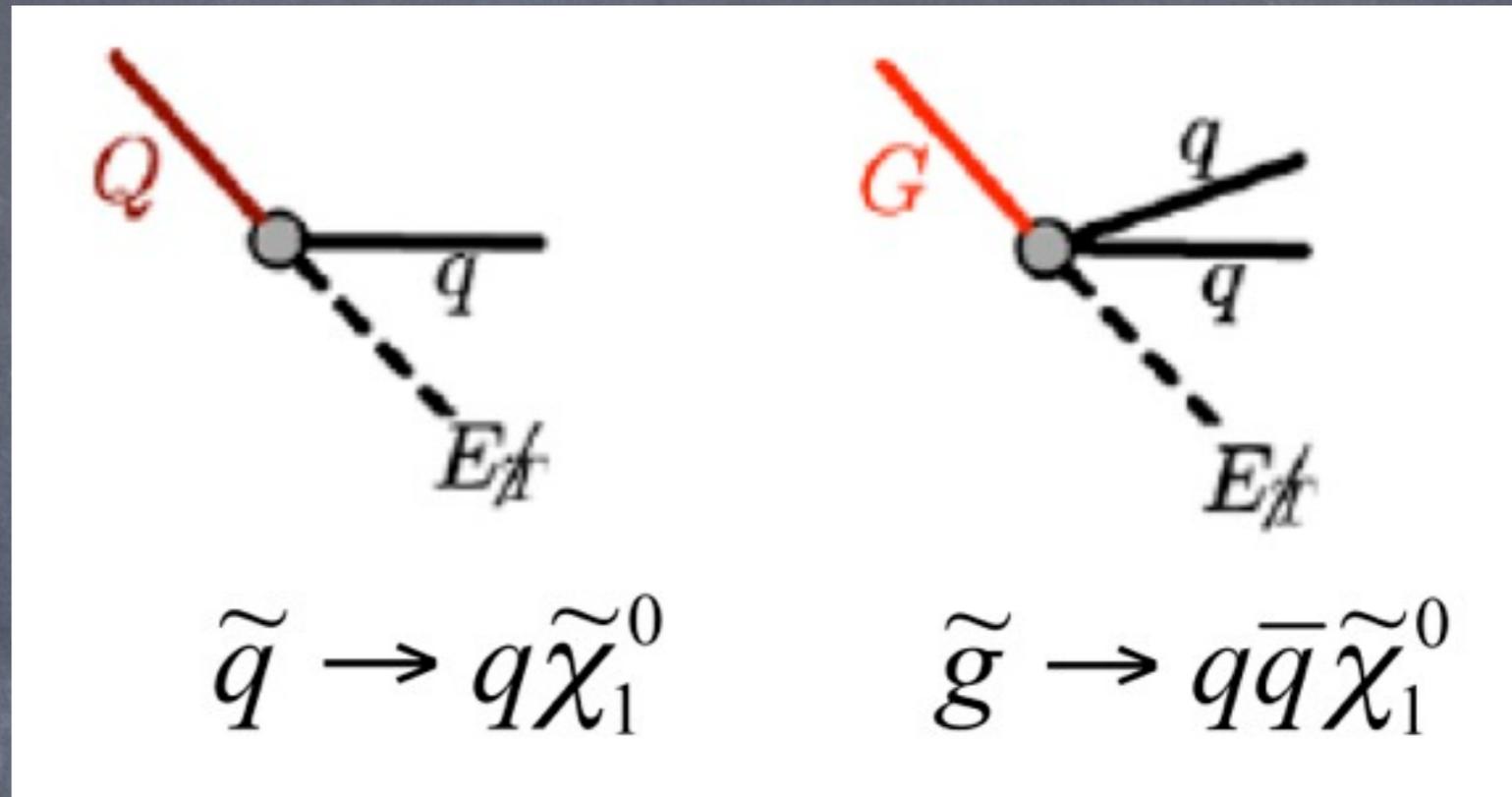
SUSY Limits: Topological Approach

- Results of SUSY searches usually presented as limits on mSUGRA parameters
- More practical and illuminating are limits on cross-section times branching fractions for given decay topologies of gluinos and squarks

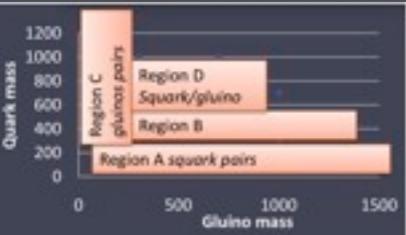


Simplest Topology: Jets + MET

- Assume R-parity conservation and neutralino LSP
- Squark can decay to 1 quark + neutralino LSP
- Glucino can decay to 2 quarks + neutralino LSP
- Most generic SUSY signatures, model independent if inclusive

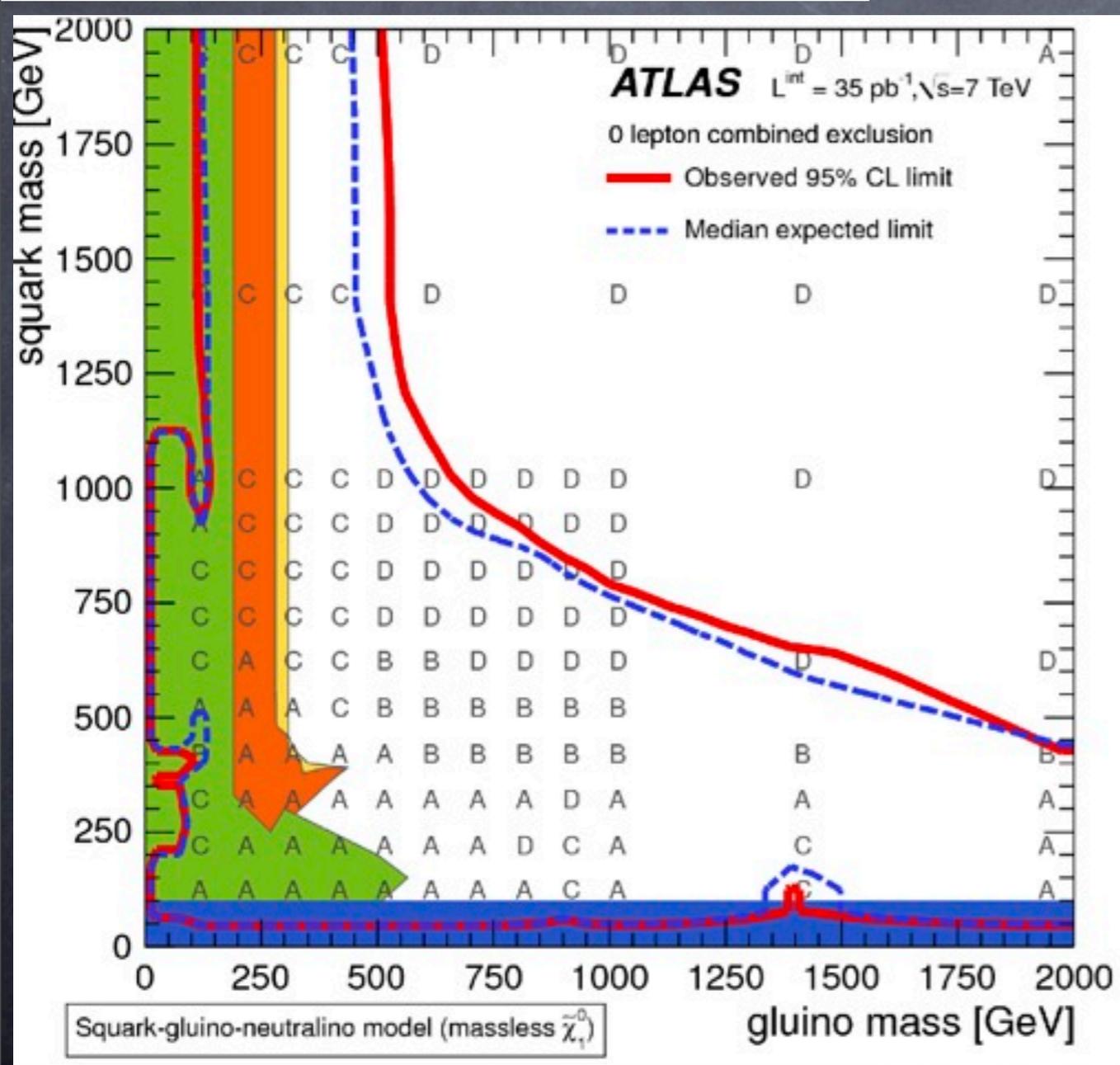


Jets+MET searches at ATLAS



	A	B	C	D
Pre-selection				
Number of required jets	≥ 2	≥ 2	≥ 3	≥ 3
Leading jet p_T [GeV]	> 120	> 120	> 120	> 120
Other jet(s) p_T [GeV]	> 40	> 40	> 40	> 40
E_T^{miss} [GeV]	> 100	> 100	> 100	> 100
Final selection				
$\Delta\phi(\text{jet}, \vec{P}_T^{\text{miss}})_{\text{min}}$	> 0.4	> 0.4	> 0.4	> 0.4
$E_T^{\text{miss}}/m_{\text{eff}}$	> 0.3	-	> 0.25	> 0.25
m_{eff} [GeV]	> 500	-	> 500	> 1000
m_{T2} [GeV]	-	> 300	-	-

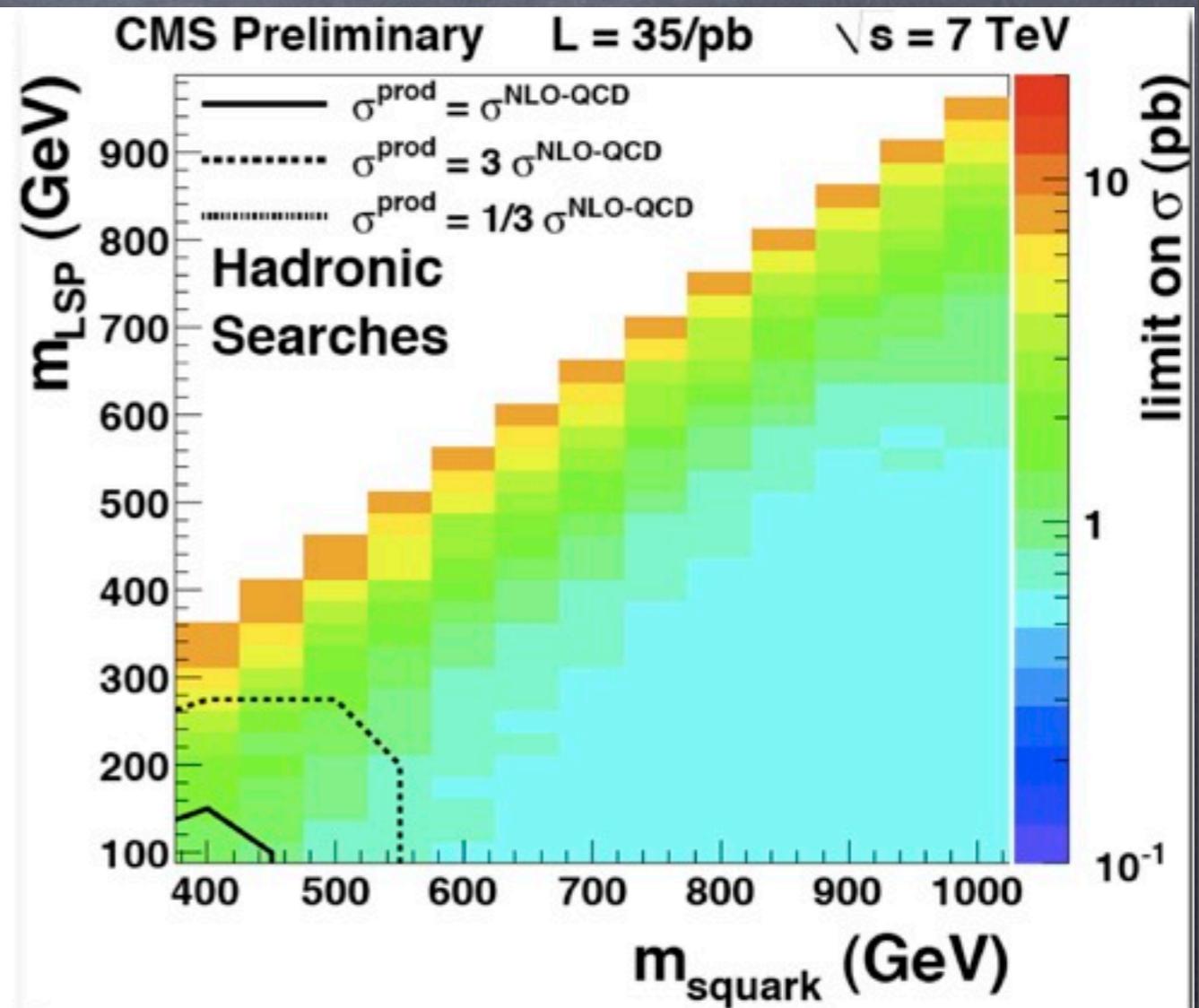
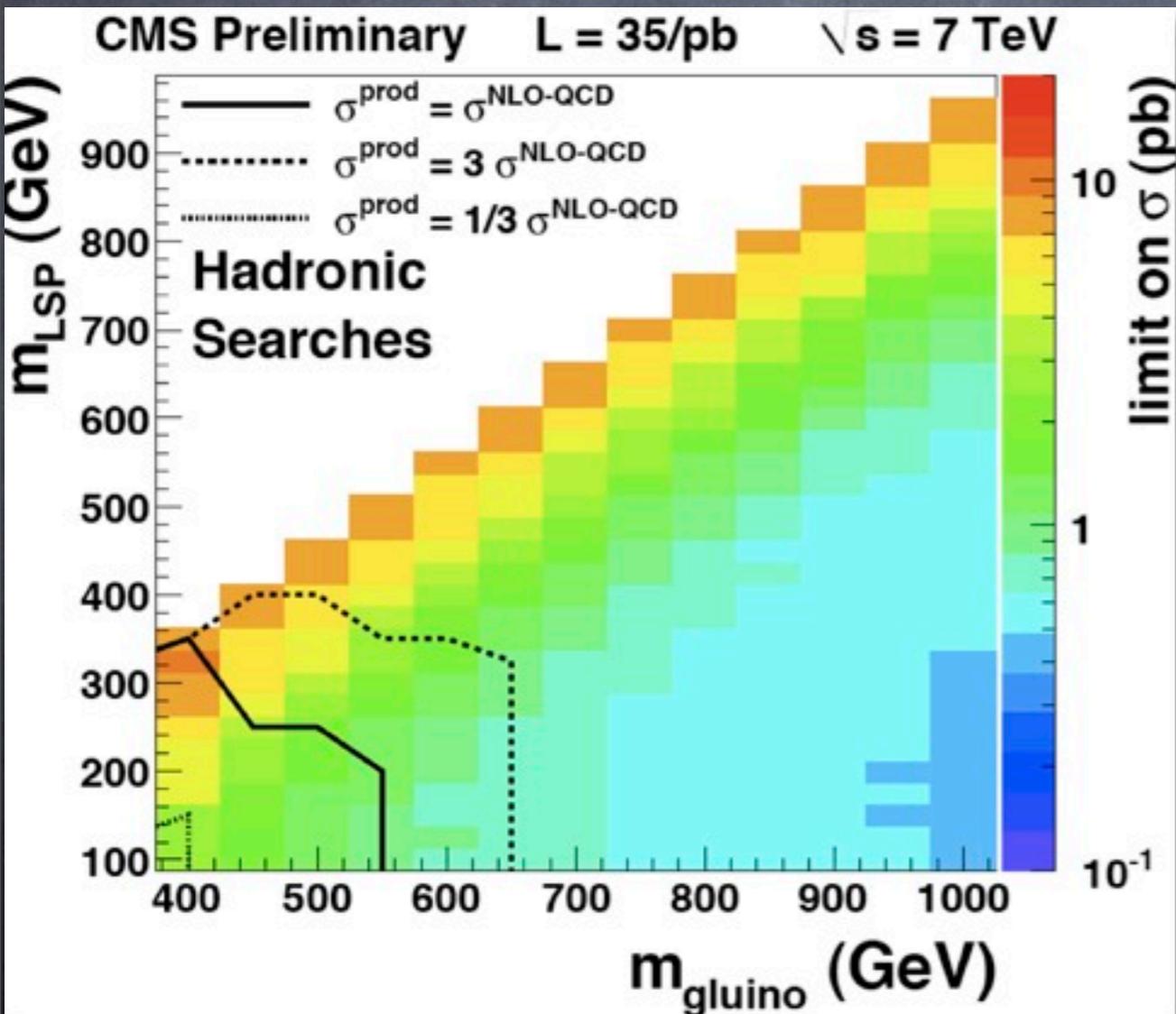
	Signal region A	Signal region B	Signal region C	Signal region D
QCD	$7_{-7}^{+8}[\text{u+j}]$	$0.6_{-0.6}^{+0.7}[\text{u+j}]$	$9_{-9}^{+10}[\text{u+j}]$	$0.2_{-0.2}^{+0.4}[\text{u+j}]$
W+jets	$50 \pm 11[\text{u}]_{-10}^{+14}[\text{j}] \pm 5[\mathcal{L}]$	$4.4 \pm 3.2[\text{u}]_{-0.8}^{+1.5}[\text{j}] \pm 0.5[\mathcal{L}]$	$35 \pm 9[\text{u}]_{-8}^{+10}[\text{j}] \pm 4[\mathcal{L}]$	$1.1 \pm 0.7[\text{u}]_{-0.3}^{+0.2}[\text{j}] \pm 0.1[\mathcal{L}]$
Z+jets	$52 \pm 21[\text{u}]_{-11}^{+15}[\text{j}] \pm 6[\mathcal{L}]$	$4.1 \pm 2.9[\text{u}]_{-0.8}^{+2.1}[\text{j}] \pm 0.5[\mathcal{L}]$	$27 \pm 12[\text{u}]_{-6}^{+10}[\text{j}] \pm 3[\mathcal{L}]$	$0.8 \pm 0.7[\text{u}]_{-0.0}^{+0.6}[\text{j}] \pm 0.1[\mathcal{L}]$
$t\bar{t}$ and t	$10 \pm 0[\text{u}]_{-3}^{+3}[\text{j}] \pm 1[\mathcal{L}]$	$0.9 \pm 0.1[\text{u}]_{-0.3}^{+0.4}[\text{j}] \pm 0.1[\mathcal{L}]$	$17 \pm 1[\text{u}]_{-4}^{+6}[\text{j}] \pm 2[\mathcal{L}]$	$0.3 \pm 0.1[\text{u}]_{-0.1}^{+0.2}[\text{j}] \pm 0.0[\mathcal{L}]$
Total SM	$118 \pm 25[\text{u}]_{-23}^{+32}[\text{j}] \pm 12[\mathcal{L}]$	$10.0 \pm 4.3[\text{u}]_{-1.9}^{+4.0}[\text{j}] \pm 1.0[\mathcal{L}]$	$88 \pm 18[\text{u}]_{-18}^{+26}[\text{j}] \pm 9[\mathcal{L}]$	$2.5 \pm 1.0[\text{u}]_{-0.4}^{+1.0}[\text{j}] \pm 0.1[\mathcal{L}]$
Data	87	11	66	2



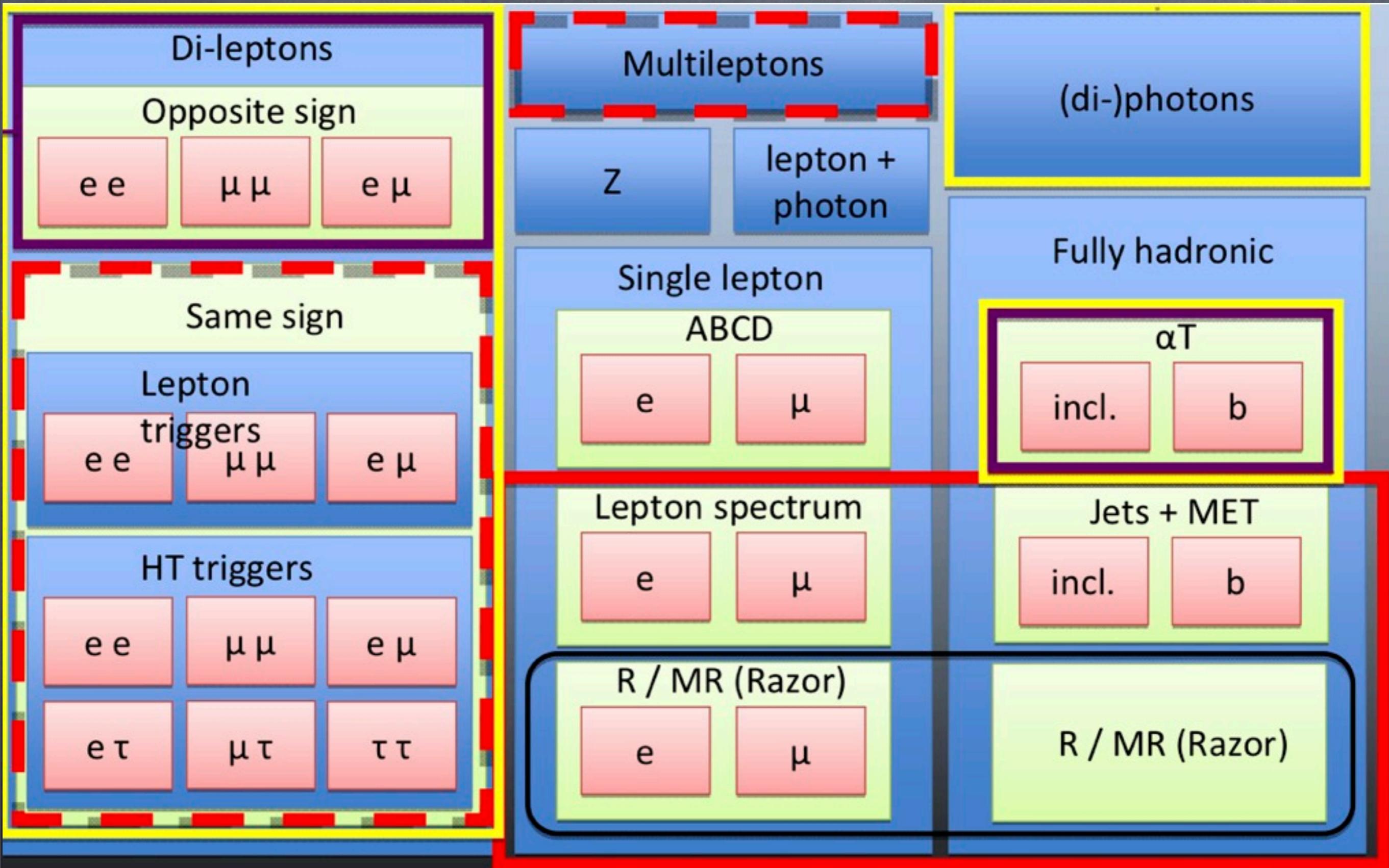
- 4 search regions targeting different squark/gluino masses
- Robust limit of 500 GeV on the gluino mass
- For equal squark and gluino masses the limit is more stringent, roughly 800 GeV

CMS JETS+MET search

- Several searches with different methods: HT and αT variables, "Razor" variable
- Model independent limits as a function of the LSP mass



SUSY is fun for experimentalists!

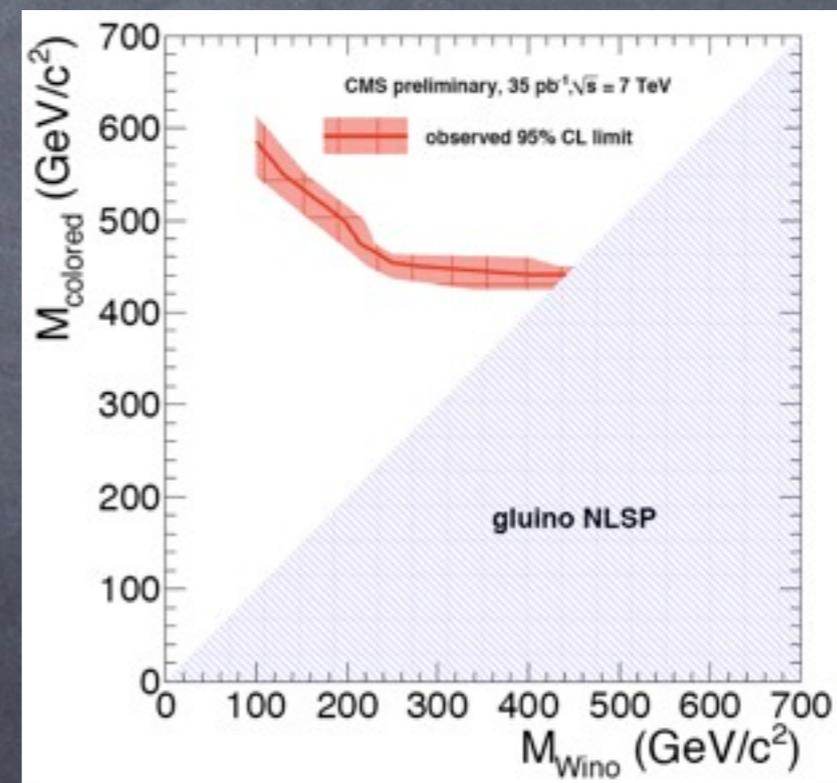
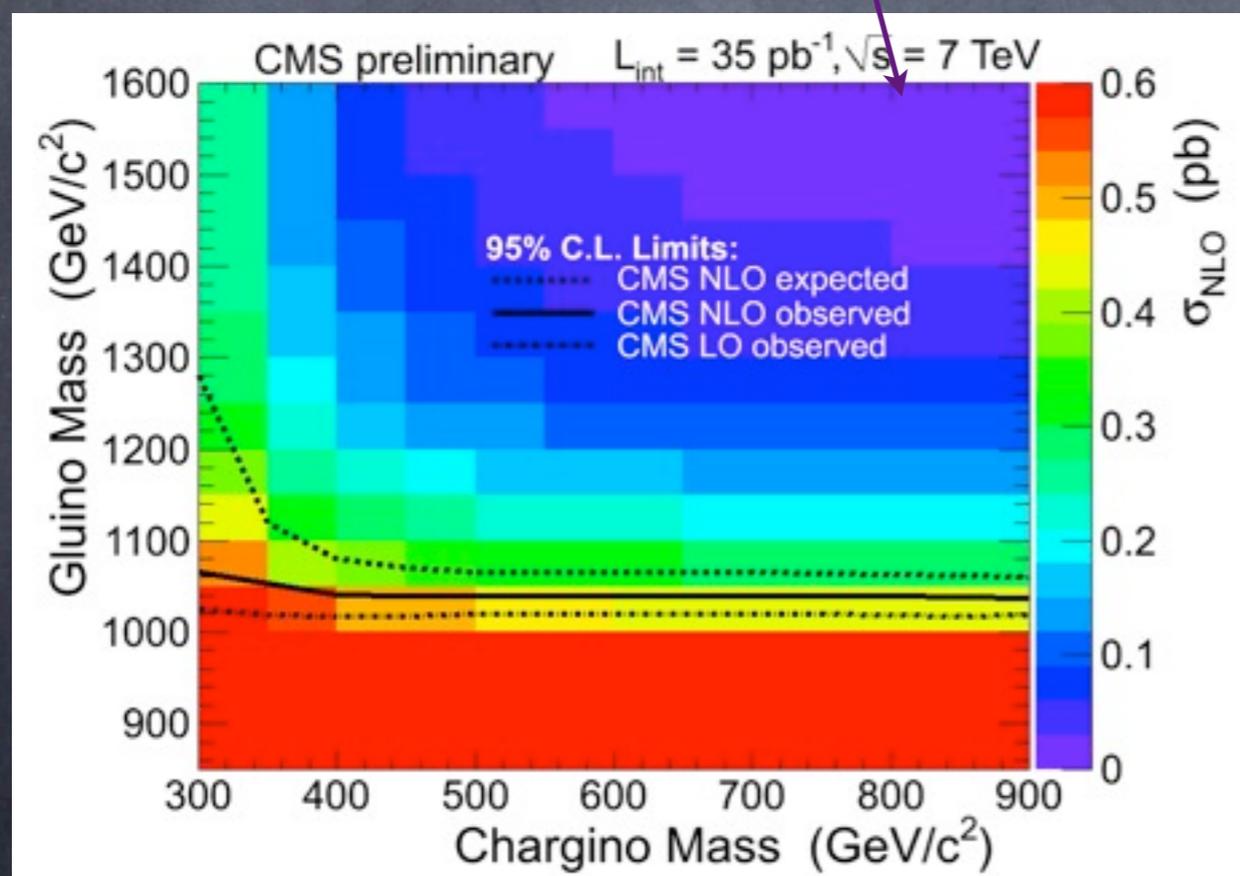
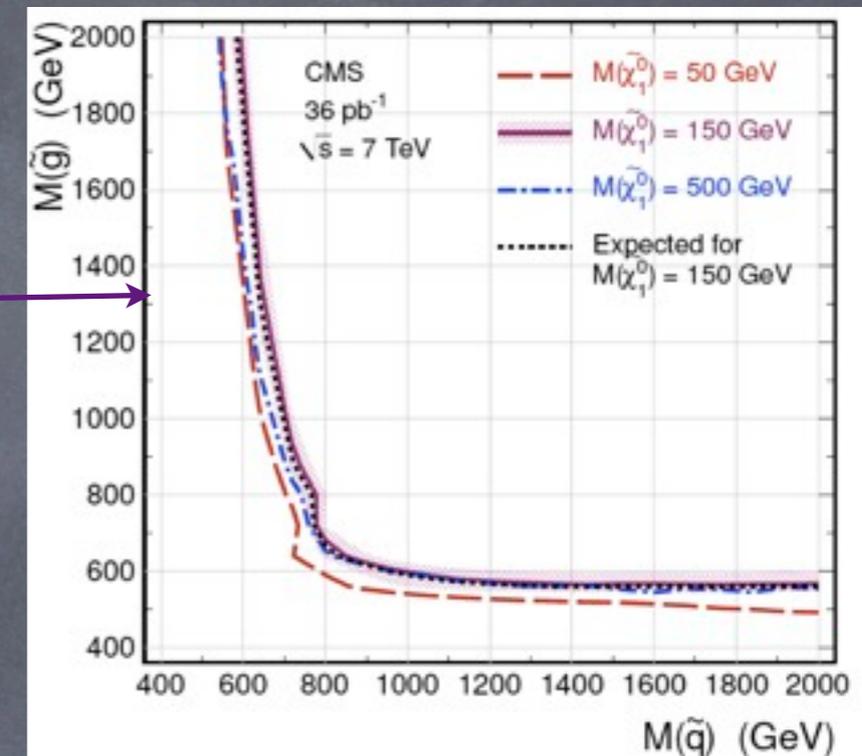


Some More Results

*Diphoton + Jets + MET

*Photon + Lepton + Jets + MET

*3 or Leptons + MET



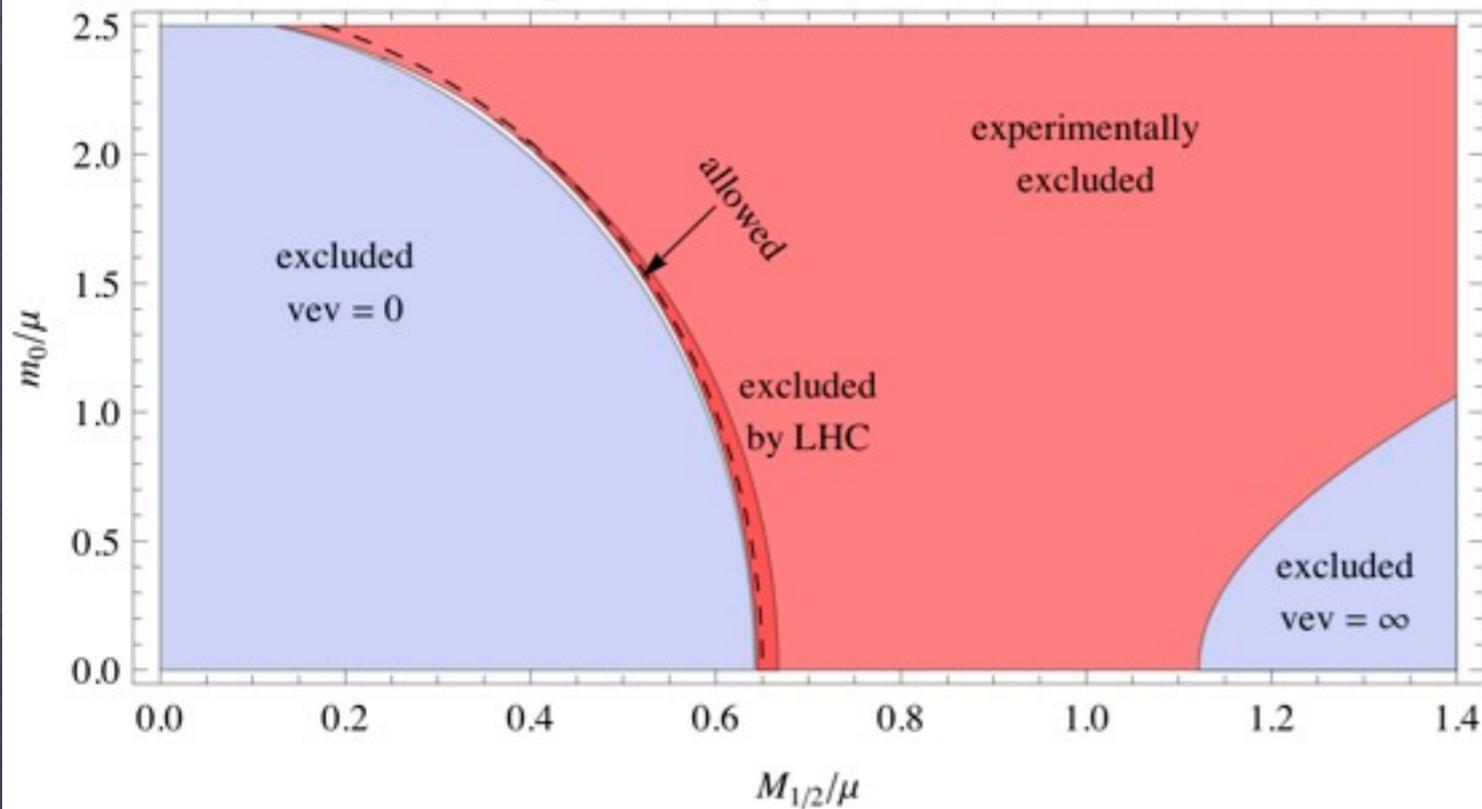
SUSY limits summary

- Robust limit of 500 GeV on gluino mass
- If squarks and gluinos have comparable masses, the limit goes up to 800 GeV
- Even stronger limits possible if the dominant decay chain produces 2 leptons
- No new limits on stops yet, currently $m_{\text{stop}} > m_{\text{top}}$

Interpretation of LHC SUSY Limits

$$M_{\frac{1}{2}}^2 \approx 0.2m_0^2 + 0.7M_3^2 - 2\mu^2 = (91 \text{ GeV})^2 \times 50 \left(\frac{M_3}{780 \text{ GeV}} \right)^2 + \dots$$

CMSSM parameter space with $\tan\beta = 3, A_0 = 0$



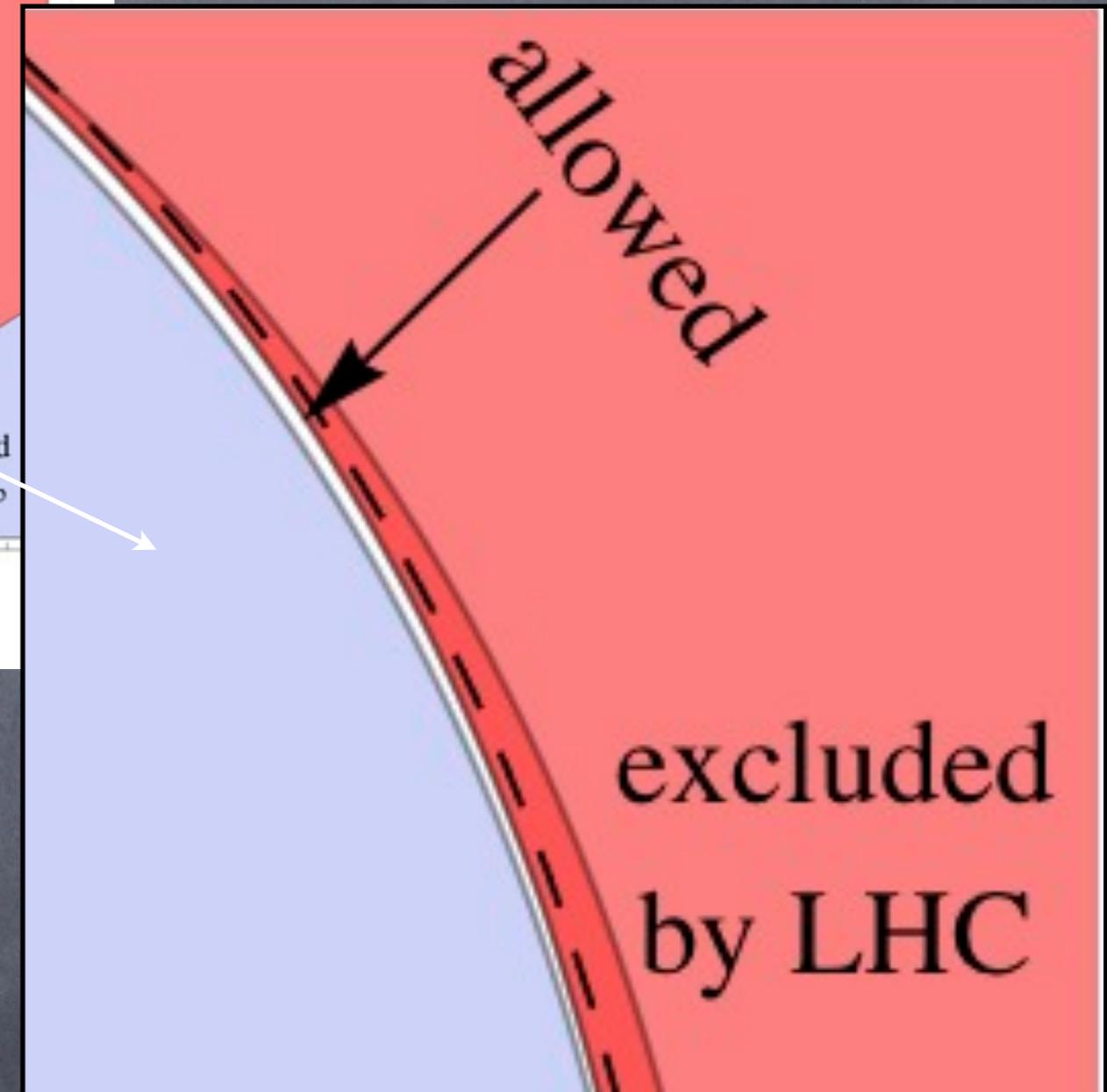
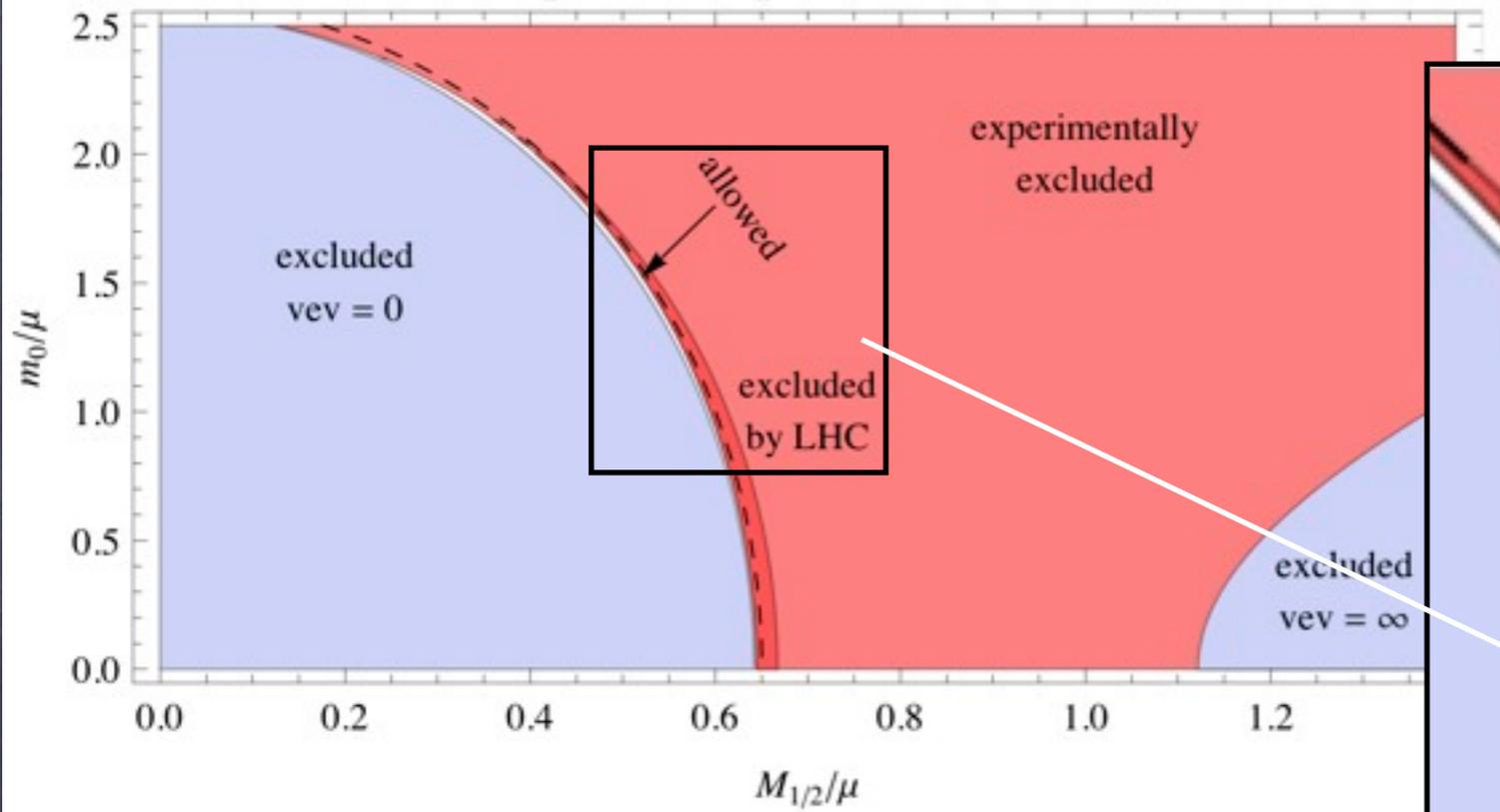
- LHC further shrinks the parameter space, kettling the parameters toward the no EW-breaking boundary

Strumia, 1101.2195

Interpretation of LHC SUSY Limits

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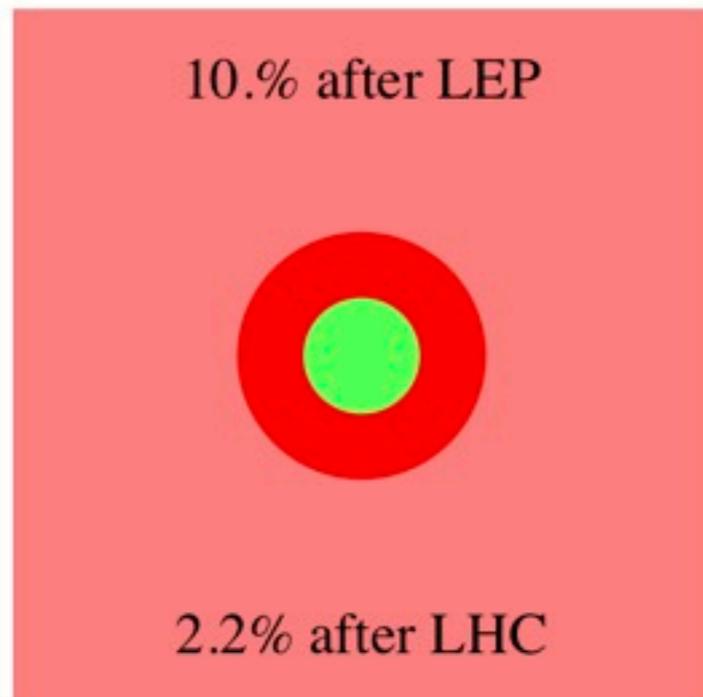
- LHC further shrinks the parameter space, kettling the parameters toward the no EW-breaking boundary.

Strumia, 1101.2195

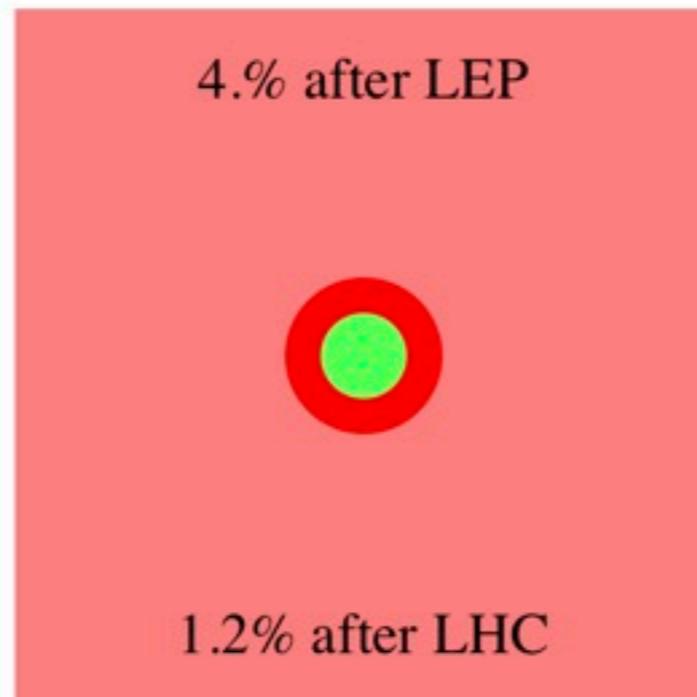
Interpretation of LHC SUSY Limits

Fraction of surviving CMSSM parameter space

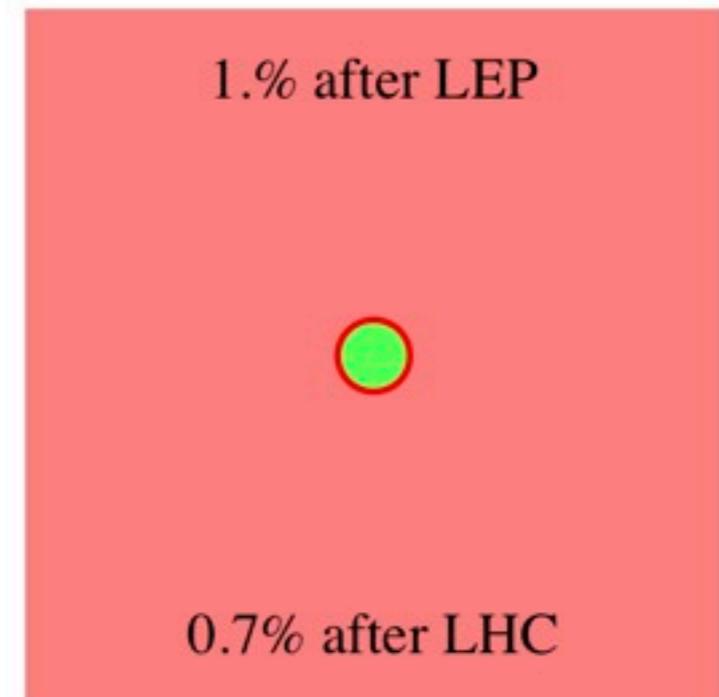
any m_h



$m_h > 100$ GeV



$m_h > 110$ GeV



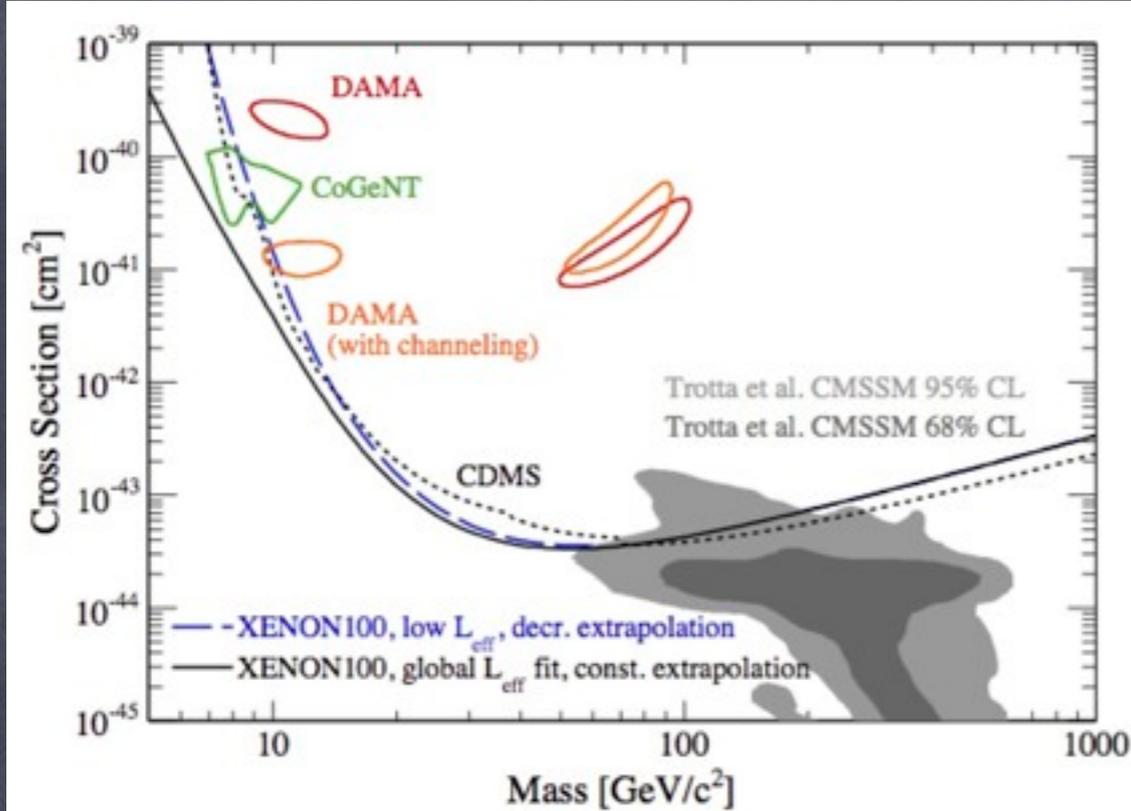
- LHC impact on MSSM-like scenarios is not significant. LEP limits on the Higgs mass already required that SUSY is heavy and fine-tuning is severe. LHC just confirmed that independently.
- But important impact on non-MSSM SUSY scenarios that avoid fine-tuning thanks to additional contributions to the Higgs mass. Now it's probably impossible to find a SUSY scenario where fine-tuning is better than 1 percent.

Back to the Frog Metaphor



The Water is Boiling...

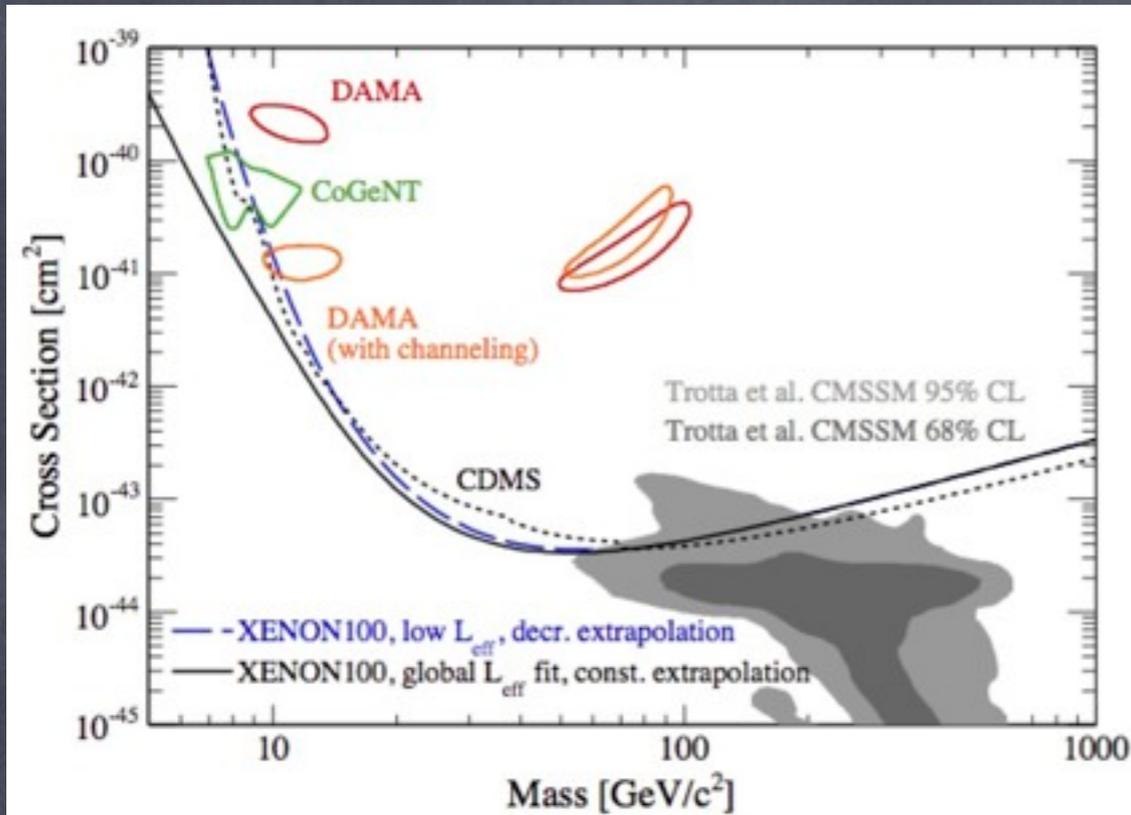
MSSM and Dark Matter



We are about to explore the most interesting region of the SUSY parameter space. Dark matter is just behind the corner!

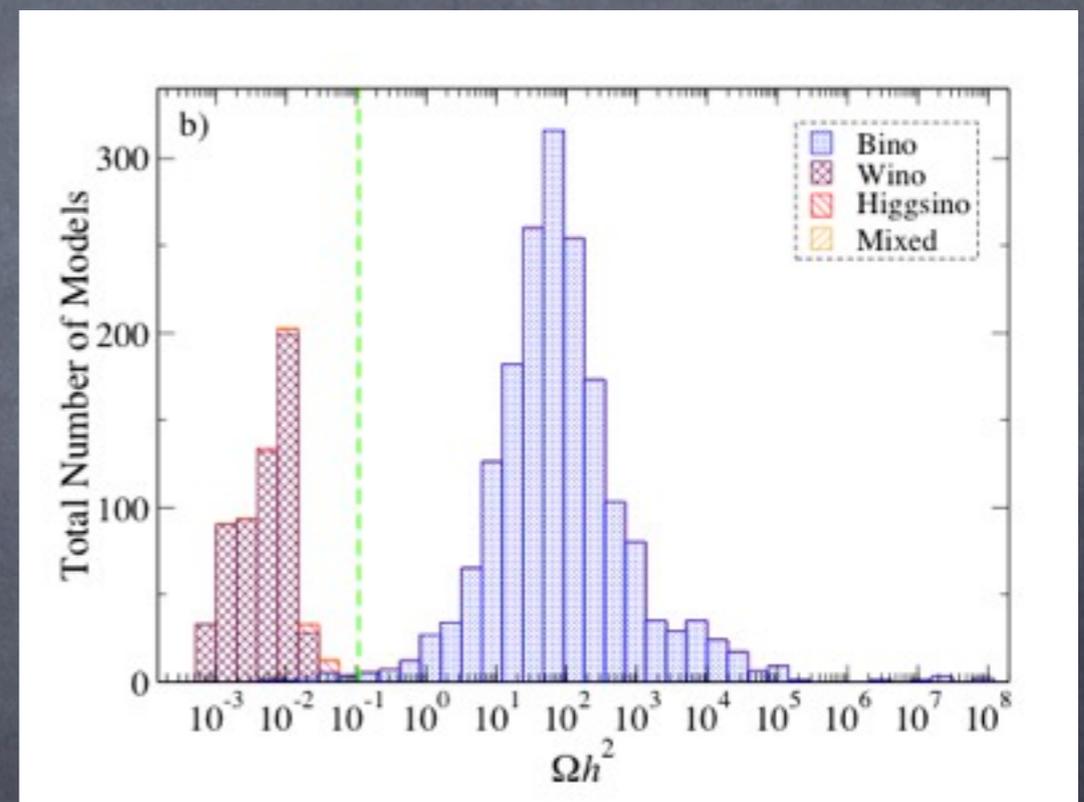
Baer, 1012.0248

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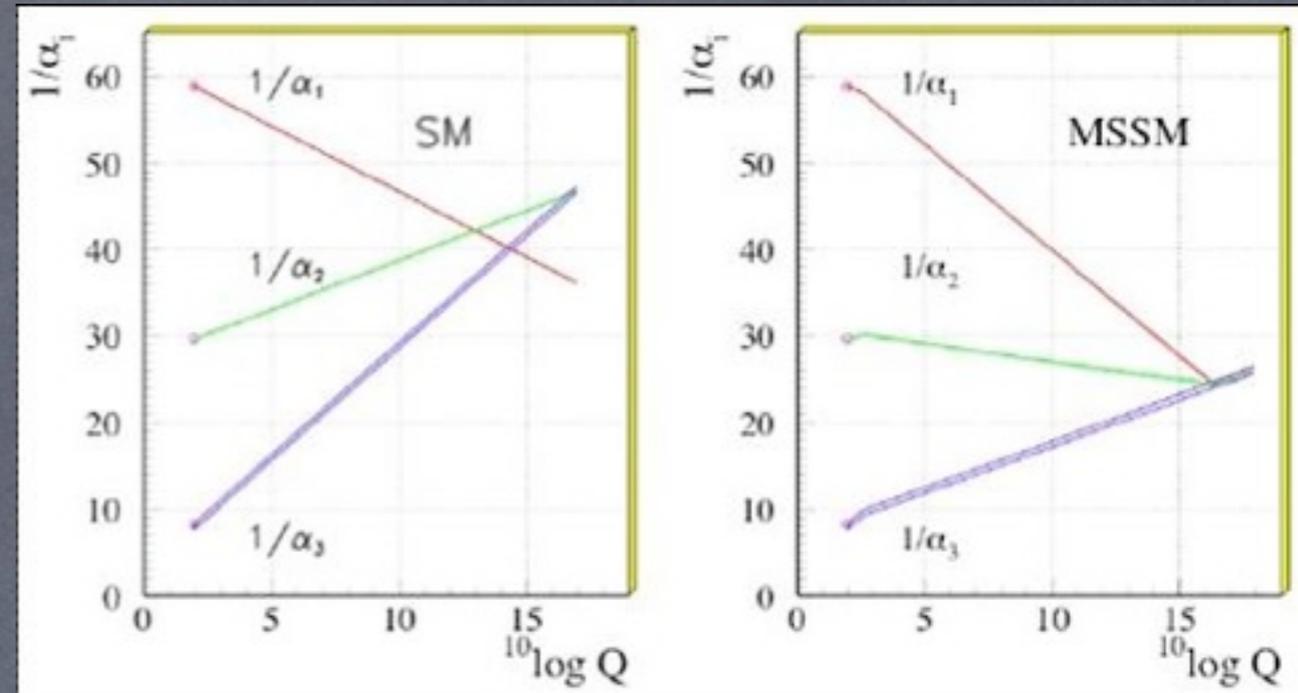
Dark matter in MSSM is fine-tuned, on top of electroweak fine-tuning. Bino dark matter typically gives too much relic abundance, whereas wino or Higgsino give too little.



Baer, 1012.0248

MSSM and Unification

- I have no snide remarks about this one...
- ...except that the strongest argument for SUSY is its connection to another hypothetical theory ;-)









90's

00's

10's



Images drawn by sister of Colin Bennett and presented by CMS at Moriond

My take on SUSY

- SUSY explains the weak scale by connecting it to the supersymmetry breaking scale:

$$m_{\text{SUSY}} \sim m_Z$$

- If SUSY was there it probably would have shown up at LEP. It would probably also show up indirectly in thousands low-energy measurements.
- Limits from Higgs searches at LEP and direct limits from the LHC imply m_{SUSY} of order 1 TeV, which corresponds to at least 1% fine-tuning
- If one can leave with 1% fine-tuning one can probably leave with 0.1% fine-tuning. The latter corresponds to $m_{\text{SUSY}} \geq 3$ TeV, that is no SUSY at the LHC.

Should we then search for SUSY?

- ✿ It predicts well defined signatures that can be searched for in colliders and in other experiments
- ✿ Signature-based SUSY searches apply to a much wider class of models (dark matter, extra dimensions, T-parity little Higgs, etc.)

Should we then search for SUSY?

YES

- ✿ It predicts well defined signatures that can be searched for in colliders and in other experiments
- ✿ Signature-based SUSY searches apply to a much wider class of models (dark matter, extra dimensions, T-parity little Higgs, etc.)

Estimated LHC Reach after 1fb-1 Alves et al. 1102.5338

