Expression of Interest: **Proposal to search for Heavy Neutral Leptons at the SPS**

(CERN-SPSC-2013-024 / SPSC-EOI-010)

On behalf of:

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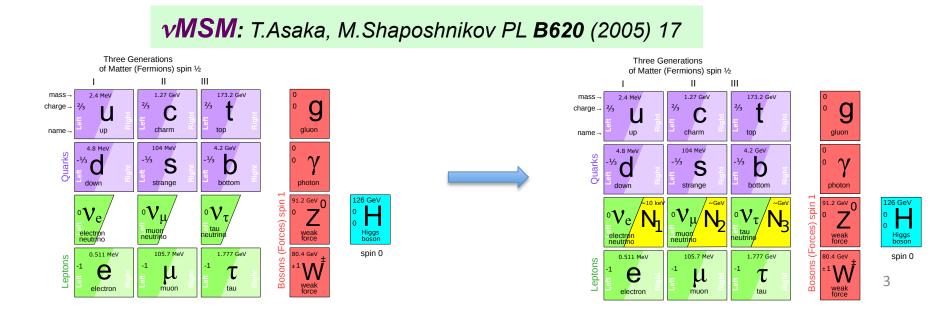
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Triumph of the Standard Model



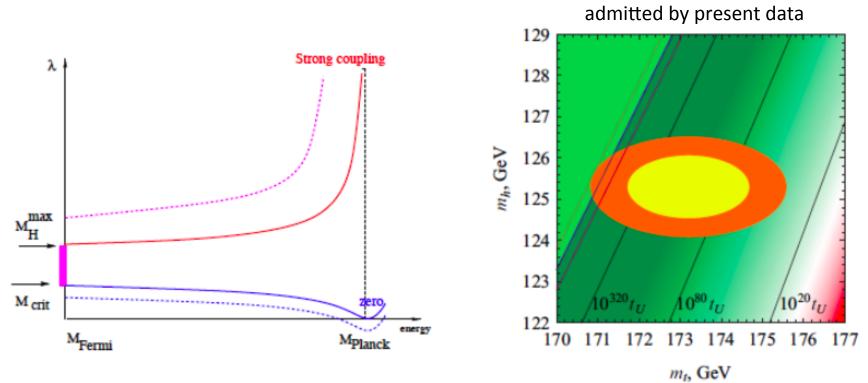
Theoretical motivation

- Discovery of the 126 GeV Higgs boson → Triumph of the Standard Model The SM may work successfully up to Planck scale !
- SM is unable to explain:
 - Neutrino masses
 - Excess of matter over antimatter in the Universe
 - The nature of non-baryonic Dark Matter
- All three issues can be solved by adding three new fundamental fermions, right-handed Majorana Heavy Neutral Leptons (HNL): N₁, N₂ and N₃



SM may well be a consistent effective theory all the way up to the Plank scale

- ✓ M_H < 175 GeV → SM is a weakly coupled theory up to the Plank energies !
- ✓ M_H > 111 GeV → EW vacuum is stable or metastable with a lifetime greatly exceeding the age of our Universe (Espinosa et al)

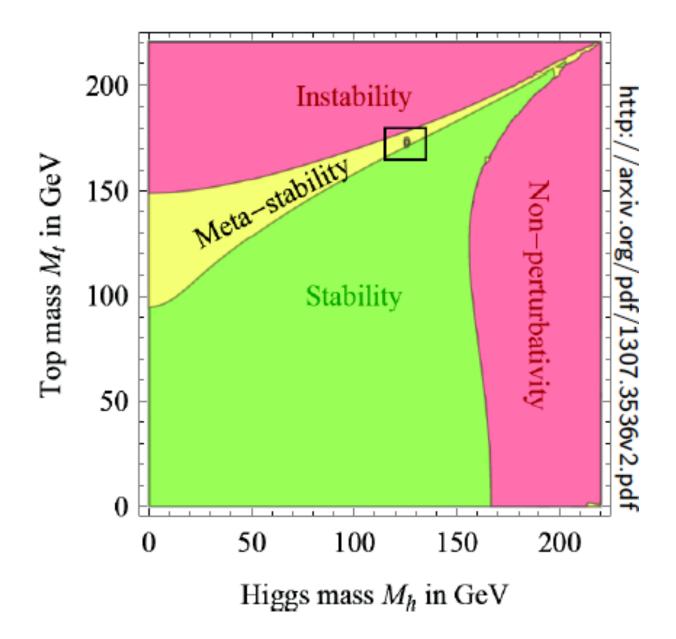


✓ No sign of New Physics seen

4

Stable vacuum is perfectly

Hard to believe that this is a pure coincidence !



No sign of New Physics seen What is not found..

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: SUSY 2013

	Model	e, μ, τ, γ	Jets	E ^{miss}	∫£ dt[fb	⁻¹] Mass limit	J	Reference
Inclusive Searches	$\begin{array}{l} \text{MSUGRA/CMSSM} \\ \text{MSUGRA/CMSSM} \\ \text{MSUGRA/CMSSM} \\ \hline \\ \tilde{q}\tilde{q}, \bar{q} \rightarrow q \tilde{x}_{1}^{D} \\ \hline \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q \tilde{x}_{1}^{D} \\ \hline \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q (\mathcal{U}/\mathcal{U}/\mathcal{W}) \tilde{x}_{1}^{D} \\ \hline \\ \tilde{g}\tilde{g}, \tilde{g} \rightarrow q q (\mathcal{U}/\mathcal{U}/\mathcal{W}) \tilde{x}_{1}^{D} \\ \hline \\ \text{GMS8} (\tilde{\ell} \text{NLSP}) \\ \hline \\ \text{GMS8} (\tilde{\ell} \text{NLSP}) \\ \hline \\ \text{GGM} (\text{bino } \text{NLSP}) \\ \hline \\ \text{GGM} (\text{higgsino-bino } \text{NLSP}) \\ \hline \\ \hline \\ \text{GGM} (\text{higgsino-bino } \text{NLSP}) \\ \hline \\ \hline \\ \text{Gravitino } \text{LSP} \\ \hline \\ \end{array}$	$\begin{matrix} 0 \\ 1 & e, \mu \\ 0 \\ 0 \\ 1 & e, \mu \\ 2 & e, \mu \\ 2 & e, \mu \\ 2 & e, \mu \\ 1.2 & \tau \\ 2 & \gamma \\ 1 & e, \mu + \gamma \\ \gamma \\ 2 & e, \mu (Z) \\ 0 \end{matrix}$	2-6 jets 3-6 jets 2-6 jets 2-6 jets 3-6 jets 0-3 jets 2-4 jets 0-2 jets 1 b 0-3 jets mono-jet	Yes Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.3 20.3 20.3 20.3 20.3 20.3 4.7 20.7 4.8 4.8 4.8 4.8 5.8 10.5	4: 2 1.7 Te' 4: 2 1.2 TeV 2 1.1 TeV 4: 2 1.1 TeV 4: 2 1.1 TeV 4: 2 1.3 TeV 2 1.3 TeV 2 1.3 TeV 2 1.18 TeV 2 1.2 TeV 2 1.4 TeV 2 619 GeV 2 600 GeV 8 600 GeV 8 600 GeV	$ \begin{array}{l} \label{eq:constraint} \textbf{V} & m(\tilde{q}) = m(\tilde{g}) \\ & any m(\tilde{q}) \\ & any m(\tilde{q}) \\ & m(\tilde{\xi}_1^2) = 0 \text{GeV} \\ & m(\tilde{\xi}_1^2) = 0 \text{GeV} \\ & m(\tilde{\xi}_1^2) = 0 \text{GeV} \\ & tar \eta < 15 \\ & tar \eta < 16 \\ & m(\tilde{\xi}_1^2) > 50 \text{GeV} \\ & m(\tilde{\xi}_1^2) > 50 \text{GeV} \\ & m(\tilde{\xi}_1^2) > 50 \text{GeV} \\ & m(\tilde{\xi}_1^2) > 200 \text{GeV} \\ & m(\tilde{\xi}_1^2) > 200 \text{GeV} \\ & m(\tilde{g}_1^2) > 10^{-4} \text{eV} \\ \end{array} $	ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 1308.1841 ATLAS-CONF-2013-047 ATLAS-CONF-2013-047 ATLAS-CONF-2013-062 ATLAS-CONF-2013-069 1208.4688 ATLAS-CONF-2013-026 1209.0753 ATLAS-CONF-2012-144 1211.1167 ATLAS-CONF-2012-152 ATLAS-CONF-2012-152
3 rd gen. <u>ĕ</u> med.	$\begin{array}{c} \hat{g} \rightarrow b \tilde{b} \tilde{b}_{1}^{2} \\ \hat{g} \rightarrow t \tilde{t} \tilde{x}_{1}^{0} \\ \hat{g} \rightarrow t \tilde{t} \tilde{x}_{1} \\ \hat{g} \rightarrow b \tilde{t} \tilde{x}_{1} \end{array}$	Ο Ο Ο-1 e,μ Ο-1 e,μ	3 b 7-10 jets 3 b 3 b	Yes Yes Yes Yes	20.1 20.3 20.1 20.1	8 1.2 TeV 8 1.1 TeV 8 1.34 TeV 8 1.34 TeV	m(ℓ ³ / ₁)<600 GeV m(ℓ ³ / ₁)<550 GeV m(ℓ ³ / ₁)<400 GeV m(ℓ ³ / ₁)<400 GeV	ATLAS-CONF-2013-061 1308-1841 ATLAS-CONF-2013-061 ATLAS-CONF-2013-061
3 rd gen. squarks direct production	$ \begin{array}{l} \overline{b}_1 \overline{b}_1, \ \overline{b}_1 \rightarrow b \overline{k}_1^0 \\ \overline{b}_1 \overline{b}_1, \ \overline{b}_1 \rightarrow t \overline{k}_1^- \\ \overline{t}_1 \overline{t}_2 (light), \ \overline{t}_1 \rightarrow t \overline{k}_1^- \\ \overline{t}_1 \overline{t}_2 (light), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (medium), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (medium), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (medium), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{k}_1^0 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{t}_1 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{t}_1 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{t}_1 \\ \overline{t}_1 \overline{t}_1 (heavy), \ \overline{t}_1 \rightarrow t \overline{t}_1 \\ \overline{t}_1 \rightarrow t \overline{t}_1 \ \overline{t}_1 \ \overline{t}_1 \rightarrow t \overline{t}_1 \\ \overline{t}_1 \rightarrow t \overline{t}_1 \ \overline{t}_1 $	$\begin{array}{c} 0 \\ 2 \ e, \mu ({\rm SS}) \\ 1-2 \ e, \mu \\ 2 \ e, \mu \\ 2 \ e, \mu \\ 0 \\ 1 \ e, \mu \\ 0 \\ 1 \ e, \mu \\ 0 \\ 3 \ e, \mu (Z) \end{array}$	2 b 0-3 b 1-2 b 0-2 jets 2 jets 2 b 1 b 2 b 0n0-jet/c-ti 1 b 1 b	Yes Yes Yes Yes Yes Yes Yes Yes Yes	20.1 20.7 4.7 20.3 20.3 20.1 20.7 20.5 20.3 20.7 20.7	5ŋ 100-620 GeV 5ŋ 275-430 GeV t̄a 110-167 GeV t̄a 130-220 GeV t̄a 130-220 GeV t̄a 225-525 GeV t̄a 200-510 GeV t̄a 200-510 GeV t̄a 320-660 GeV t̄a 90-200 GeV t̄a 500 GeV t̄a 500 GeV	$\begin{split} m(\tilde{k}_1^3) &\!$	1308.2631 ATLAB-CONF-2013-007 1208.4305, 1209.2102 ATLAB-CONF-2013-048 ATLAB-CONF-2013-048 ATLAS-CONF-2013-037 ATLAS-CONF-2013-047 ATLAS-CONF-2013-048 ATLAS-CONF-2013-045 ATLAS-CONF-2013-045
EW direct	$ \begin{array}{l} \tilde{t}_{\perp,R}\tilde{\ell}_{\perp,R}, \tilde{t} \rightarrow \ell \tilde{x}_{1}^{0} \\ \tilde{x}_{1}^{-1}\tilde{\chi}_{1}^{-}, \tilde{\chi}_{1}^{-} \rightarrow \tilde{\ell} \nu(\ell \bar{\nu}) \\ \tilde{x}_{1}^{-1}\tilde{\chi}_{1}^{-}, \tilde{k}_{1}^{-} \rightarrow \tilde{\tau} \nu(\ell \bar{\nu}) \\ \tilde{x}_{1}^{-1}\tilde{k}_{0}^{0} \rightarrow \tilde{\ell}_{\nu}\nu^{\tilde{\ell}}_{\ell}(\bar{\nu} \bar{\nu}), (\bar{\nu}\tilde{\ell}_{\perp}\ell(\bar{\nu}\nu)) \\ \tilde{x}_{1}^{-1}\tilde{\chi}_{0}^{0} \rightarrow W \tilde{\chi}_{1}^{0}\ell \tilde{x}_{1}^{0} \\ \tilde{x}_{1}^{-1}\tilde{\chi}_{0}^{0} \rightarrow W \tilde{\chi}_{1}^{0}\ell h \tilde{\chi}_{1}^{0} \end{array} $	2 θ, μ 2 θ, μ 2 τ 3 θ, μ 3 θ, μ 1 θ, μ	0 - 0 2 b	Yes Yes Yes Yes Yes	20.3 20.3 20.7 20.7 20.7 20.7 20.3	λ 85-315 GeV X [±] ₁ 125-450 GeV X [±] ₁ 180-330 GeV X [±] ₁ , X [±] ₂ 600 GeV X [±] ₁ , X [±] ₂ 315 GeV X [±] ₁ , X [±] ₂ 285 GeV	$\begin{split} m(\tilde{\epsilon}_{1}^{2}) &= 0 \; \text{GeV} \\ m(\tilde{\epsilon}_{1}^{2}) &= 0 \; \text{GeV}, \; m(\tilde{\ell}, \tilde{\nu}) = 0.5 (m(\tilde{\epsilon}_{1}^{2}) + m(\tilde{\epsilon}_{1}^{2})) \\ m(\tilde{\epsilon}_{1}^{2}) &= 0 \; \text{GeV}, \; m(\tilde{\ell}, \tilde{\nu}) = 0.5 (m(\tilde{\epsilon}_{1}^{2}) + m(\tilde{\epsilon}_{1}^{2})) \\) &= m(\tilde{\epsilon}_{1}^{2}), \; m(\tilde{\epsilon}_{1}^{2}) = 0, \; m(\tilde{\ell}, \tilde{\nu}) = 0.5 (m(\tilde{\epsilon}_{1}^{2}) + m(\tilde{\epsilon}_{1}^{2})) \\ m(\tilde{\epsilon}_{1}^{2}) = m(\tilde{\epsilon}_{1}^{2}), \; m(\tilde{\epsilon}_{1}^{2}) = 0, \; \text{sleptons decoupled} \\ m(\tilde{\epsilon}_{1}^{2}) = m(\tilde{\epsilon}_{1}^{2}), m(\tilde{\epsilon}_{1}^{2}) = 0, \; \text{sleptons decoupled} \end{split}$	ATLAS-CONF-2013-049 ATLAS-CONF-2013-049 ATLAS-CONF-2013-028 ATLAS-CONF-2013-035 ATLAS-CONF-2013-035 ATLAS-CONF-2013-083
Long-Eved particles	Direct $\tilde{\chi}_1^+ \tilde{\chi}_1^-$ prod., long-lived $\tilde{\chi}_1^+$ Stable, stopped \tilde{g} R-hadron GMSB, stable $\tilde{\tau}, \tilde{\chi}_1^0 \rightarrow \tilde{\tau}(\tilde{e}, \tilde{\mu}) + \tau(e$ GMSB, $\tilde{\chi}_1^0 \rightarrow \gamma \tilde{G}$, long-lived $\tilde{\chi}_1^0$ $\tilde{q}\tilde{q}, \tilde{\chi}_1^0 \rightarrow qq\mu$ (RPV)	Disapp. trk Q t, μ) 1-2 μ 2 γ 1 μ, displ. vtx	1 jet 1-5 jets - -	Yes Yes Yes	20.3 22.9 15.9 4.7 20.3	X ⁺ 1 270 GeV 822 GeV B 875 GeV 822 GeV X ⁺ 1 475 GeV 1.0 TeV	$\begin{split} m(\tilde{r}_1^1) &= m(\tilde{r}_1^0) = 160 \text{ MeV}, r(\tilde{r}_1^1) = 0.2 \text{ ns} \\ m(\tilde{r}_1^1) &= 160 \text{ GeV}, 10 \ \mu \text{ s} \neq r(\tilde{g}) \neq 1000 \text{ s} \\ 10 < tan \beta < 0 \\ 0.4 < r(\tilde{r}_1^1) < 2 \text{ ns} \\ 1.5 < cr < 156 \text{ mm}, \text{ BR}(\mu) = 1, m(\tilde{r}_1^0) = 108 \text{ GeV} \end{split}$	ATLAS-CONF-2013-069 ATLAS-CONF-2012-057 ATLAS-CONF-2013-058 1304,6310 ATLAS-CONF-2013-082
RPV	$ \begin{array}{l} LFV\;\rho p \! \rightarrow \! \bar{v}_{\tau} + X, \bar{v}_{\tau} \! \rightarrow \! e + \mu \\ LFV\;\rho p \! \rightarrow \! \bar{v}_{\tau} + X, \bar{v}_{\tau} \! \rightarrow \! e(\mu) + \tau \\ Bilinear\;RPV\;CMSSM \\ \bar{\mathcal{K}}_1^+ \bar{\mathcal{K}}_1^-, \bar{\mathcal{K}}_1^+ \! \rightarrow \! W \bar{\mathcal{K}}_1^0, \bar{\mathcal{K}}_1^0 \! \rightarrow \! e \bar{v}_{\rho}, e \mu \bar{\nu} \\ \bar{\mathcal{K}}_1^- \bar{\mathcal{K}}_1^-, \bar{\mathcal{K}}_1^+ \rightarrow W \bar{\mathcal{K}}_1^0, \bar{\mathcal{K}}_1^0 \! \rightarrow \! e \bar{\tau} \bar{\nu}_{\rho}, e r \bar{\nu}, \\ \bar{\mathcal{g}}_1^0 \! \rightarrow \! \bar{q} q \\ \bar{\mathcal{g}} \! \rightarrow \! \bar{q} 1, \bar{\mathcal{K}}_1^- \! \rightarrow \! b s \end{array} $	$\begin{array}{c} 2 \ e, \mu \\ 1 \ e, \mu + \tau \\ 1 \ e, \mu \\ e \end{array} \\ \begin{array}{c} 4 \ e, \mu \\ \tau \end{array} \\ \begin{array}{c} 3 \ e, \mu + \tau \\ 0 \\ 2 \ e, \mu \ (SS) \end{array}$	7 jets - - 6-7 jets 0-3 b	Yes Yes Yes Yes	4.6 4.7 20.7 20.7 20.3 20.7	7, 1.61 TeV 7, 1.1 TeV 6,8 1.2 TeV \tilde{x}_1^+ 760 GeV \tilde{x}_1^+ 350 GeV	$\begin{split} &\mathcal{X}_{111}=0.10, \ \mathcal{X}_{122}=0.05\\ &\mathcal{X}_{211}=0.10, \ \mathcal{X}_{1223}=0.05\\ &m(\tilde{q})=m(\tilde{g}), \ \sigma_{2,2P}<1\ mm\\ &m(\tilde{q}_1^2)=300\ GeV, \ \mathcal{X}_{212}>0\\ &m(\tilde{q}_1^2)=300\ GeV, \ \mathcal{X}_{212}>0\\ &m(\tilde{q}_1^2)=300\ GeV, \ \mathcal{X}_{212}>0\\ &BR(c)=BR(b)=BR(c)=0\% \end{split}$	1212.1272 1212.1272 ATLAS-CONF-2012-140 ATLAS-CONF-2013-036 ATLAS-CONF-2013-036 ATLAS-CONF-2013-091 ATLAS-CONF-2013-007
Other	Scalar gluon pair, sgluon $\rightarrow q\bar{q}$ Scalar gluon pair, sgluon $\rightarrow t\bar{t}$ WIMP interaction (D5, Dirac χ)	0 2 e,µ (SS) 0	4 jets 1 b mono-jet	Yes Yes	4.6 14.3 10.5	sgluan 100-287 GeV sgluan 800 GeV M' scale 704 GeV	inal. limit fram 1110.2683 m(z)<80 GeV, limit a1<887 GeV for D8	1210.4826 ATLAS-CONF-2013-051 ATLAS-CONF-2012-147
		√s = 8 TeV artial data	$\sqrt{s} = 0$ full of			10 ⁻¹ 1	Mass scale [TeV]	

ATLAS Preliminary

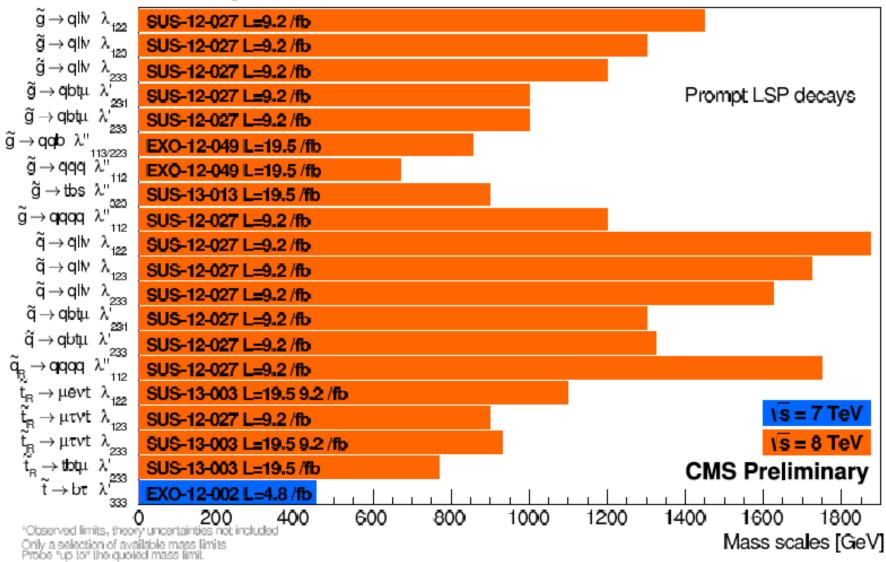
 $\int \mathcal{L} dt = (4.6 - 22.9) \text{ fb}^{-1}$ $\sqrt{s} = 7, 8 \text{ TeV}$

No sign of New Physics seen

What is not found..

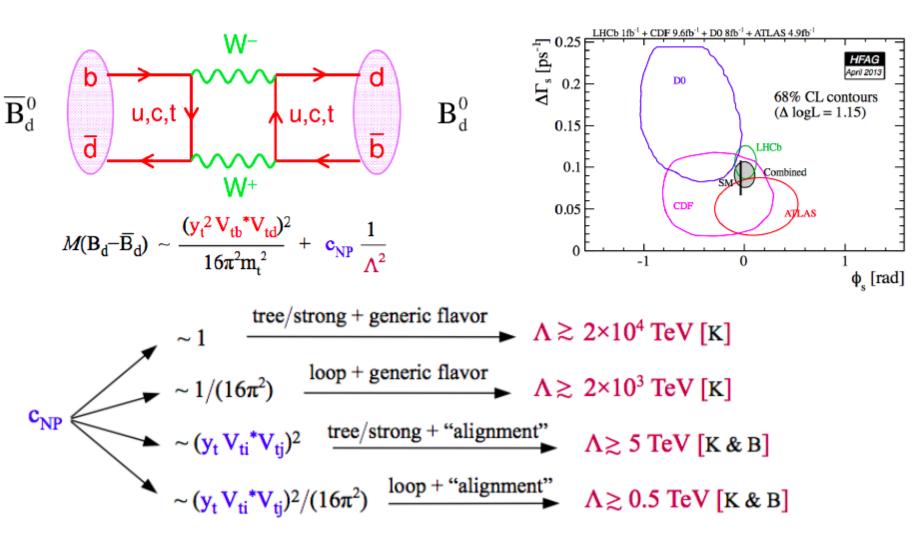
Summary of CMS RPV SUSY Results*

EPSHEP 2013



Bounds on the scale of New Physics

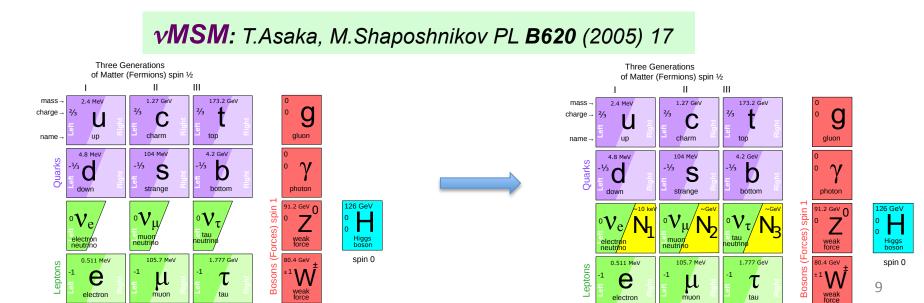
Most stringent limits come from observables in BB mixing



8

Theoretical motivation

- Discovery of the 126 GeV Higgs boson → Triumph of the Standard Model The SM may work successfully up to Planck scale !
- SM is unable to explain:
 - Neutrino masses & oscillations
 - Excess of matter over antimatter in the Universe
 - The nature of non-baryonic Dark Matter
- All three issues can be solved by adding three new fundamental fermions, right-handed Majorana Heavy Neutral Leptons (HNL): N₁, N₂ and N₃



See-saw generation of neutrino masses

Most general renormalisable Lagrangian of all SM particles (+3 singlets wrt the SM gauge group):

$$L_{\text{singlet}} = i\bar{N}_I\partial_\mu\gamma^\mu N_I - Y_{I\alpha}\bar{N}_I^c\tilde{H}L^c_\alpha - M_I\bar{N}_I^cN_I + \text{h.c.}$$

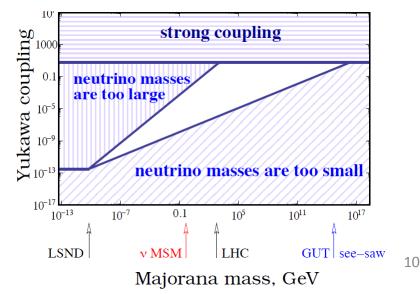
Yukawa term: mixing of N_I with active neutrinos to explain oscillations

Majorana term which carries no gauge charge

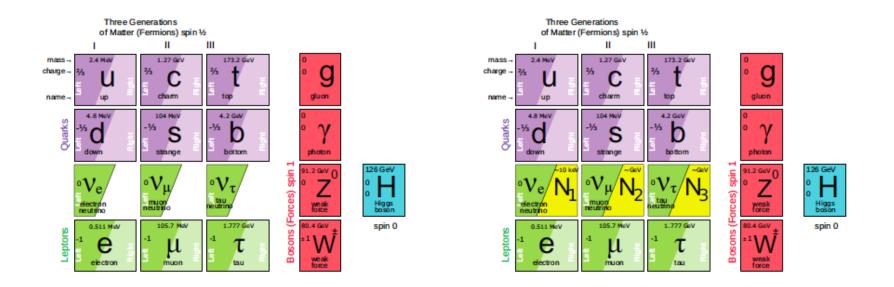
The scale of the active neutrino mass is given by the see-saw formula: $m_{\nu} \sim where m_D \sim Y_{I\alpha}v$ - typical value of the Dirac mass term

Example:

For $M \sim 1$ GeV and $m_v \sim 0.05$ eV it results in $m_D \sim 10$ keV and Yukawa coupling $\sim 10^{-7}$



The vMSM model



N = Heavy Neutral Lepton - HNL

Role of N_1 with mass in keV region: dark matter Role of N_2 , N_3 with mass in 100 MeV – GeV region: "give" masses to neutrinos and produce baryon asymmetry of the Universe Role of the Higgs: give masses to quarks, leptons, Z and W and inflate the Universe.

Masses and couplings of HNLs

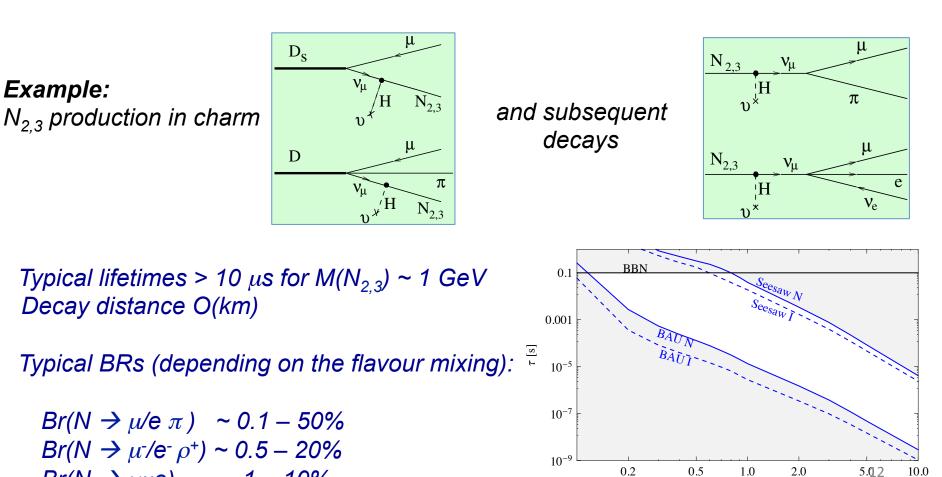
• N_1 can be sufficiently stable to be a DM candidate, $M(N_1) \sim 10 \text{ keV}$

•

 $Br(N \rightarrow v\mu e) \sim 1 - 10\%$

 M(N₂) ≈ M(N₃) ~ a few GeV → CPV can be increased dramatically to explain Baryon Asymmetry of the Universe (BAU)

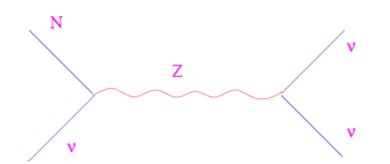
Very weak $N_{2,3}$ -to-v mixing (~ U^2) $\rightarrow N_{2,3}$ are much longer-lived than the SM particles



M [GeV]

Dark Matter candidate HNL N₁

Yukawa couplings are small \rightarrow *N* can be very stable.



Main decay mode: $N \rightarrow 3\nu$. Subdominant radiative decay channel: $N \rightarrow \nu\gamma$. For one flavour:

$$au_{N_1} = 10^{14}\, ext{years} \left(rac{10\ ext{keV}}{M_N}
ight)^5 \left(rac{10^{-8}}{ heta_1^2}
ight)$$

$$heta_1 = rac{m_D}{M_N}$$

Dark Matter candidate HNL N₁

DM particle is not stable. Main decay mode $N_1 \rightarrow 3\nu$ is not observable. Subdominant radiative decay channel: $N \rightarrow \nu\gamma$. Photon energy: $E_{\gamma} = \frac{M}{2}$

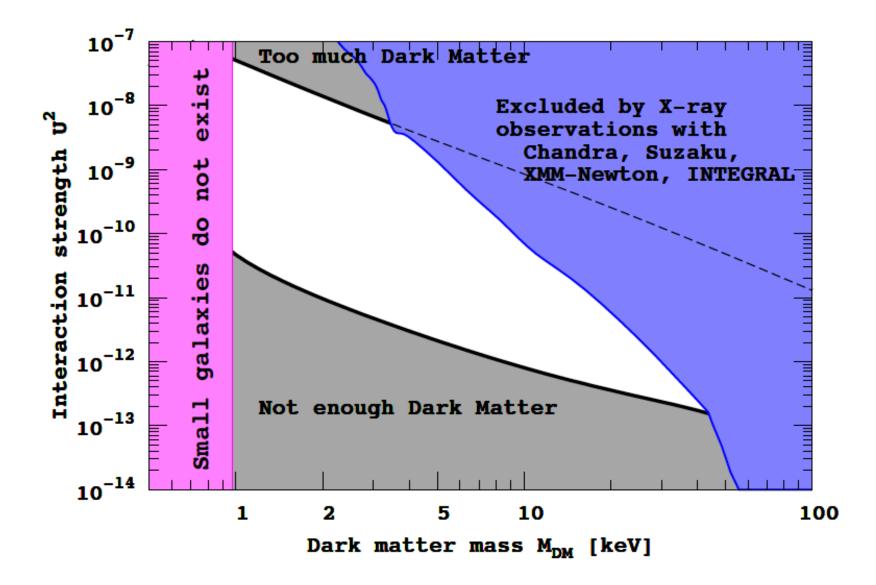
Radiative decay width:

$$\Gamma_{
m rad} = rac{9\,lpha_{ extsf{EM}}\,G_F^2}{256\cdot 4\pi^4}\,\sin^2(2 heta)\,{M_{ extsf{N}}}^5$$

Constraints on DM HNL N₁

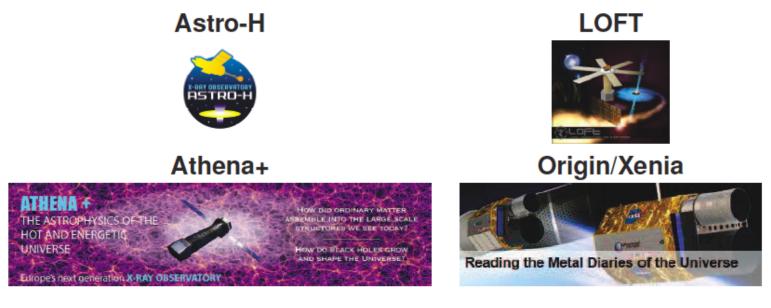
- ✓ **Stability** → N_1 must have a lifetime larger than that of the Universe
- ✓ **Production** → N_1 are created in the early Universe in reactions $l\bar{l} \rightarrow vN_1$, $q\bar{q} \rightarrow vN_1$ etc. Need to provide correct DM abundance
- ✓ Structure formation → N_1 should be heavy enough ! Otherwise its free streaming length would erase structure non-uniformities at small scales (Lyman- α forest spectra of distant quasars and structure of dwarf galaxies)
- ✓ X-ray spectra → Radiative decays N₁→_γv produce a mono-line in photon galaxies spectrum. This line has not yet been seen by X-ray telescopes (such as Chandra or XMM-Newton)

Allowed parameter space for DM HNL N₁



Searches for DM HNL N₁ in space

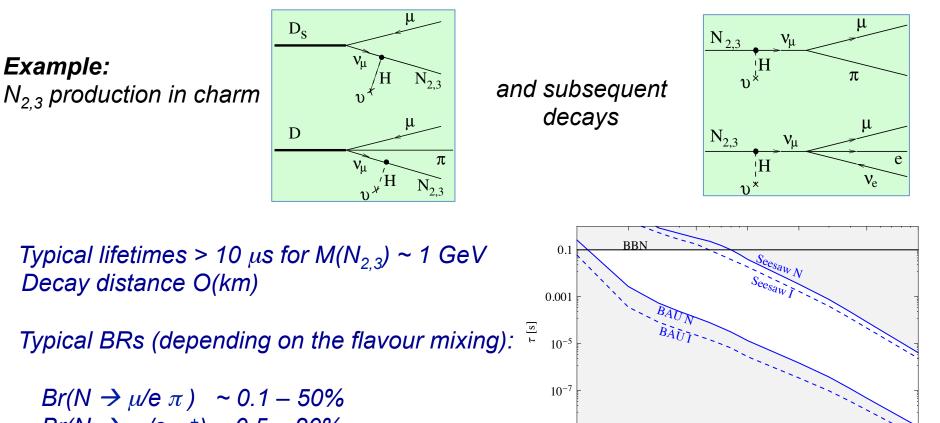
- Has been previously searched with XMM-Newton, Chandra, Suzaku, INTEGRAL
- Spectral resolution is not enough (required $\Delta E/E \sim 10^{-3}$)
- Proposed/planned X-ray missions with sufficient spectral resolution:



Masses and couplings of HNLs

- N_1 can be sufficiently stable to be a DM candidate, $M(N_1) \sim 10 \text{keV}$
- $M(N_2) \approx M(N_3) \sim a$ few GeV \rightarrow CPV can be increased dramatically to explain Baryon Asymmetry of the Universe (BAU)

Very weak $N_{2,3}$ -to-v mixing (~ U^2) $\rightarrow N_{2,3}$ are much longer-lived than the SM particles



 10^{-9}

0.2

0.5

1.0

M [GeV]

2.0

5.08

10.0

 $Br(N \rightarrow \mu^{-}/e^{-} \rho^{+}) \sim 0.5 - 20\%$ $Br(N \rightarrow \nu \mu e) \sim 1 - 10\%$

•

Baryon asymmetry

• CP is not conserved in vMSM

6 CPV phases in the lepton sector and 1 CKM phase in the quark sector (to be compared with only one CKM phase in the SM)

• Deviations from thermal equilibrium

- ✓ HNL are created in the early Universe
- ✓ CPV in the interference of HNL mixing and decay
- ✓ Lepton number goes from HNL to active neutrinos
- ✓ Then lepton number transfers to baryons in the equilibrium sphaleron processes

Constraints on BAU HNL N_{2,3}

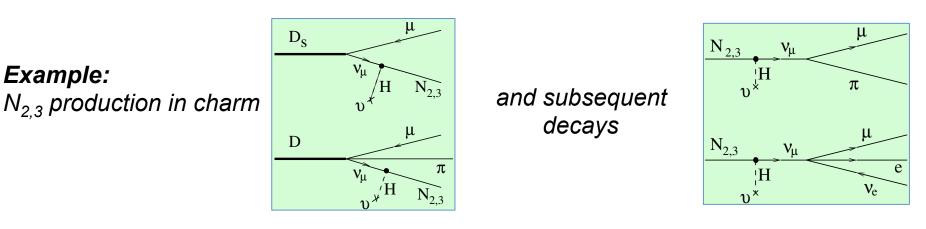
Baryon asymmetry is generated by CPV in HNL mixing and decays + sphalerons

- ✓ BAU generation requires out of equilibrium → mixing angle of $N_{2,3}$ can not be large
- ✓ To generate correct order of the active neutrino masses the mixing angle of $N_{2,3}$ to active neutrino can not be too small
- ✓ Decays of $N_{2,3}$ should keep BBN scenario working
- ✓ *Experimental constraints*
- **PS** Explanation of DM with N_1 reduces a number of free parameters \rightarrow Degeneracy of $N_{2,3}$ masses is required to ensure sufficient CPV

Masses and couplings of HNLs

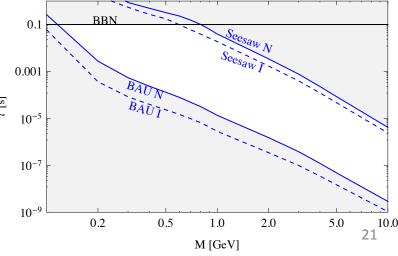
- N_1 can be sufficiently stable to be a DM candidate, $M(N_1) \sim 10 \text{keV}$
- M(N₂) ≈ M(N₃) ~ a few GeV → CPV can be increased dramatically to explain Baryon Asymmetry of the Universe (BAU) using sphaleron lepton-to-baryon number transformation

Very weak $N_{2,3}$ -to-v mixing (~ U^2) $\rightarrow N_{2,3}$ are much longer-lived than the SM particles

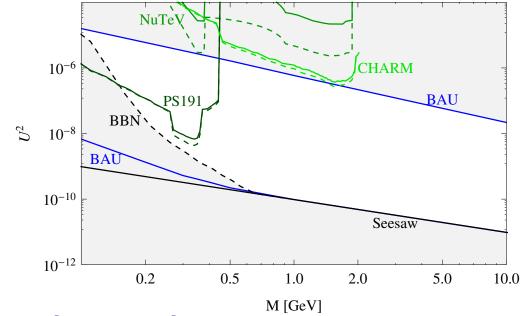


- Typical lifetimes > 10 μ s for $M(N_{2,3}) \sim 1$ GeV Decay distance O(km)
- Typical BRs (depending on the flavour mixing): [™]

Br(N → μ/e π) ~ 0.1 - 50% Br(N → μ⁻/e⁻ ρ⁺) ~ 0.5 - 20% Br(N → νμe) ~ 1 - 10%



Experimental and cosmological constraints



Recent progress in cosmology

- The sensitivity of previous experiments did not probe the interesting region for HNL masses above the kaon mass

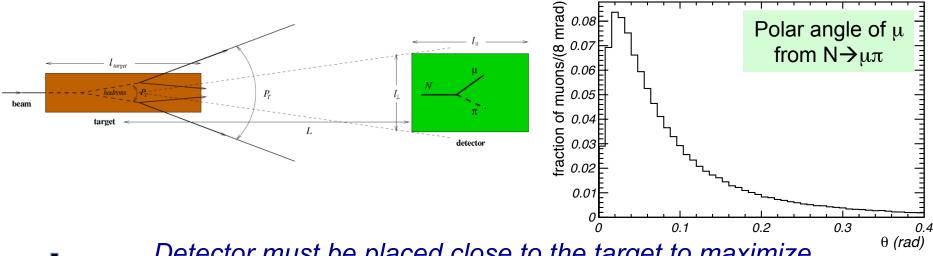
Strong motivation to explore cosmologically allowed parameter space **Proposal for a new experiment at the SPS to search for New long-lived Particles produced in charm decays** Experimentally this domain has not been very well explored ! 22

Experimental requirements

• Search for HNL in Heavy Flavour decays

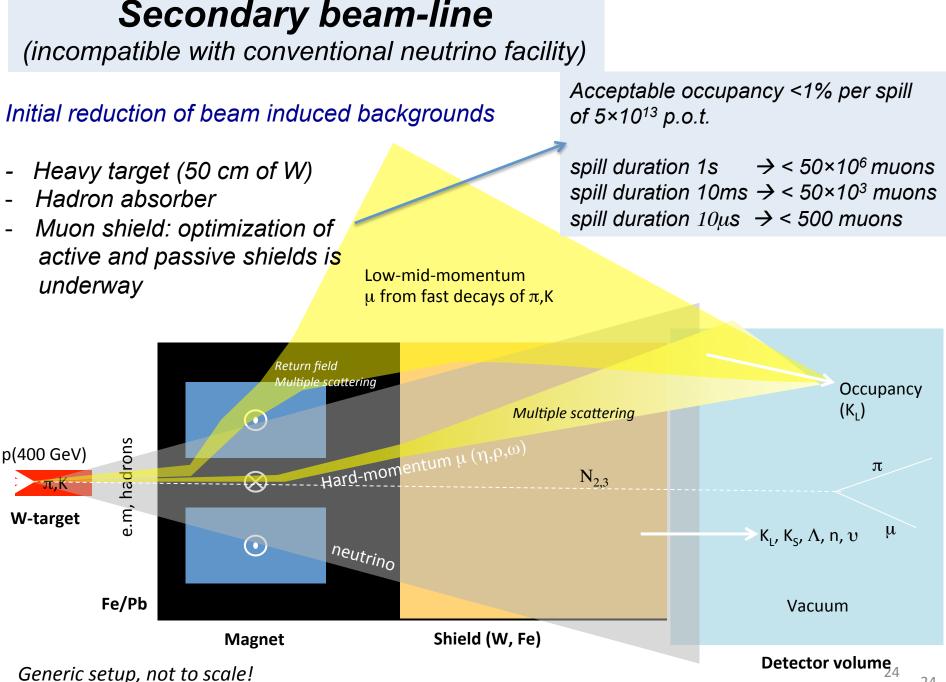
Beam dump experiment at the SPS with a total of 2×10²⁰ protons on target (pot) to produce large number of charm mesons

• HNLs produced in charm decays have significant P_{T}



Detector must be placed close to the target to maximize geometrical acceptance

Effective (and "short") muon shield is essential to reduce
 muon-induced backgrounds (mainly from short-lived resonances accompanying charm production)

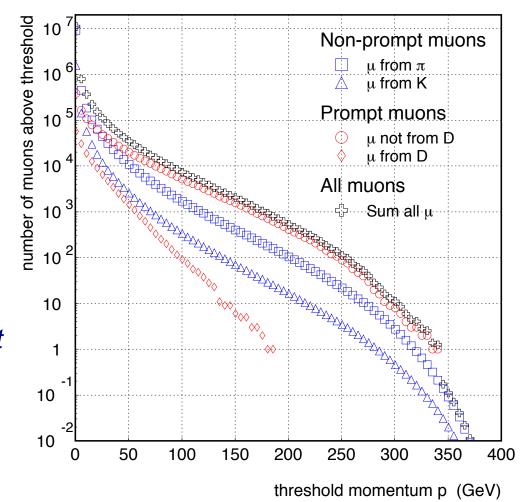


Secondary beam-line (cont.)

Muon shield

Main sources of the muon flux (estimated using PYTHIA with 10⁹ protons of 400 GeV energy)

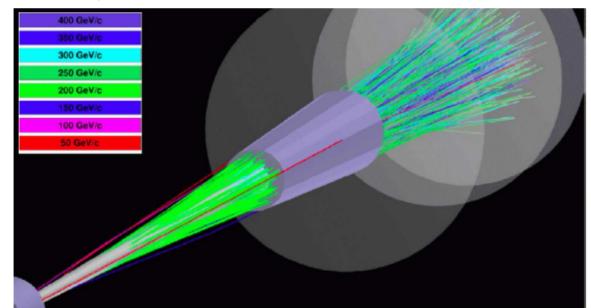
- A muon shield made of ~55 m W(U) should stop muons with energies up to 400 GeV
- Cross-checked with results from CHARM beam-dump experiment
- Detailed simulations will define the exact length and radial extent of the shield

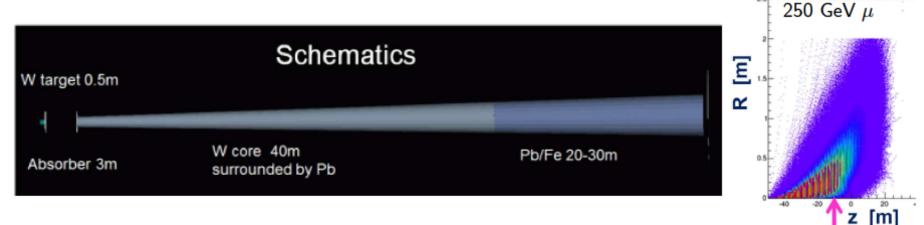


Muon shield optimization

Passive μ -filter

- Geant studies to estimate flux.
- MS and €: limit W-length to 40 m.
- High-p at small θ : WØ12-50 cm
- +20-30 m of Pb/Fe :
- reduction of 10^7 possible
- Robust/easy to operate

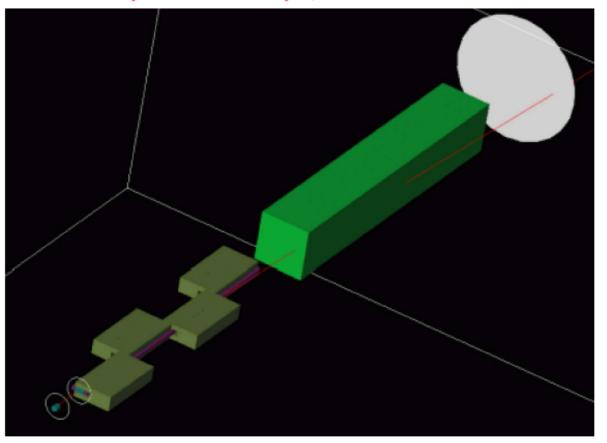




Muon shield optimization

Alternative: Active (+passive) μ -filter

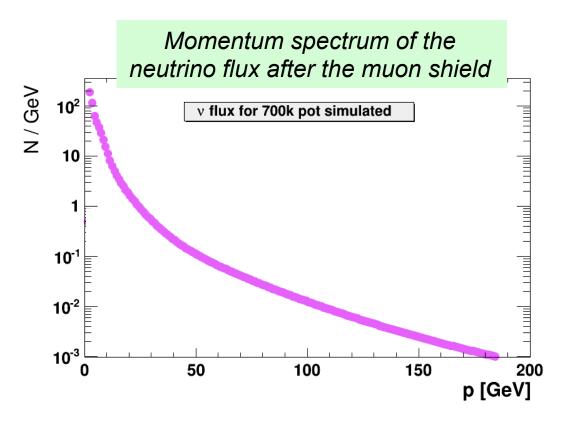
- Use 6 m long C-shaped magnets.
- Produces 40 Tm total field with 4 magnets: high-p swept out.
- Problem: return-B of low-p μ :
- alternate return-B left/right
- Add passive Fe-shield
- reduction of $10^7\ {\rm possible}$



Work in progress, need to optimize together with SPS-spill length, and induced background.

Experimental requirements (cont.)

- Minimize background from interactions of active neutrinos in the detector decay volume
 - Requires evacuation of the detector volume

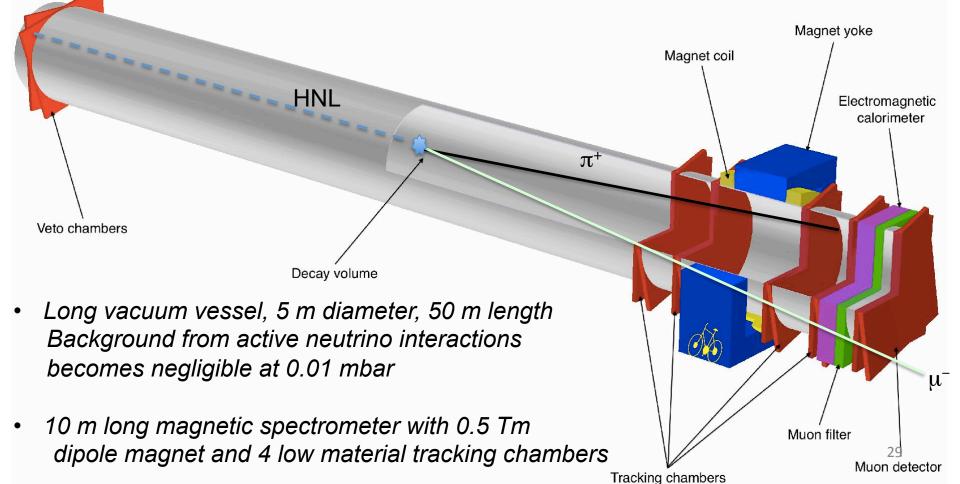


2×10⁴ neutrino interactions per 2×10²⁰ pot in the decay volume at atmospheric pressure \rightarrow becomes negligible at 0.01 mbar

Detector concept (based on existing technologies)

• Reconstruction of the HNL decays in the final states: $\mu^-\pi^+$, $\mu^-\rho^+$ & $e^-\pi^+$

Requires long decay volume, magnetic spectrometer, muon detector and electromagnetic calorimeter, preferably in surface building

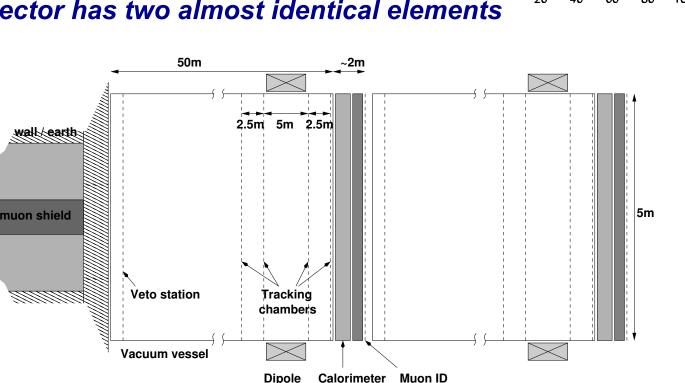


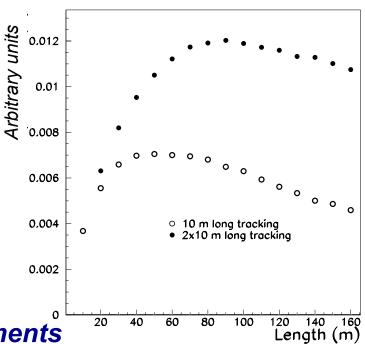
Detector concept (cont.)

Geometrical acceptance

- Saturates for a given HNL lifetime as a • function of detector length
- The use of two magnetic spectrometers • increases the acceptance by 70%

Detector has two almost identical elements

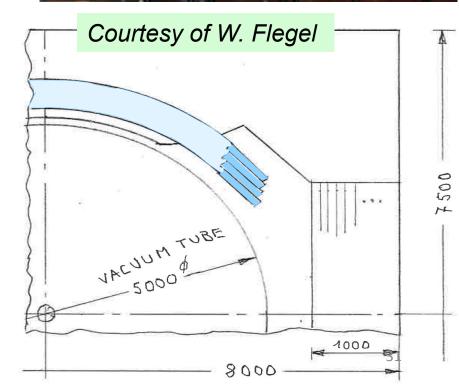




Detector apparatus based on existing technologies

- Experiment requires a dipole magnet similar to LHCb design, but with ~40% less iron and three times less dissipated power
- Free aperture of ~ 16 m² and field integral of ~ 0.5 Tm
 - Yoke outer dimension: 8.0×7.5×2.5 m³
 - Two Al-99.7 coils
 - Peak field ~ 0.2 T
 - Field integral ~ 0.5 Tm over 5 m length

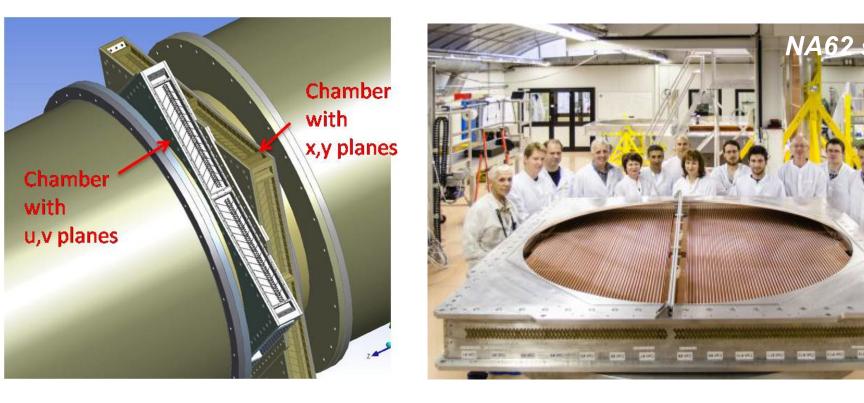




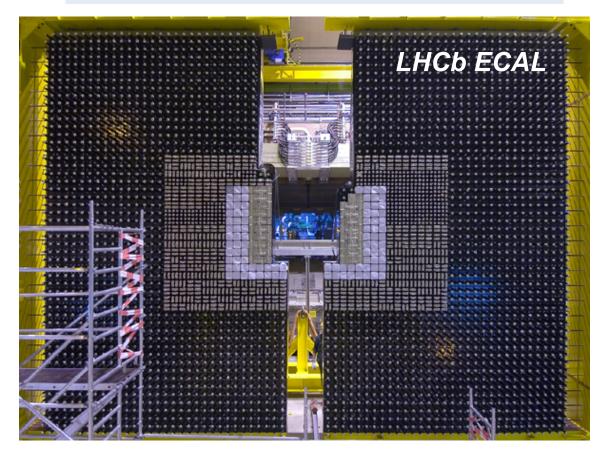
Detector apparatus (cont.) based on existing technologies

NA62 vacuum tank and straw tracker

- < 10⁻⁵ mbar pressure in NA62 tank
- Straw tubes with 120 μm spatial resolution and 0.5% X₀/X material budget Gas tightness of NA62 straw tubes demonstrated in long term tests



Detector apparatus (cont.) based on existing technologies



LHCb electromagnetic calorimeter

- Shashlik technology provides economical solution with good energy and time resolution

Residual backgrounds

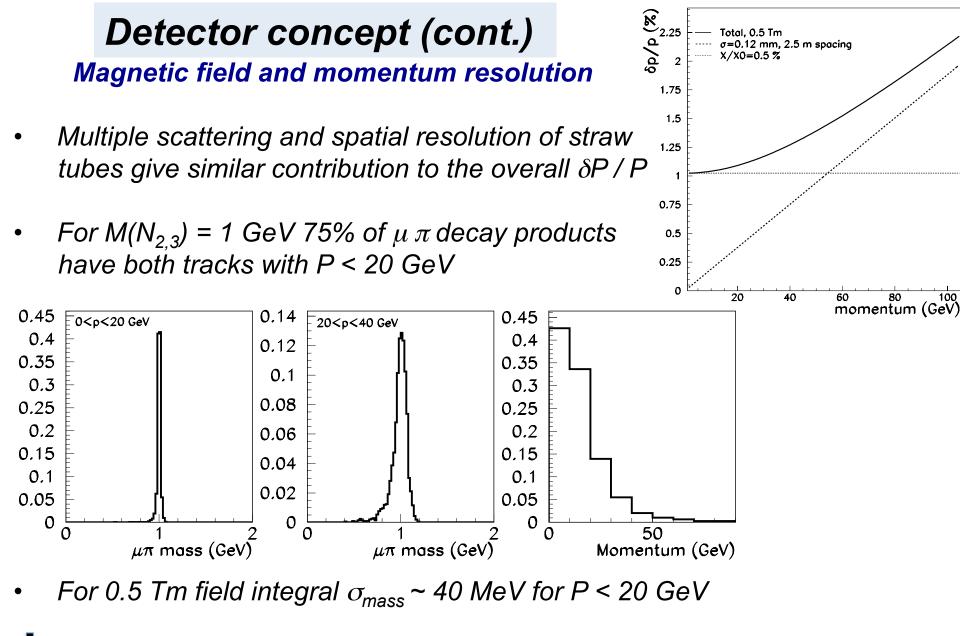
Use a combination of GEANT and GENIE to simulate the Charged Current and Neutral Current neutrino interaction in the final part of the muon shield (cross-checked with CHARM measurement)

yields CC(NC) rate of $\sim 6(2) \times 10^5$ per int. length per 2×10^{20} pot

Instrumentation of the end-part of the muon shield would allow the rate of CC + NC to be measured and neutrino interactions to be tagged

- ~10% of neutrino interactions in the muon shield just upstream of the decay volume produce Λ or K⁰ (as follows from GEANT+GENIE and NOMAD measurement)
- *Majority of decays occur in the first 5 m of the decay volume*
- Requiring μ -id. for one of the two decay products

 \rightarrow 150 two-prong vertices in 2×10²⁰ pot



Ample discrimination between high mass tail from small number of residual $K_{I} \rightarrow \pi^{+}\mu^{-}v$ and 1 GeV HNL

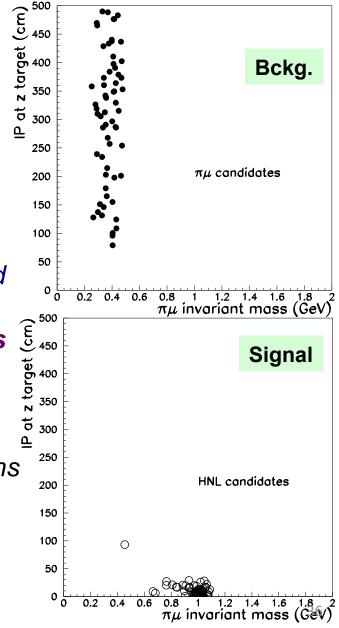
Detector concept (cont.)

Impact Parameter resolution

K_L produced in the final part of the muon shield have very different pointing to the target compared to the signal events

> Use Impact Parameter (IP) to further suppress K_L background

- IP < 1 m is 100% eff. for signal and leaves only a handful of background events
- The IP cut will also be used to reject backgrounds induced in neutrino interactions in the material surrounding the detector



Expected event yield

- Integral mixing angle U^2 is given by $U^2 = U_e^2 + U_{\mu}^2 + U_{\tau}^2$
- A conservative estimate of the sensitivity is obtained by considering only the decay $N_{2,3} \rightarrow \mu^- \pi^+$ with production mechanism $D \rightarrow \mu^+ NX$, which probes U_{μ}^{2}
- $U^2 \longleftrightarrow U_{\mu}^2$ depends on flavour mixing
- Expected number of signal events:

 $N_{signal} = n_{pot} \times 2\chi_{cc} \times BR(U_{\mu}^{2}) \times \varepsilon_{det}(U_{\mu}^{2})$

$$n_{pot} = 2 \times 10^{20}$$

 $\chi_{cc} = 0.45 \times 10^{-3}$

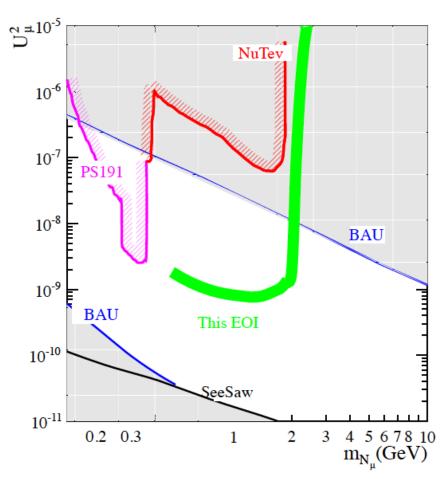
 $BR(U_{\mu}^{2}) = BR(D \rightarrow N_{2,3}X) \times BR(N_{2,3} \rightarrow \mu\pi)$ BR(N_{2,3} $\rightarrow \mu^{-}\pi^{+}$) is assumed to be 20%

 ε_{det} (U_{μ}^{2}) is the probability of the $N_{2,3}$ to decay in the fiducial volume and μ , π are reconstructed in the spectrometer

Expected event yield (cont.)

Assuming $U_{\mu}^{2} = 10^{-7}$ (corresponding to the strongest experimental limit currently for $M_{N} \sim 1$ GeV) and $\tau_{N} = 1.8 \times 10^{-5}$ s

~12k fully reconstructed N $\rightarrow \mu^{-}\pi^{+}$ events are expected for M_{N} = 1 GeV



120 events for cosmologically favoured region: $U_{\mu}^{2} = 10^{-8} \& \tau_{N} = 1.8 \times 10^{-4} s_{38}$

Expected event yield (cont.)

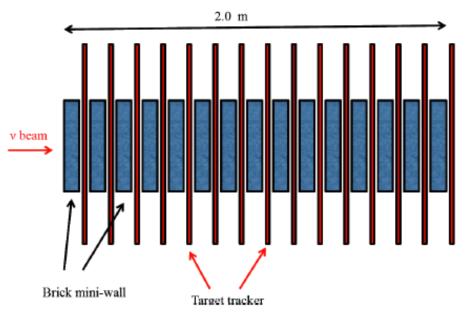
- ECAL will allow the reconstruction of decay modes with π^0 such as $N \rightarrow \mu^- \rho^+$ with $\rho^+ \rightarrow \pi^+ \pi^0$, doubling the signal yield
- Study of decay channels with electrons such as $N \rightarrow e\pi$ would further increase the signal yield and constrain U_e^2

In summary, for $M_N < 2$ GeV the proposed experiment has discovery potential for the cosmologically favoured region with $10^{-7} < U_{\mu}^{2} < a$ few × 10⁻⁹

SM physics

 ν_{τ} Physics with 2×10^{20} pot

- Scaling from DONUT: 20 times more CC with same ν -target mass.
- But can increase u-target mass "easily", lets say to 3~% of OPERA emulsion surface:



- Only requires limited space along beam-line, hence "no" loss for HNL acceptance.
- HNL spectrometer is forward spectrometer of ν -physics program.
- ν -target allows to tag K_L which coincide with ν -interactions.
- Expect 1500-2000 CC ν_{τ} interactions.
- In addition: $5 \times \nu_{\mu}$ CC charm production than CHORUS (2k).

Status of the SPSC review

- Oct 2013: submitted our EOI: CERN-SPSC-2013-024 ; arXiv:1310.1762 ; SPSC-EOI-010. 2013
- SPSC assigned 4 referees, who came with a list of questions.
- 3/1/2014: answers to questions: snoopy.web.cern.ch/snoopy/EOI/SPSC-EOI-010_ResponseToReferees.pdf
- 15/1/2014: SPSC discussed our proposal.

17/1/2014: The official feedback from the Committee is as follows :

"The Committee received with interest the response of the proponents to the questions raised in its review of EOI010.

The SPSC **recognises** the interesting physics potential of searching for heavy neutral leptons and investigating the properties of neutrinos.

Considering the large cost and complexity of the required beam infrastructure as well as the significant associated beam intensity, such a project should be designed as a general purpose beam dump facility with the broadest possible physics programme, including maximum reach in the investigation of the hidden sector.

To further review the project the Committee **would need** an extended proposal with further developed physics goals, a more detailed technical design and a stronger collaboration."

Cheers,

Gavin, Lau, Matthew and Thierry

(for the SPS Committee).

Conclusion and Next steps

- The proposed experiment will search for NP in the largely unexplored domain of new, very weakly interacting particles with masses below the Fermi scale
- Detector is based on existing technologies
 Ongoing discussions of the beam lines with experts
- The impact of HNL discovery on particle physics is difficult to overestimate !
- The proposed experiment perfectly complements the searches for NP at the LHC and in neutrino physics

A collaboration is currently being setup with aim for the first collaboration meeting in June. Let us know if you are interested to join !

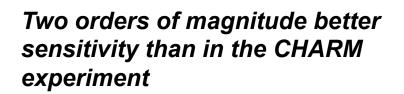
BACK - UP

Other BSM physics

Search for light, very weakly interacting, yet unstable New Particles

Massive paraphotons, p (in secluded Dark Matter models), e.g. $\Sigma \rightarrow pV$ with $V \rightarrow \mu\mu$

M. Pospelov, A. Ritz, M.B. Voloshin (2008)



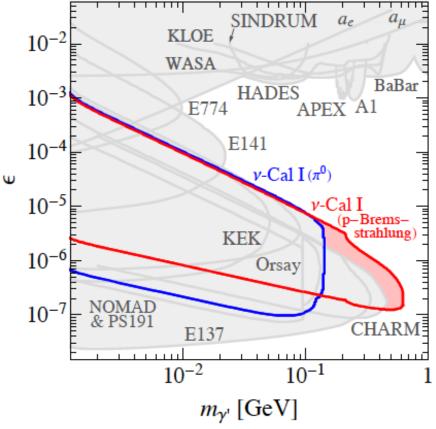


Figure 1: Present direct limits on the model parameter space $(\epsilon, m_{\gamma'})$, for details and original references see [36].

Other BSM physics

Search for light, very weakly interacting, yet unstable New Particles

Light s-goldstinos (super-partners of SUSY goldstinos), e.g. $D \rightarrow \pi X$ with $X \rightarrow \mu \mu$

D.S. Gorbunov (2001)

$$N_{\pi^+\pi^-} \simeq 2 \times \left(\frac{1000 \text{ TeV}}{\sqrt{F}}\right)^8 \left(\frac{M_{\lambda_g}}{3 \text{ TeV}}\right)^4 \left(\frac{m_X}{1 \text{ GeV}}\right)^2$$

R-parity violating neutralinos in SUSY goldstinos, e.g. $D \rightarrow \mu \bar{\chi}_0$ with $\bar{\chi}_0 \rightarrow \mu^+ \mu^- \nu$

A. Dedes, H.K. Dreiner, P. Richardson (2001)

$$N \simeq 20 \times \left(\frac{m_{\chi_0}}{1 \text{ GeV}}\right)^6 \left(\frac{\lambda}{10^{-8}}\right)^2 \left(\frac{\text{Br}\left(D \to \chi_0 + \ldots\right)}{10^{-10}}\right)$$

The mass of the Higgs boson is very close to the stability bound of the Higgs mass *

$$M_{crit} = [129.3 + \frac{y_t(M_t) - 0.9361}{0.0058} \times 2.0 - \frac{\alpha_s(M_Z) - 0.1184}{0.0007} \times 0.5] \text{ GeV}$$

 $y_t(M_t)$ - top Yukawa in $\overline{\mathrm{MS}}$ scheme

Matching at EW scale	Central value	theor. error
Bezrukov et al, $\mathcal{O}(lpha lpha_s)$	129.4 GeV	1.0 GeV
Degrassi et al, $\mathcal{O}(lpha lpha_s, y_t^2 lpha_s, \lambda^2, \lambda lpha_s)$	129.6 GeV	0.7 GeV
Buttazzo et al, complete 2-loop	129.3 GeV	0.07 GeV

Chetyrkin et al, Mihaila et al, Bednyakov et al, 3 loop running to high energies

