

# High resolution timing with Silicon detectors

Is it possible to build a tracker with concurrent excellent time and position resolution?

Can we provide in one detector, or in combination

**Timing resolution ~ 10 ps**  
**Space resolution ~ 10's of  $\mu\text{m}$**

- Tracking in 4 Dimensions
- HL-LHC conditions
- CMS Barrel and Endcap detectors
- Review of Ultra-fast silicon detectors
- How to build a large detector



# The effect of timing information

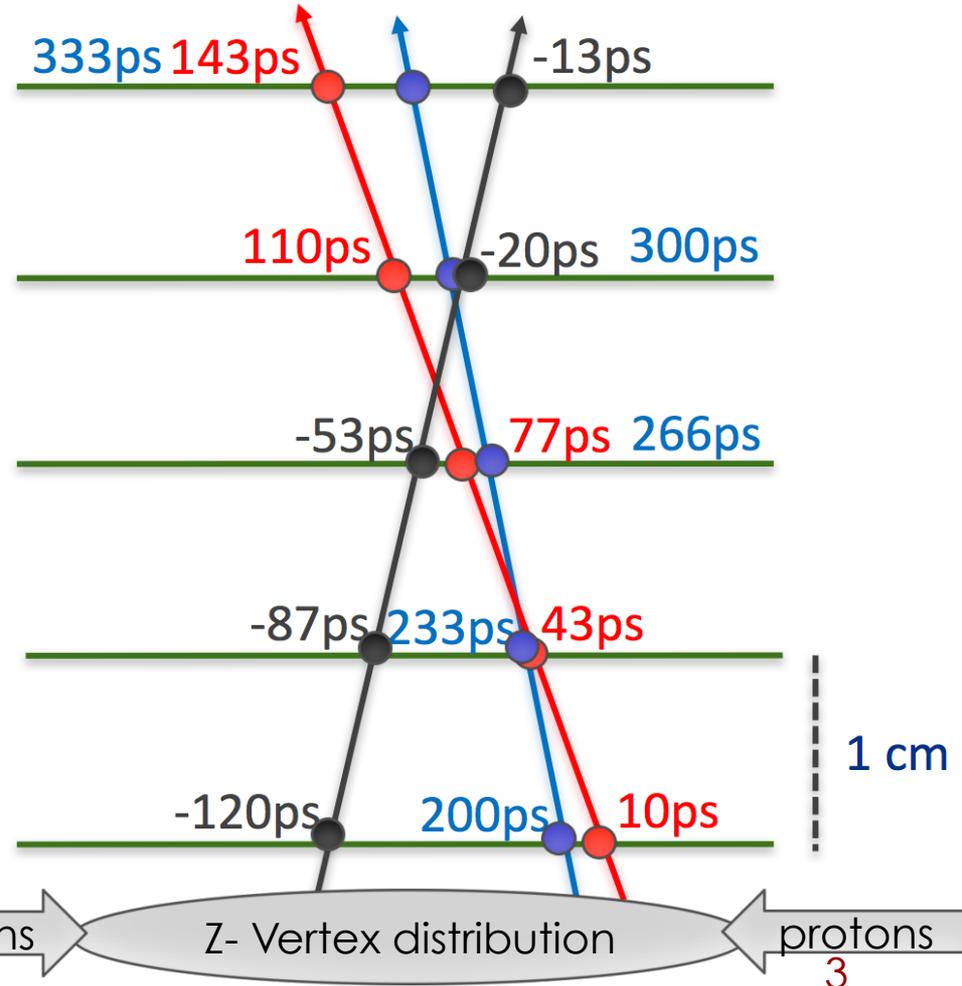
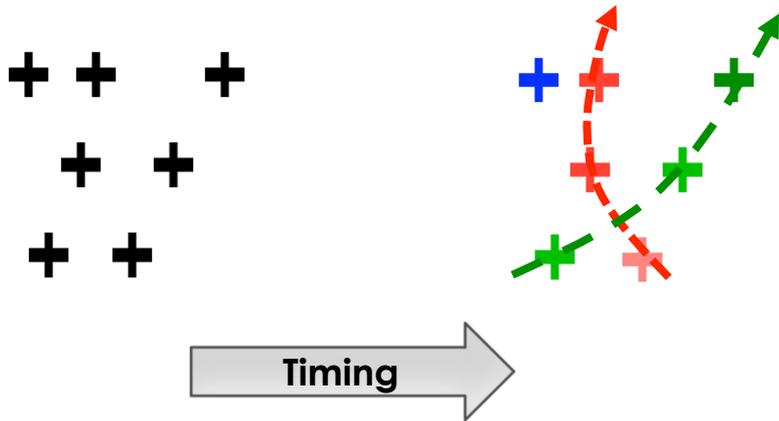
**The inclusion of track-timing in the event information has the capability of changing radically how we design experiments.**

**Timing can be available at different levels of the event reconstruction.**

- 1) Timing at each point along the track
- 2) Timing in the event reconstruction
- 3) Timing at the trigger level

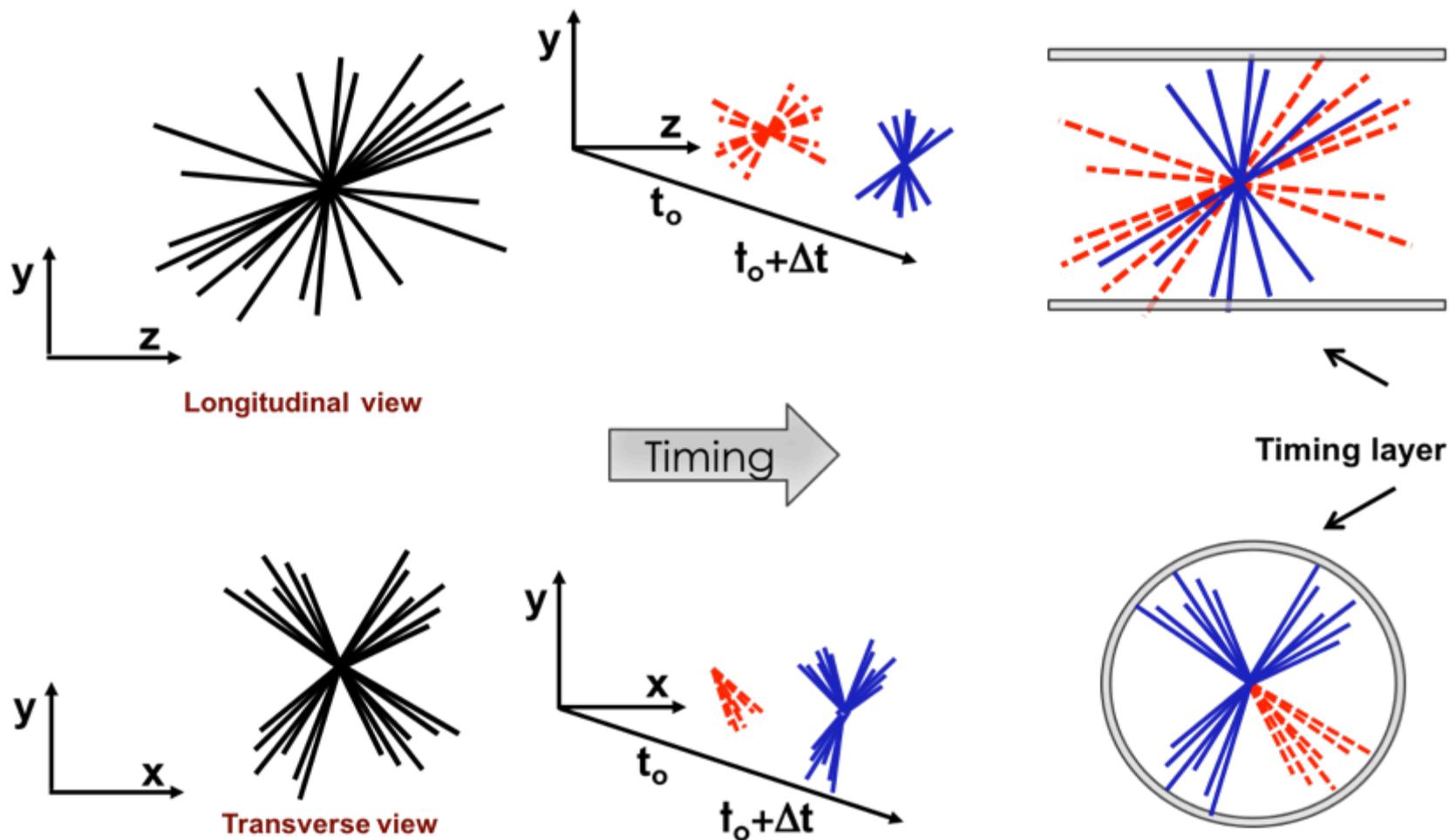
# Timing at each point along the track

- Massive simplification of pattern recognition, new tracking algorithms will be faster even in very dense environments
- Use only “time compatible points”



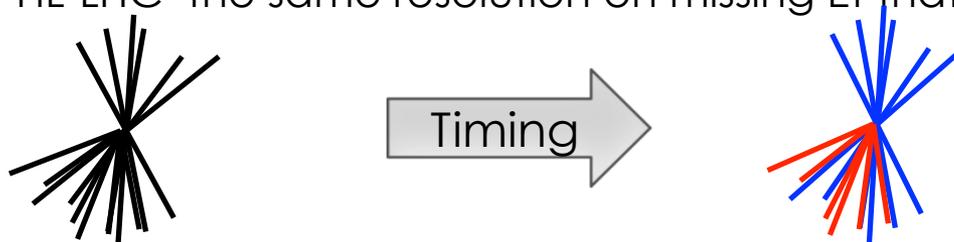
# Timing in the event reconstruction - I

Timing allows distinguishing overlapping events by means of an extra dimension.

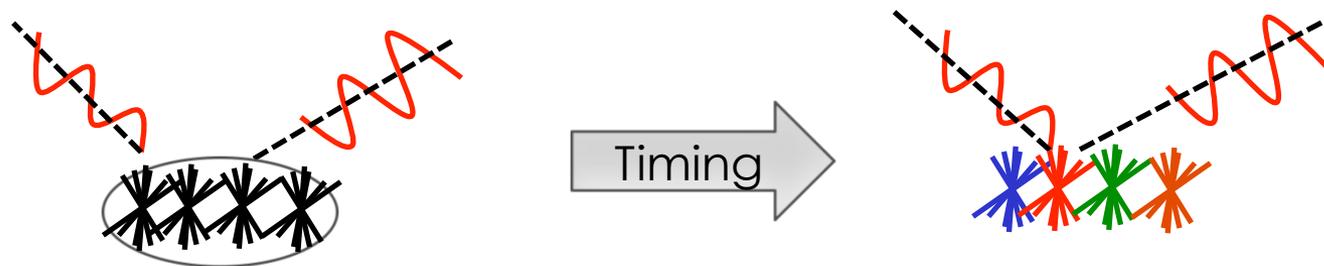


# Timing in the event reconstruction - II

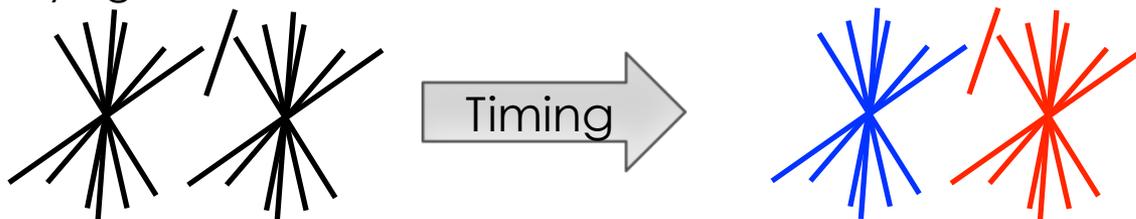
**Missing Et:** consider overlapping vertexes, one with missing Et: Timing allows obtaining at HL-LHC the same resolution on missing Et that we have now



$H \rightarrow \gamma\gamma$ : The timing of the  $\gamma\gamma$  allows to select an area (1 cm) where the vertex is located. The vertex timing allows to select the correct vertex within this area

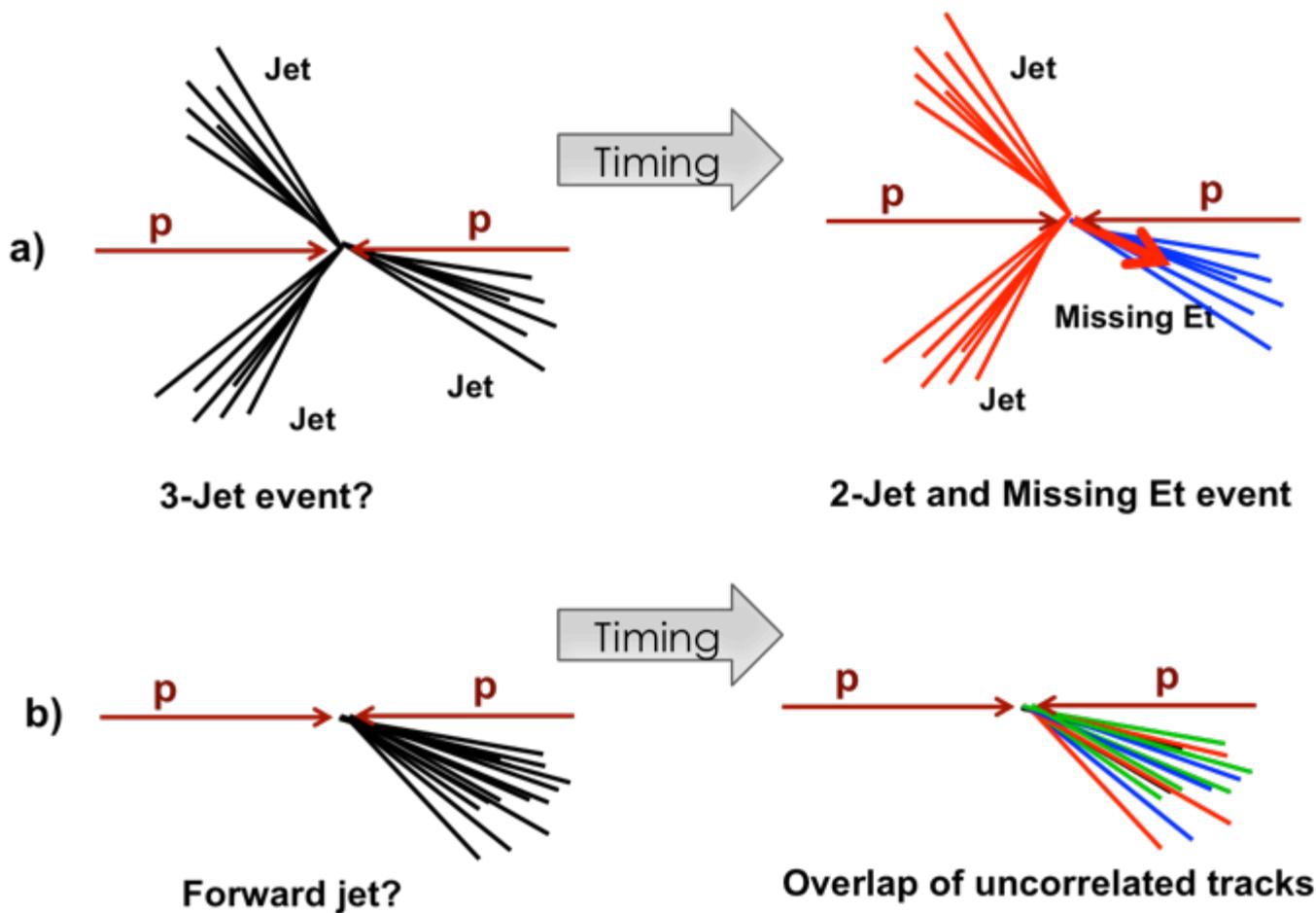


**Displaced vertexes:** The timing of the displaced track and that of each vertex allow identifying the correct vertex



# The effect of timing information:

**Timing at the trigger decision:** it allows reducing the trigger rate, rejecting topologies that look similar, but they are actually different.



# Why do we want a timing layer?

HL-LHC limit:  $2 \times 10^{35} \text{ cm}^{-2}\text{s}^{-1}$  at the beginning of each fill.  
Limited by luminosity leveling at  $5.2$  or  $7.2 \times 10^{34} \text{ cm}^{-2}\text{s}^{-1}$ .

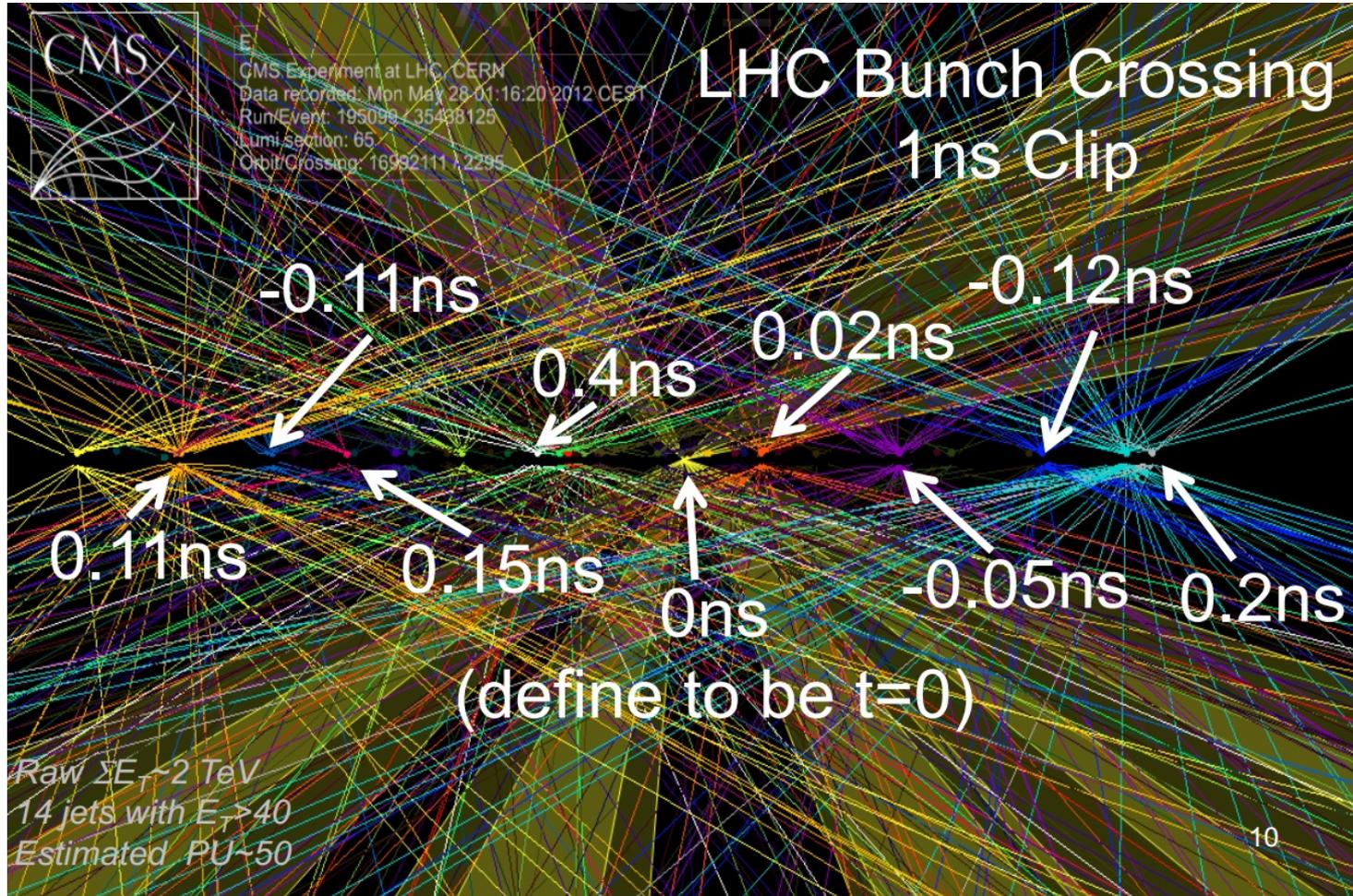
→ Possibly LHC will be able to deliver luminosity in excess of what the experiments can take: we need to be able to take data efficiently at very high instantaneous luminosity

The purpose of a timing upgrade of the CMS detector is to consolidate the particle-flow performance at a multiplicity of 140 pileup events and to extend it up to 200 pileup events, exploiting the additional information provided by the precision timing of both tracks and energy deposits in the calorimeters.

# Vertexes in space and time

Current situation, pile-up  $\sim 50$ :

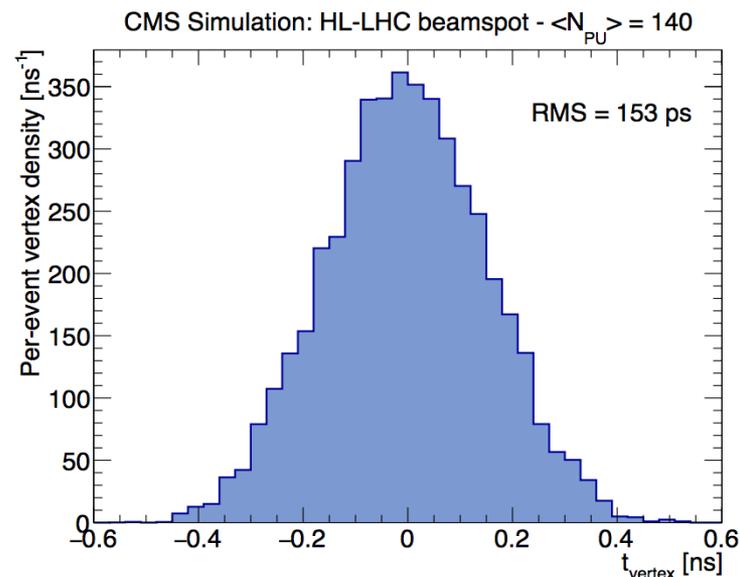
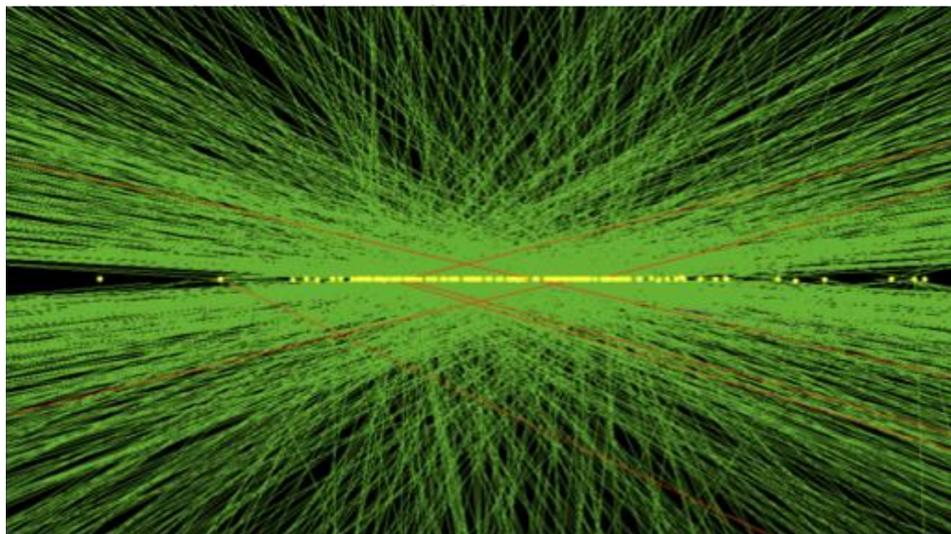
Vertexes do not overlap in space  $\rightarrow$  tracking resolves the vertexes



# Vertexes in space and time

HL-LHC situation, pile-up  $\sim 150 - 200$ :

Vertexes **overlap** in space  $\rightarrow$  tracking **does not resolves** all vertexes



There are between 15-20% of tracking vertexes (longitudinal resolution  $\sim 200$  micron) that are actually composed by 2 or more interactions

$\rightarrow$  Loss of events  $\rightarrow$  loss of luminosity

# Pileup and event density

**Pile-up:** number of concurrent scattering processes (140 – 200).



**Density of events:** number of events 1 mm (0.2 – 2 event/mm)



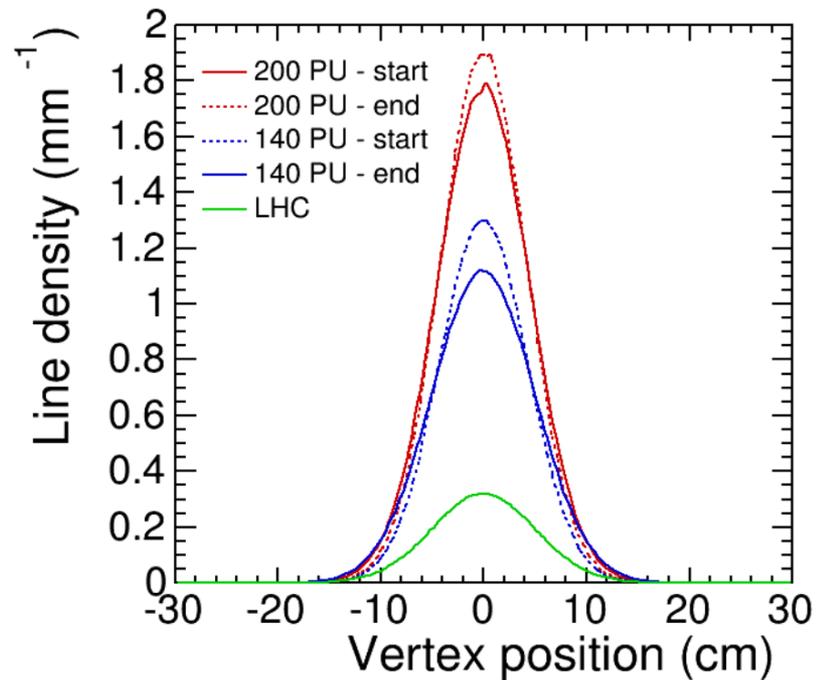
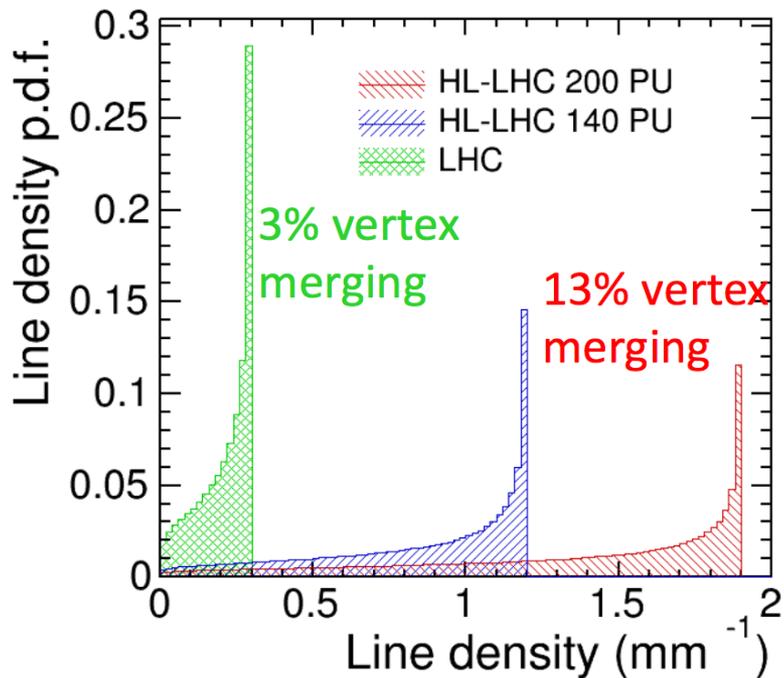
Why are they different?

**Pile-up** is a global quantity, and it can be fought with very high granularity. It influences, for example, the total amount of tracks and neutral clusters

**Density of events:** it can be fought with longitudinal resolution and timing.

→ Charge particles

# Vertex merging

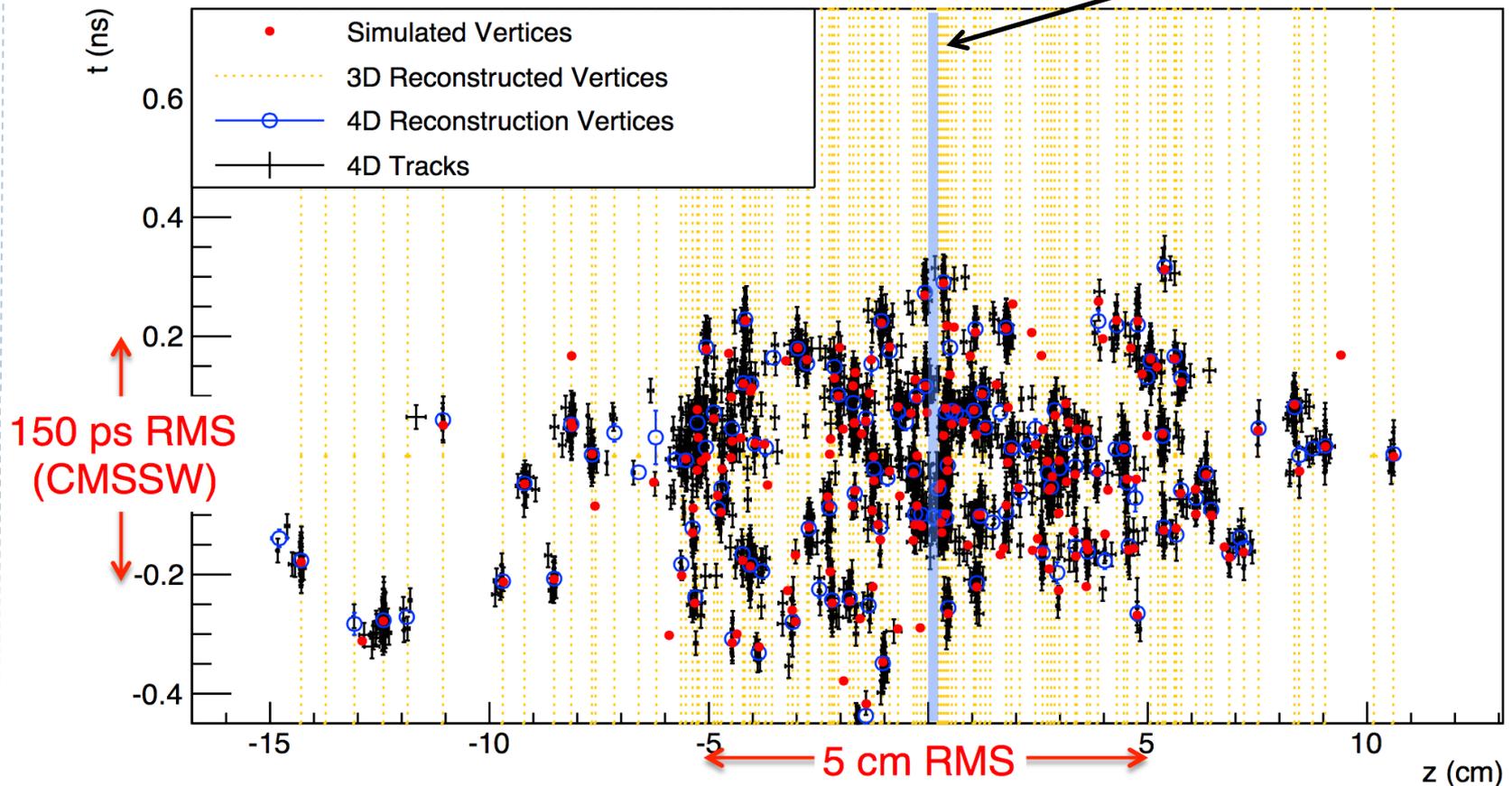


**Vertexes are clustered in high density region, producing large overlaps**

# Position – time of each vertex

## ▶ 200 pileup collisions

Longitudinal resolution  $\sim 300$  micron

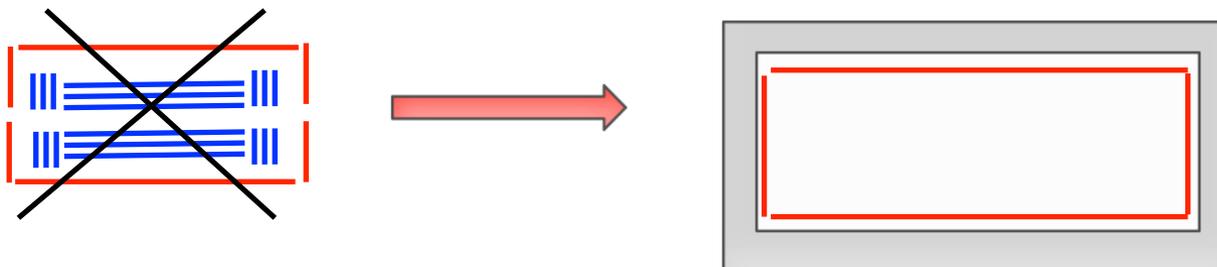


## ▶ HL-LHC baseline (as of ECFA): $t_{\text{RMS}} = 180$ ps, $z_{\text{RMS}} = 4.8$ cm

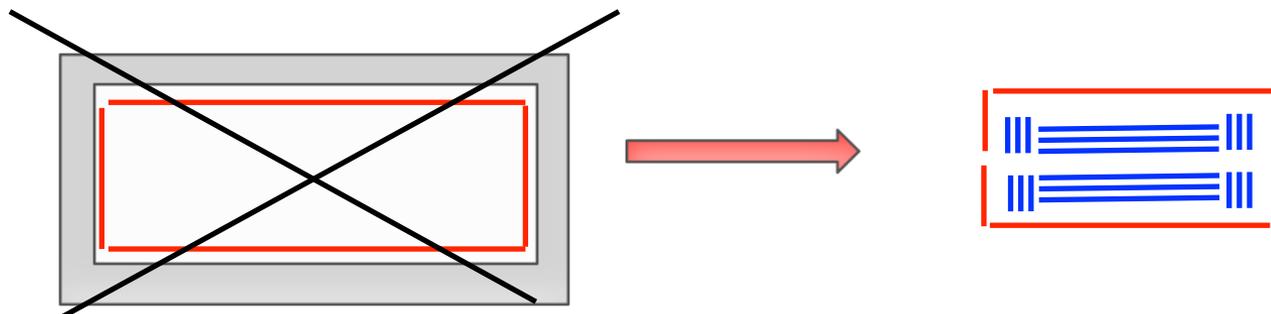
At pileup levels of 140–200, a fraction as high as 15-20% of independent vertices merges, in the absence of time information.

# Where do we stand?

**The tracking community** thinks it is a wonderful idea, clearly to be implemented **outside the tracker volume**, in front of the calorimeter



**The calorimeter community** thinks it is a wonderful idea, clearly to be implemented **far from the calorimeter**, in the tracker volume

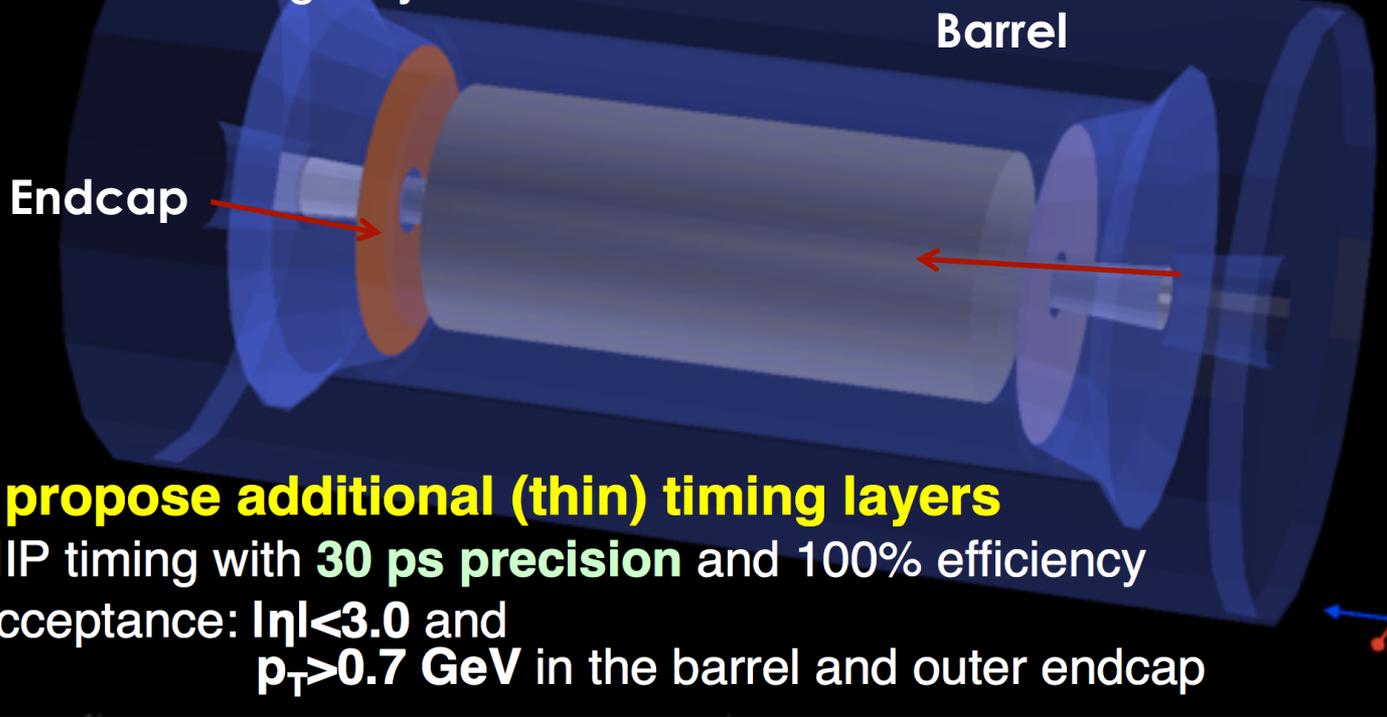


We are now in contact with **the muon community....**

# A timing layer for charge particles in CMS

## Calorimeter upgrades:

- ▶ Provide precision timing ( $\sim 30$  ps) on high energy photons in ECAL, on photons and high energy hadrons in HGCal
- ▶ Precision timing only for showers



## We propose additional (thin) timing layers

- ▶ MIP timing with **30 ps precision** and 100% efficiency
- ▶ Acceptance:  $|\eta| < 3.0$  and  $p_T > 0.7$  GeV in the barrel and outer endcap

### ECAL Crystal Barrel:

Current resolution: 150 ps for  $E > 30$  GeV

With new electronics: 30 ps for  $E > 30$  GeV

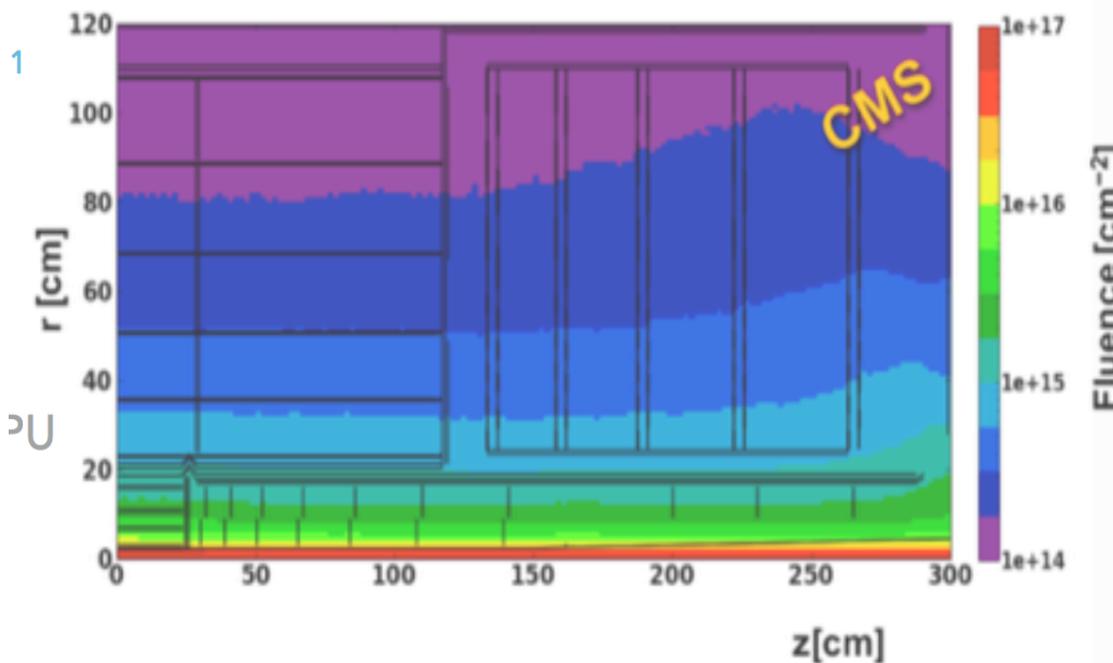
### HGCAL:

50 ps time resolution from each silicon pad for  $\#MIPS > 30$

- ➔ Each showers covers many planes ( $> 5$ )
- ➔ Shower resolution limited by systematics

# Radiation levels

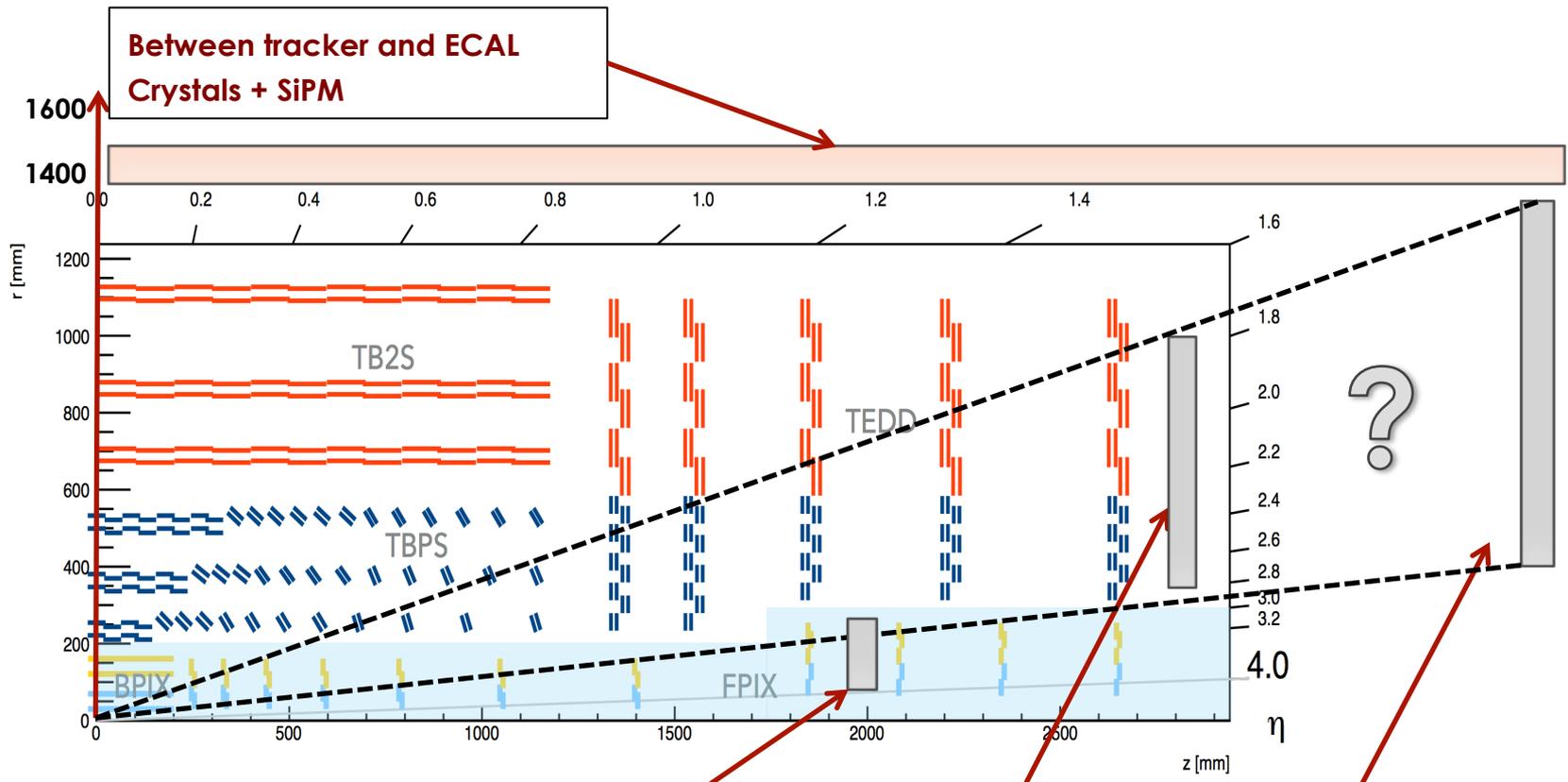
Region	$\eta$	R (cm)	z (cm)	Fluence ( $\text{cm}^{-2}$ )	Charged hardons ( $\text{cm}^{-2}$ )	Dose (kGy)
barrel	0.0	130	0	$1.5 \times 10^{14}$	$6.7 \times 10^{12}$	11.9
barrel	1.0	130	153	$1.9 \times 10^{14}$	$6.9 \times 10^{12}$	15.9
transition	1.5	130	300	$1.7 \times 10^{14}$	$8.9 \times 10^{12}$	15.0
endcap	2.0	82.7	300	$4.3 \times 10^{14}$	$5.7 \times 10^{13}$	79.4
endcap	2.5	49.6	300	$1.1 \times 10^{15}$	$2.3 \times 10^{14}$	245.
endcap	2.7	40.5	300	$1.2 \times 10^{15}$	$3.3 \times 10^{14}$	337.
endcap	3.0	29.9	300	$1.8 \times 10^{15}$	$5.8 \times 10^{14}$	566.



**Barrel is a much easier environment**

**Endcap will be driving the radiation aspects of the project**

# CMS Timing layer position



Between tracker and ECAL  
Crystals + SiPM

Silicon detector

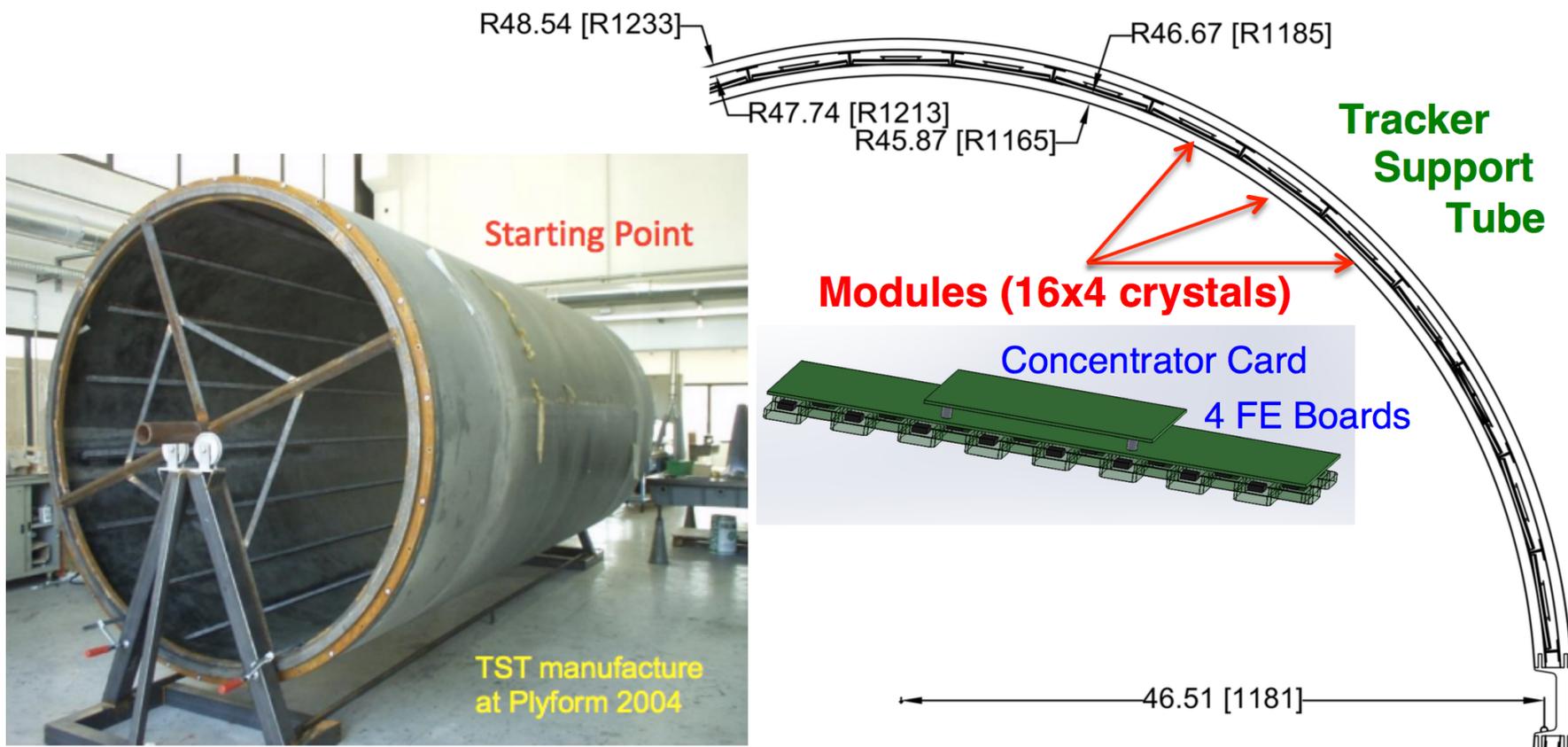
**High rapidity Option:**  
Inside Forward pixel  
3.5 m<sup>2</sup>  
0.7 ml Channel

**Option 1:**  
Inside Tracker  
3.5 m<sup>2</sup>  
0.7 ml Channel

**Option 2:**  
In front HGCAL  
7.2 m<sup>2</sup>  
1.2 ml Channel

# Barrel Timing layer: on the Tracker Support Tube

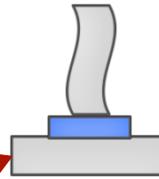
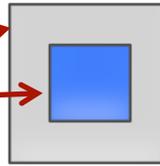
- ▶ **LYSO crystals + SiPM embedded in the TST**
  - ▶ Be ready before TK installation: 2022
  - Select production ready sensors and electronics



# Barrel design: crystals read-out by SiPM

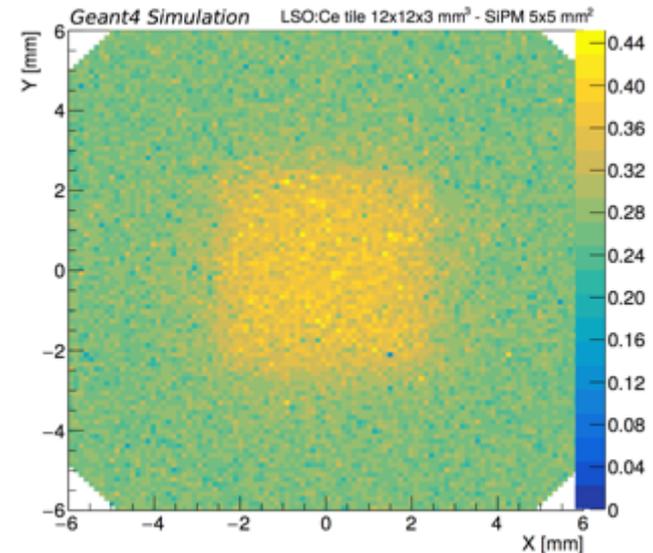
Crystal:  $1.2 \times 1.2 \text{ cm}^2$

SiPM =  $0.5 \times 0.5 \text{ cm}^2$



Crystal:  $0.3 \text{ cm}^2$  thick

Single photon efficiency  
as a function of position



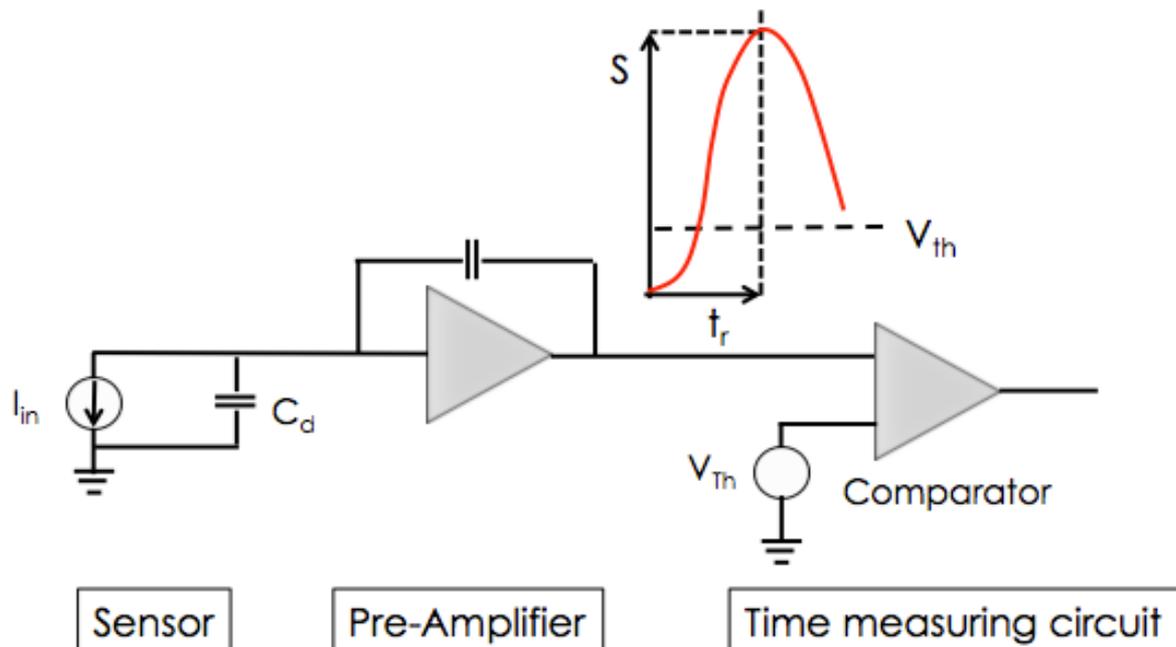
- Based on LYSO:Ce crystals read-out with silicon photomultipliers (SiPMs)
- $1.2 \times 1.2 \text{ cm}^2$  crystal, 3 mm thick, read-out by a  $0.5 \times 0.5 \text{ cm}^2$  SiPM
- time resolution in the order of 10-20 ps.

➔ A MIP traversing 3mm of LYSO produces 90,000 photons and with a 5mm x 5mm SiPM per  $\text{cm}^2$  tile yields a S/N of above 100 throughout the entire HL-LHC program.

- Use the design of the present Tracker Support Tube (TST) and rails to instrument the region outer tracker ring and the ECAL front-end cooling plates with a thin standalone detector
- Both the crystals and the SiPM are proven to be radiation tolerant up to neutron equivalent fluence  $3 \times 10^{14} \text{ cm}^{-2}$ , when cooled to  $-30^\circ\text{C}$ .
- We do not foresee to use time information in the level-1 trigger decision.

# Silicon time-tagging detector

(a simplified view)



**Time is set when the signal crosses the comparator threshold**

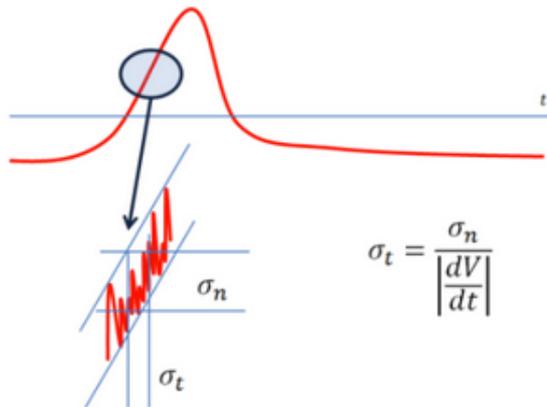
The timing capabilities are determined by the characteristics of the signal at the output of the pre-Amplifier and by the TDC binning.

**Strong interplay between sensor and electronics**

# Time resolution

$$\sigma_t = \left( \frac{N}{dV/dt} \right)^2 + (\text{Landau Shape})^2 + \text{TDC}$$

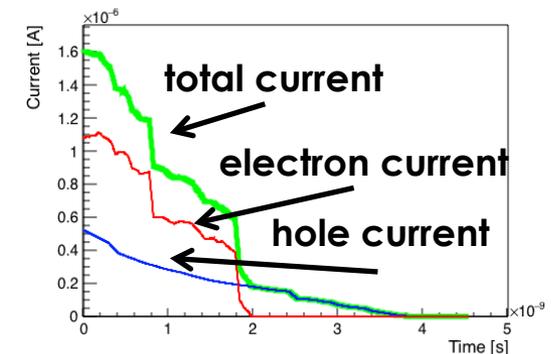
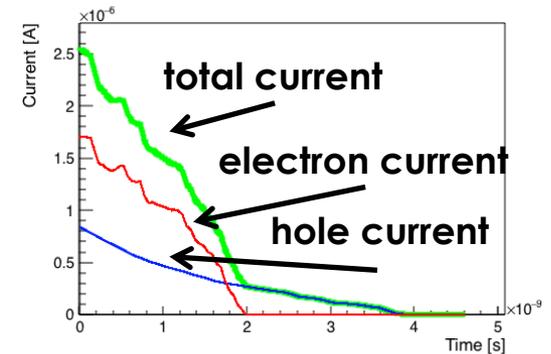
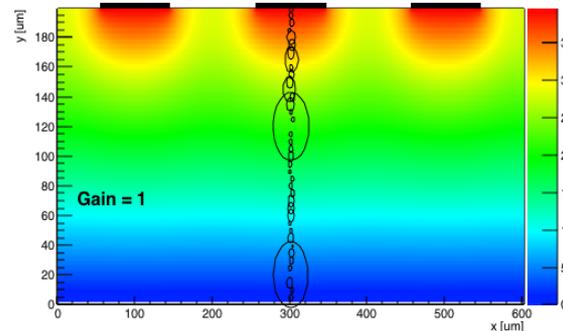
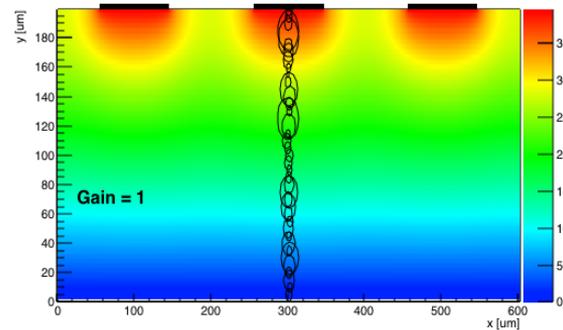
Usual "Jitter" term  
 Here enters everything that  
 is "Noise" and the  
 steepness of the signal



$$\sigma_t = \frac{\sigma_n}{\left| \frac{dV}{dt} \right|}$$

**Time walk:** Amplitude variation, corrected in electronics

**Shape variations:** non homogeneous energy



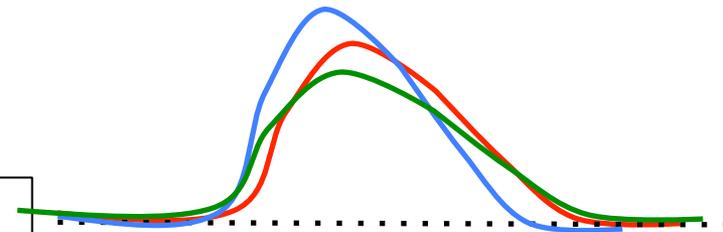
# Not all geometries are possible

Signal shape is determined by Ramo's Theorem:

$$i \propto qvE_w$$

Drift velocity

Weighting field

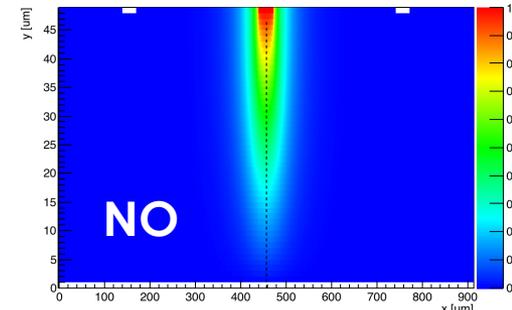
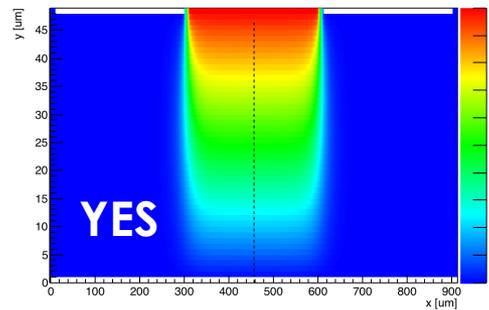


The key to good timing is the uniformity of signals:

**Drift velocity** and **Weighting field** need to be **as uniform as possible**

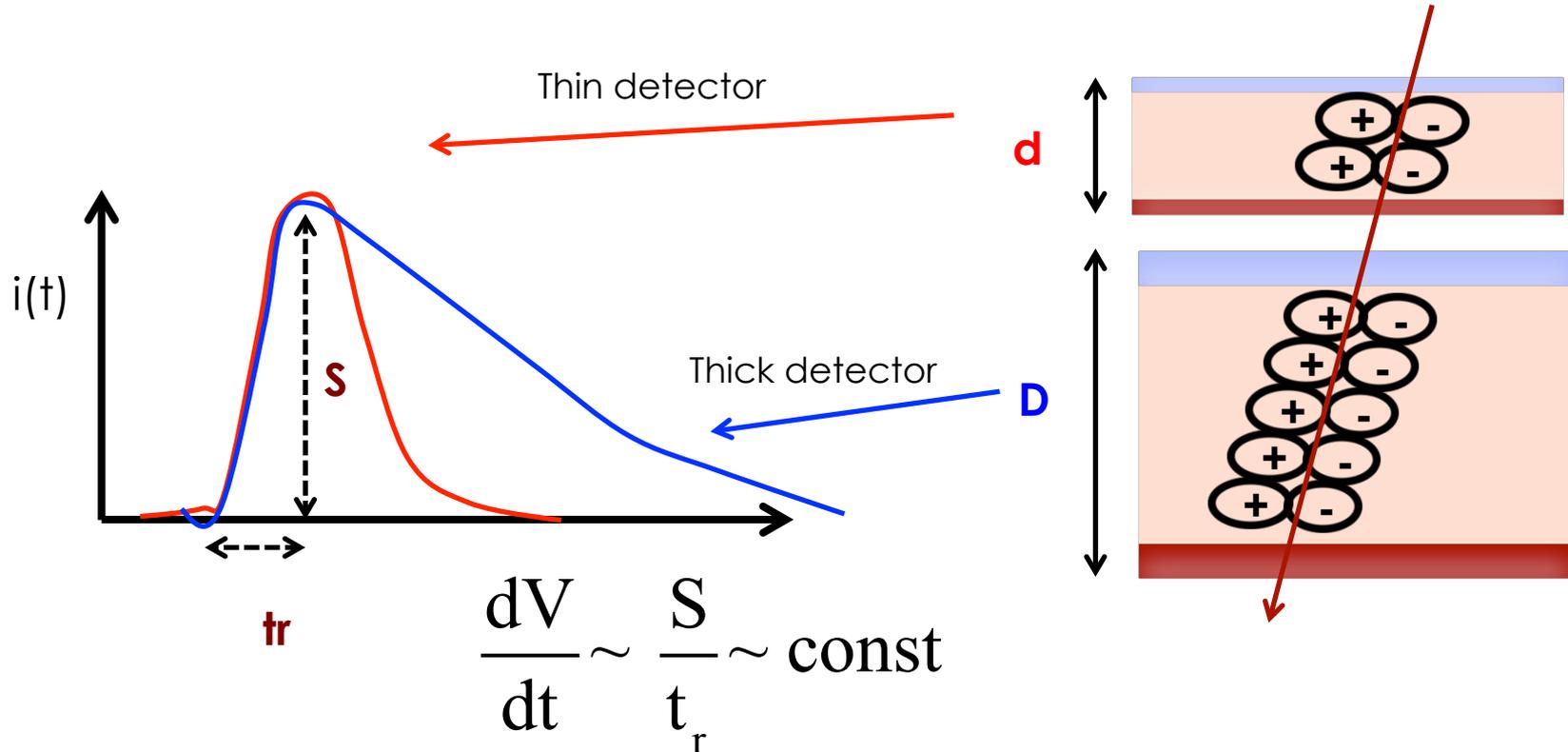
**Basic rule: parallel plate geometry: strip implant ~ strip pitch >> thickness**

Everything else does not work



# dV/dt in Thin and Thick silicon detectors

(Simplified model for pad detectors)



Thick detectors have longer signals, not higher signals

Best result : NA62, 150 ps on a 300 x 300 micron pixels

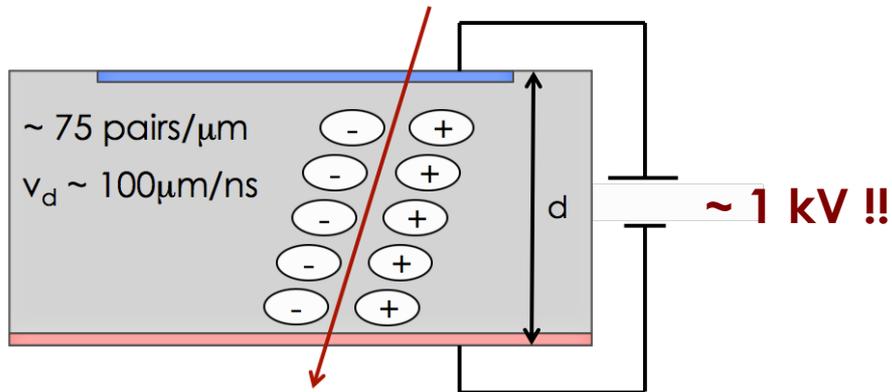
**How can we do better?**

**➔ Add internal gain**

# Gain need $E \sim 300\text{kV/cm}$ . How can we do it?

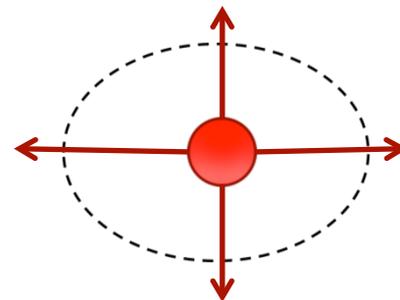
1) Use external bias: assuming a 50 micron silicon detector, we need  $V_{\text{bias}} = \sim 1\text{ kV}$

**Possible, but really difficult**



2) Use Gauss Theorem:

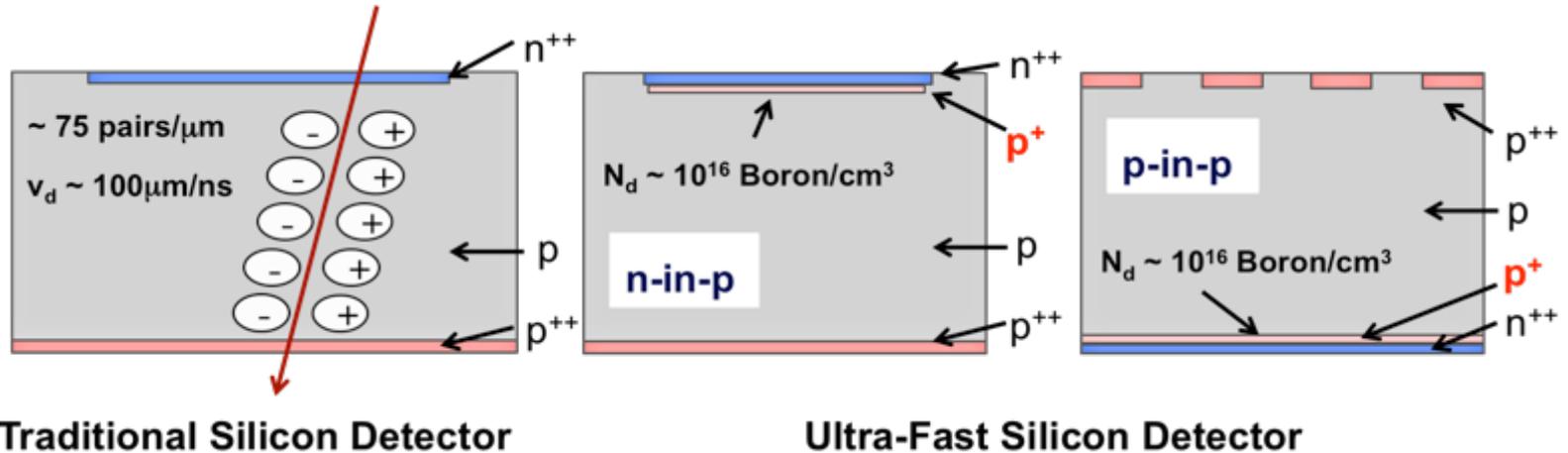
$$\sum q = 2\pi r * E$$



$$E = 300\text{ kV/cm} \rightarrow q \sim 10^{16} / \text{cm}^3$$

**Need to have  $10^{16}/\text{cm}^3$  charges !!**

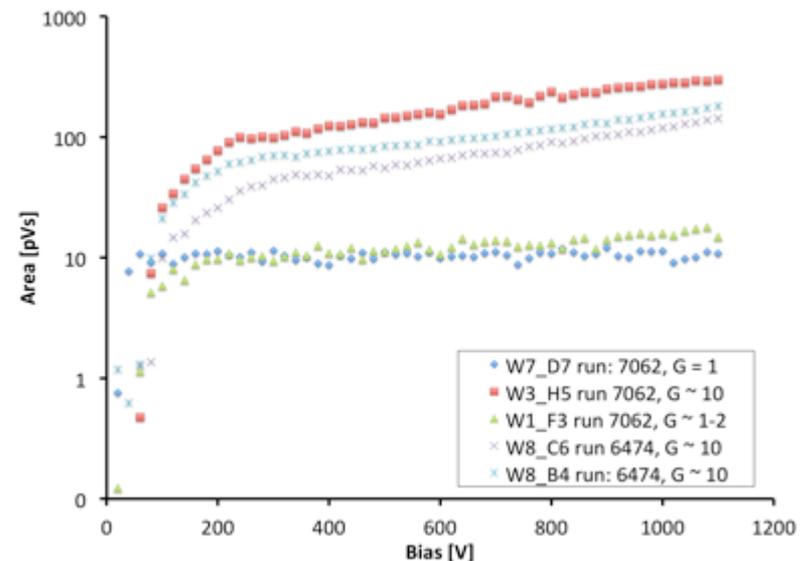
# LGAD - Ultra-Fast Silicon Detector



Adding a highly doped, thin layer of **p-implant** near the p-n junction creates a high electric field that accelerates the electrons enough to start multiplication. Same principle of APD, but with much lower gain.

**Gain changes very smoothly with bias voltage.**

**Easy to set the value of gain requested.**



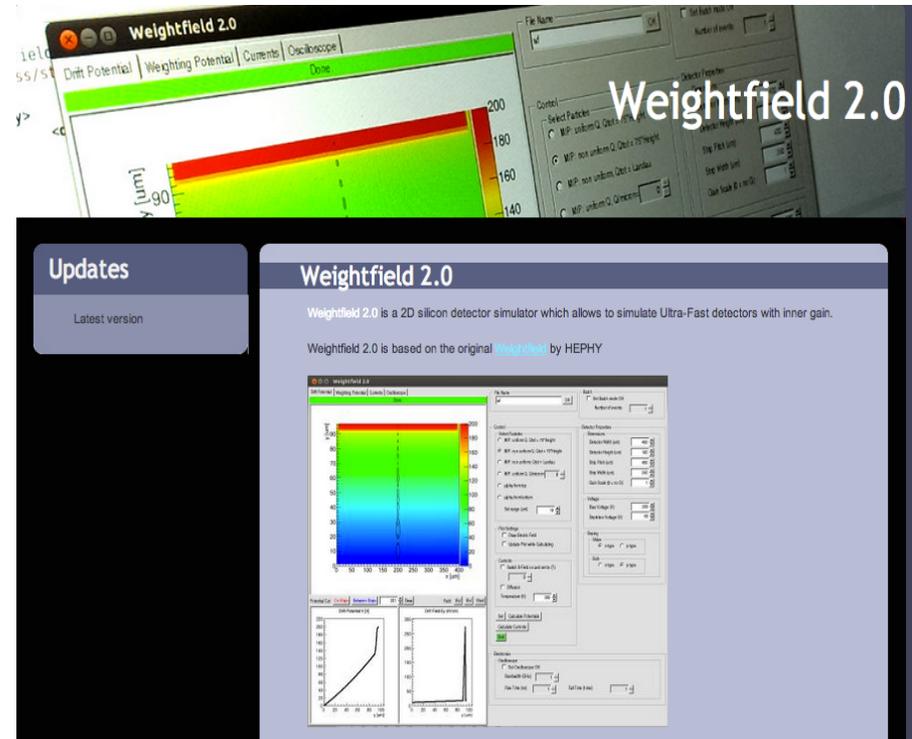
# Simulation

We developed a full sensor simulation to optimize the sensor design

WeightField2, F. Cenna, N. Cartiglia 9<sup>th</sup> Trento workshop, Genova 2014  
 Available at <http://personalpages.to.infn.it/~cartigli/weightfield2>

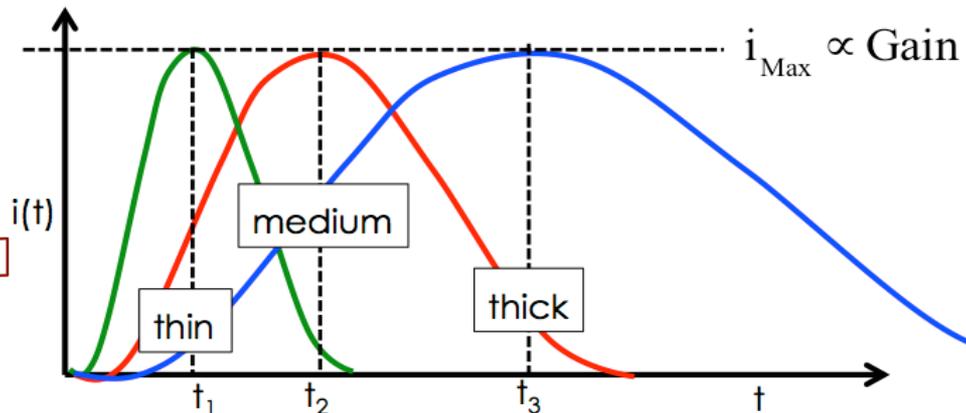
## It includes:

- Custom Geometry
- Calculation of drift field and weighting field
- Currents signal via Ramo's Theorem
- Gain
- Diffusion
- Temperature effect
- Non-uniform deposition
- Electronics



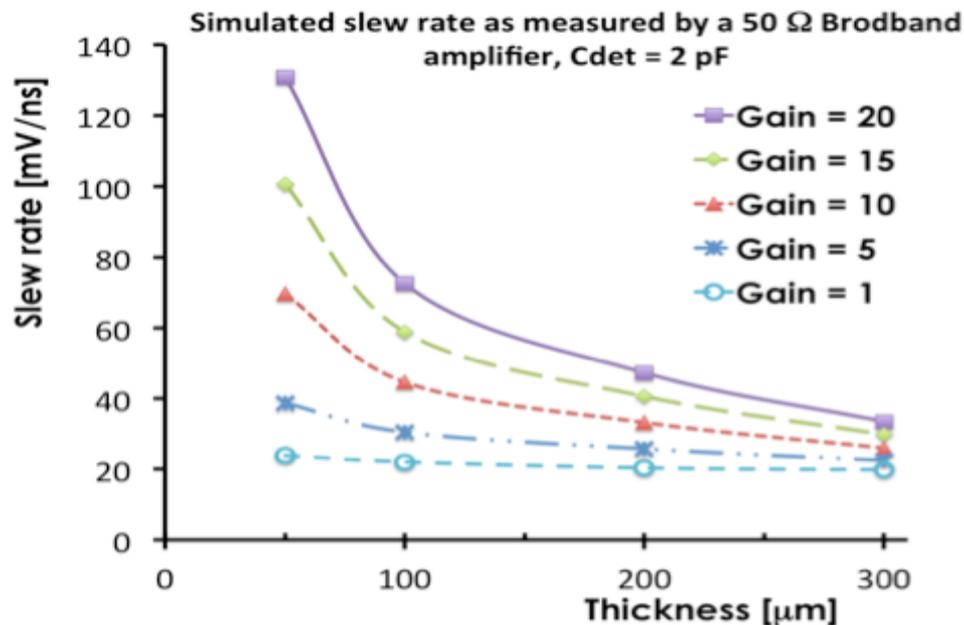
**For each event, it produces a file with the current output that can be used as input in the simulation of the electronic response.**

# Gain and slew rate vs thickness



**For a fixed gain:**

- amplitude = constant
- rise time  $\sim 1/\text{thickness}$



**The slew rate:**

- Increases with gain
- Increases  $\sim 1/\text{thickness}$

**→ Go thin!!**

**Significant improvements in time resolution require thin detectors**

# Ultra Fast Silicon Detectors

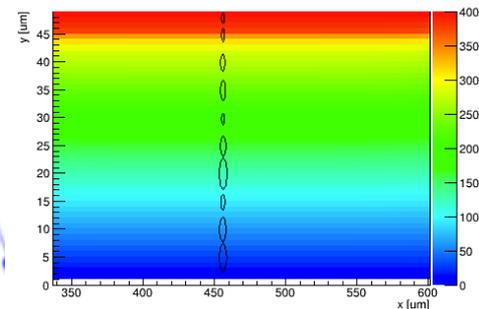
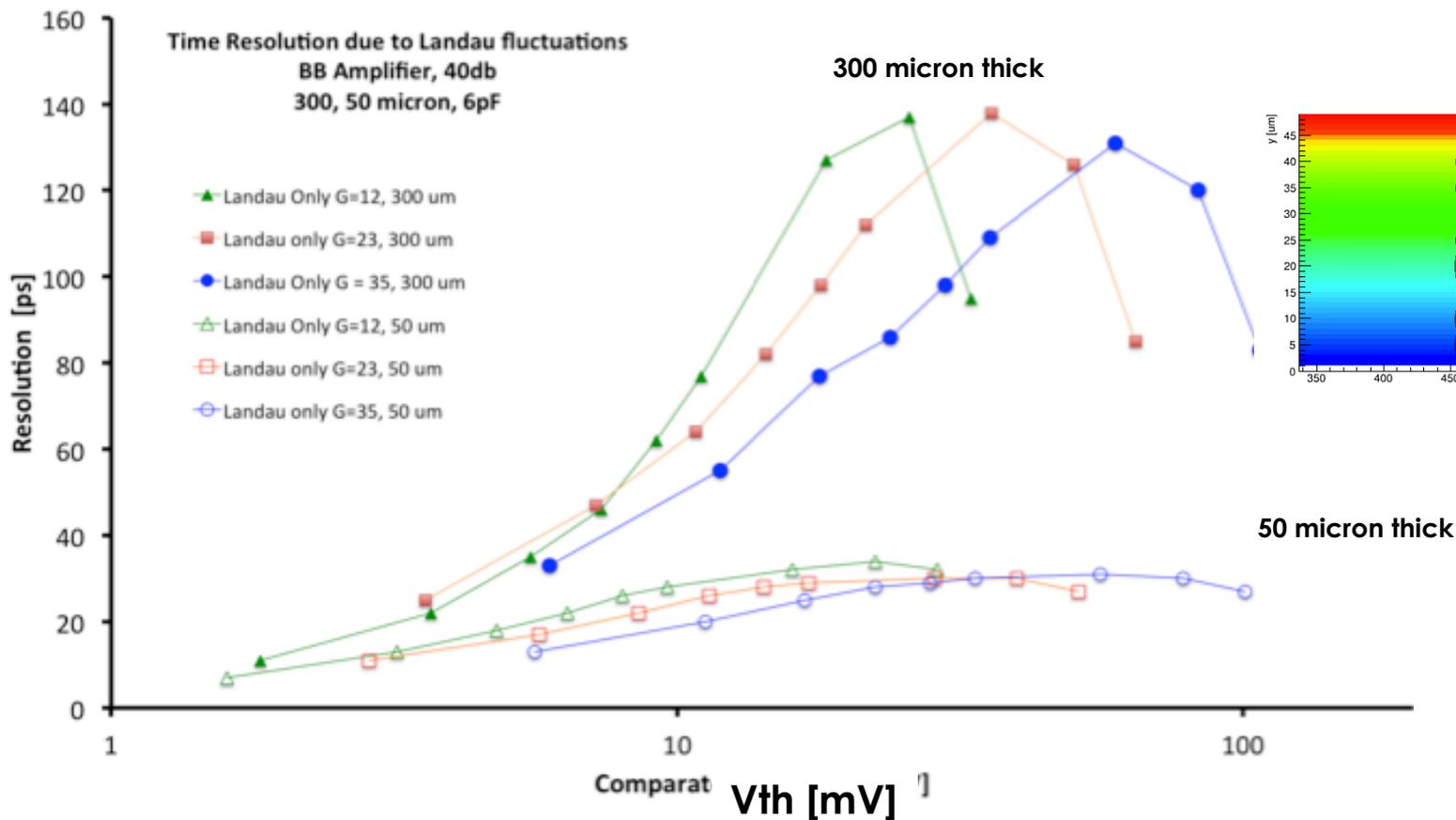
**UFSD are LGAD detectors optimized to achieve the best possible time resolution**

**Specifically:**

1. Thin to maximize the slew rate ( $dV/dt$ )
2. Parallel plate – like geometries (pixels..) for most uniform weighting field
3. High electric field to maximize the drift velocity
4. Highest possible resistivity to have uniform E field
5. Small size to keep the capacitance low
6. Small volumes to keep the leakage current low (shot noise)

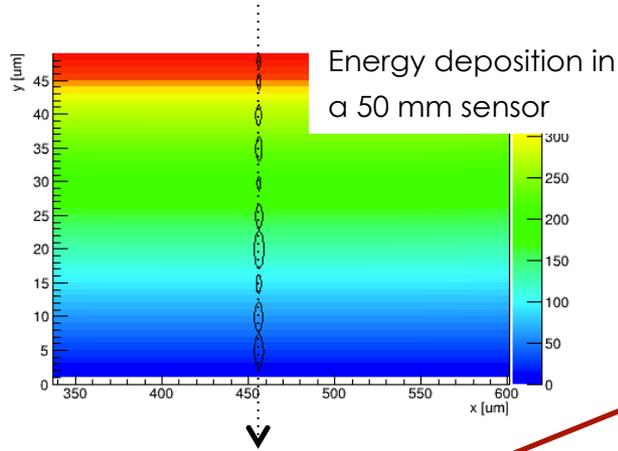
# Non uniform charge deposition along the track

This is a physical limit to time resolution: beat it with thin detectors and low comparator threshold.

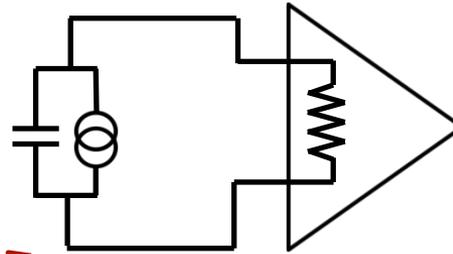


- ➔ Set the comparator threshold as low as you can
- ➔ Use thin sensors

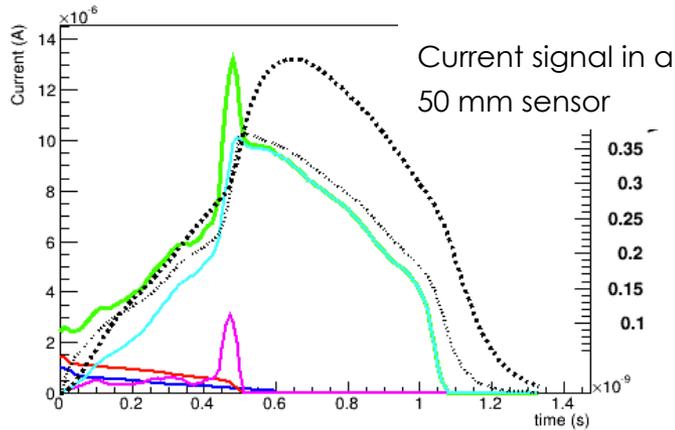
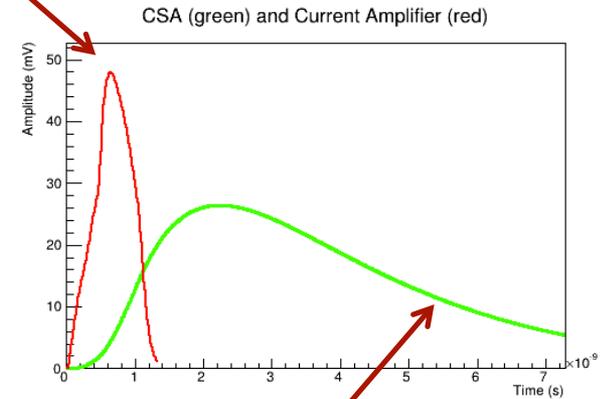
# What is the best pre-amp choice?



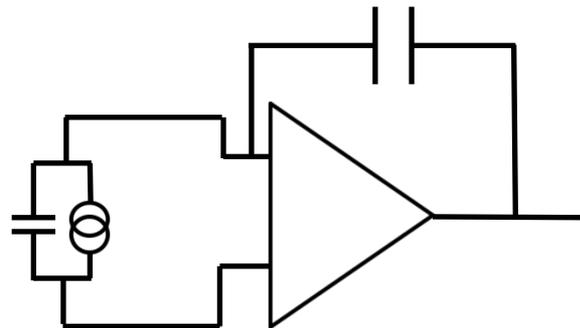
## Current Amplifier



- Fast slew rate
- Higher noise
- Sensitive to Landau bumps
- More power

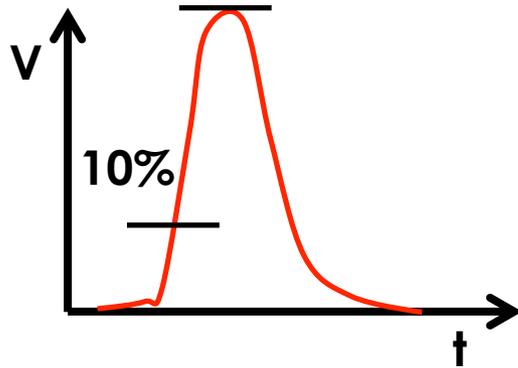


## Integrating Amplifier



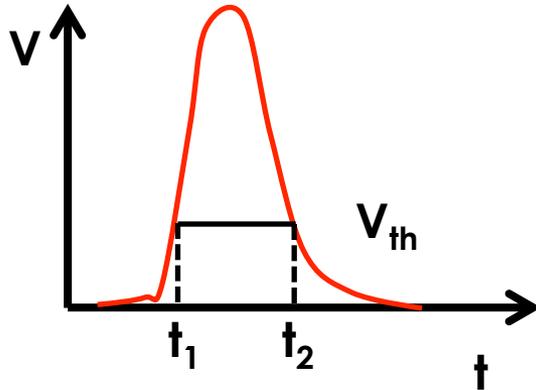
- Slower slew rate
- Lower noise
- Signal smoothing
- Less power

# What is the best “time measuring” circuit?



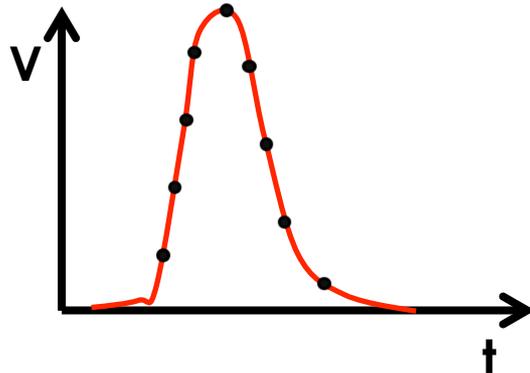
## Constant Fraction Discriminator

The time is set when a fixed fraction of the amplitude is reached



## Time over Threshold

The amount of time over the threshold is used to correct for time walk



## Multiple sampling

Most accurate method, needs a lot of computing power.

Possibly too complicated for large systems

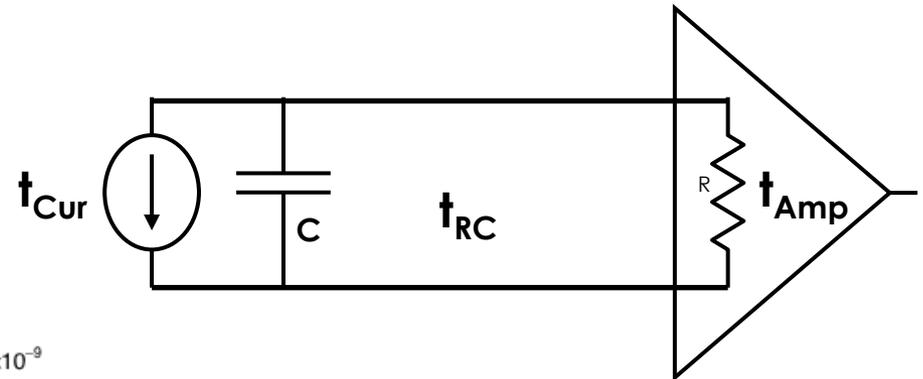
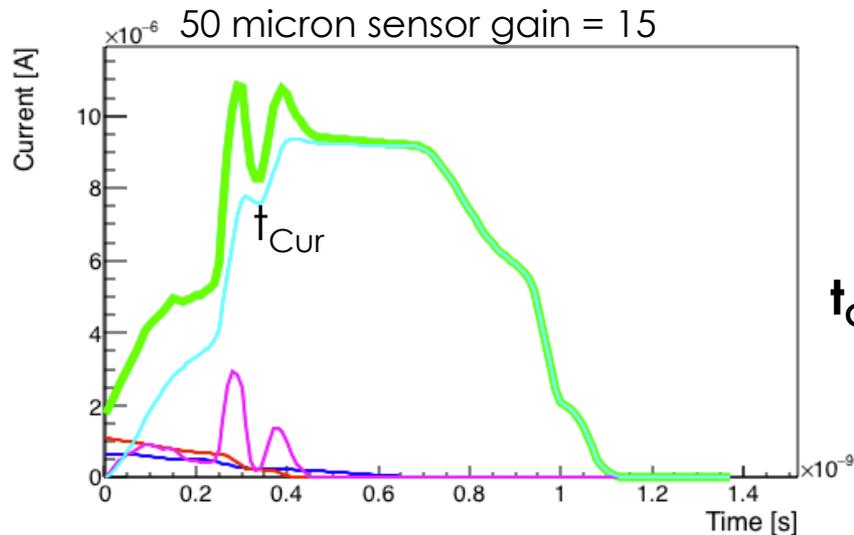
# The players: signal, noise and slope

Signal  $dV/dt$

Landau Noise

Shot Noise

Electronic Noise



The current rise time ( $t_{Cur}$ )

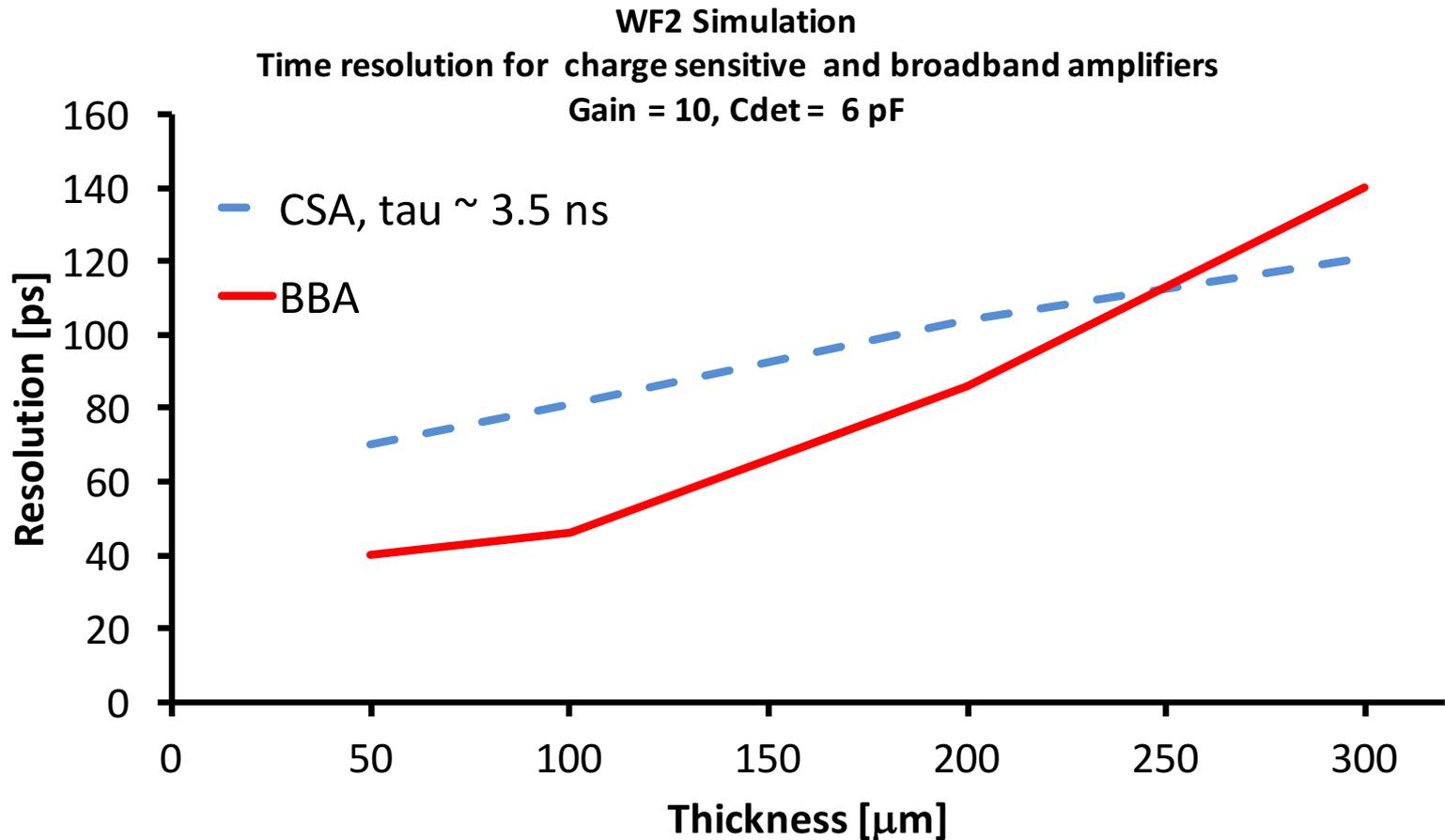
The RC circuit ( $t_{RC}$ )

Amplifier rise time ( $t_{Amp}$ )

There are 3 quantities determining the output rise time after the amplifier:

1. The signal rise time ( $t_{Cur}$ )
2. The RC circuit formed by the detector capacitance and the amplifier input impedance ( $t_{RC}$ )
3. The amplifier rise time ( $t_{Amp}$ )

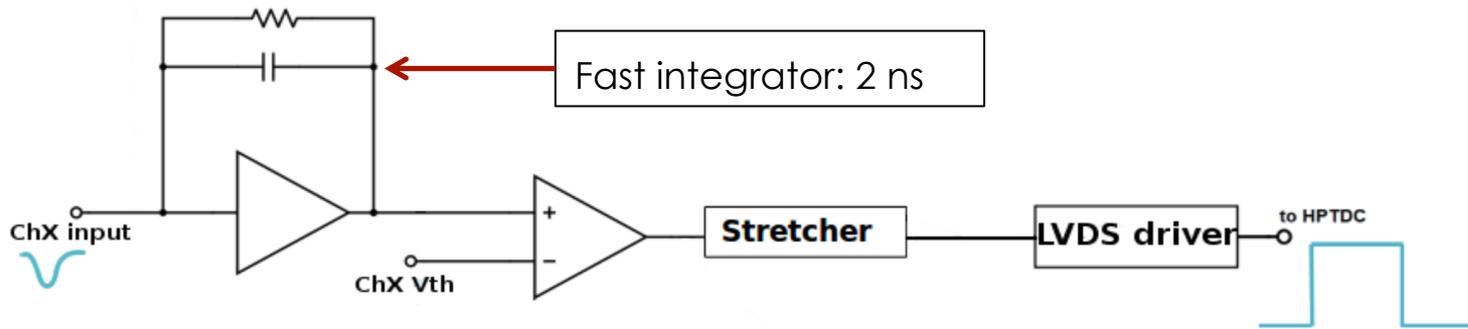
# Integrator or current amplifier?



- integrators work best with signals that are longer than their integration time
- Current amplifiers work best with very fast signals

# TOFFEE: Time Of Flight Front-End Electronics

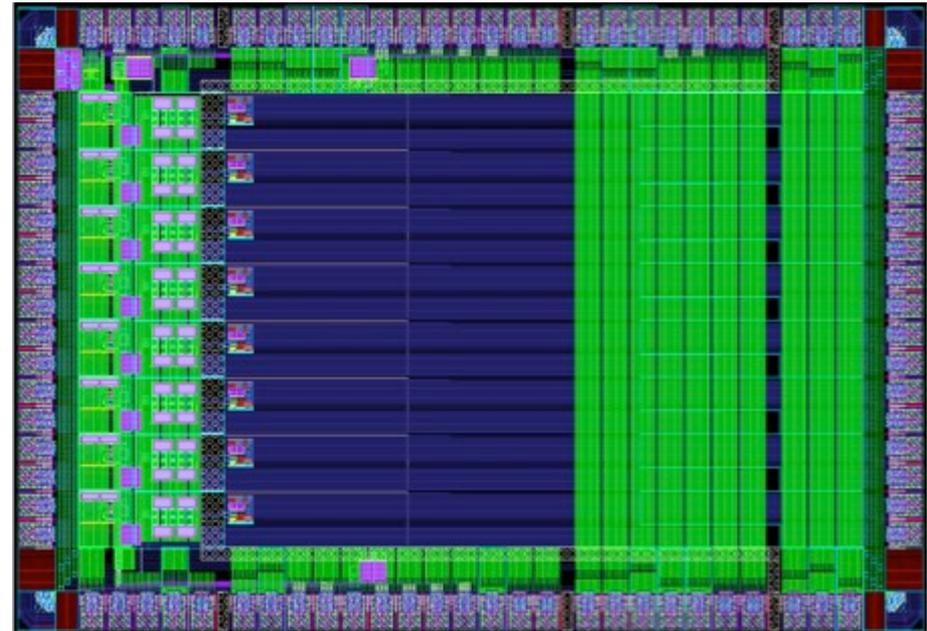
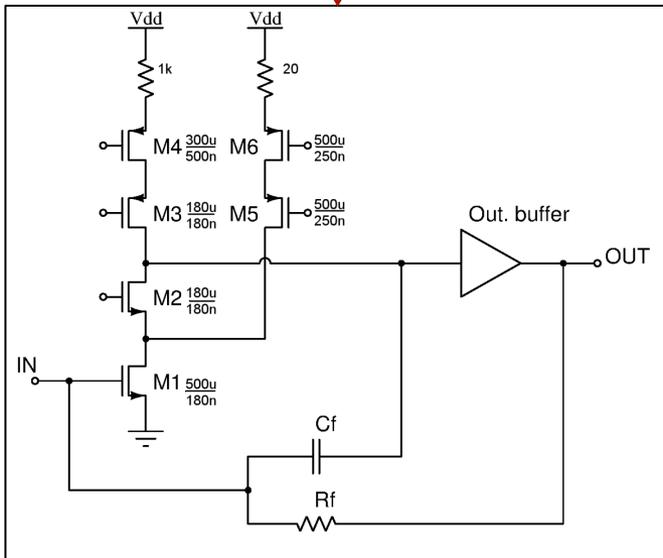
TOFFEE: Fully custom, 8 channel, 110 nm TMC chip to read-out UFSD



AMPLIFIER

DISCRIMINATOR  
first stage +  
second stage

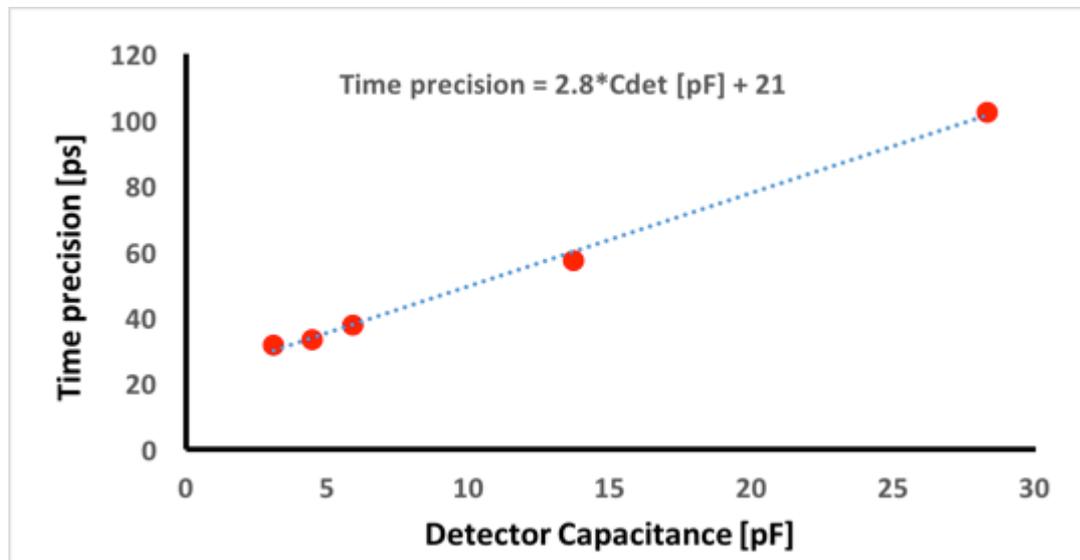
STRETCHER =  
10x delay blocks



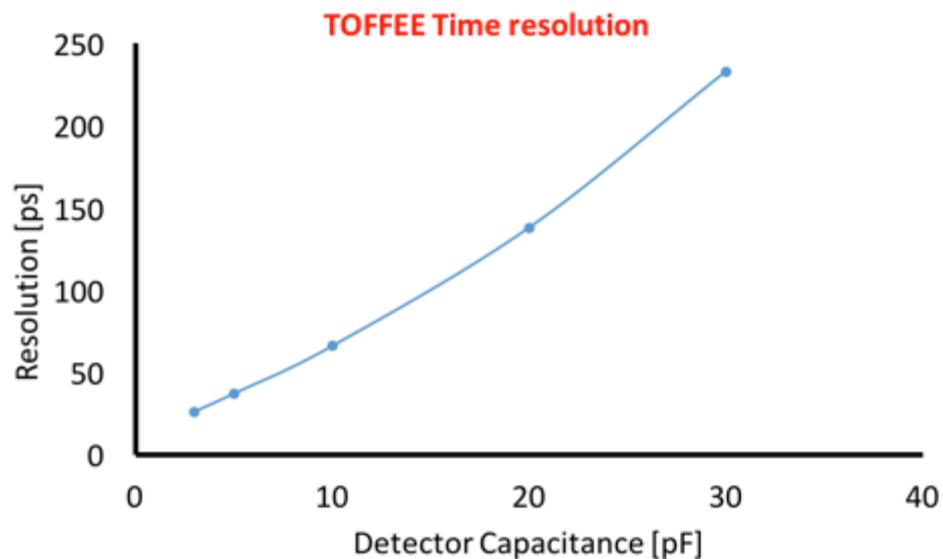
# Keep your capacitance low

Two examples of the effect of capacitance on the time resolution

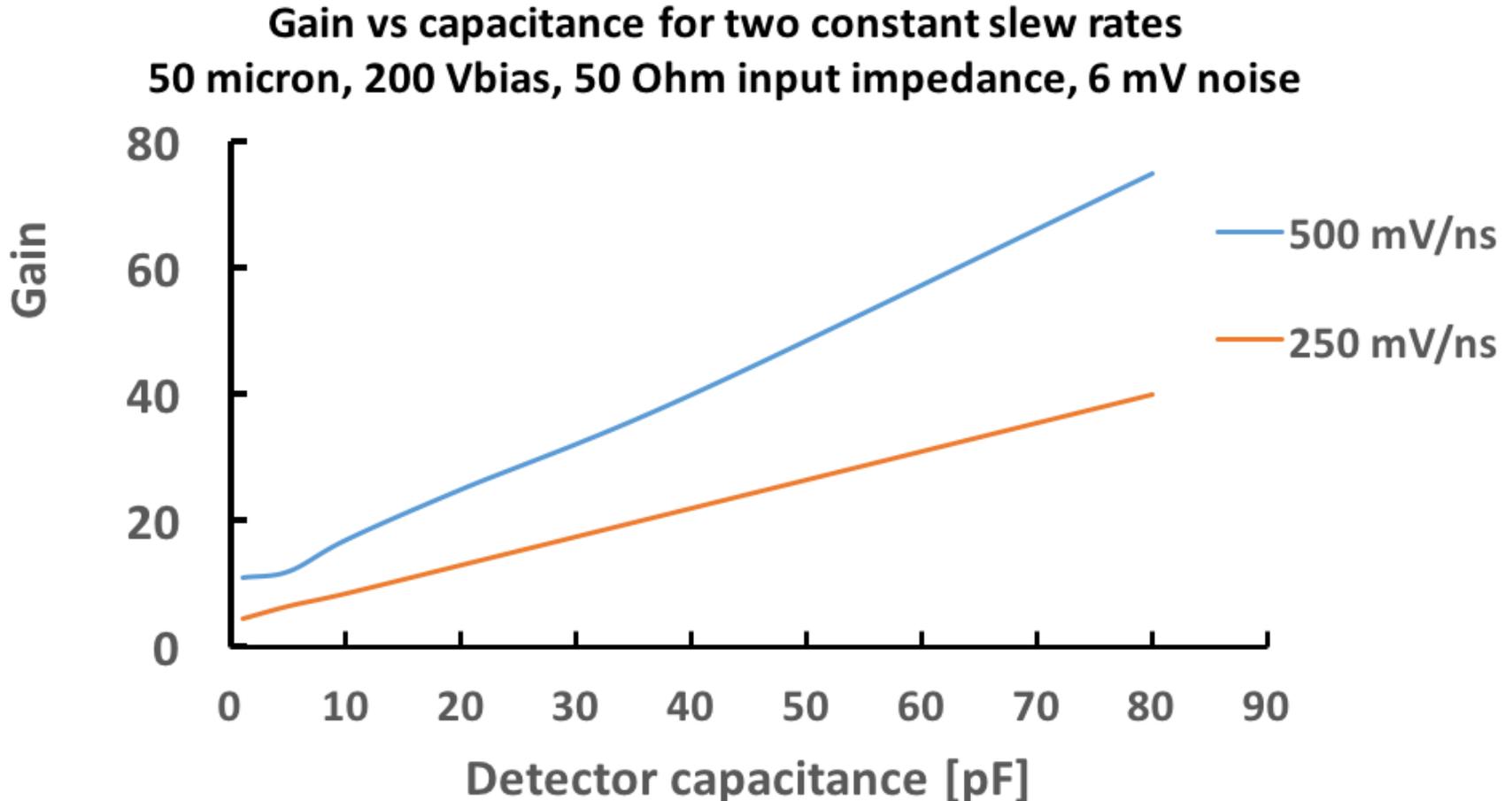
TOTEM current amplifier



TOFEE fast integrator



# How to read large sensor: more gain

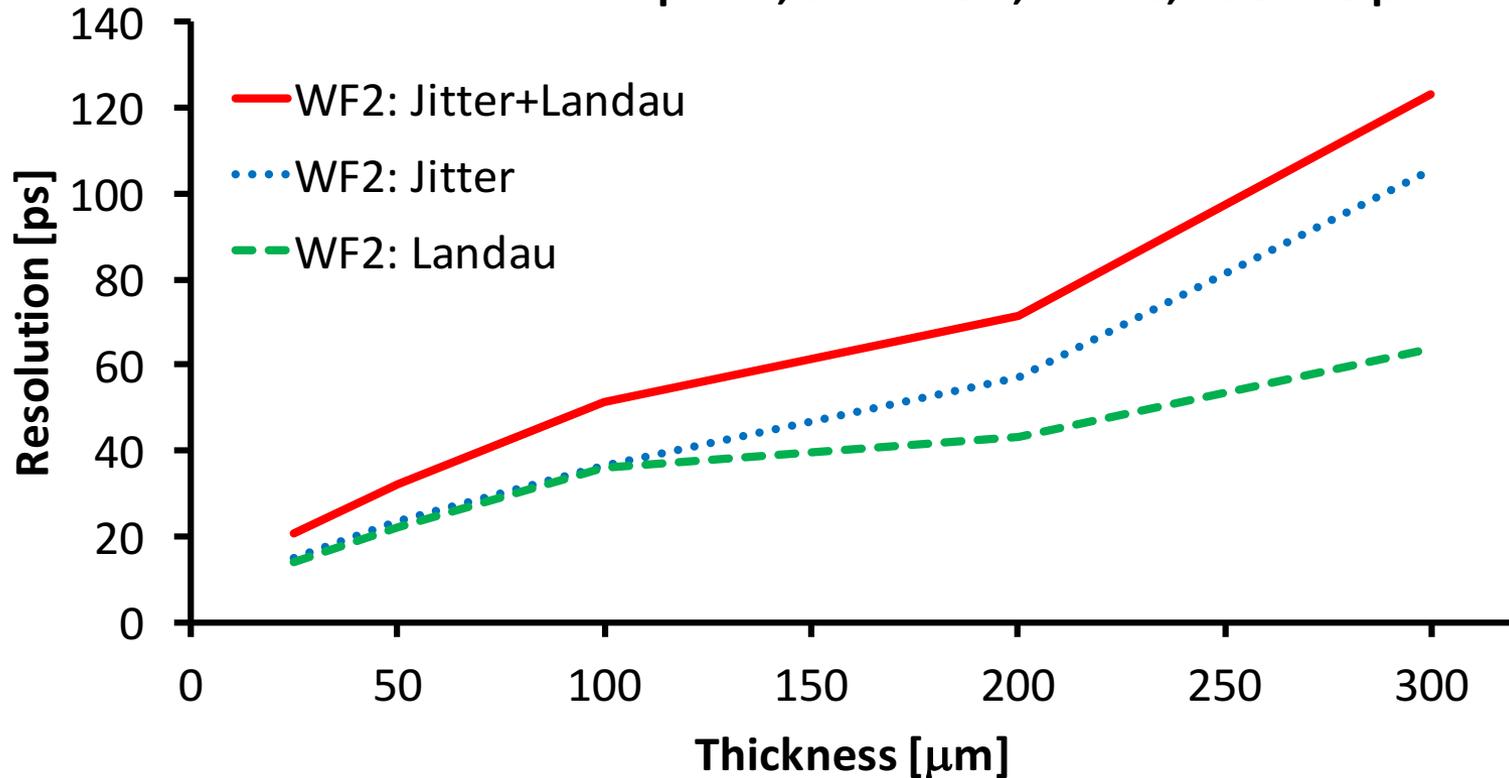


**But be aware: more gain is bad for noise, and terrible for radiation damage**

# How precise can we be?

## WF2 Simulation

BBA amplifier, CFD = 15%,  $G \sim 20$ ,  $C_{det} = 6$  pF



- ➔ Two main contributions: charge non uniformity and Jitter
- ➔ The time resolution has a lower limit due to charge non uniformity

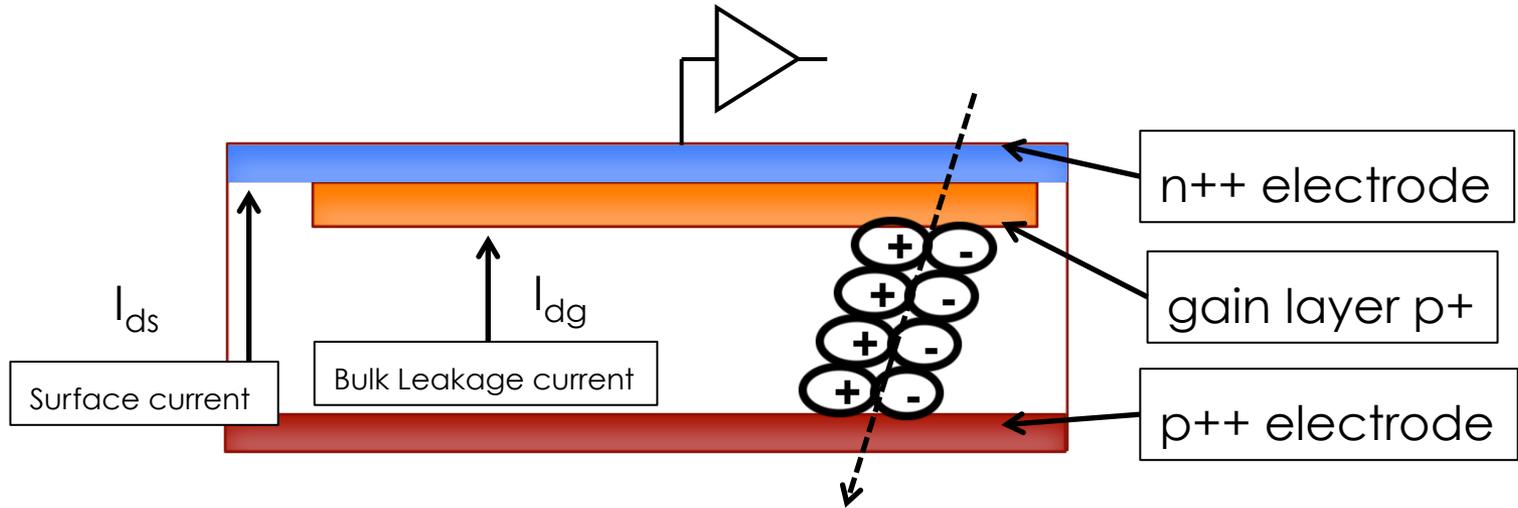
# Irradiation effects

---

## **Irradiation causes 3 main effects:**

- Decrease of charge collection efficiency due to trapping  
→ Very small in thin sensor
- Increased leakage current, shot noise → next slides
- Gain layer disappearance → following slides

# Shot noise in LGAD - APD



$$i_{Shot}^2 = 2eI_{Det} = 2e \left[ I_{Surface} + (I_{Bulk}) M^2 F \right]$$

Current density, nA/sqrt(f)

$$F = Mk + \left( 2 - \frac{1}{M} \right) (1 - k)$$

$$F \sim M^x$$

$k = e/h$  ionization rate

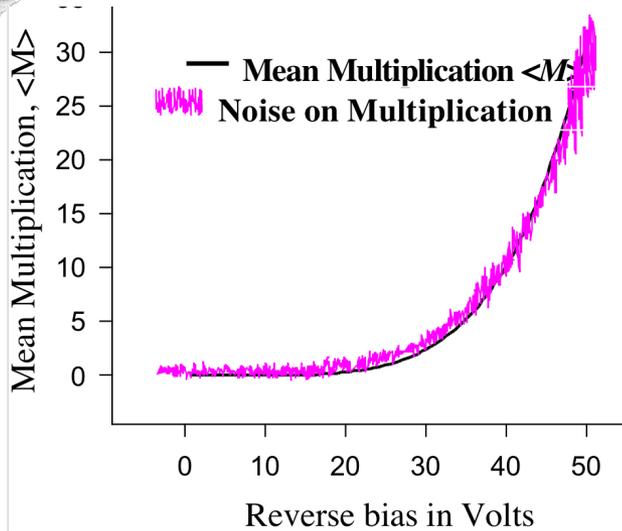
$x =$  excess noise index

$M =$  gain

Correction factor to the standard Shot noise, due to the noise of the multiplication mechanism

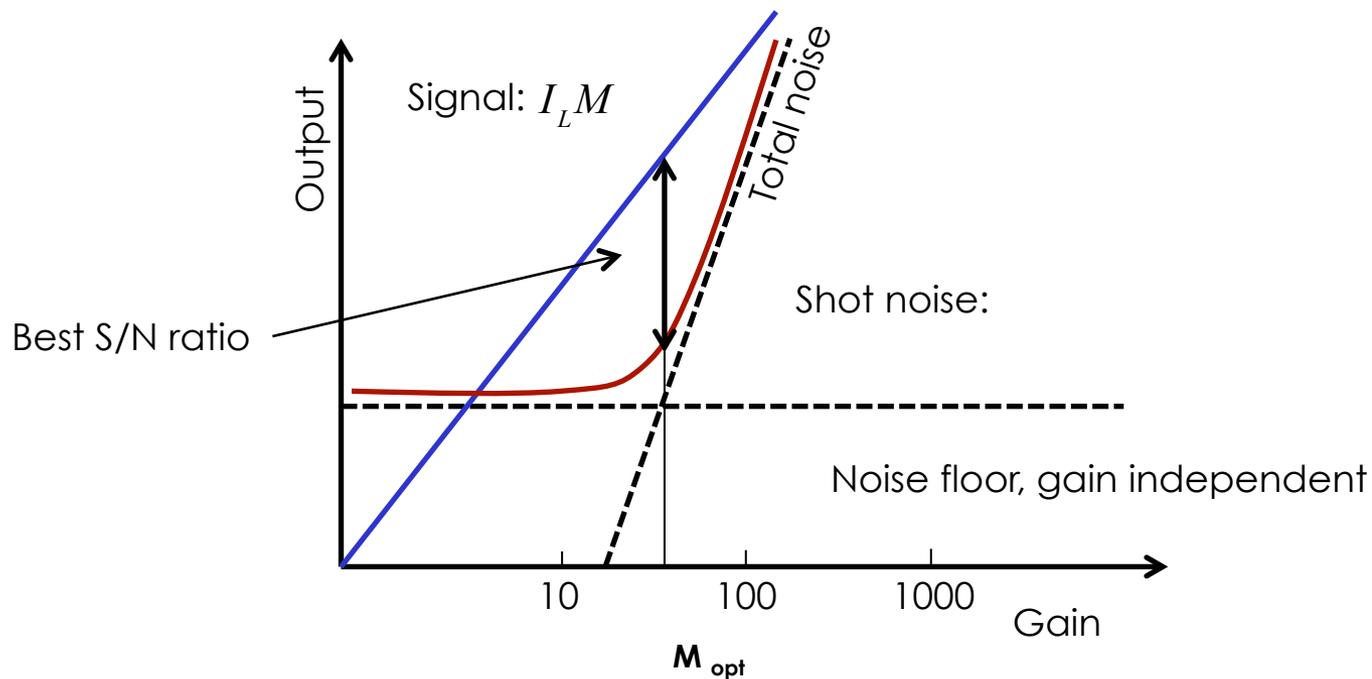
$$F = \frac{\langle M^2 \rangle}{\langle M \rangle^2} \Rightarrow \langle M^2 \rangle = \langle M \rangle^2 F$$

# Noise in LGAD & APD – Aide Memoire



**Noise increases faster than then signal → the ratio S/N becomes worse at higher gain.**

**There is an Optimum Gain value: 10-20?**



# Shot noise

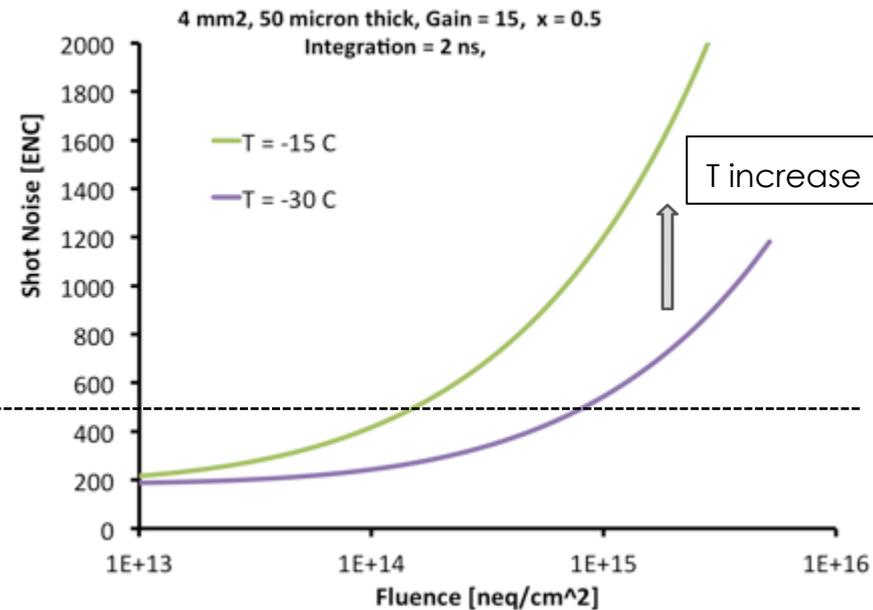
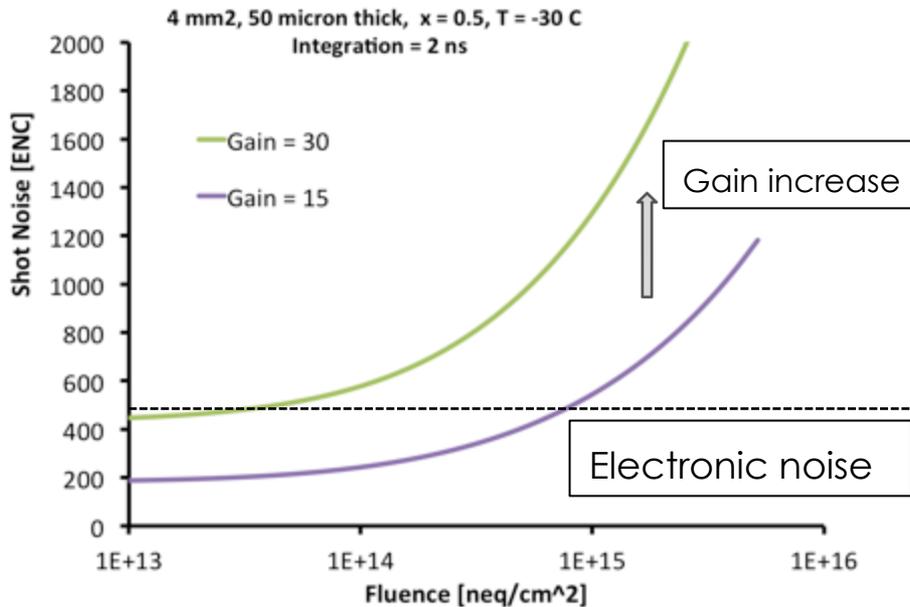
Let's assume a 4 mm<sup>2</sup> pad, 50 micron thick, and a electronic noise of 500 ENC

**What is the effect of shot noise as a function of radiation?**

Steep dependence on gain

$$I = \alpha * \Phi * \text{Volume} \quad \alpha = 3 \cdot 10^{-17} / \text{cm}$$

$$\text{Shot noise: } ENC = \sqrt{\int i_{\text{Shot}}^2 df} = \sqrt{\frac{I * (\text{Gain})^{2+x}}{2e}} * \tau_{\text{Int}}$$



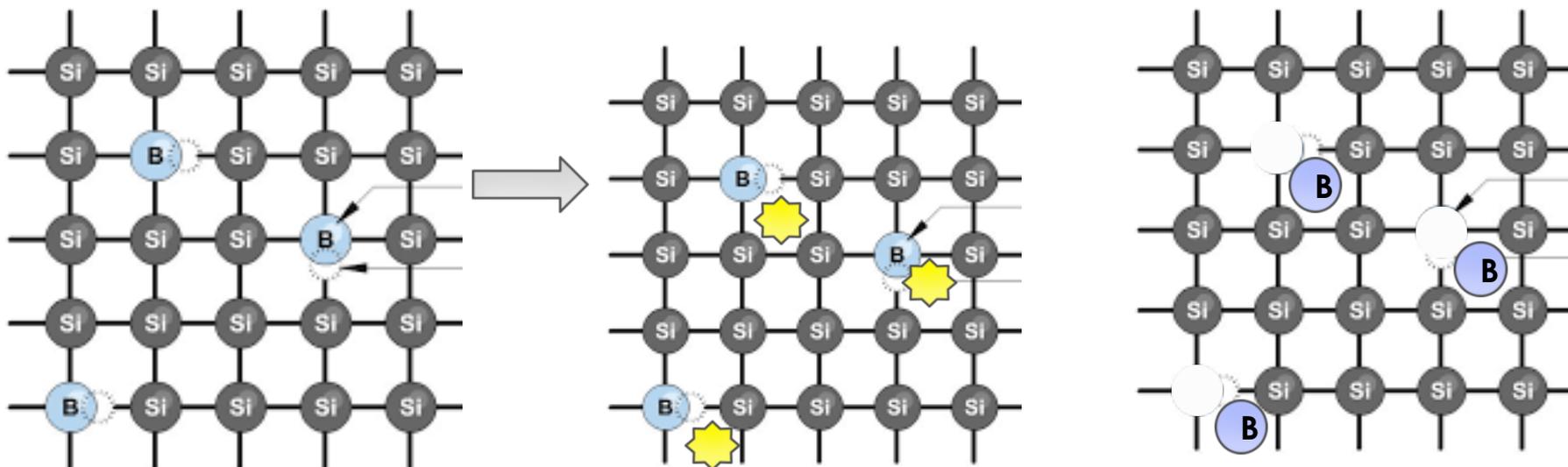
**To minimize Shot noise:**

- ➔ Low gain!! Keep the gain below ~ 20
- ➔ Cool the detectors
- ➔ Use small pads to have less leakage current

# Radiation issue: Initial acceptor removal

This term indicates the “removal” of the initially present p-doping.  
 For UFSD this is particularly problematic as it removes the gain layer

**Irradiation → Defects → Boron becomes interstitial**

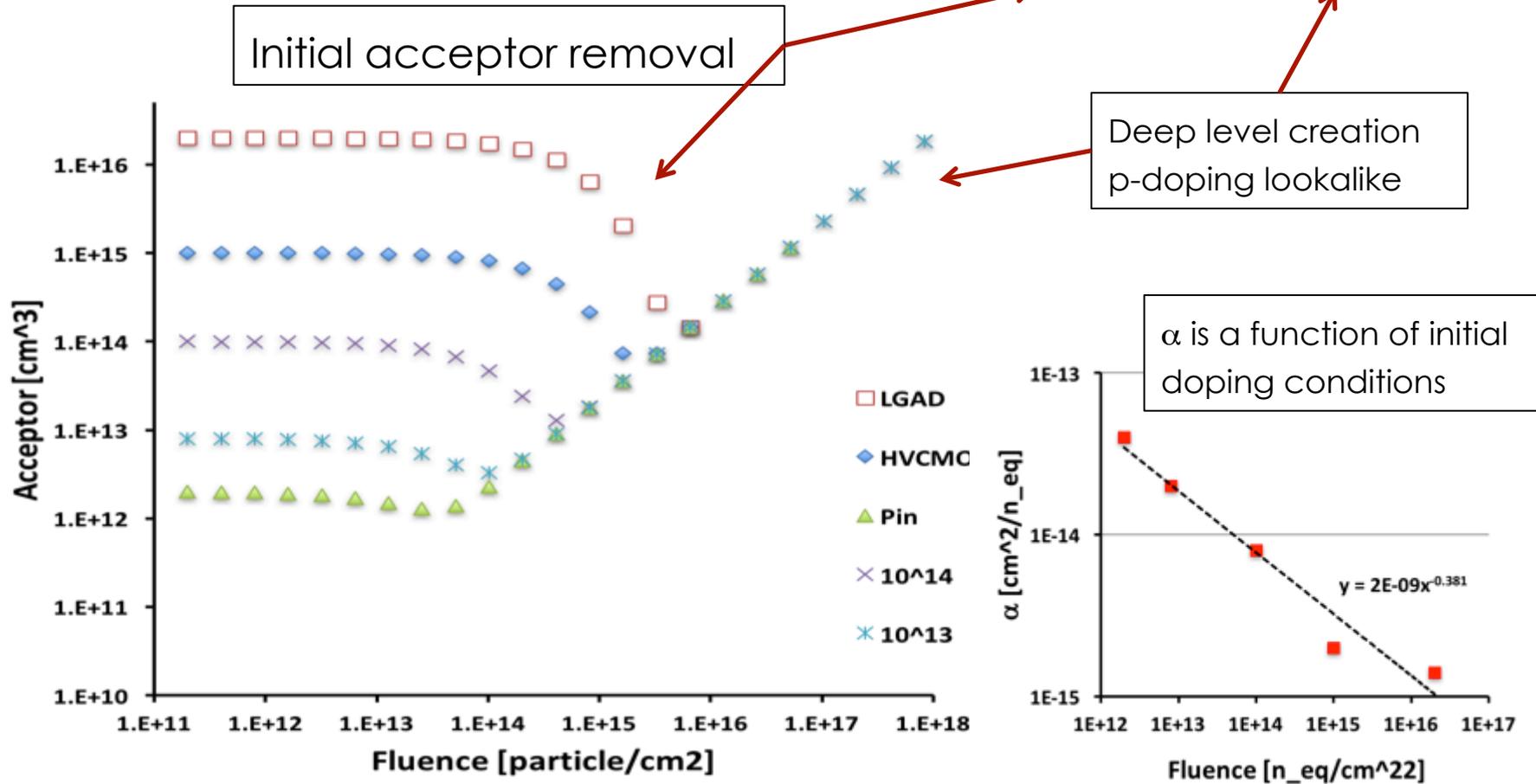


The boron doping is still there, only it has been moves into a different position and it does not contribute to the doping profile, it is inactive

# Irradiation main problem: gain layer disappearance

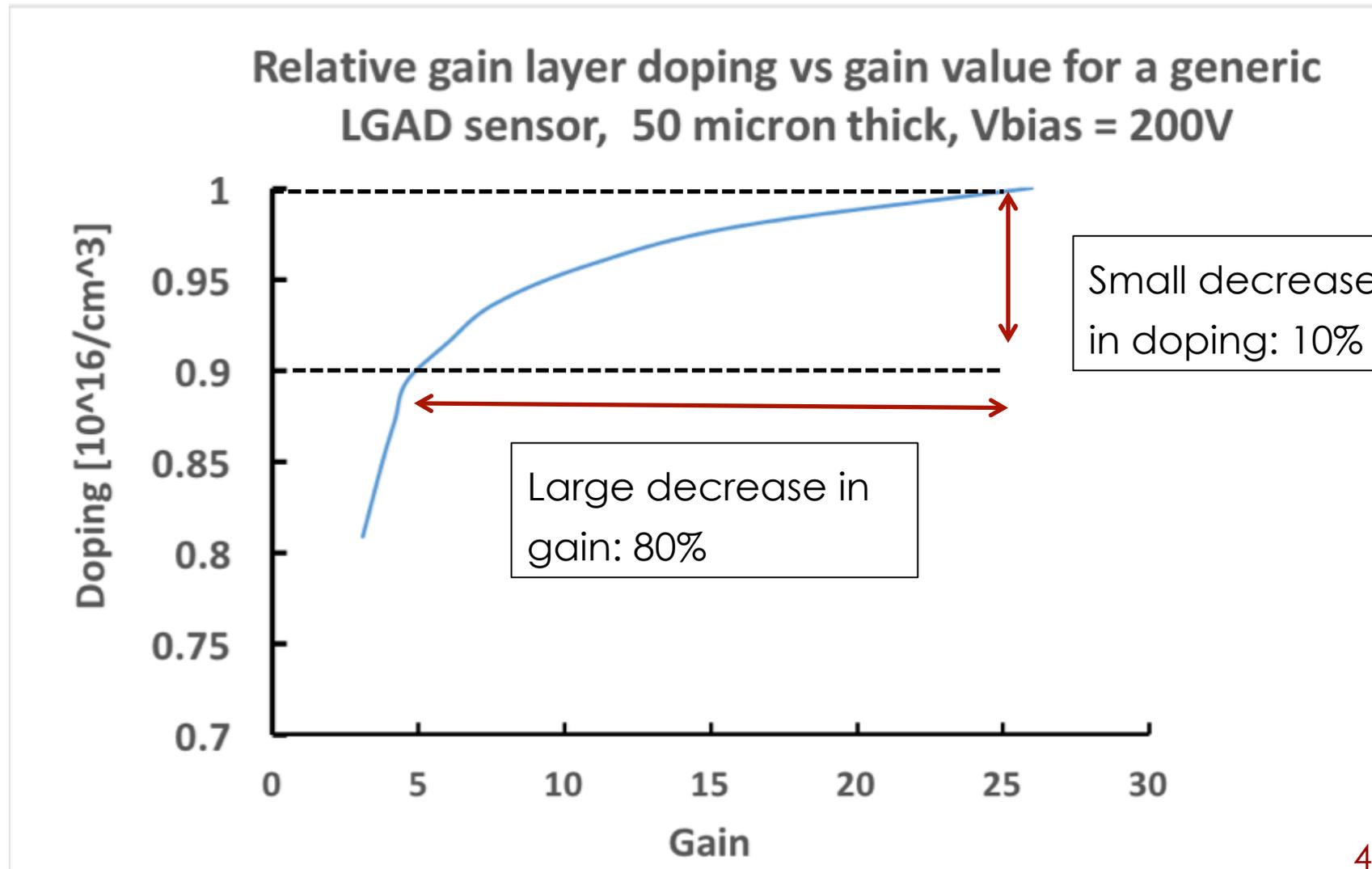
Density of the doping vs irradiation:

$$N_D = N_0 e^{-\alpha\phi} + \beta\phi$$



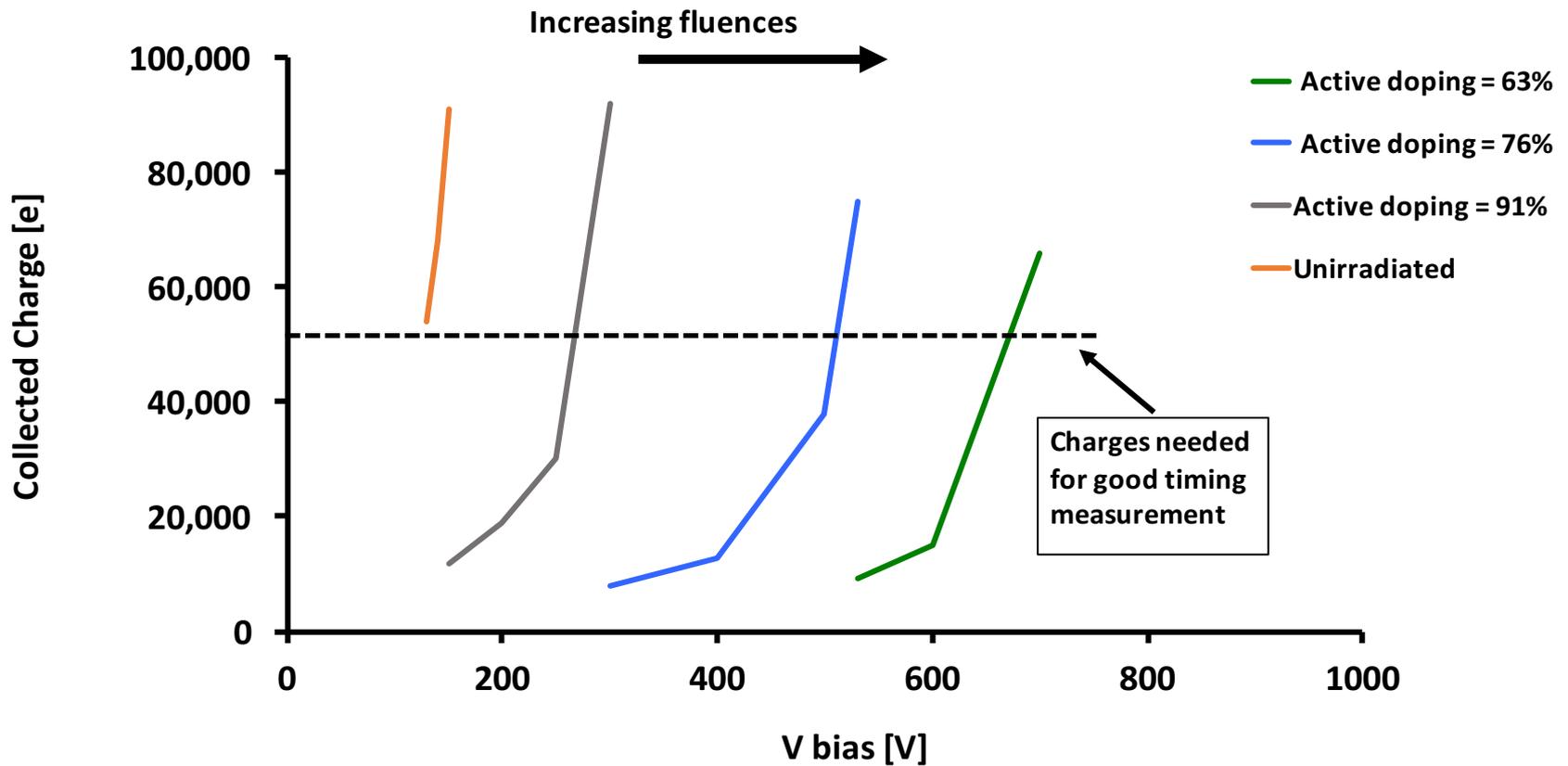
# Gain vs gain layer doping

Unfortunately, the gain is very sensitive to the doping level



# Compensation with Vbias

The necessary field can be recovered by increasing the external Vbias: proven to work up to  $5 \cdot 10^{14} \text{ n}^{\text{eq}}/\text{cm}^2$



## How can we sustain more radiation?

There is the understanding that Boron is not “gone”, but it is simply inactive, it has gone from being sub-stitutional to being interstitial.

➔ The Boron presence has been measured after irradiation

What can we do?

➔ Try different manufactures

- CNM has been tested
- FBK and HPK are being tested as we speak, results at TREDI 2017, end of February

➔ Try different dopant: use Gallium instead of Boron

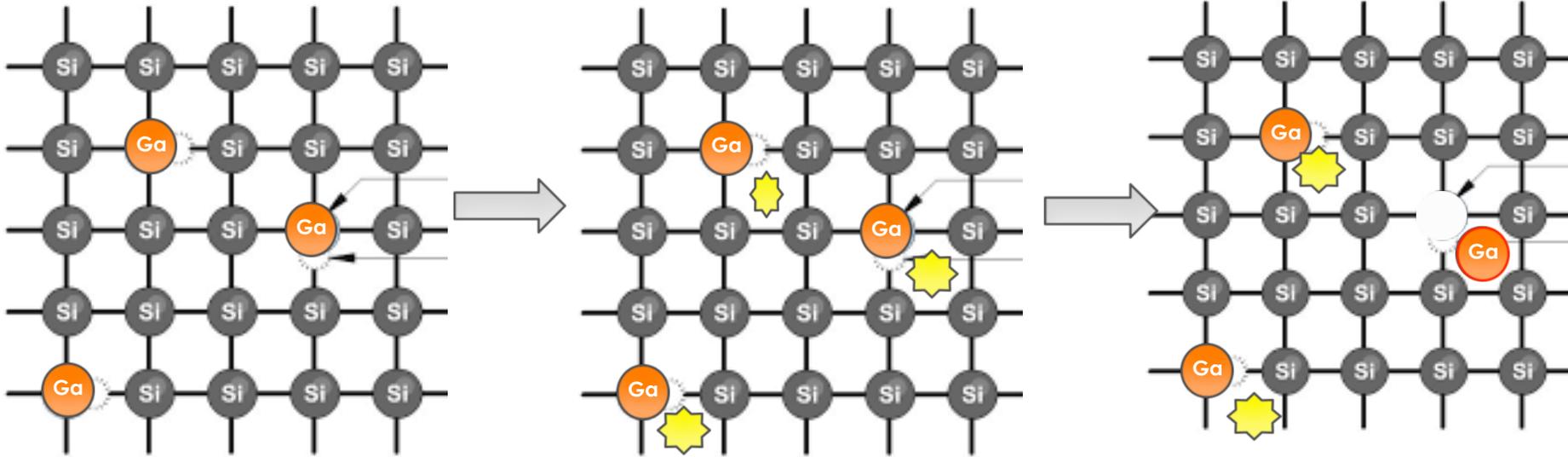
- Being manufactured right now

➔ Add carbon to the gain layer

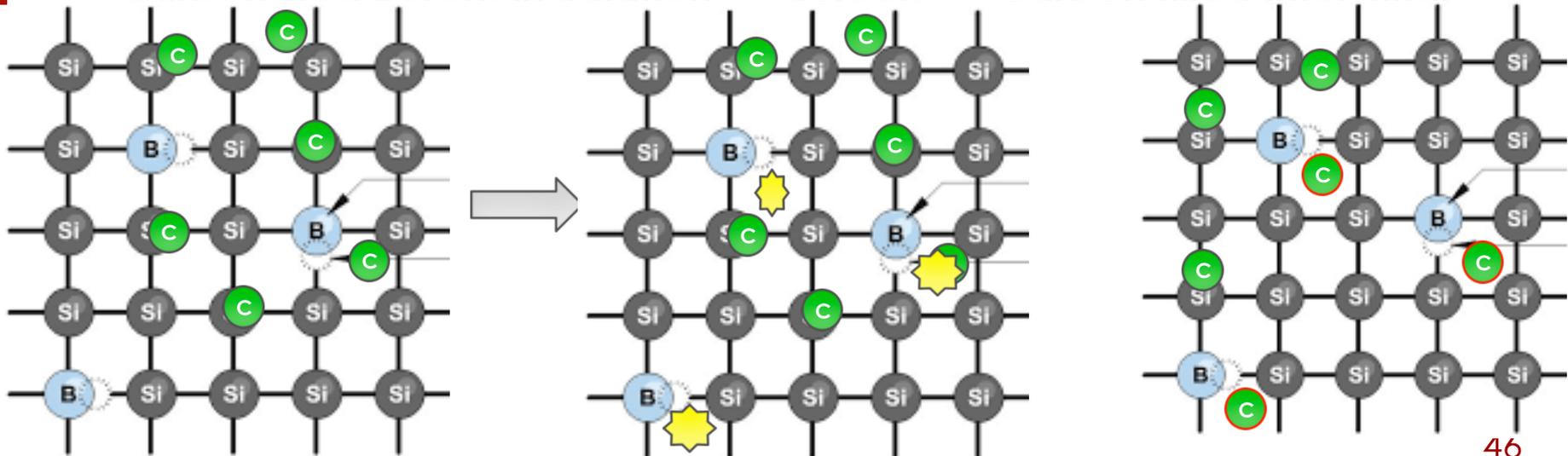
- It might protect the Boron...

# Initial acceptor removal: mitigation

**Gallium doping:** Irradiation  $\rightarrow$  defects  $\rightarrow$  Lower diffusivity



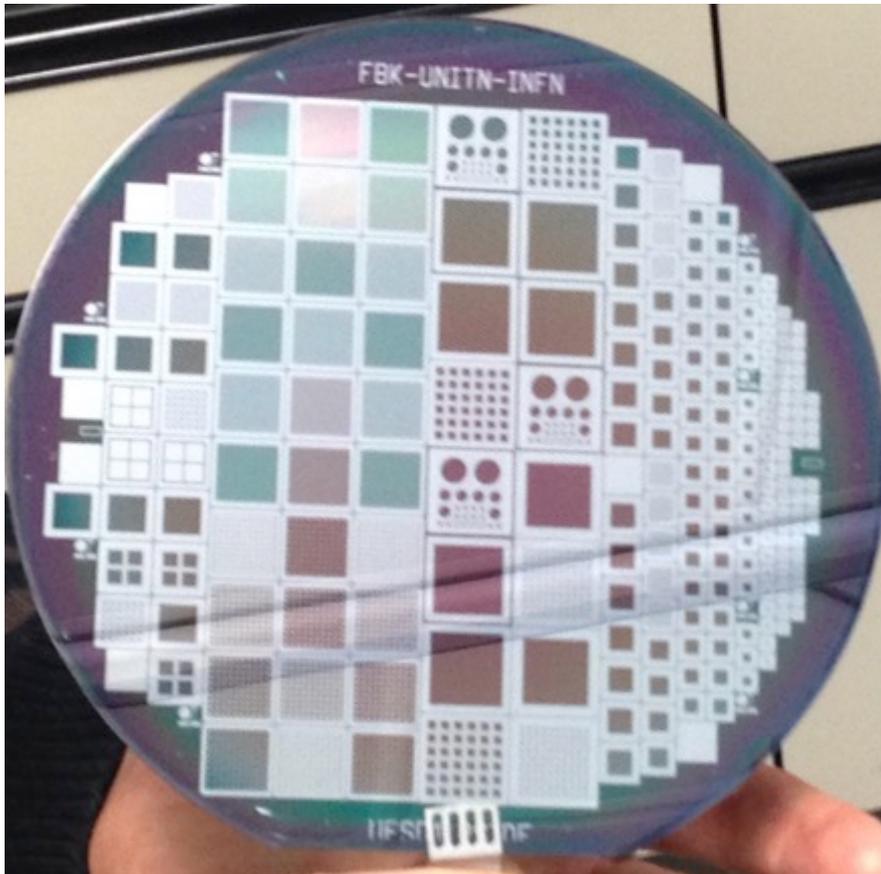
**Carbonated Boron:** Irradiation  $\rightarrow$  defects  $\rightarrow$  Carbon filled interstitial



# Sensors: FBK & CNM

FBK 300-micron production  
Very successful, good gain and overall behavior

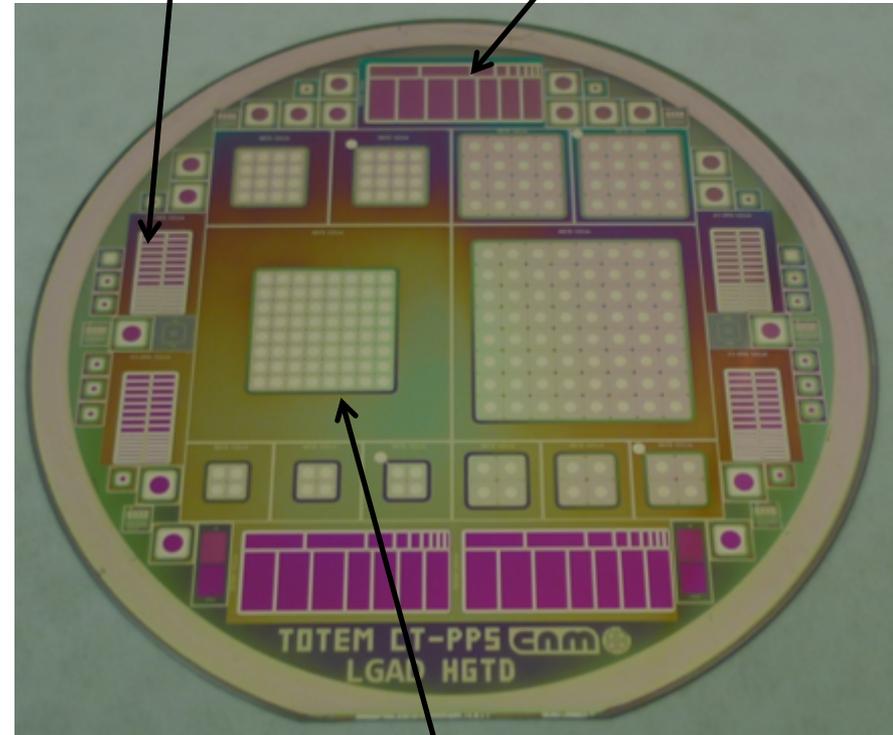
→ We have now a second producer



CNM 75-micron  
CNM 50-micron production

x4 CT-PPS

x3 TOTEM



ATLAS High Granularity Timing Det.

# Sensors for the CMS CT-PPS detectors

New production of 50 micron thick, segmented UFSD sensors.

Gain  $\sim 15$

32 fat strip array for CT-PPS

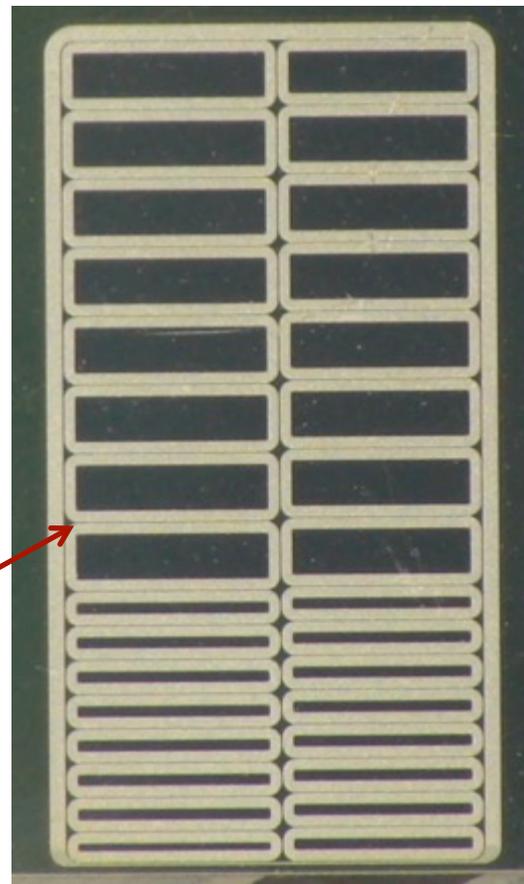
Strips:

3 mm x 0.5 mm

3 mm x 1 mm

Distance between pads: 50 micron

→ Able to produce segmented UFSD

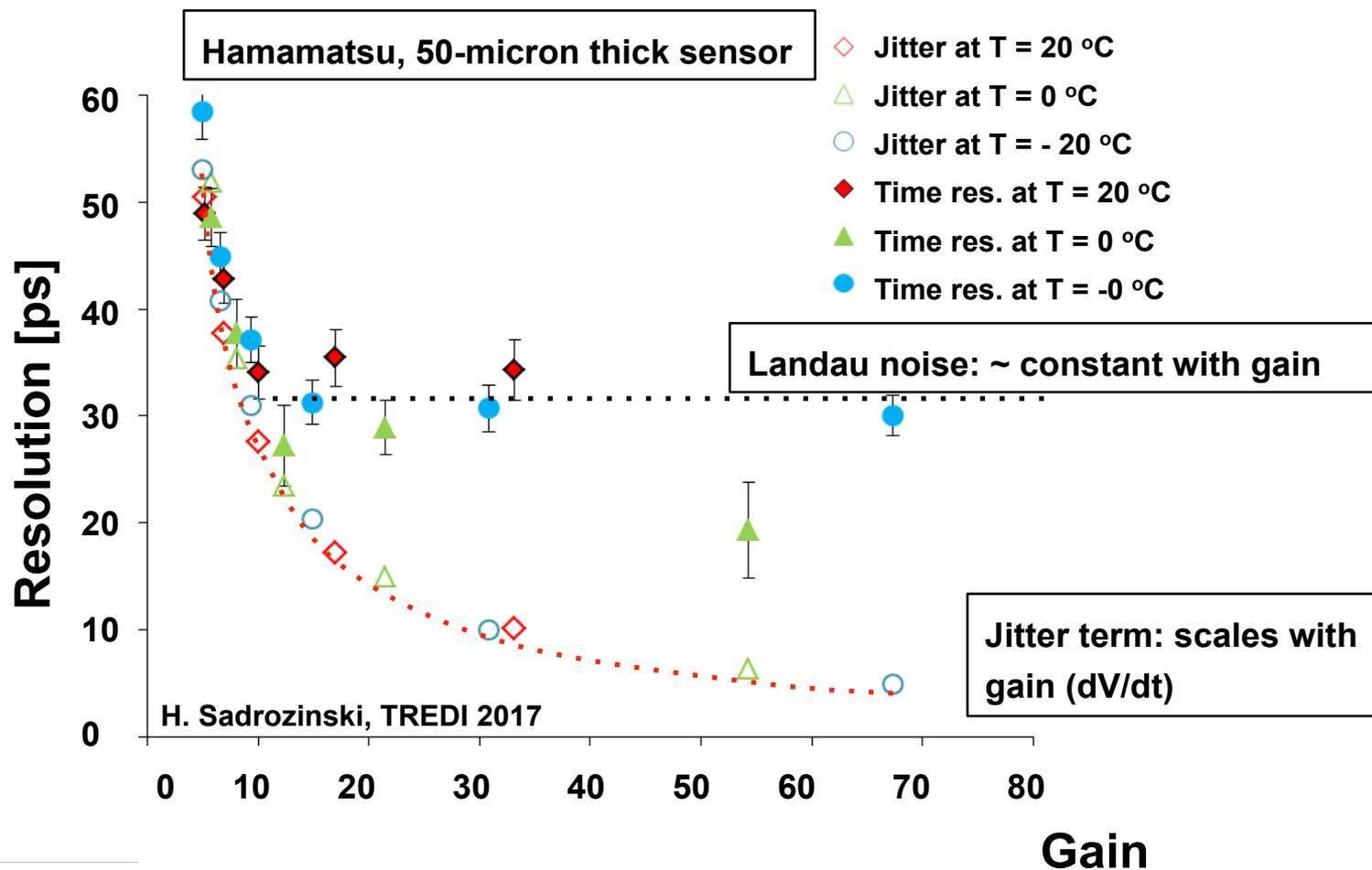


12 mm

6 mm

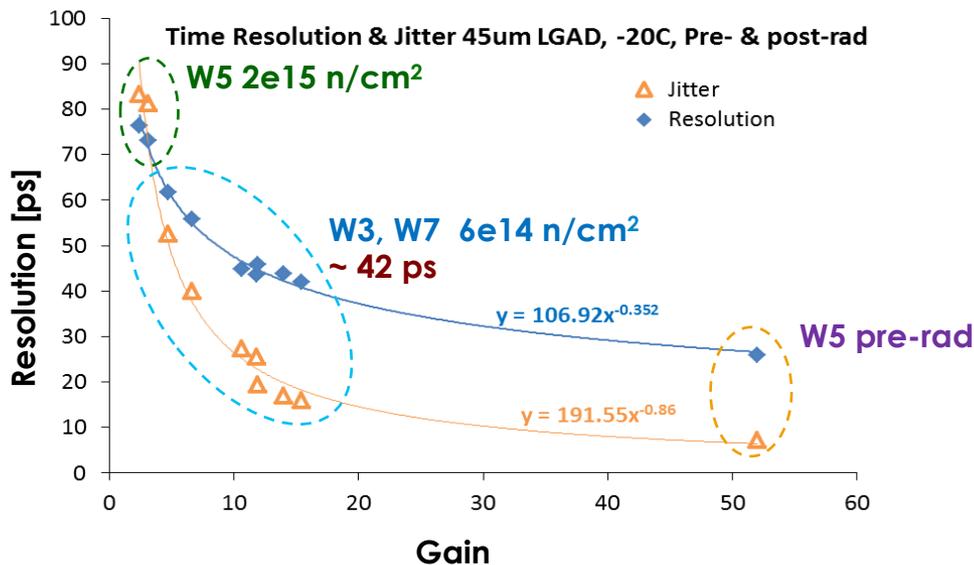
# Latest results on UFSD time resolution

UFSD from Hamamatsu confirm our simulation: 30 ps time resolution,  
Value of gain  $\sim 20$



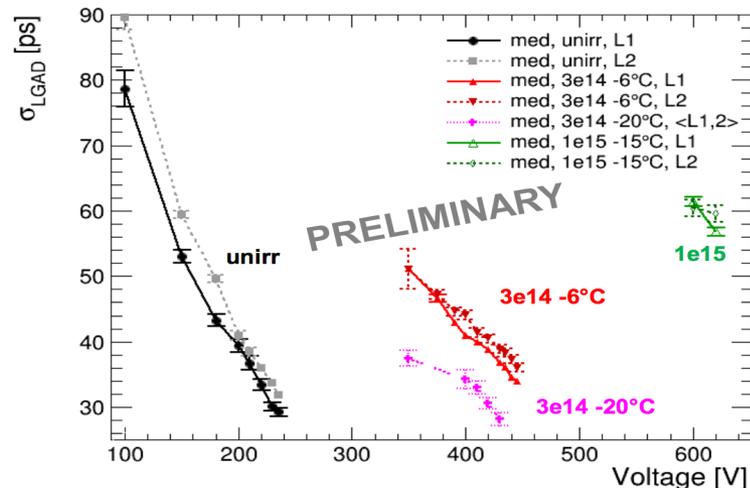
# Time resolution for irradiated sensors

H. Sadrozinski, TREDI 2017



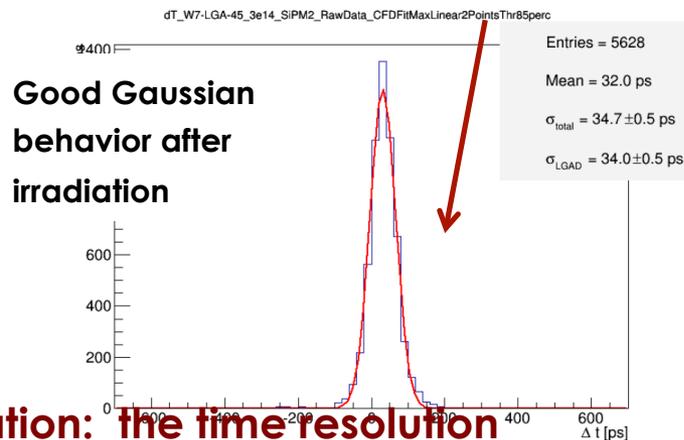
J. Lange, TREDI 2017

Irradiated, Medium Dose



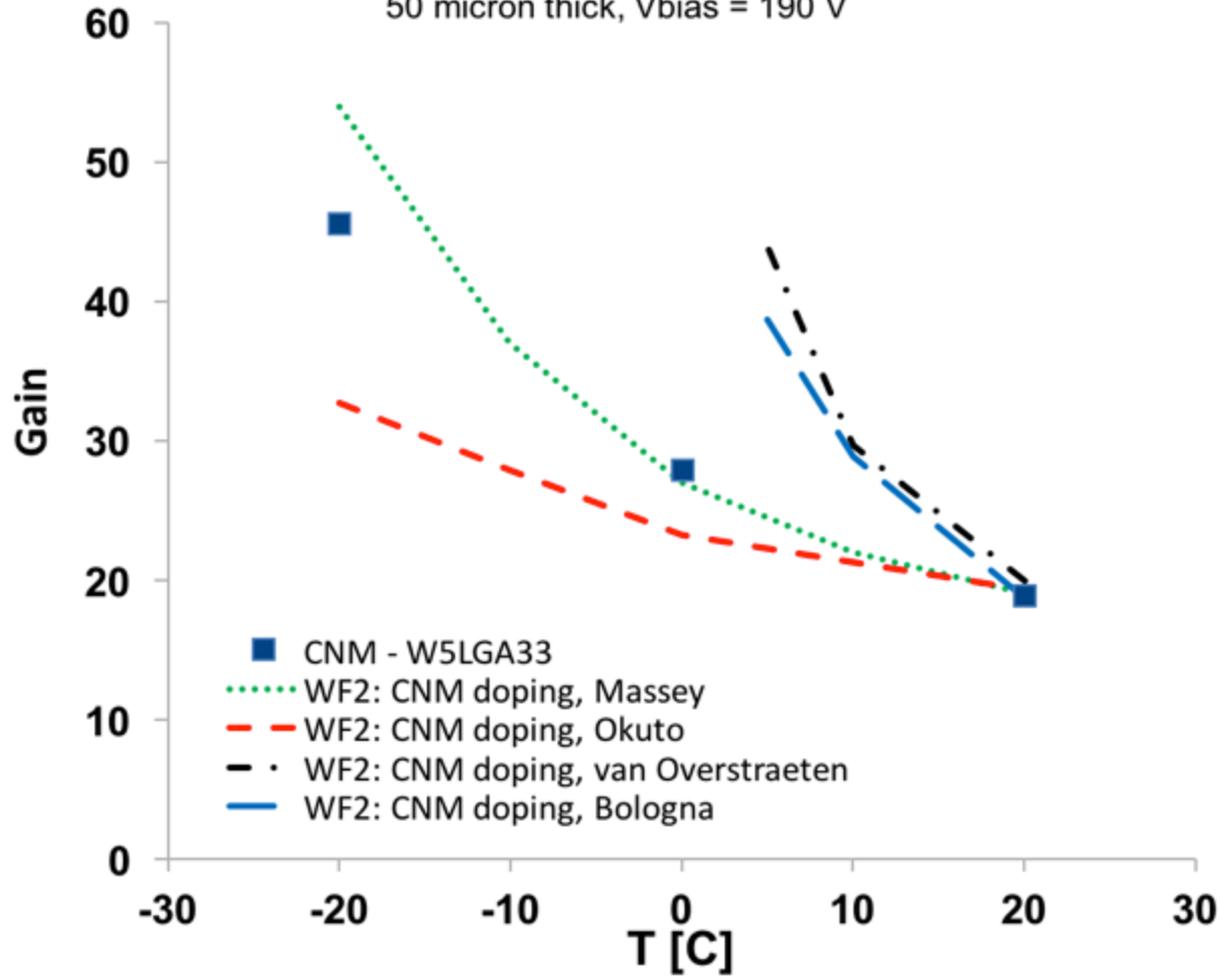
**No unexpected features**, the signals are still large and the leakage current does not prevent to reach good time resolution.

**No difference in behavior before and after irradiation: the time resolution scales with gain. → Keep the gain high**



# Running cold

CNM W5 LGA33, 1.2 x 1.2 mm pad  
50 micron thick, Vbias = 190 V

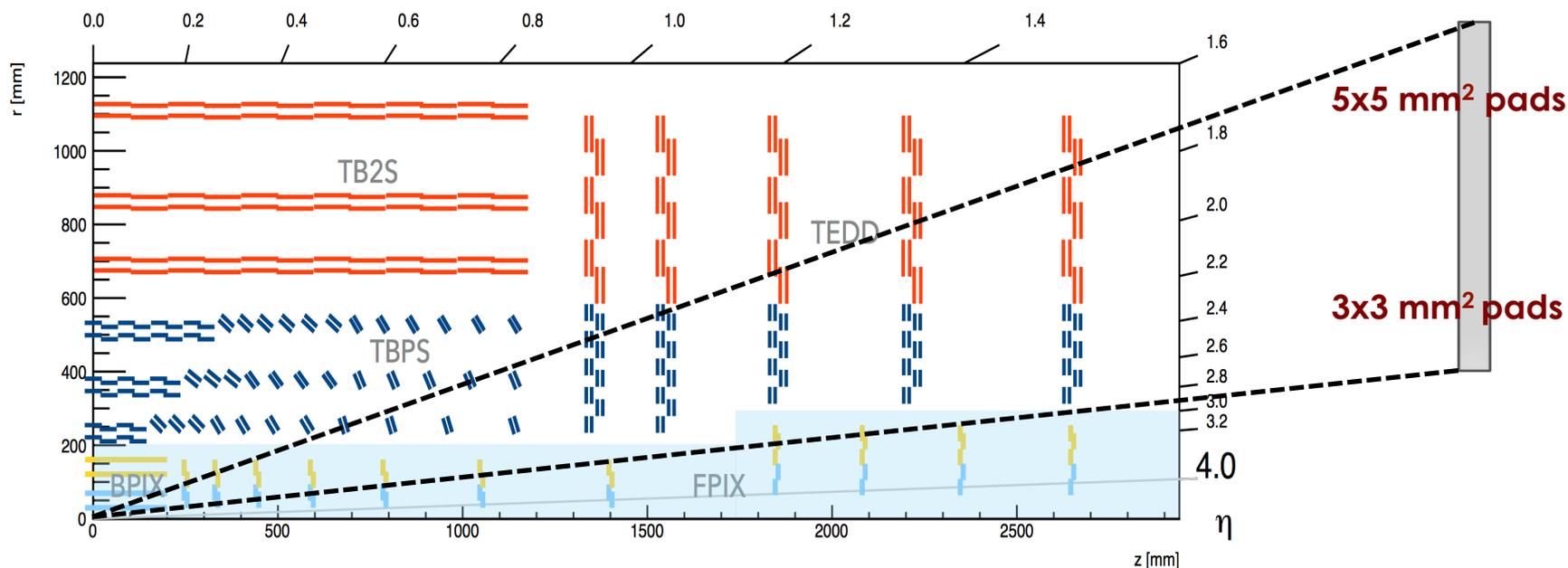


# And now we need to build the detector..

**Position?** We don't quite know, but we assume outside the tracker

**Granularity?** If we want to keep 1% occupancy, the sensor should be about  $3 \times 3 \text{ mm}^2$  in the inner part and  $5 \times 5 \text{ mm}^2$  in the out

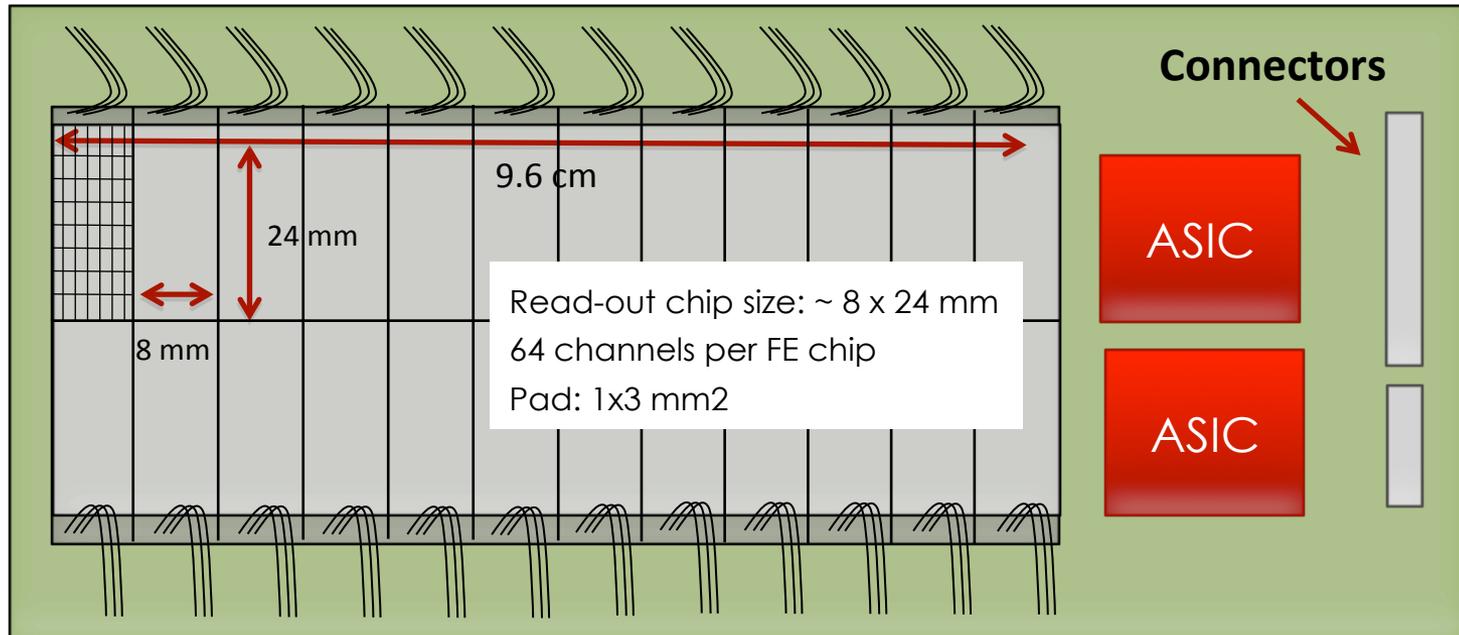
**Not in the trigger**



**And: mechanics, cooling, HV & LV distribution, High precision clock, data transmission**

# CMS Timing Module for the endcap

Use the geometry 5x10 cm<sup>2</sup> as baseline



**From physics:** what is the granularity? 1x1 mm<sup>2</sup>, 1x2 mm<sup>2</sup>, 1x3 mm<sup>2</sup>.  
Shall we use hexagon to minimize the perimeter?

30  $\mu$ m  
dead area



**Dead area:**

**12%**

**9%**

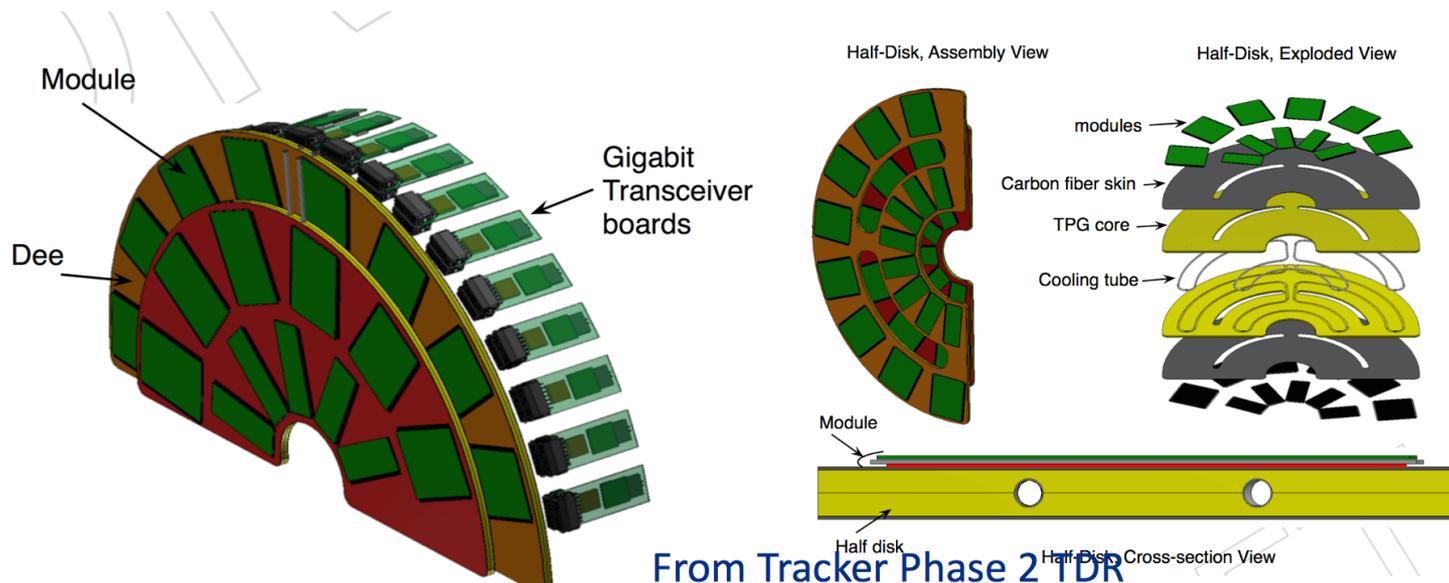
**8%**

**9%**

**6%**

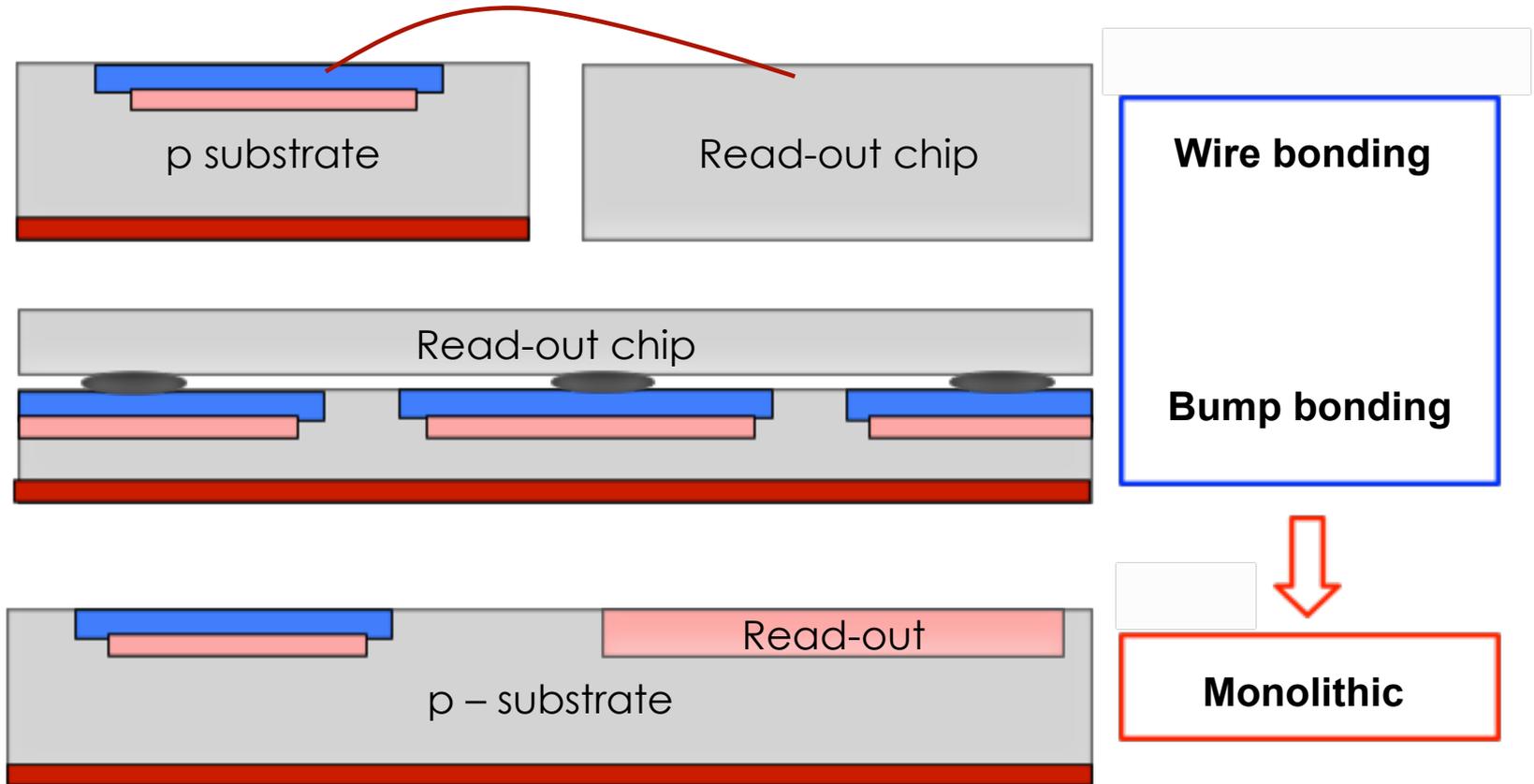
# Endcap disks

- ETL support structure similar to pixel upgrade
  - Modules are mounted on planar half-disc structures
  - Modules placed on both sides of the dees to provide hermetic coverage
- CO<sub>2</sub> flows in embedded tubes under the ROCs
  - Sensors and ROC running at -25°C
- Many common aspects with the tracker: a lot of expertise and experience to borrow
  - Manual jig-based assembly technique, will consider robotic assembly as well
- The routing of services, cooling, and support designed by FNAL engineer



From Tracker Phase 2 TDR

# Can we use Monolithic technology?



# Summary

---

4Dimensional timing can be achieved using special silicon sensors with internal gain

In the proposed Timing Detector, CMS foresees to use crystals in the barrel and UFSD sensors in the endcap

UFSD sensors are currently the only technology available able to maintain good gain up to  $\sim 10^{15} n_{eq}/\text{cm}^2$

CMS timing layer will be discussed **TODAY** at the CMS management board, and then, eventually, by the full collaboration during the first week of April.

# References

---

1. R. Mulargia et al, "Temperature dependence of the response of Ultra Fast Silicon Detectors" Pixel 2016, to be published in JINST 11 (2016) C12016
2. N. Cartiglia et al, "The 4D pixel challenge", JINST 11 (2016) C12016
3. N. Cartiglia et al, "Beam test results of a 16 ps timing system based on ultra-fast silicon detectors", NIMA 850 (2017) 83 – 88
4. B. Baldassarri et al, Signal formation in irradiated silicon detectors, NIMA 845 (2017) 20 – 23
5. N. Cartiglia et al, Tracking in 4 Dimensions, NIMA 845 (2017) 47 – 51
6. N. Cartiglia et al, Design Optimization of Ultra-Fast Silicon Detector, NIMA 796 (2015) 141-148, doi:10.1016/j.nima.2015.04.025i
7. F. Cenna et al, Weightfield2: A fast simulator for silicon and diamond solid state detector, NIMA 796 (2015) 149-153, doi:10.1016/j.nima.2015.04.015
8. G-F Dalla Betta et al, Design and TCAD simulation of double-sided pixelated low gain avalanche detectors, NIMA 796 (2015) 54-157, doi:10.1016/j.nima.2015.03.039
9. N. Cartiglia et al., Timing Capabilities of Ultra-Fast Silicon Detectors, doi:10.5506/APhysPolBSupp.7.657
10. N. Cartiglia, et al., Performance of Ultra-Fast Silicon Detectors, JINST 9 (2014) C02001
11. H.-W. Sadrozinski et al., Sensors for ultra-fast silicon detectors, NIMA 765 (2014) 7
12. H.-W. Sadrozinski, et al., Ultra-fast silicon detectors, NIMA 730 (2013) 226-231

# Acknowledgement

---

This research was carried out with the contribution of the Ministero degli Affari Esteri, “Direzione Generale per la Promozione del Sistema Paese” of Italy.



*Ministero degli Affari Esteri  
e della Cooperazione Internazionale*

DIREZIONE GENERALE  
PER LA PROMOZIONE DEL SISTEMA PAESE  
*Unità per la cooperazione scientifica  
e tecnologica bilaterale e multilaterale*

This work is currently supported by INFN Gruppo V, UFSD project (Torino, Trento Univ., FBK).

This work was developed in the framework of the CERN RD50  
The work is supported by HORIZON2020 Grants and UFSD ERC grant  
UFSD669529

# Backup

---

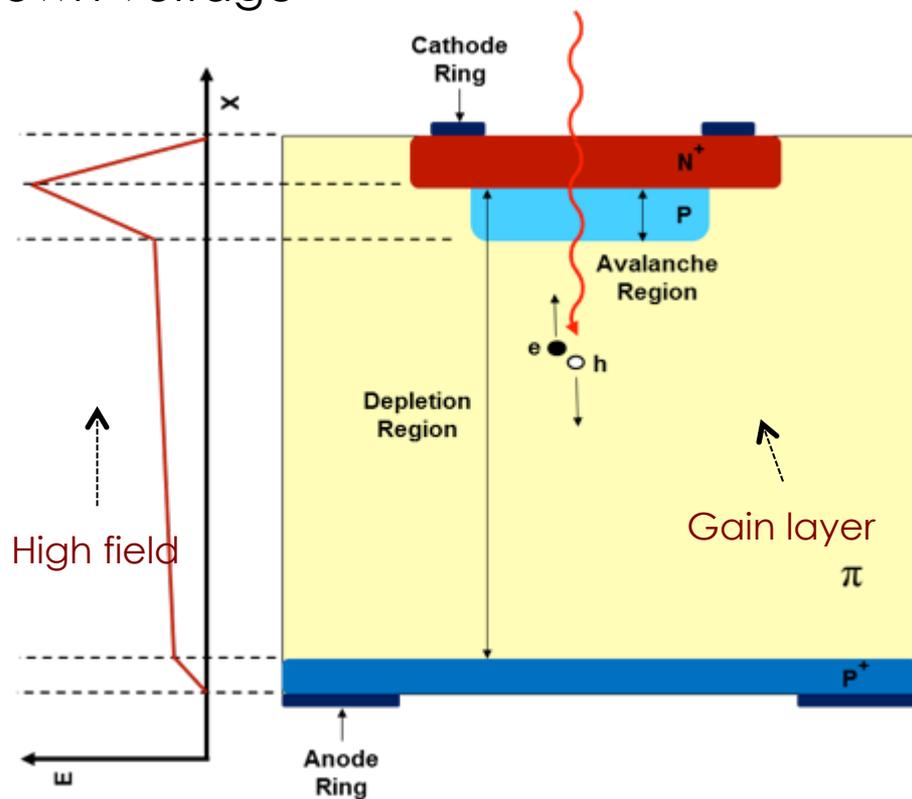
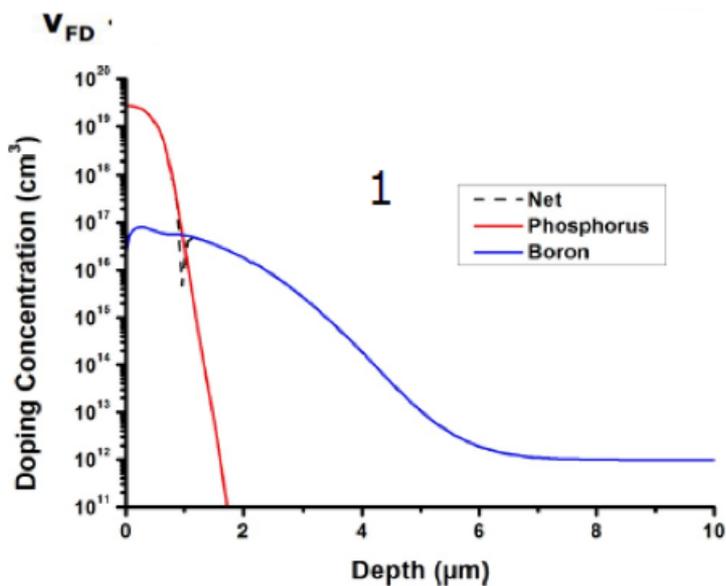
# Low Gain Avalanche Detectors (LGADs)

The LGAD sensors, as proposed and manufactured by CNM

(National Center for Micro-electronics, Barcelona):

**High field obtained by adding an extra doping layer**

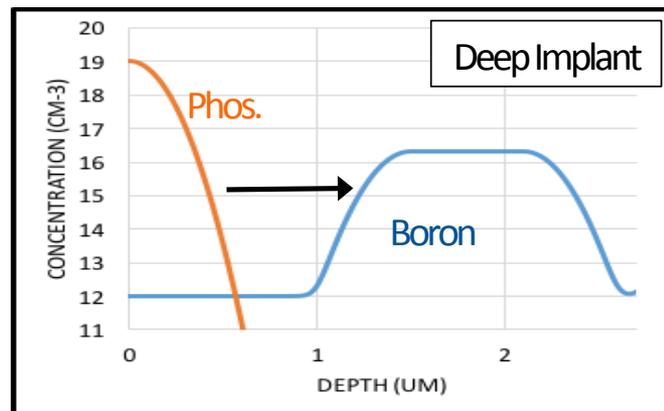
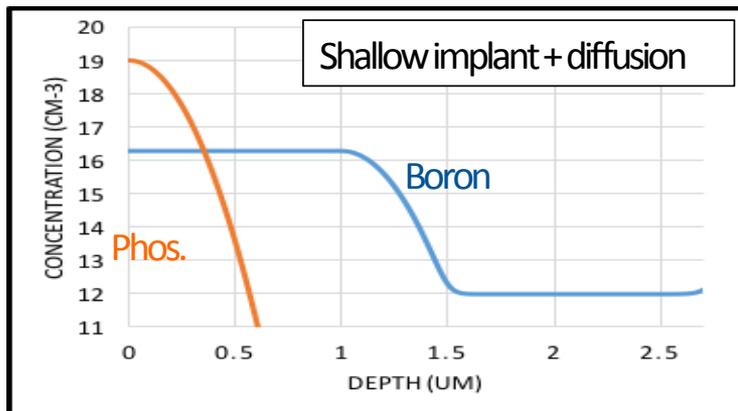
$E \sim 300 \text{ kV/cm}$ , closed to breakdown voltage



# Gain layer design

The doping profile of the Gain layer controls the shape of the Electric Field

2 technological approaches are possible:

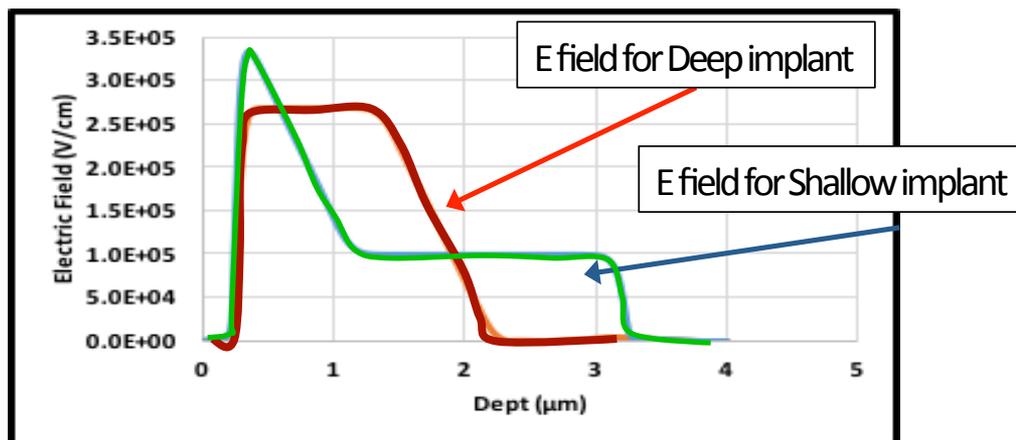


The deep implant approach has several advantages:

- Avoid peaked Electric Field -> less noise
- Is more reliable (independent of thermal diffusion and of doping compensation effect)

CMM

CENTRE FOR MATERIALS AND MICROSYSTEMS



# When do we decide?

CMS plans to reach a decision on whether or not to proceed with the R&D phase by LHCC in May

... circumstances, technical readiness, timeline and cost

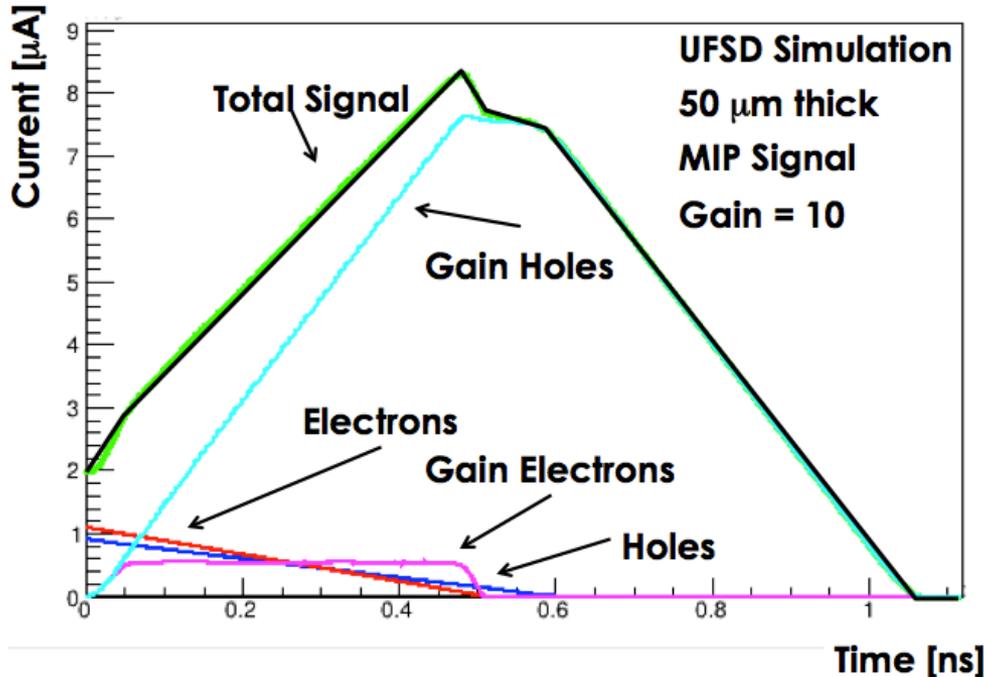
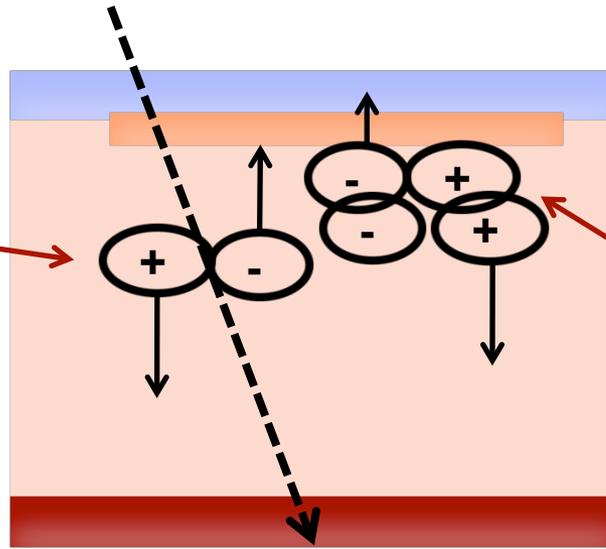
JAN				
	<b>2. Present an updated document at the CMS Week</b>	<b>Jan 30<sup>th</sup></b>		
	<ul style="list-style-type: none"><li>From Jan 30th to LHCC to finalize content → can report at WGM meetings</li></ul>			
	<b>3. Give an information talk to the LHCC</b>	<b>Feb 21<sup>st</sup></b>		
	<b>4. Review report at WGM</b>	<b>Mar 16<sup>th</sup></b>		
	→ End of the review process			
	<b>5. CMS decision at CMS week</b>	<b>Apr 3<sup>rd</sup></b>		
	<b>6. Formally present CMS position to LHCC</b>	<b>May 9<sup>th</sup></b>		
MAY	APR	MAR	FEB	JAN

# How gain shapes the signal

**Gain electron:**  
absorbed immediately

**Gain holes:**  
long drift home

Initial electron, holes



Electrons multiply and produce additional electrons and holes.

- **Gain electrons have almost no effect**
- **Gain holes dominate the signal**

➔ **No holes multiplications**

# WeightField2: a program to simulate silicon detectors

Drift Potential | Weighting Potential | Currents and Oscilloscope | Electronics
Weightfield 2.6

Done.

Plotting at: On Strips Between Strips 465 Draw

Drift Potential V [V]

Drift Field E (kV/cm)

Control

Precision (1=best, 10=fastest): 10

Sampling (GigaSample): 100

File Name  
 ON wf

Batch  
 ON # of events: 1

Select Particles

MIP: uniform Q, Qtot = 75\*Height

MIP: non uniform Q, Qtot = 75\*Height

MIP: non uniform, Qtot = Landau

MIP: uniform Q, Q/micron = 75

alpha from top (E = 5 MeV)

alpha from bottom (E = 5 MeV)

Set range (Max = 30 um): 10

Plot Settings

Draw Electric Field

No 1D Plots  No 1D & 2D

Currents

Switch B-Field on and set to (T): 0

Diffusion

Temperature (K): 300

Set
Calculate Potentials

Calculate Currents
Stop
Exit

Detector Properties

Type  
 Si  Diamond  Free

Strips  
 n-type  p-type

Bulk  
 n-type  p-type

Dimensions

# of strips (1,3,5,...): 3

Detector Height (um): 285

Strip Pitch (um): 300

Strip Width (um): 290

Gain Scale (1 = no G): 1

Force Fixed Gain:  ON

h/e Gain ratio: 0

Gain layer recess (um): 0

Voltage

Bias Voltage (V): 800

Depletion Voltage (V): 40

Electronics

ON

Detector Cap (pF): 1

Oscilloscope BW (GHz): 2.5

Shaper T<sub>r</sub> - T<sub>f</sub> (ns): 3.5 8

Shaper Trans Imp. (mV/IQ): 4

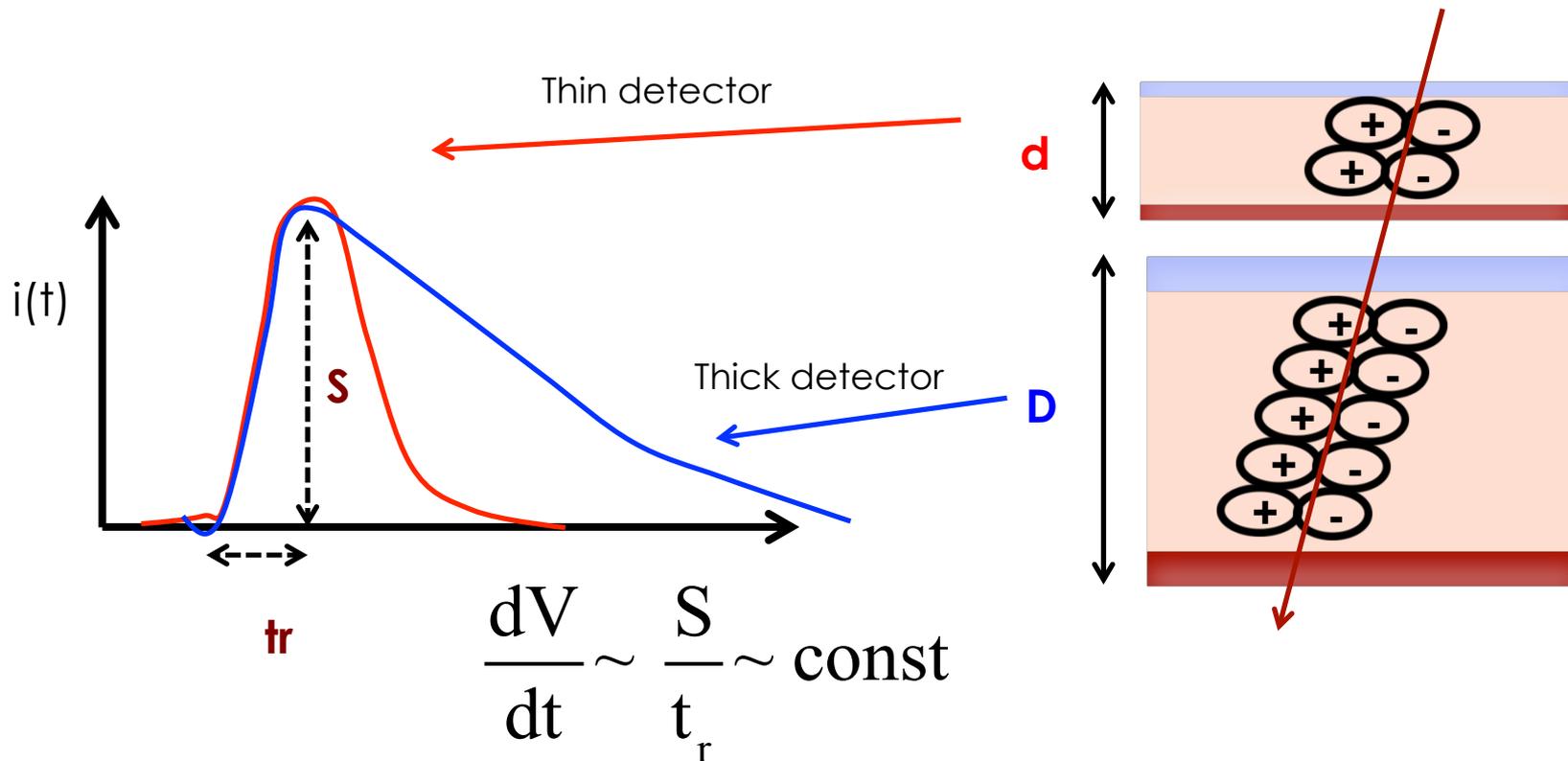
Shaper Noise & V<sub>th</sub> (mV): 1 10

PreAmp input Imp. (Ohm): 50

64

# Thin vs Thick detectors

(Simplified model for pad detectors)



Thick detectors have longer signals, not higher signals

Best result : NA62, 150 ps on a 300 x 300 micron pixels

**How can we do better?**

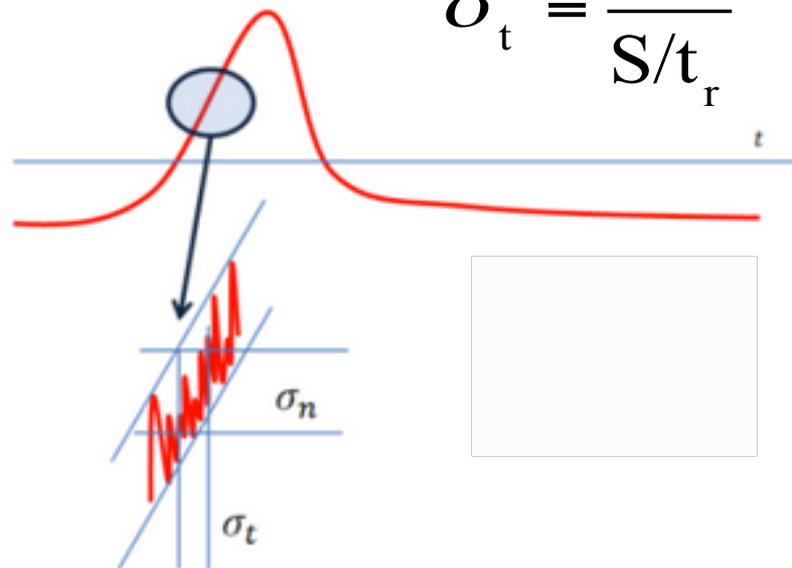
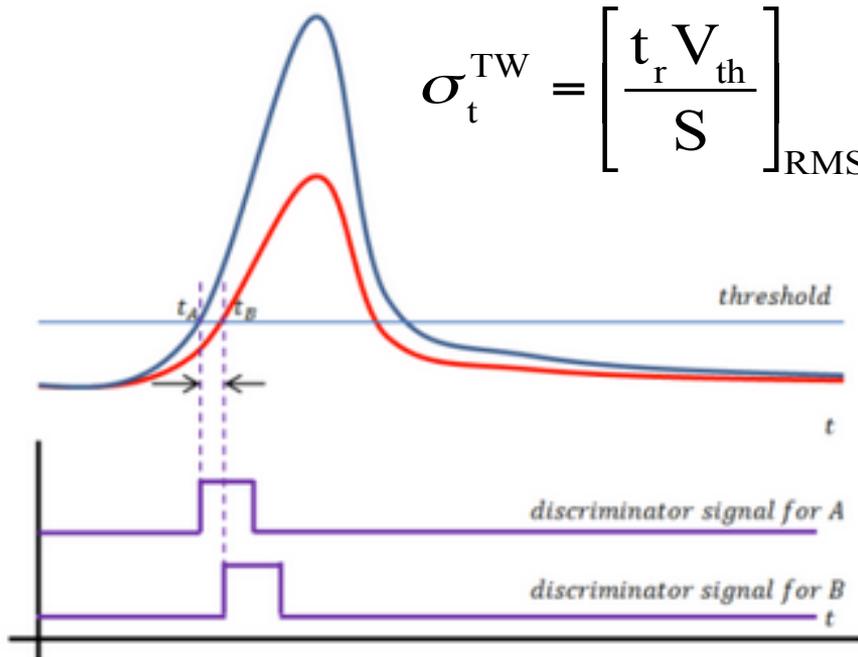
# 2 important effects: Time walk and Time jitter

**Time walk:** the voltage value  $V_{th}$  is reached at different times by signals of different amplitude

**Jitter:** the noise is summed to the signal, causing amplitude variations

$$\sigma_t^{TW} = \left[ \frac{t_r V_{th}}{S} \right]_{RMS}$$

$$\sigma_t^J = \frac{N}{S/t_r}$$



Due to the physics of signal formation

Mostly due to electronic noise

Time walk and jitter  $\sim N / (S/t_r) = N / (dV/dt)$