

Mergers, Gamma-Ray Bursts and Gold

Tsvi Piran

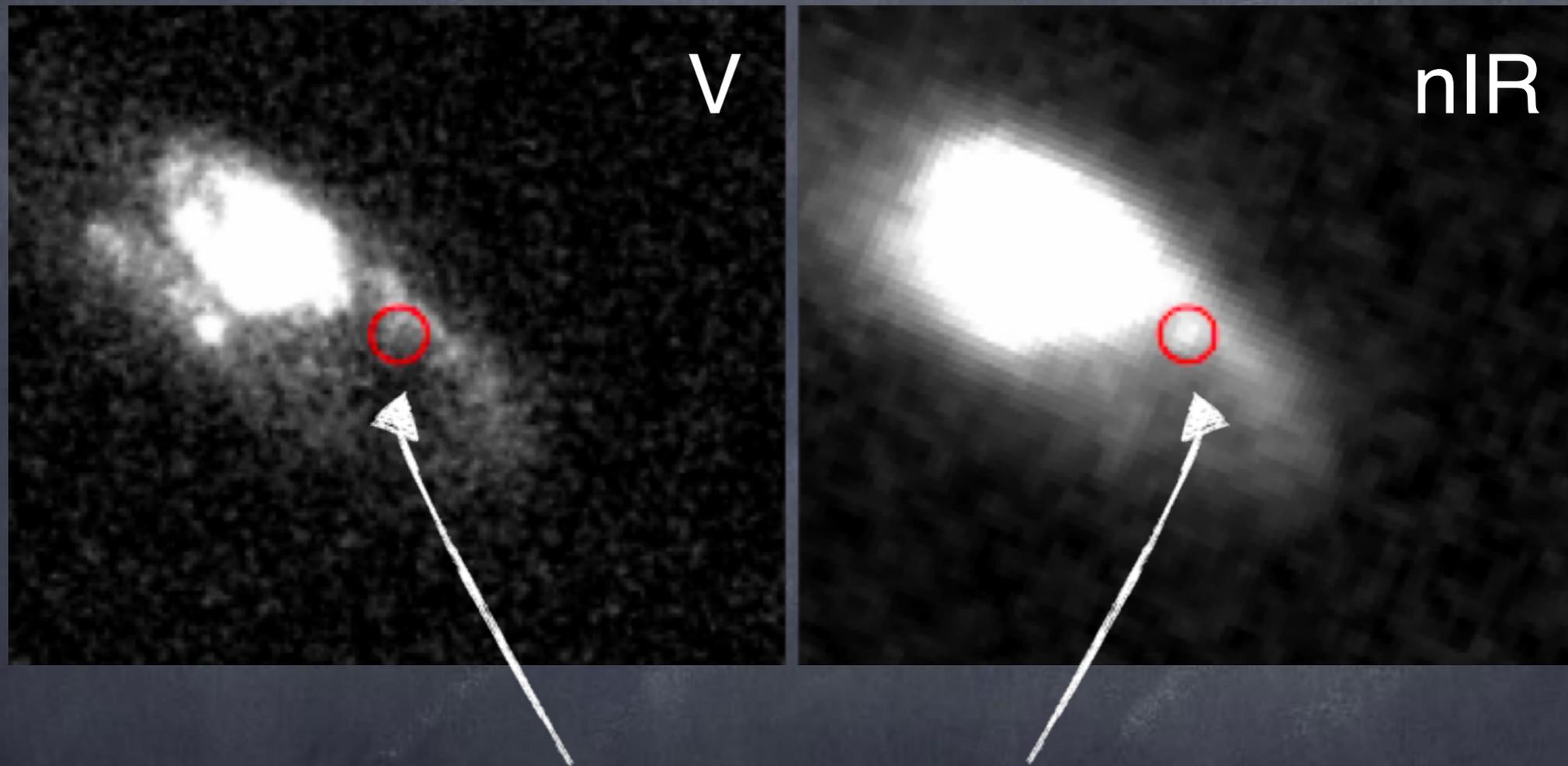
The Hebrew University

**Doron Grossmam, Kenta Hotokezaka,
Paz Beniamini, Reetanjali Moharana**



The Hubble Space Telescope

June 13th 2013

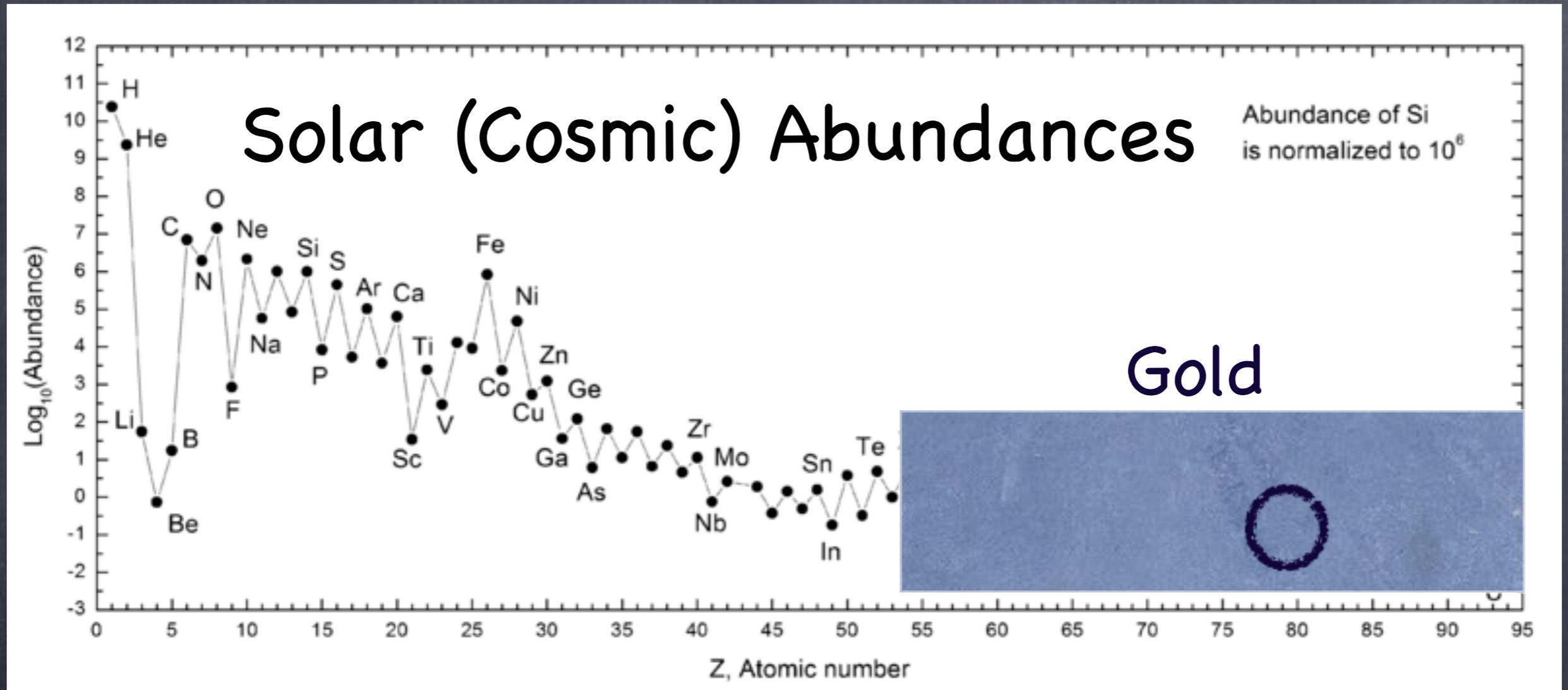


Is this the "smoking gun" proving the origin of Gold
(and other heavy elements) in the Universe?

Outline

1. Nucleosynthesis 101
2. Neutron Stars and Mergers
3. Gamma-Ray Bursts
4. The Li-Paczynski Macronova (kilonova)
5. Putting it all together - GRB 130603B
6. Additional support
7. The origin of Gold

1. Nucleosynthesis 101



How are these elements produced?

BB (Big Bang) Nucleosynthesis

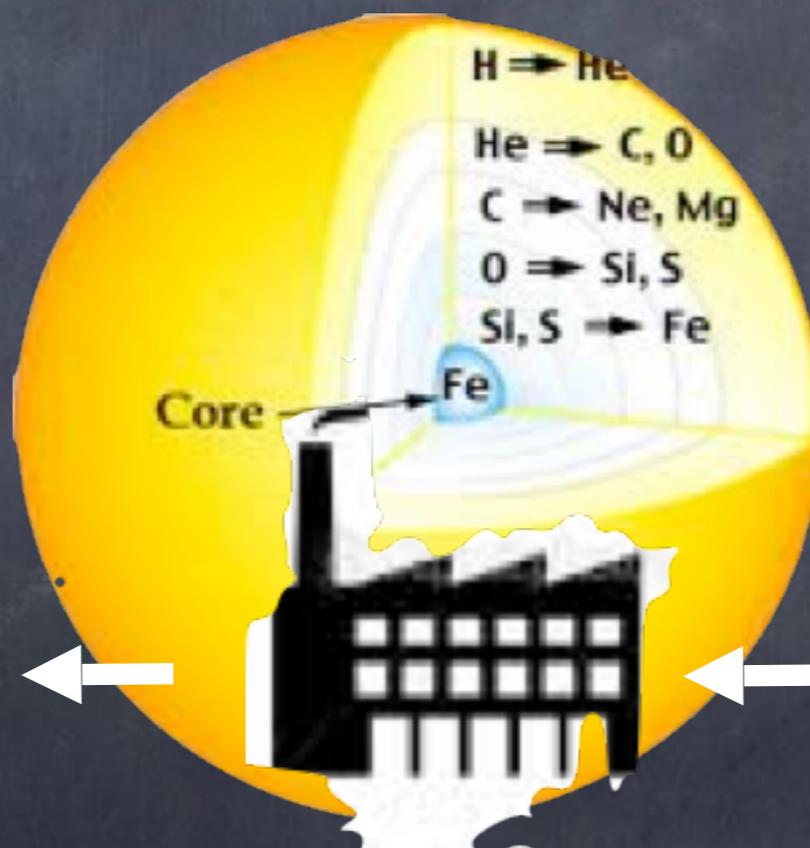
- 24% of the Universe is He.
- This He is produced in the big Bang.



George Gamow



Burbidge, Burbidge, Fowler and Hoyle
 B^2FH 1957

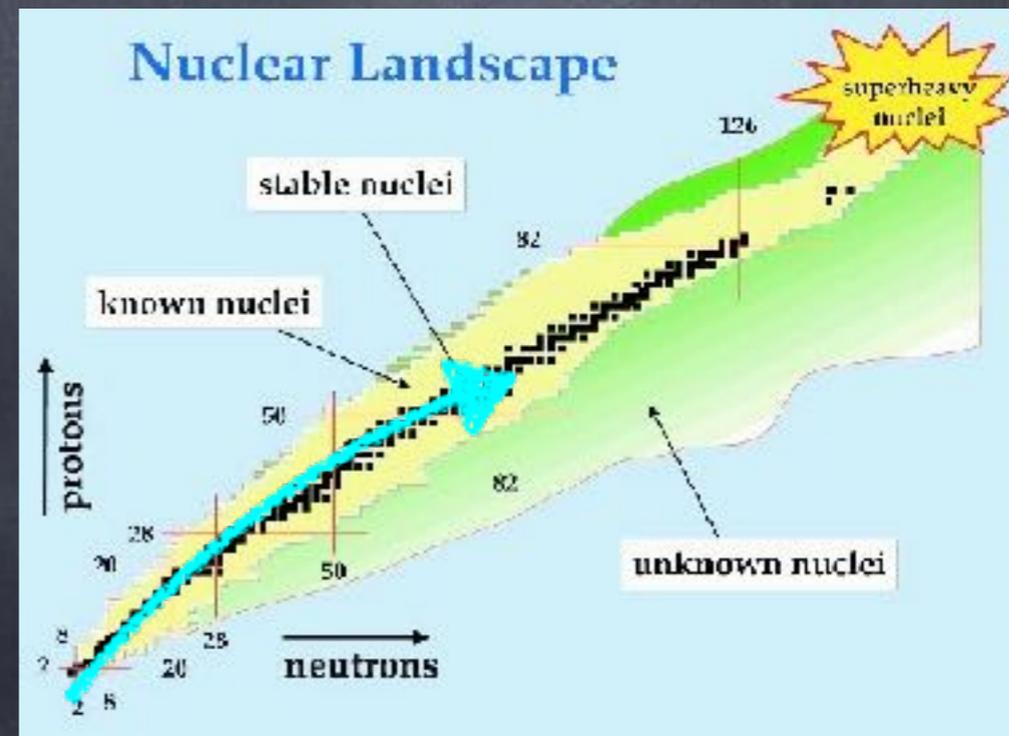
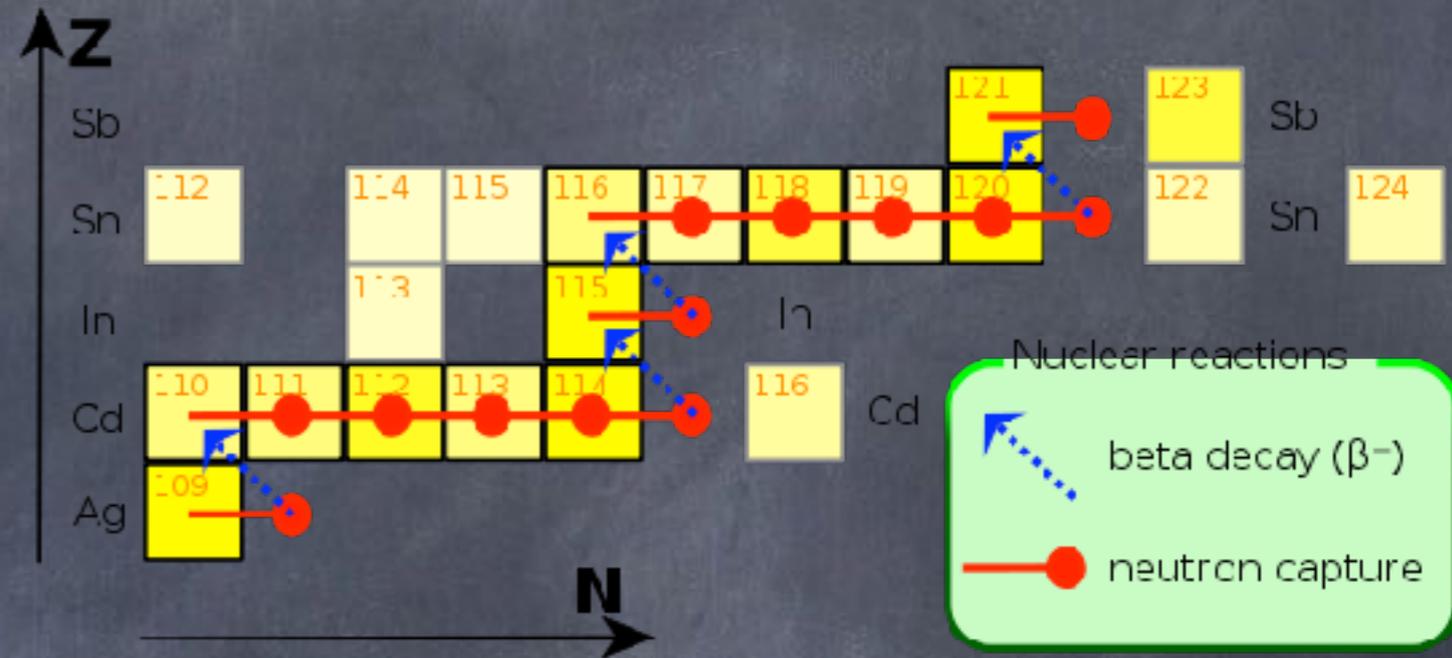


He, C, O, Ne, Mg
 Si, S, Fe, Ni....

Elements up to Iron are produced in stars

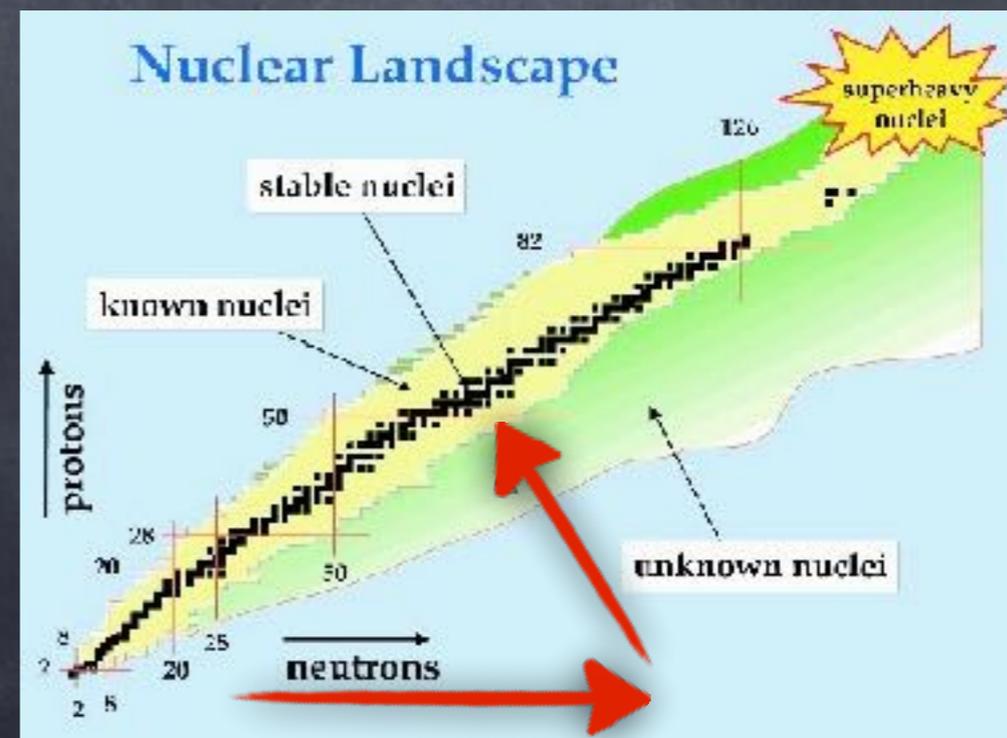
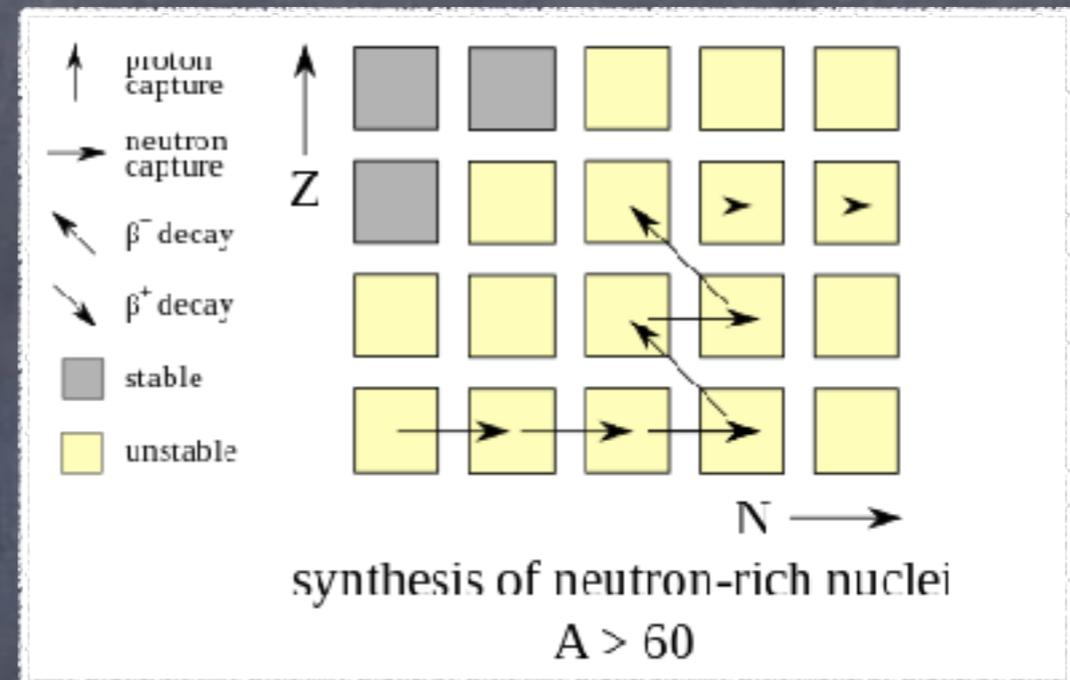
S (slow) Process

- Neutron capture slower than beta decay.
- Low neutron densities.
- time scale - years.
- Moves along the valley of nuclear stability.
- Final abundances depend on the conditions within the site.

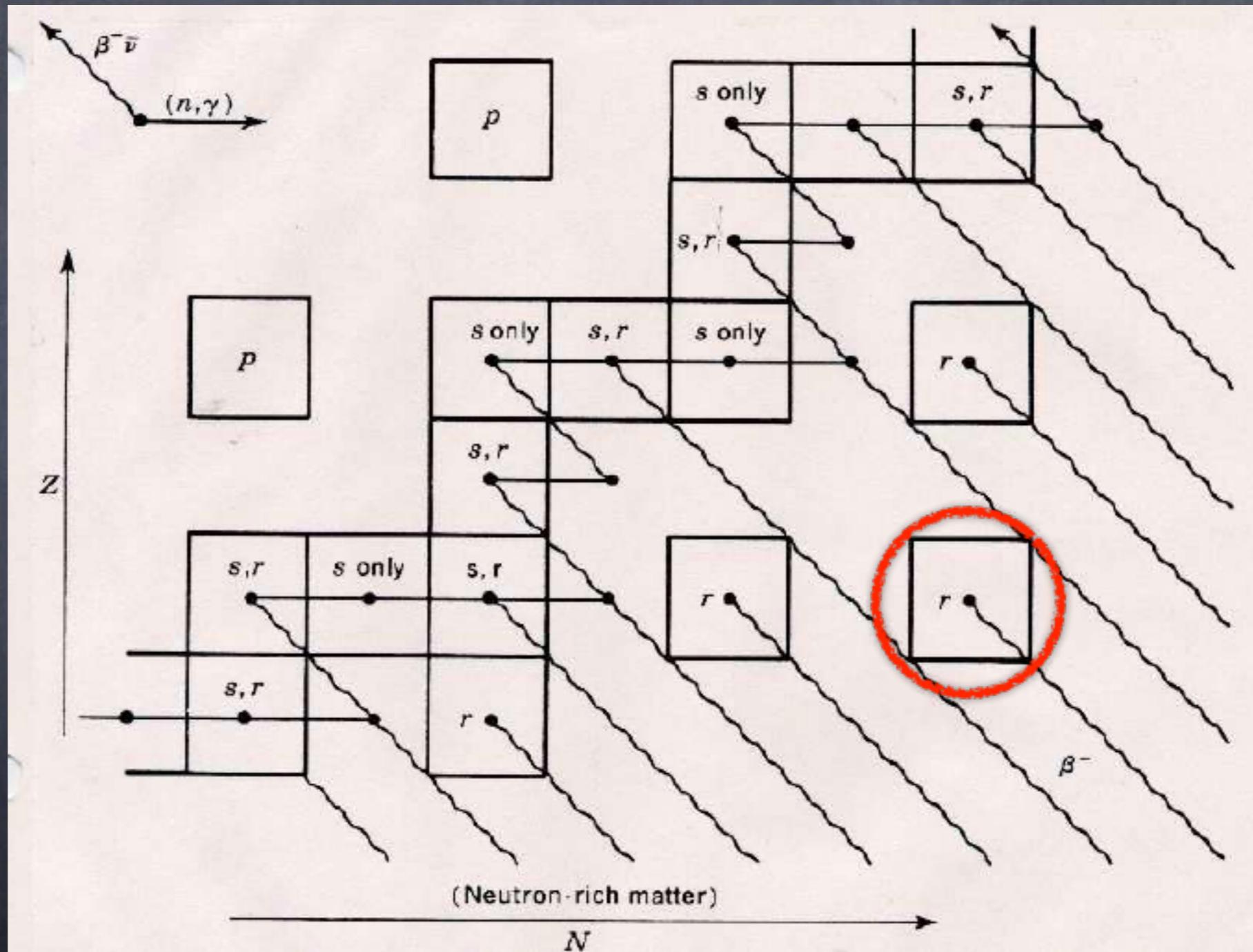


r (rapid) Process

- Neutron capture faster than beta decay.
- High neutron densities.
- Time scales – seconds.
- On the neutron rich side of nuclear stability.
- Uniform final abundances.

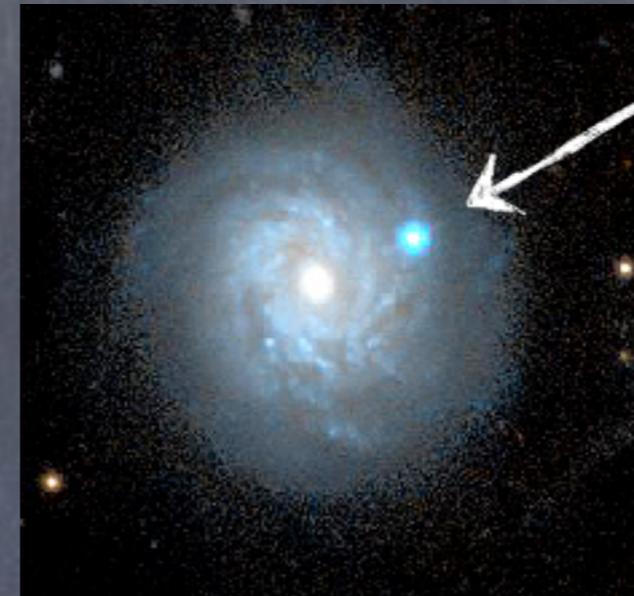


s and r processes

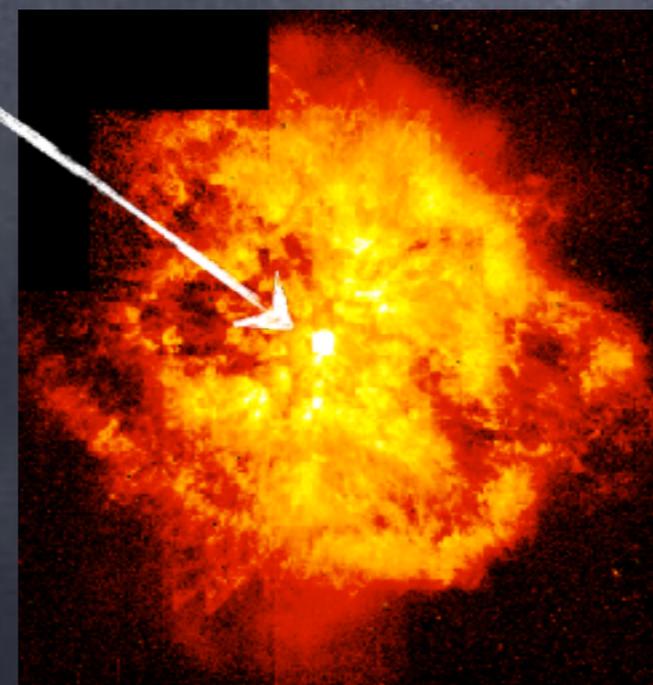


Explosive r-process

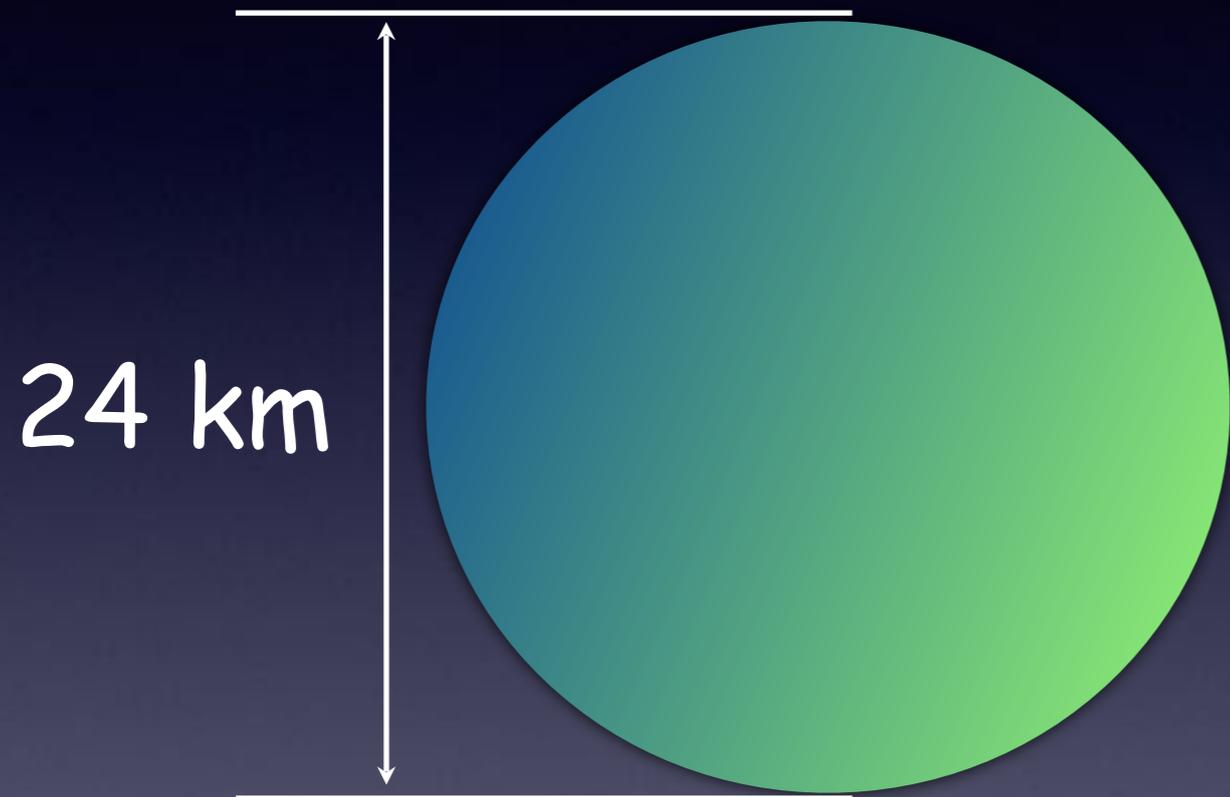
- ν flux from the newborn neutron star produce excess of neutrons in Supernova explosion.



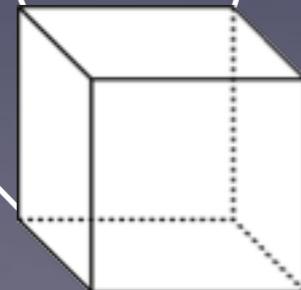
Supernova



2. Neutron stars and mergers

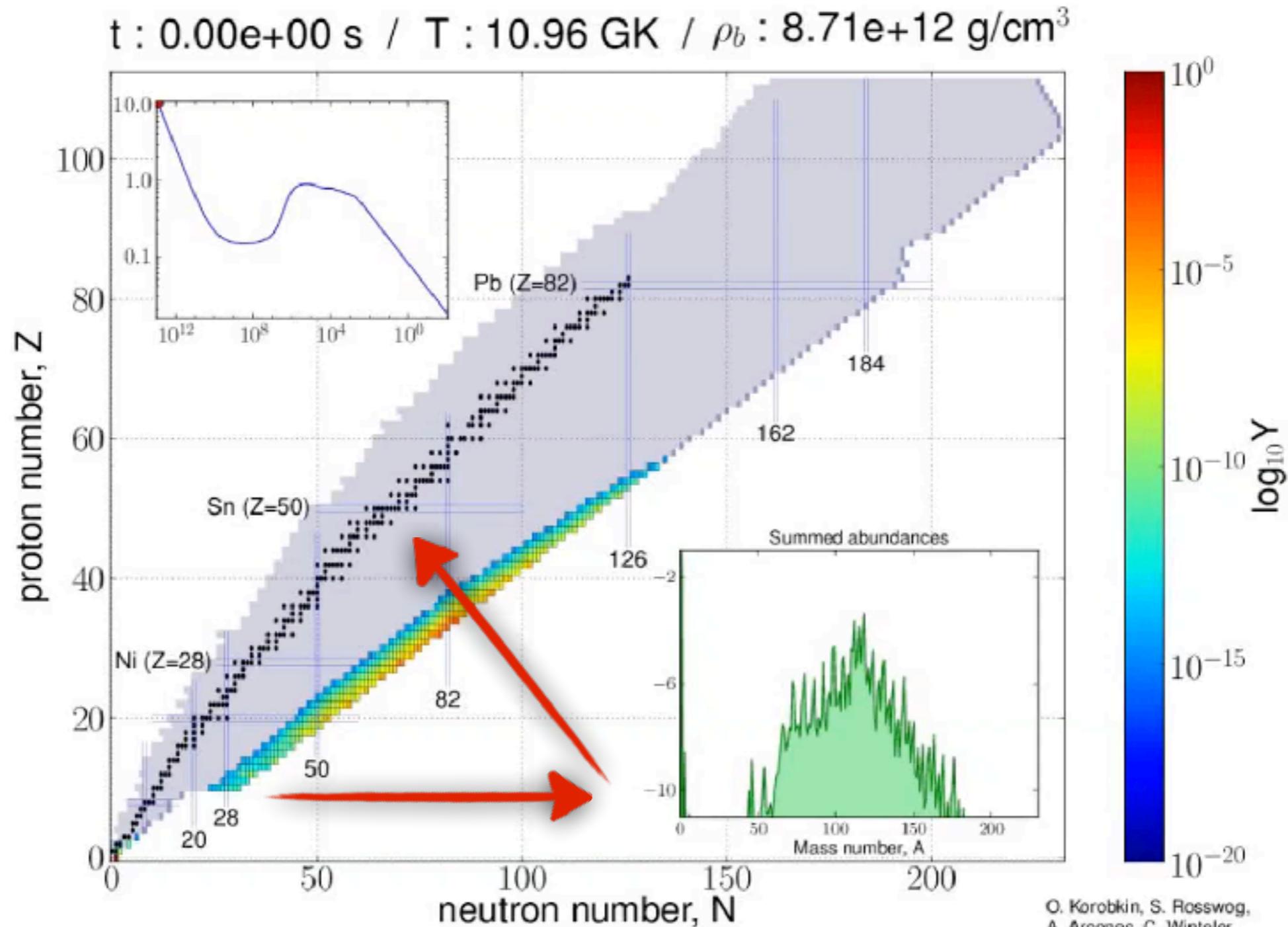


95% neutrons!





Decay of neutron star matter

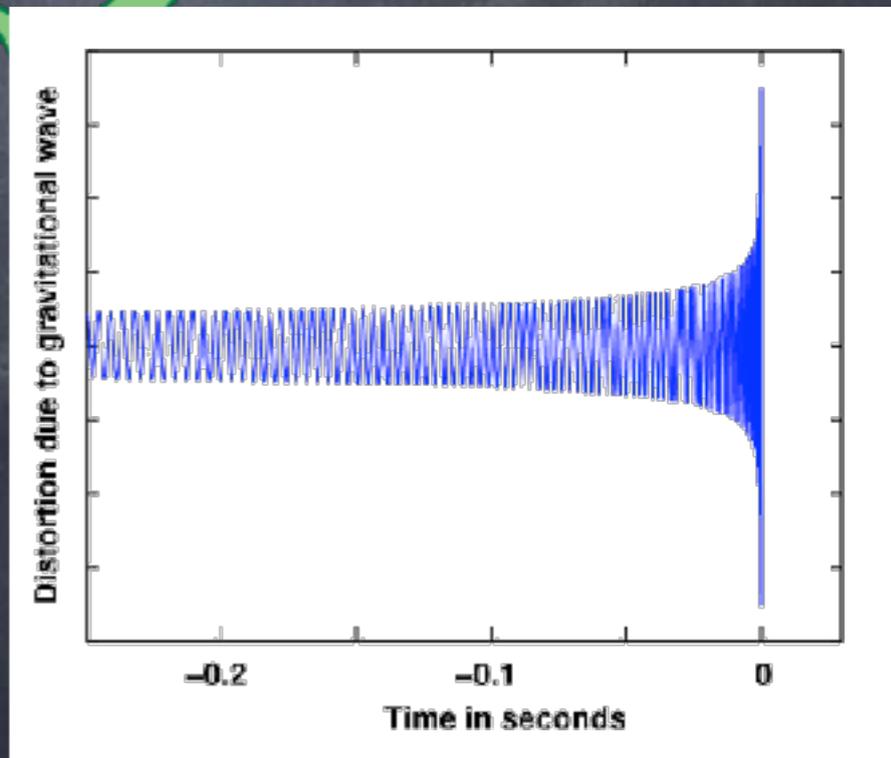


O. Korobkin, S. Rosswog,
A. Arcones, C. Winteler,
arXiv:1206.2379

Binary Neutron Stars



$$\frac{dr}{dt} = -\frac{64 G^3 (m_1 m_2)(m_1 + m_2)}{5 c^5 r^3}$$



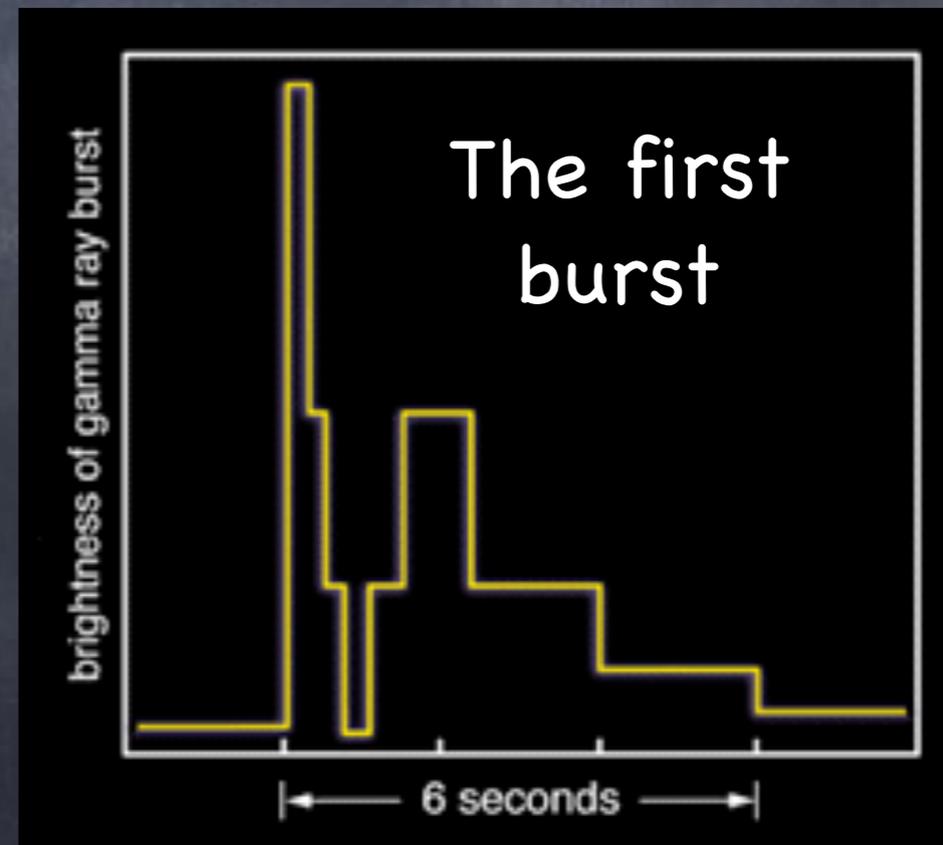
Neutron Star Mergers

3. Gamma Ray Bursts

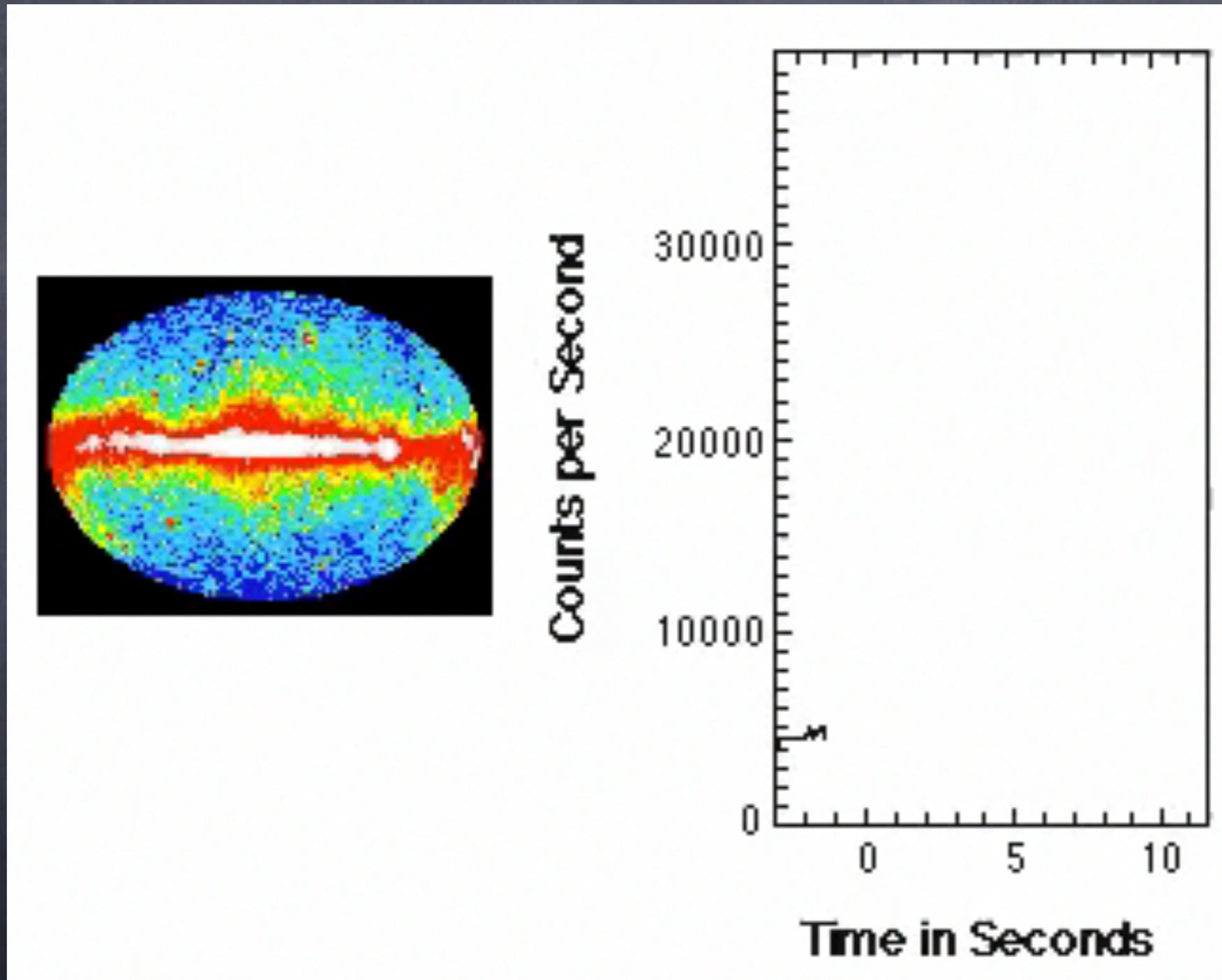


The Vela Satellites

Gamma-ray bursts



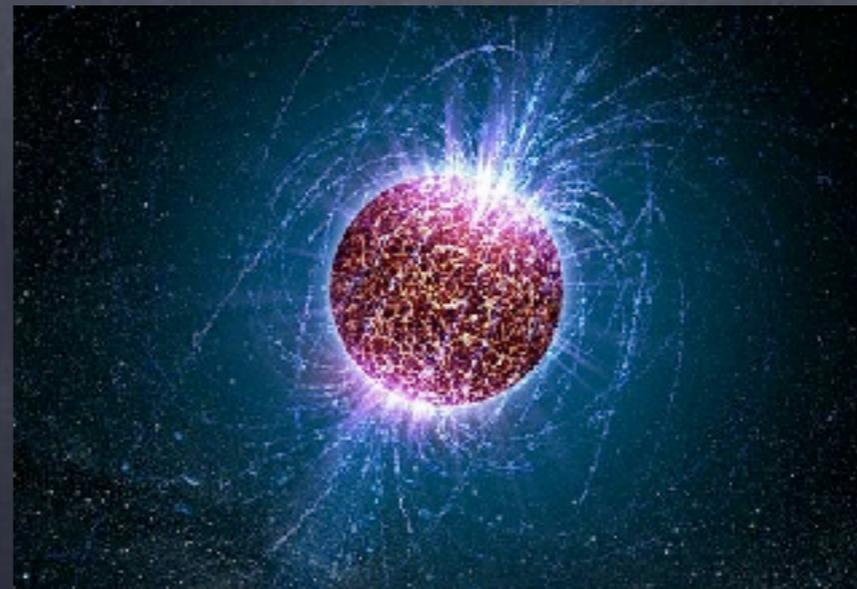
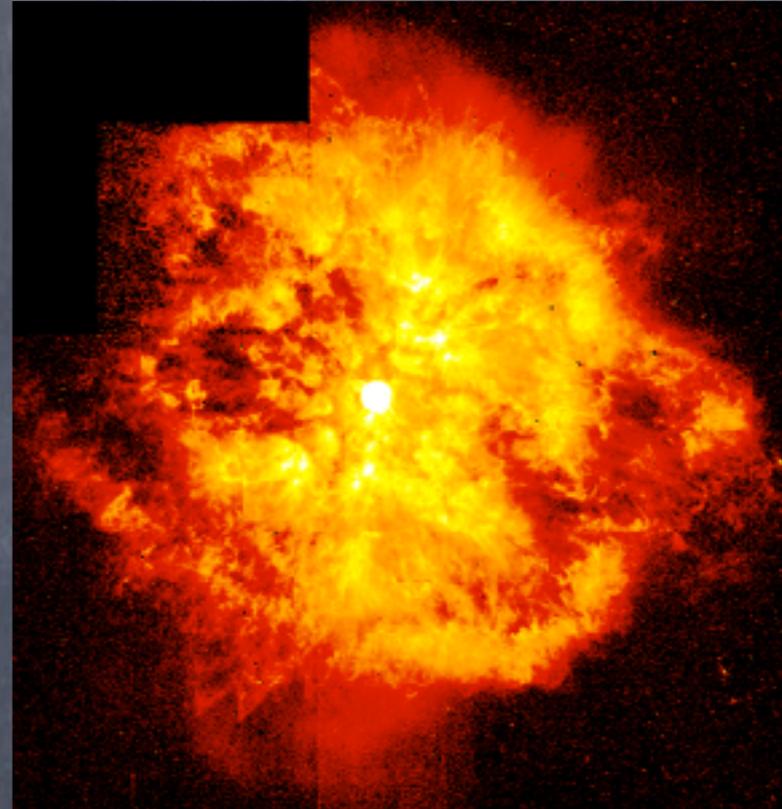
The sky in gamma-Rays



Gamma-ray bursts

The late 80ies

- r-process material from Supernovae
- GRBs from magnetic flares on galactic neutron stars ($E \sim 10^{40}$ ergs).



Two provocative ideas

LETTERS TO NATURE

Nucleosynthesis, neutrino bursts and γ -rays from coalescing neutron stars

**David Eichler*, Mario Livio†, Tsvi Piran‡
& David N. Schramm§**

NEUTRON-STAR collisions occur inevitably when binary neutron stars spiral into each other as a result of damping of gravitational radiation. Such collisions will produce a characteristic burst of gravitational radiation, which may be the most promising source of a detectable signal for proposed gravity-wave detectors¹. Such signals are sufficiently unique and robust for them to have been proposed as a means of determining the Hubble constant². However, the rate of these neutron-star collisions is highly uncertain³. Here we note that such events should also synthesize neutron-rich heavy elements, thought to be formed by rapid neutron capture (the r-process)⁴. Furthermore, these collisions should produce neutrino bursts⁵ and resultant bursts of γ -rays; the latter should comprise a subclass of observable γ -ray bursts. We argue that observed r-process abundances and γ -ray-burst rates predict rates for these collisions that are both significant and consistent with other estimates.

Gravitational Waves

Gamma-Ray Bursts

Nucleosynthesis

LETTERS TO NATURE

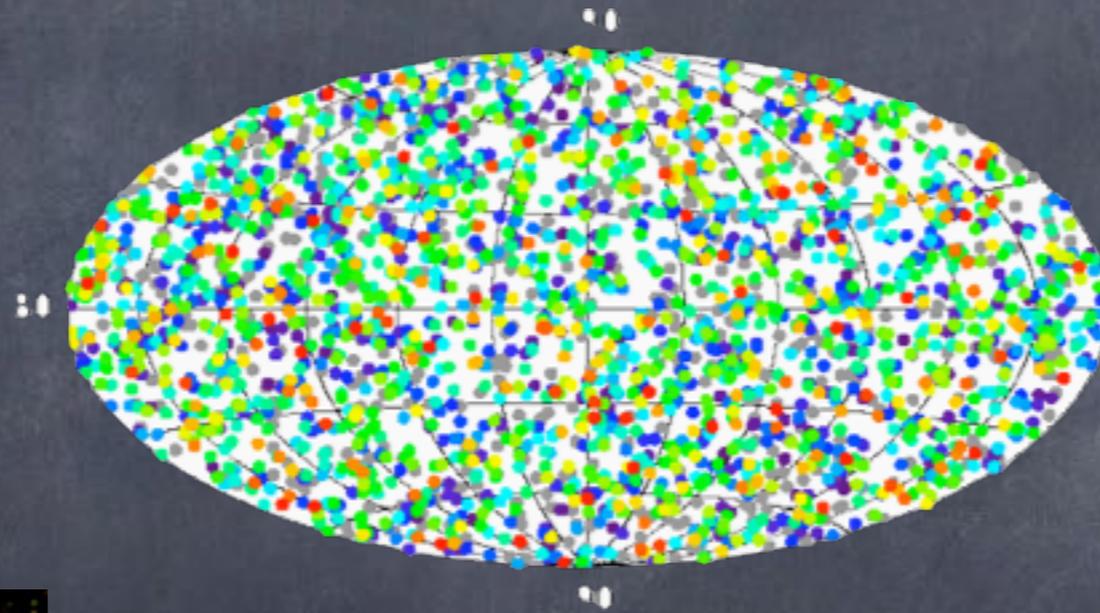
**Nucleosynthesis, neutrino bursts
and γ -rays from coalescing
neutron stars**

David Eichler*, Mario Livio†, Tsvi Piran‡
& David N. Schramm§

NEUTRON-STAR collisions occur inevitably when binary neutron stars spiral into each other as a result of damping of gravitational radiation. Such collisions will produce a characteristic burst of gravitational radiation, which may be the most promising source of a detectable signal for proposed gravity-wave detectors¹. Such signals are sufficiently unique and robust for them to have been proposed as a means of determining the Hubble constant². However, the rate of these neutron-star collisions is highly uncertain³. Here we note that such events should also synthesize neutron-rich heavy elements, thought to be formed by rapid neutron capture (the *r*-process)⁴. Furthermore, these collisions should produce neutrino bursts⁵ and resultant bursts of γ -rays; the latter should comprise a subclass of observable γ -ray bursts. We argue that observed *r*-process abundances and γ -ray-burst rates predict rates for these collisions that are both significant and consistent with other estimates.

90ies: GRBs are cosmological

1992: BATSE – GRBs have a cosmological distribution



1997: BeppoSAX – GRBs' afterglow that enables redshift measurements confirming the cosmological origin

Gamma-Ray Bursts

1988

- ~~r-process from Supernovae~~

- ~~GRBs from magnetic flares in galactic neutron stars (E $\sim 10^{40}$ ergs).~~

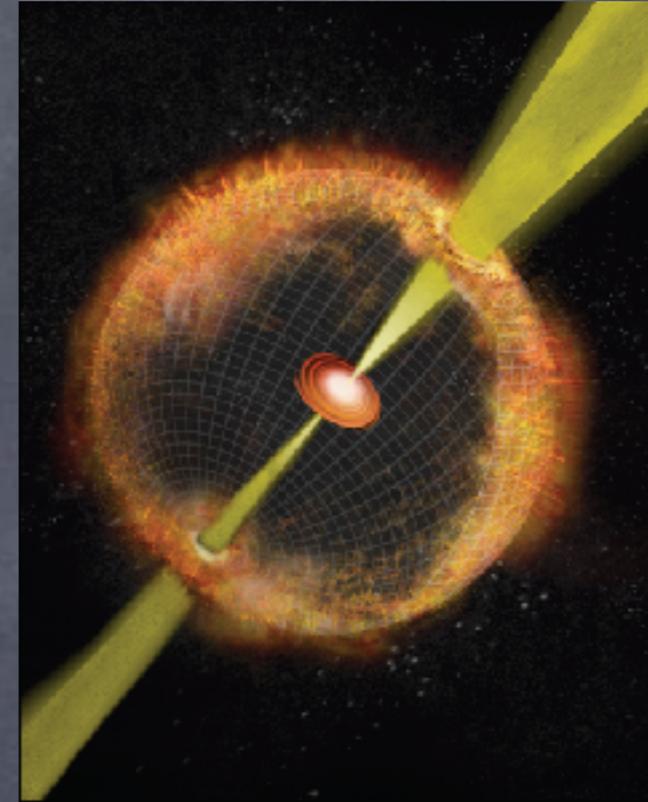
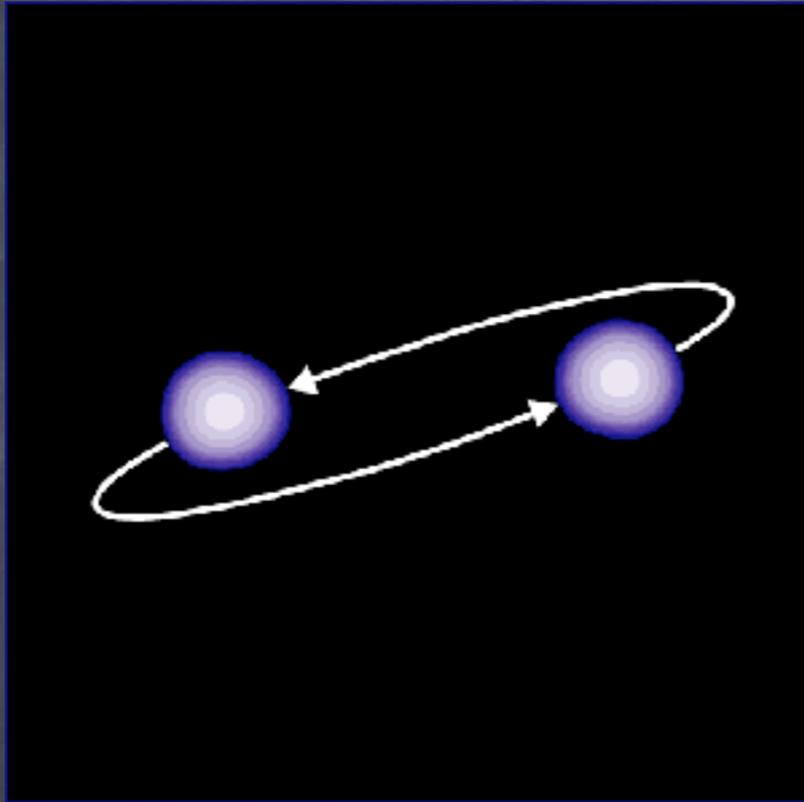
2015

- Supernovae cannot produce $A > 130$

- GRBs are cosmological (E $\sim 10^{51}$ ergs).

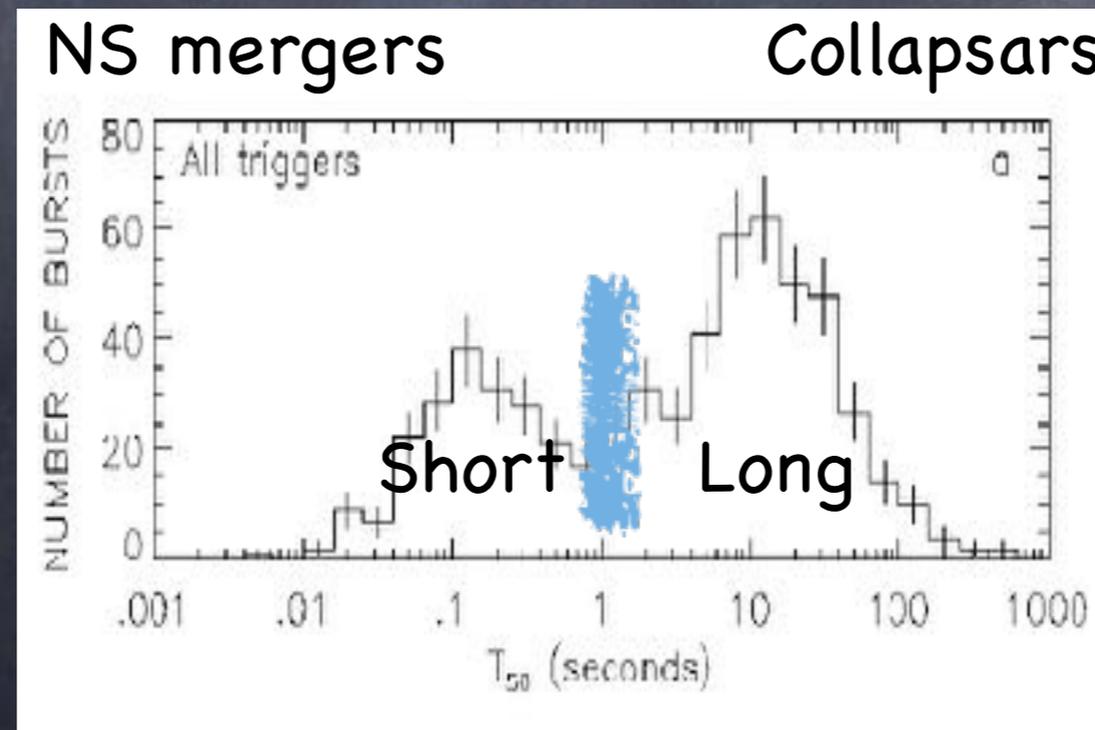
Eichler, Livio, TP,
Schramm, 88

MacFadyen & Woosley,
98



Indirect
Evidence

Direct
Evidence



Gamma-ray bursts

Mergers ejects $0.01-0.04 M_{\text{sun}}$

with $E_k \sim 10^{50}-10^{51}$ ergs



4. Macronova* (Li & Paczynski 1997)

- Radioactive decay of the neutron rich matter.
- $E_{\text{radioactive}} \approx 0.001 Mc^2 \approx 10^{50}$ erg
- A weak short Supernova like event.
- Macronovae follow short GRBs but could appear without a short GRB as those are beamed.



Bohdan Paczynski

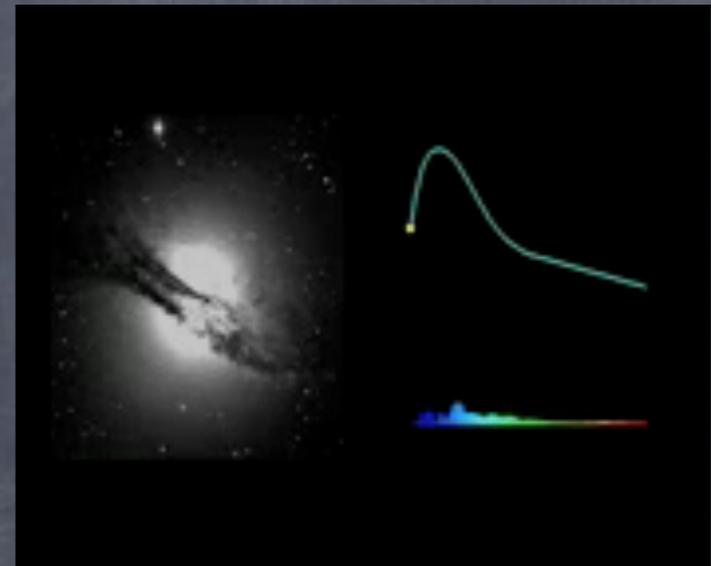
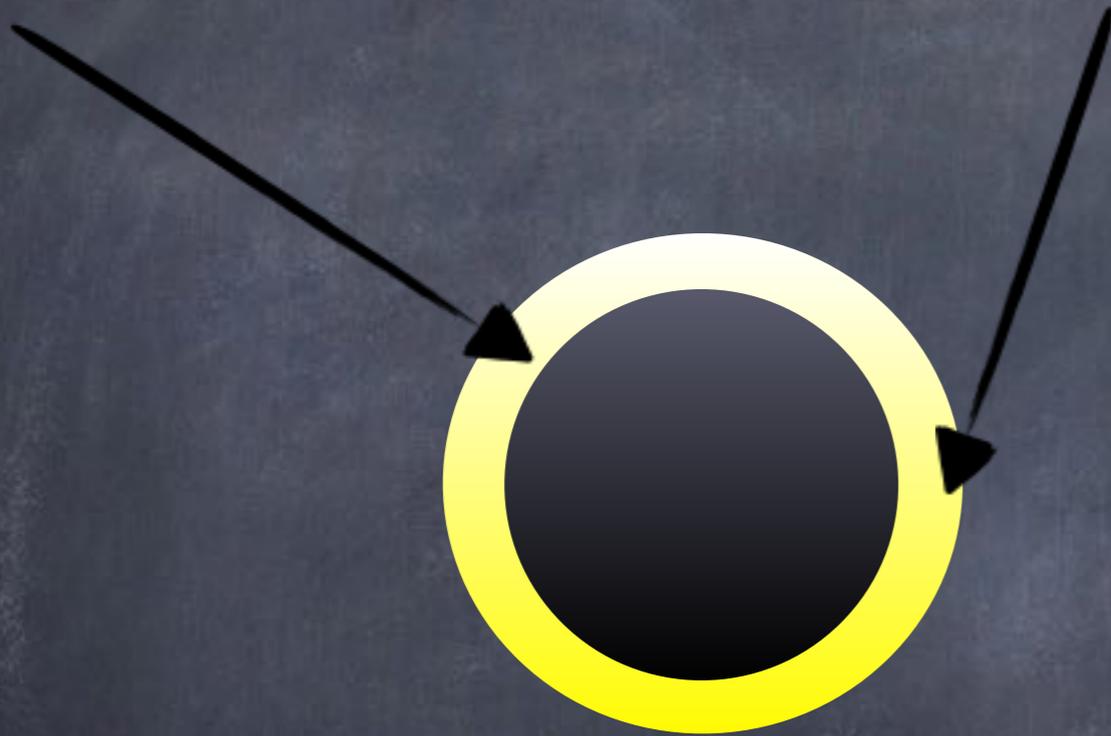


*Also called Kilonova

Supernova

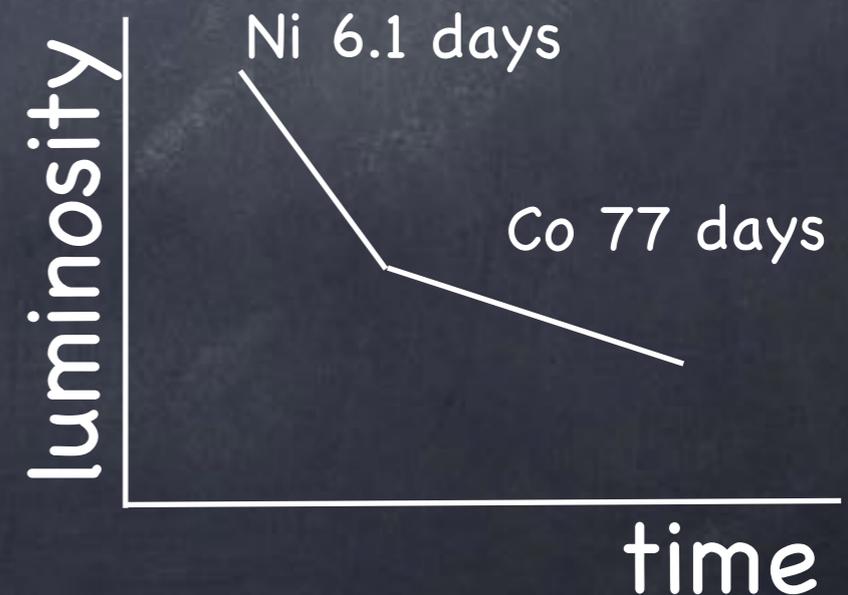
Photosphere

Photons escape



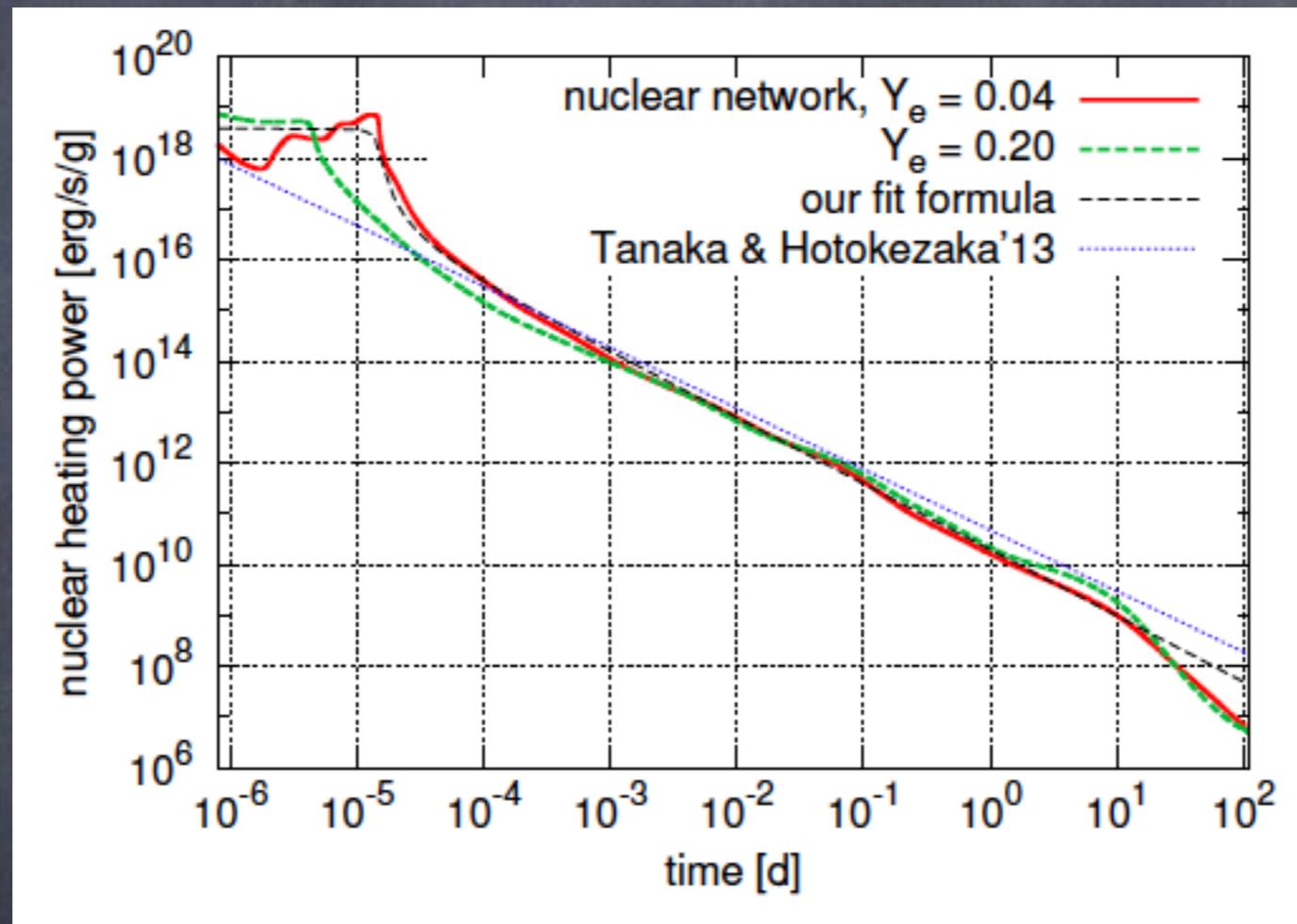
Powered by radioactive
decay of $^{56}\text{Ni} \rightarrow ^{56}\text{Co} \rightarrow ^{56}\text{Fe}$

Macronova



Radioactive Decay

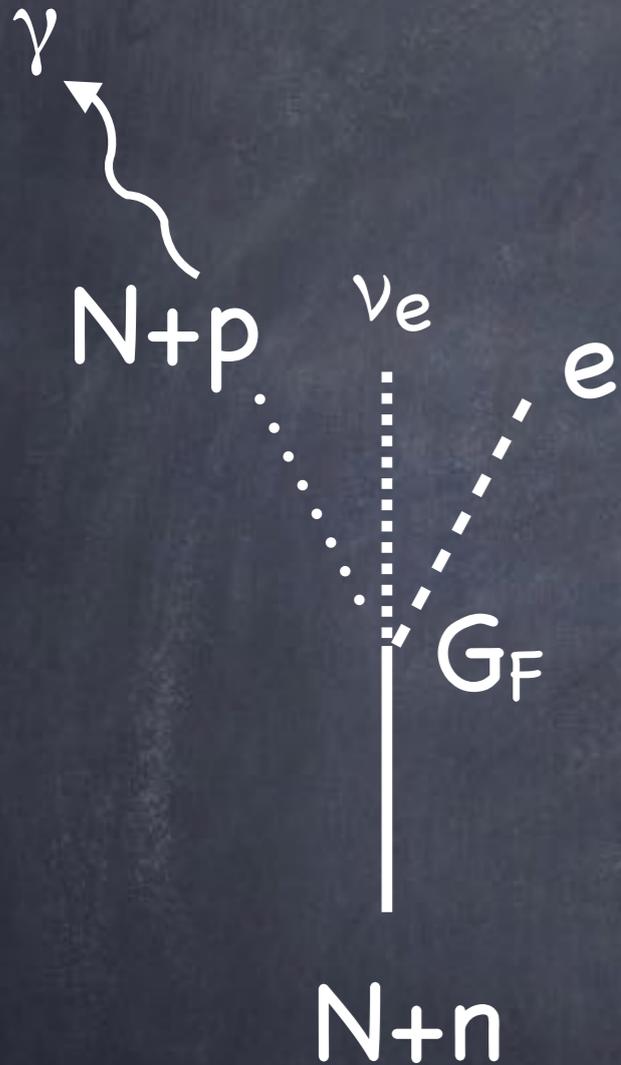
Korobkin + 13; Rosswog, Korobkin + 13



- After a second $dE/dt \propto t^{-1.3}$ (Freiburghaus + 1999; Korobkin + 2013)

Energy Generation

Hotokezaka, Sari & TP + 16



$$t_f = \frac{2\pi^3}{G_F^2} \frac{\hbar^7}{m_e^5 c^4} \approx 10^4 \text{ sec}$$

$$\dot{E} = \epsilon_e \frac{m_e c^2}{t_f} \left(\frac{t}{t_F} \right)^{-\alpha}$$

$$\frac{1}{\tau} \propto \frac{d}{dE} \int d^3 p_e \int d^3 p_\nu$$

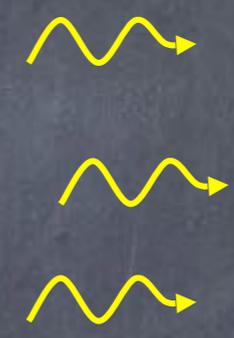
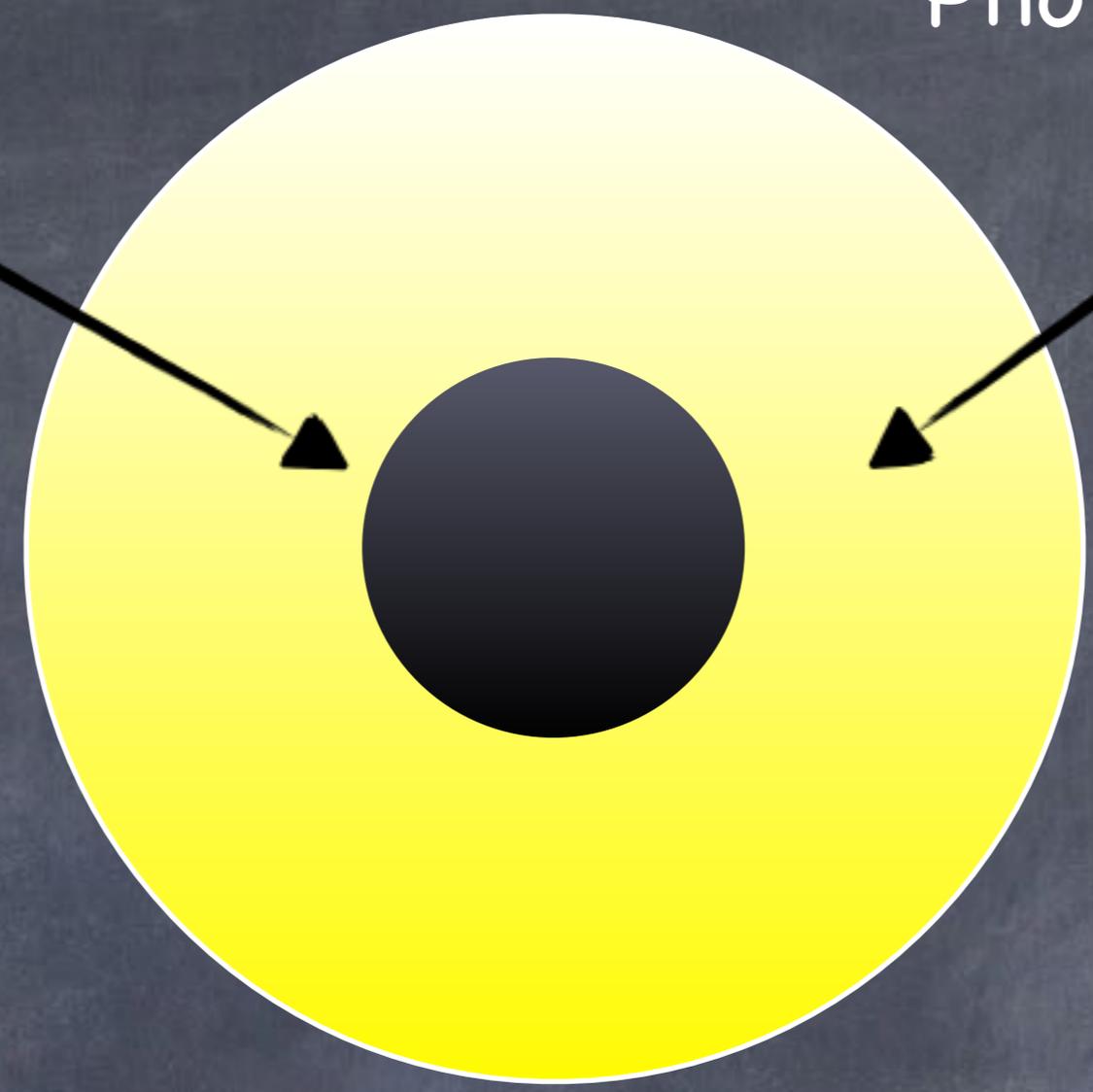
$$E^3 \text{ or } E^{3/2} \quad E^3$$

Relativistic $\frac{1}{\tau} \propto E^5 \rightarrow \alpha = 6/5$

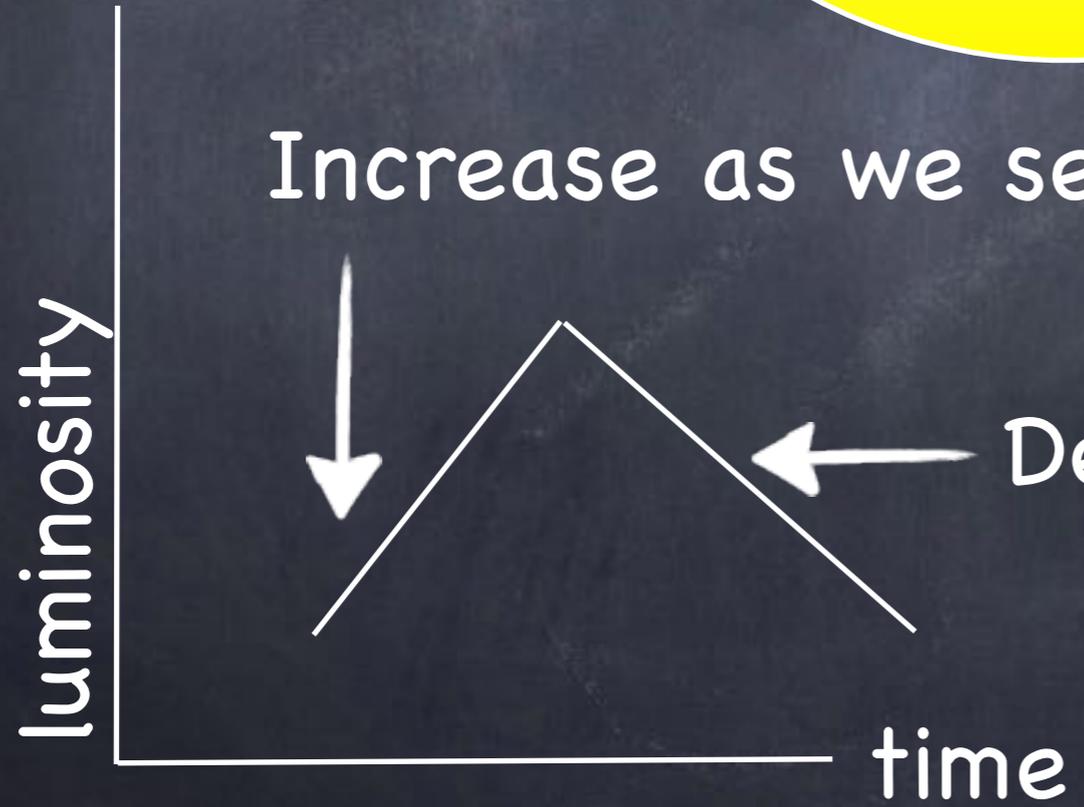
Newtonian $\frac{1}{\tau} \propto E^{7/2} \rightarrow \alpha = 9/7$

$$\tau = c/v$$

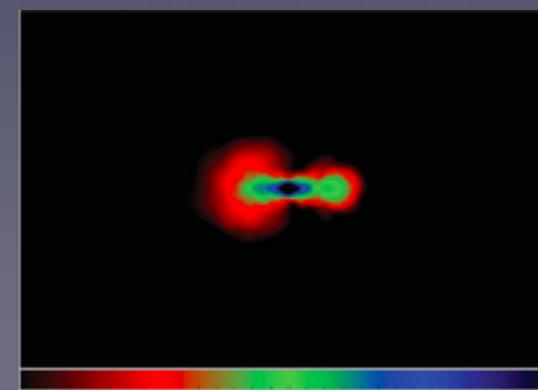
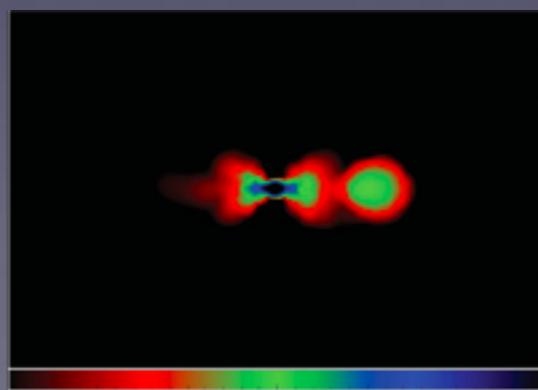
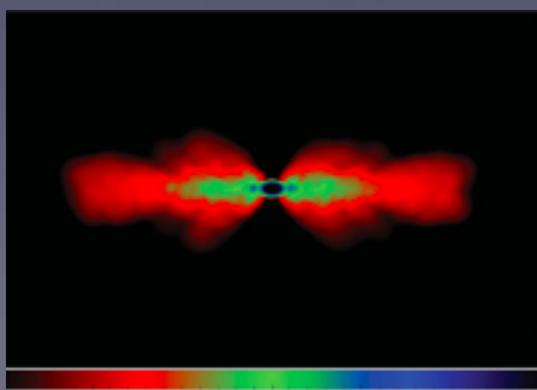
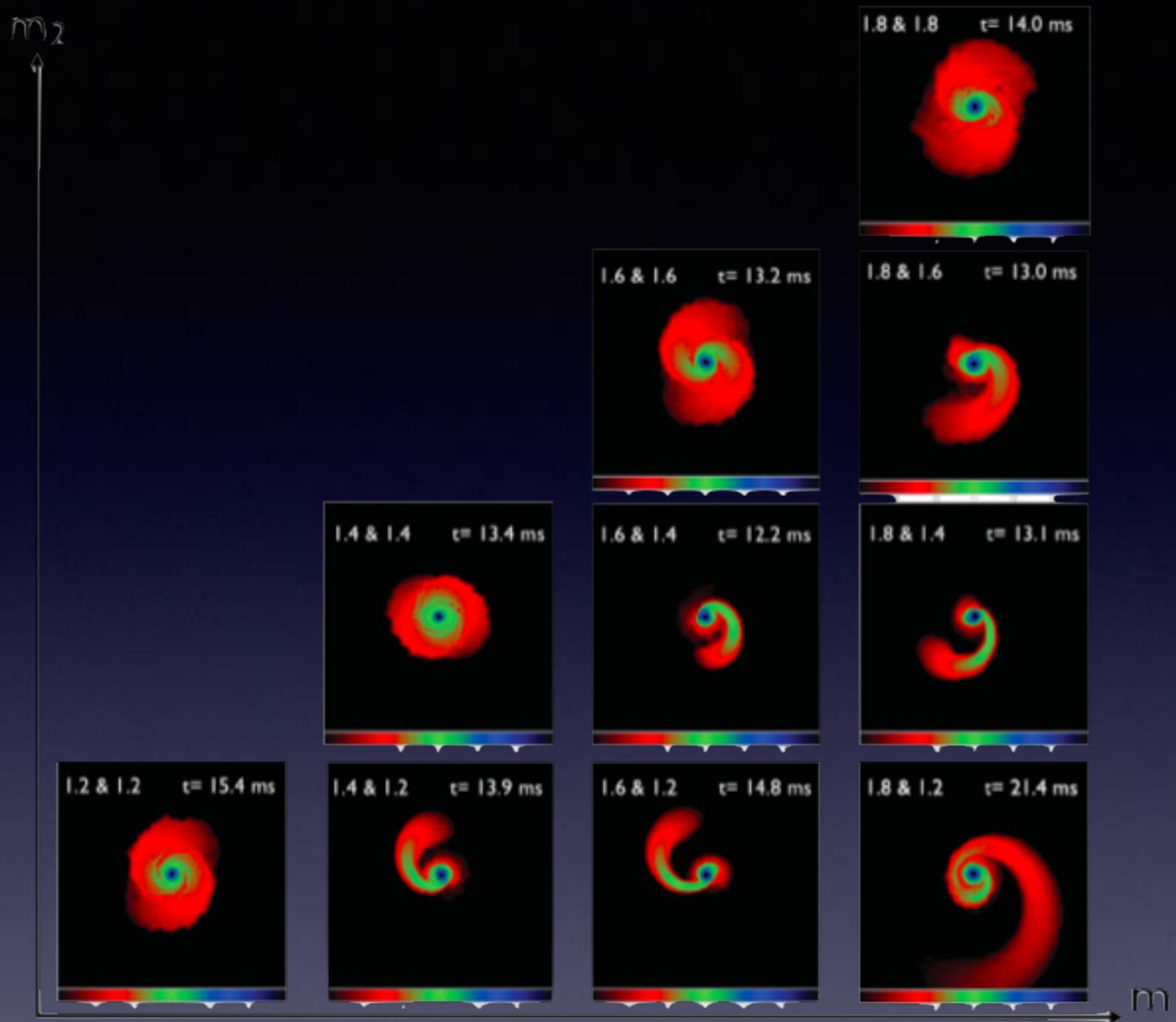
Photons escape from this region



Increase as we see a large fraction of the matter.



Decrease due to radioactive decay



Macronova

Peak time and peak luminosity

Diffusion time = expansion time \Leftrightarrow

Mass of the "emitting region"

$$\frac{m(v)}{v} = \frac{4\pi ct^2}{\kappa}$$

Luminosity

$$L(t) = \dot{\epsilon}(t)m(v) = \dot{\epsilon}_0(t/t_0)^{-\alpha}m(v)$$

Radioactive heating rate

The peak time

$$\tilde{t}_p \approx \sqrt{\frac{\kappa m_{ej}}{4\pi c\bar{v}}} = 4.9 \text{ days} \left(\frac{\kappa_{10} m_{ej,-2}}{\bar{v}_{-1}} \right)^{1/2}$$

The peak luminosity

$$\tilde{L}_p \approx \dot{\epsilon}_0 m_{ej} \left(\frac{\kappa m_{ej}}{4\pi c\bar{v}t_0^2} \right)^{-\alpha/2} = 2.5 \times 10^{40} \frac{\text{erg}}{\text{s}} \left(\frac{\bar{v}_{-1}}{\kappa_{10}} \right)^{\alpha/2} m_{ej,-2}^{1-\alpha/2}$$

Lanthanides dominate the Opacity

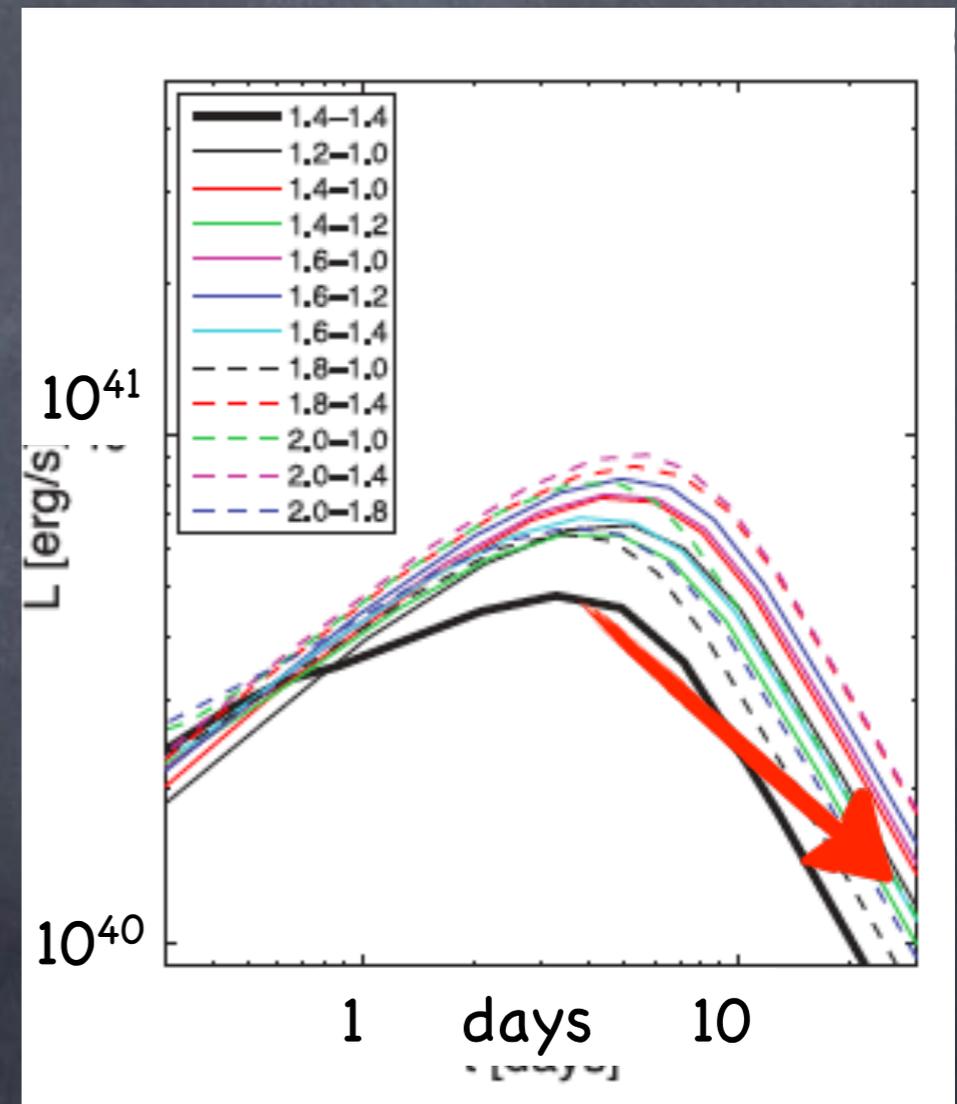
(Kassen & Barnes 13; Tanak & Hotokezaka 2013)

• $\kappa = 10 \text{ cm}^2/\text{gm}$

• $t_{\text{max}} \propto \kappa^{1/2} \Rightarrow$ longer

• $L_{\text{max}} \propto \kappa^{-0.65} \Rightarrow$ weaker

• $T \propto \kappa^{-0.4} \Rightarrow$ redder

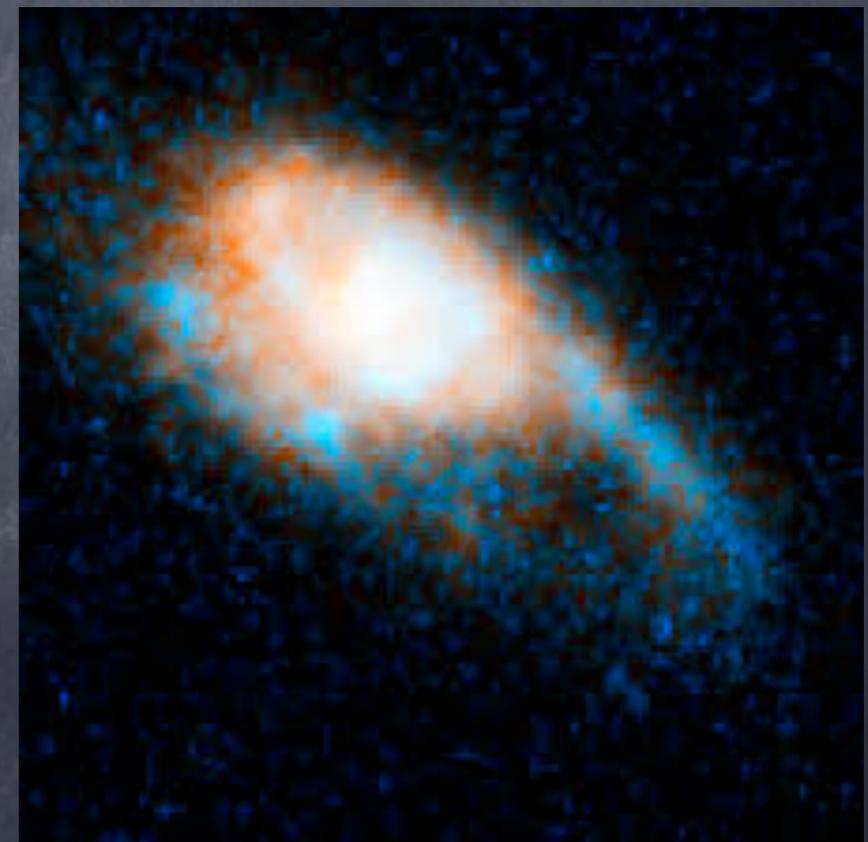
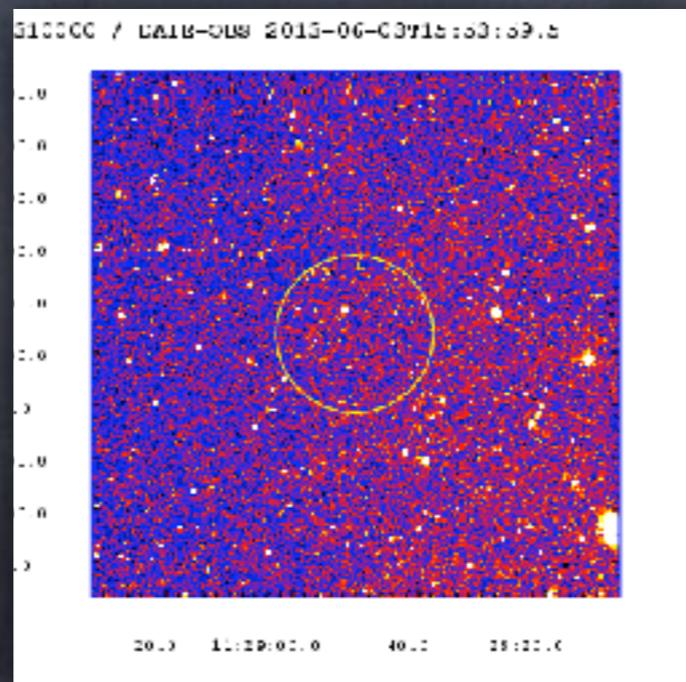
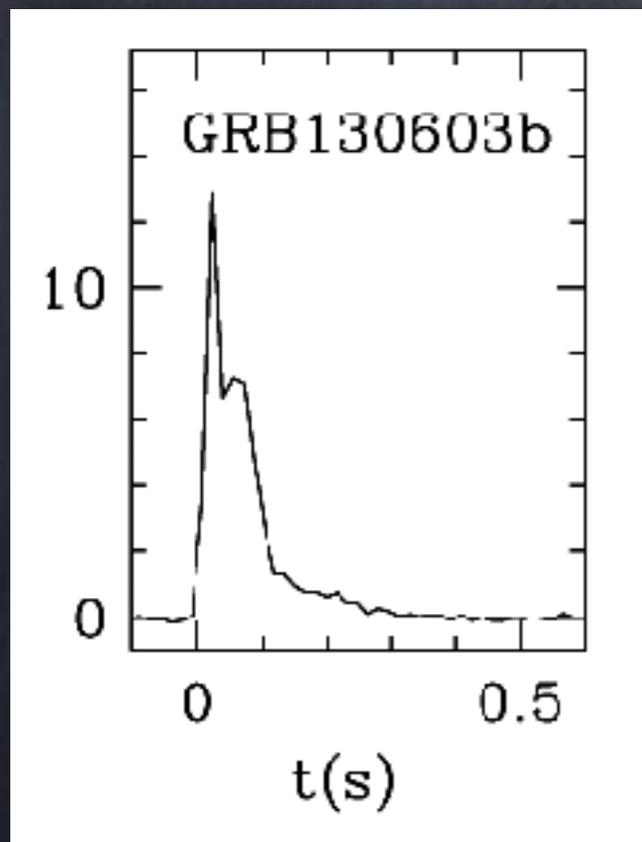


uv or optical \rightarrow IR

Putting it all together

5. Gamma-Ray Burst (GRB)

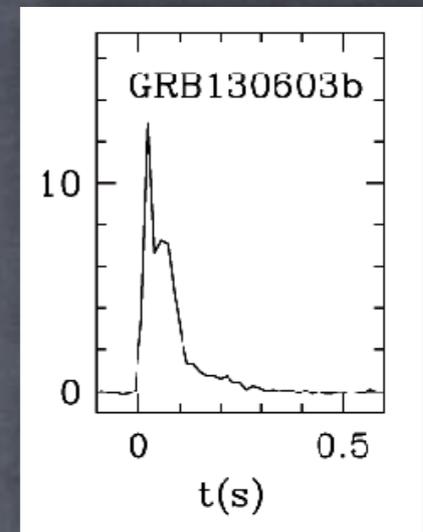
130603B



GRB 130603B

$z=0.356 \Leftrightarrow 1 \text{ Gpc} = 3 \text{ Glyr}$

GRB 130603B



At 15:49:14 UT, the Swift Burst Alert Telescope (BAT) triggered and located GRB 130603B (trigger=557310). Swift slewed immediately to the burst.

The BAT on-board calculated location is RA, Dec 172.209, +17.045 which is

RA(J2000) = 11h 28m 50s

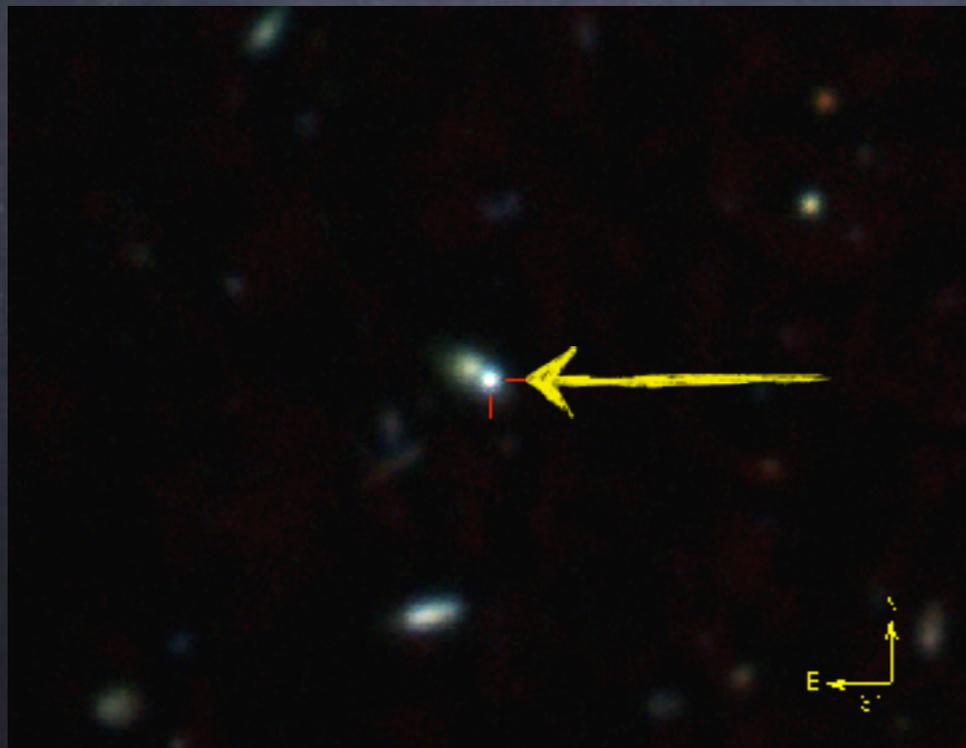
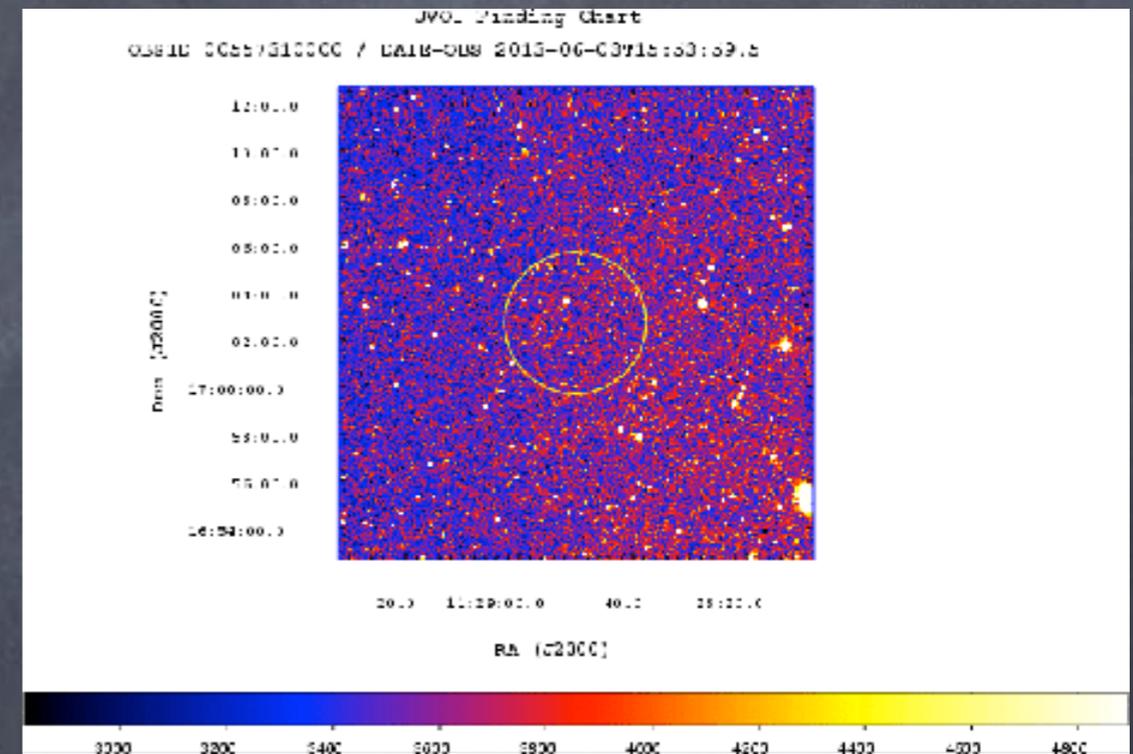
Dec(J2000) = +17d 02' 42"

with an uncertainty of 3 arcmin (radius, 90% containment, including systematic uncertainty).

The BAT light curve showed a single spike structure with a duration of about 0.4 sec.

The peak count rate was 60000 counts/sec (15-350 keV), at ~0 sec after the trigger.

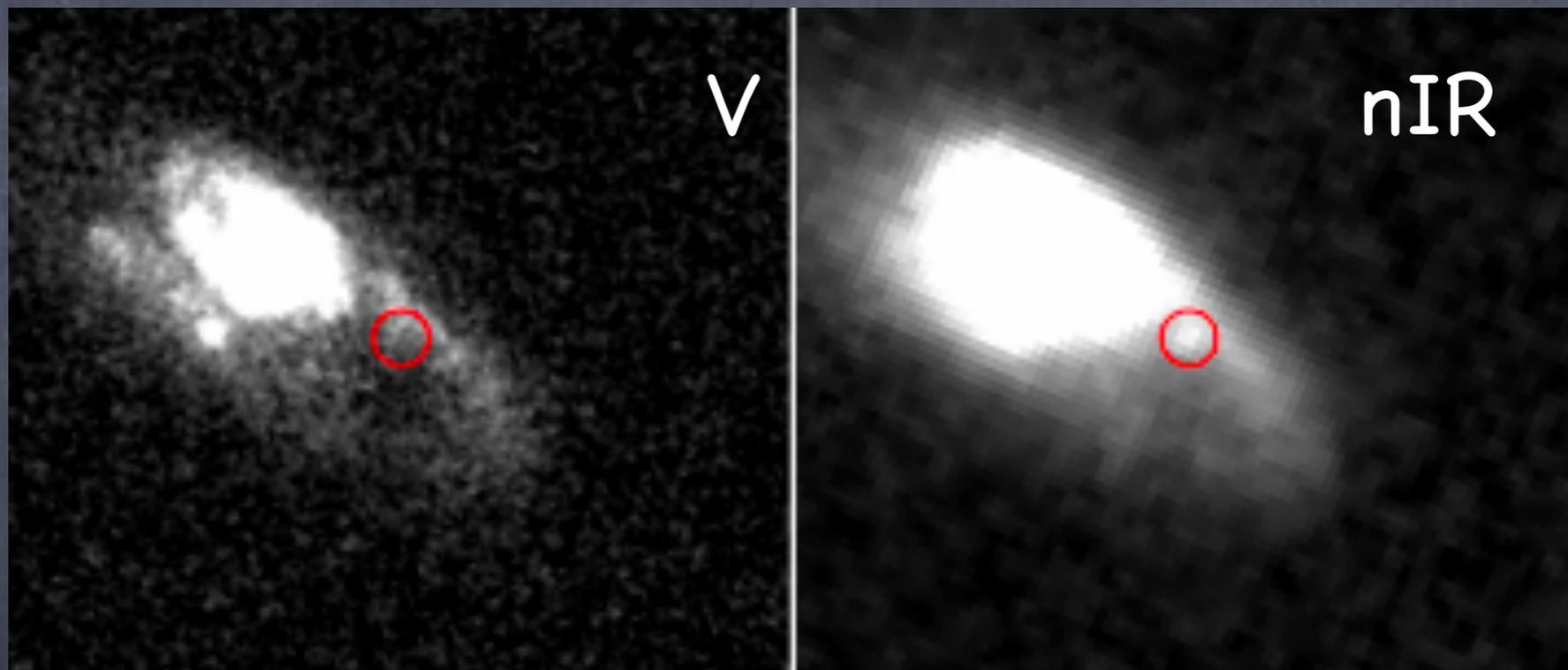
A short burst



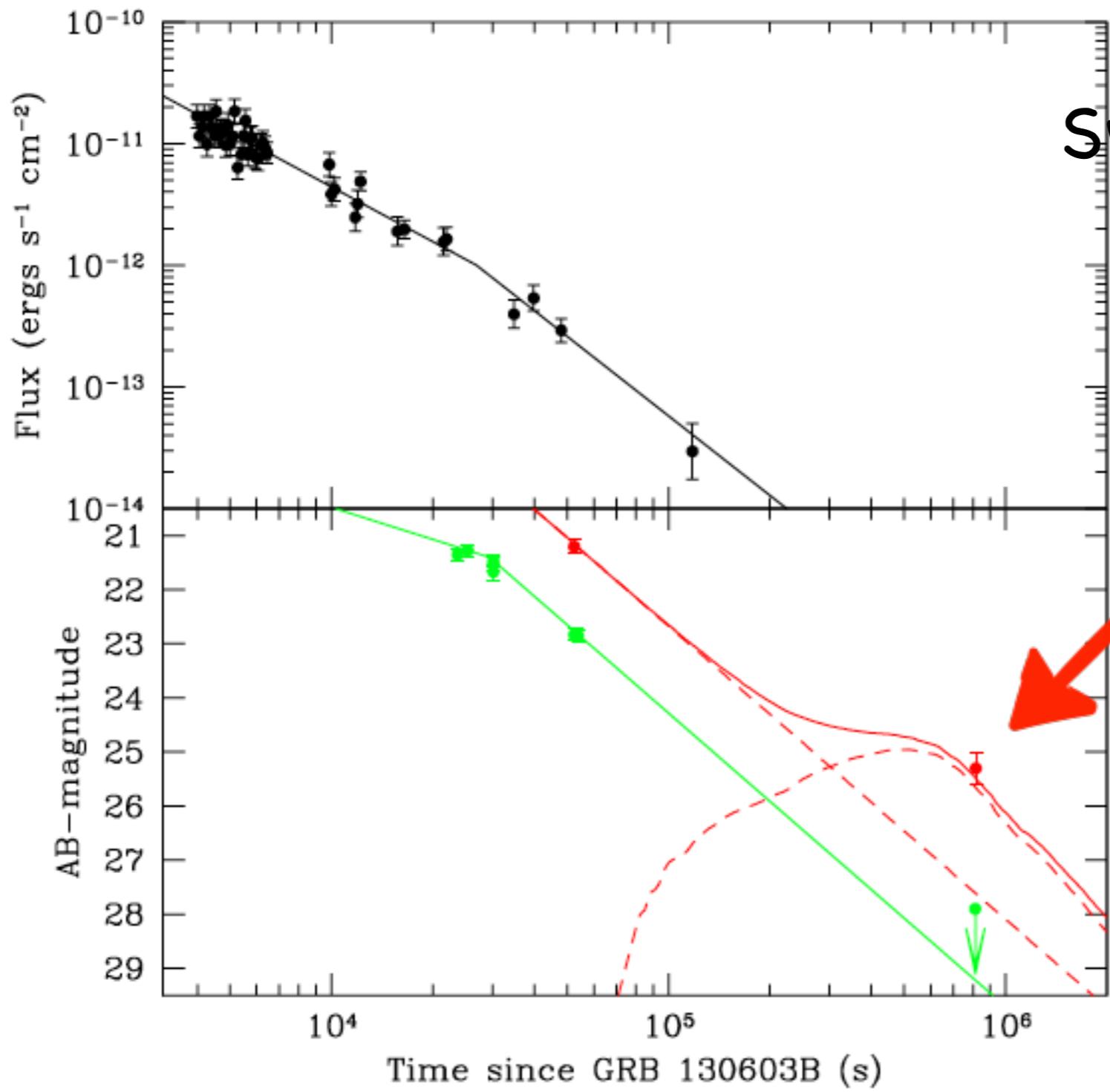
$z=0.356 \Leftrightarrow 1 \text{ Gpc} = 3 \text{ Glyr}$

GRB130603B @ 9 days AB

(6.6 days at the source frame)



HST image (Tanvir + 13)



Swift

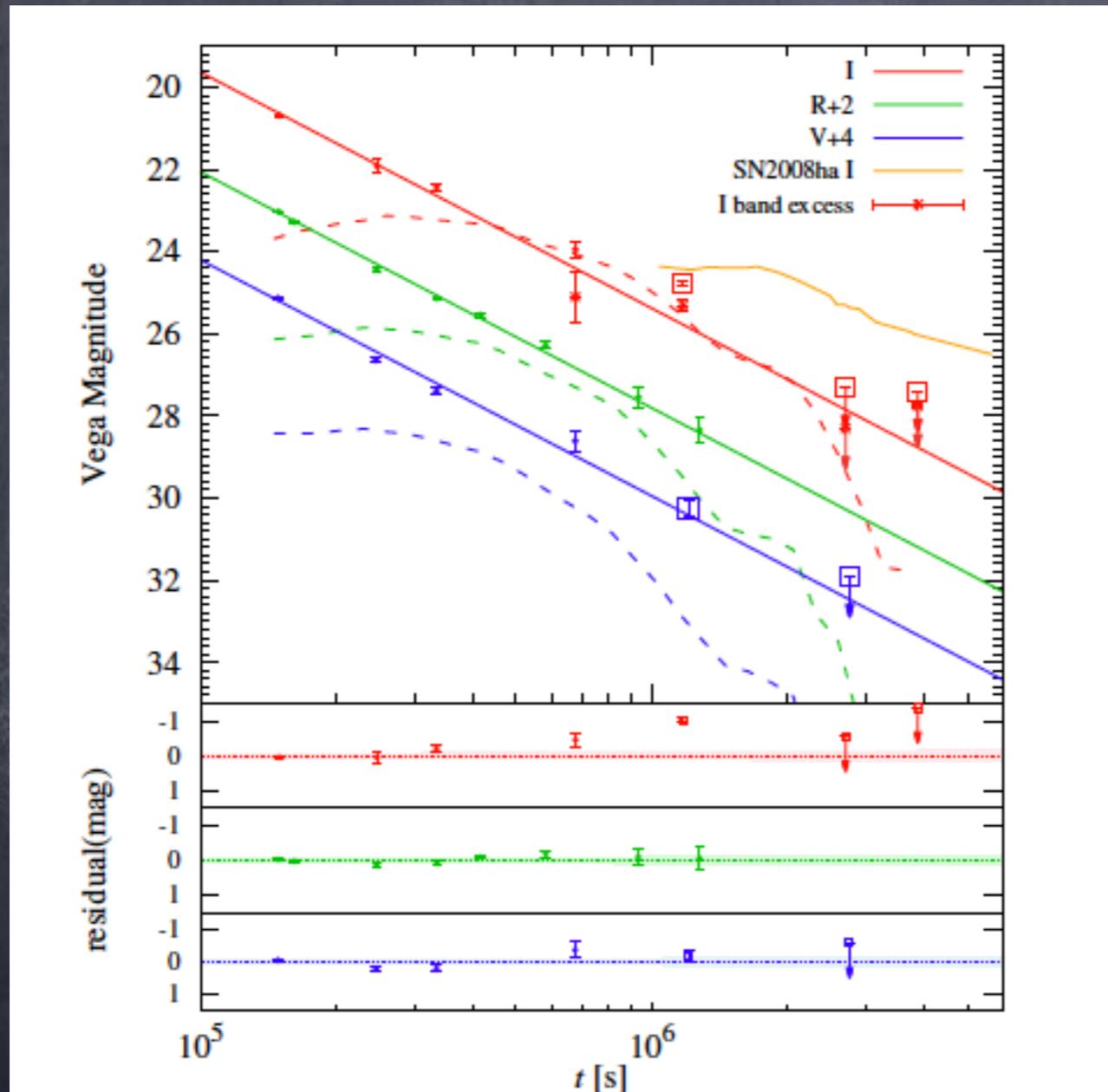
Macronova?

0.01-0.05 M_{\odot}



Tanvir + 13 (see also Berger + 13)
GRB 130603B

6. Additional Evidence

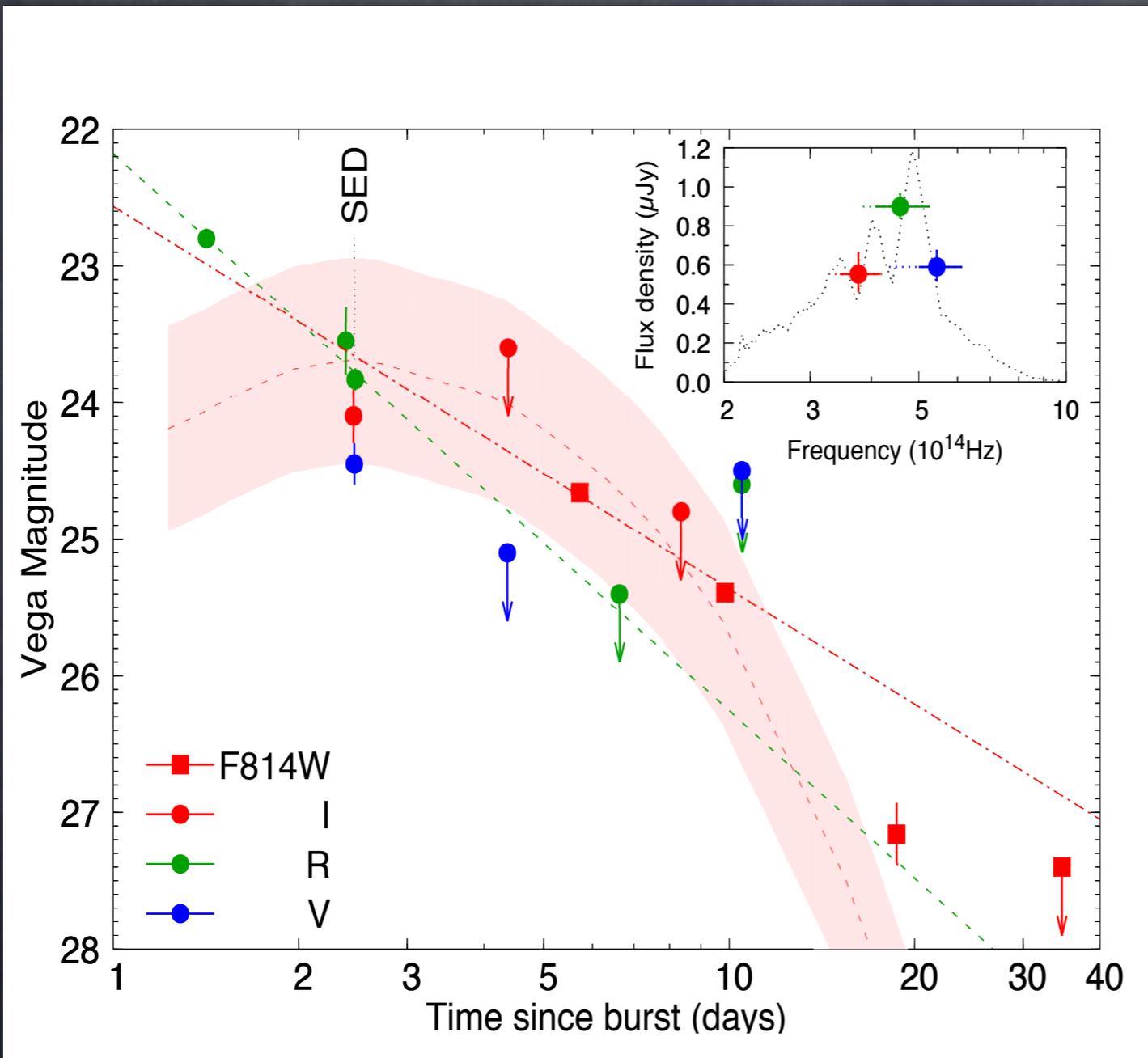


GRB 060614

Need $M \approx 0.1 M_{\odot}$

\Rightarrow BH-NS ?

Yang et al., 2015



GRB 050709

Need $M \approx 0.05 M_{\odot}$

\Rightarrow BH-NS ?

Jin et al., 2016

Additional evidence

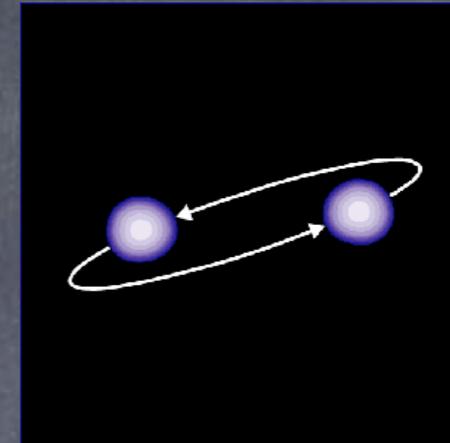
Are Macronova Frequent?

- There are 3 (6) possible (nearby) historical candidates with a good enough data
- In 3/3 (3/6) there are possible Macronovae

If correct



Confirmation of the GRB neutron star merger model (Eichler, Livio, TP & Schramm 1989).



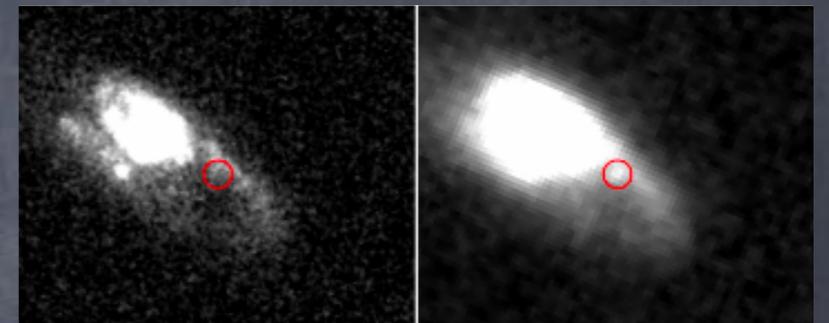
Confirmation of the Li-Paczynski Macronova.



Confirmation that compact binary mergers are the source of heavy ($A > 130$) r-process material (Gold, Silver, Platinum, Plutonium, Uranium etc...).



7. The Origin of GOLD



Implications

Mass ejected in a merger

Observed luminosity =
 10^{41} erg/sec @ 6.6 days

$$m_{ej} > 0.02(\epsilon/0.5)^{-1} m_{\odot}$$

of mergers

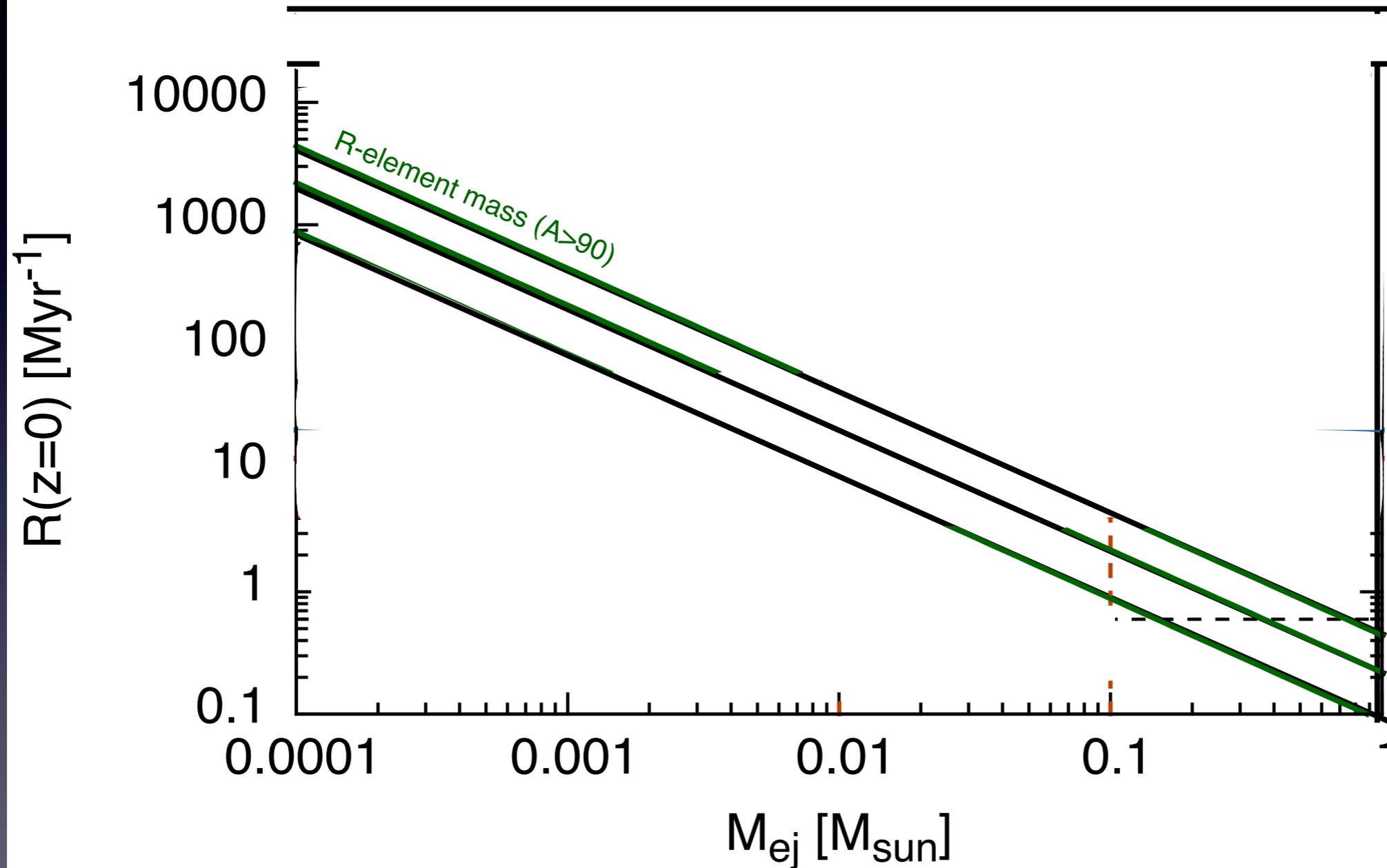
$$N = 2.5 \times 10^5 \left(\frac{M^{A>130}}{10^4 m_{\odot}} \right) \left(\frac{m_{ej}}{0.04 m_{\odot}} \right)^{-1}$$

$A>130$ r-process material in the Galaxy

Mergers' Rate

$$\begin{aligned} R_{merger} &= 20 \left(\frac{m_{ej}}{0.04 m_{\odot}} \right)^{-1} \left(\frac{M^{A>130}}{10^4 m_{\odot}} \right) \text{Myr}^{-1} \\ &= 200 \left(\frac{m_{ej}}{0.04 m_{\odot}} \right)^{-1} \left(\frac{M^{A>130}}{10^4 m_{\odot}} \right) \text{Gpc}^{-3} \text{yr}^{-1} \end{aligned}$$

R-Process

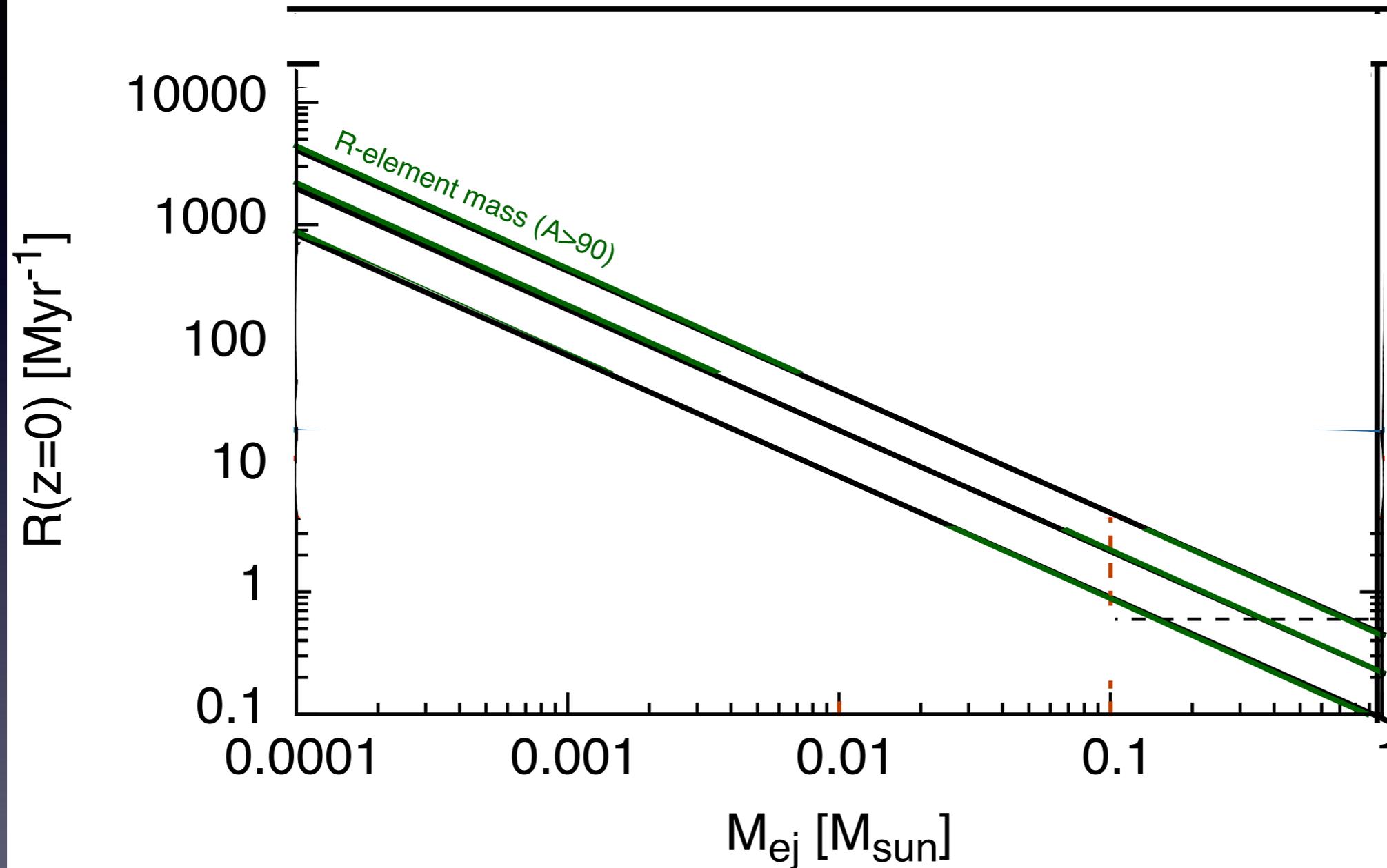


lines of R-mass: Current event rate is lower than the average one by a factor of 5 (lower line), 3 (middle line).

lines of SGRB: beaming factor $f_b^{-1} = 10, 30, 70$ (Wanderman & Piran 2015)

lines of NSNS: 95% confidence level (Kim et al 2015)

R-Process



Can we break the yield - rate degeneracy?
Hotokezaka, TP Paul, Nature Phys 2015

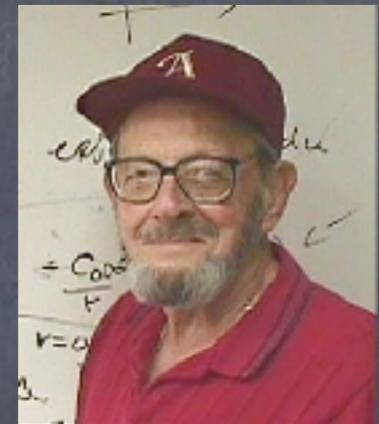
One cannot give a talk in Astronomy these days without a reference to the Solar System and life.

- The early Solar System had ^{244}Pu ($\tau = 117$ Myr) Wasserburg et al, (2006).

No evidence for ^{244}Pu deposition in deep-sea crust and sediment accumulated over the last ~ 25 Myr (M. Paul et al., 2001; A. Wallner et al., in preparation).

^{244}Pu is NOT from the Inter Stellar Medium! => Actinides production near the early Solar System just prior to formation.

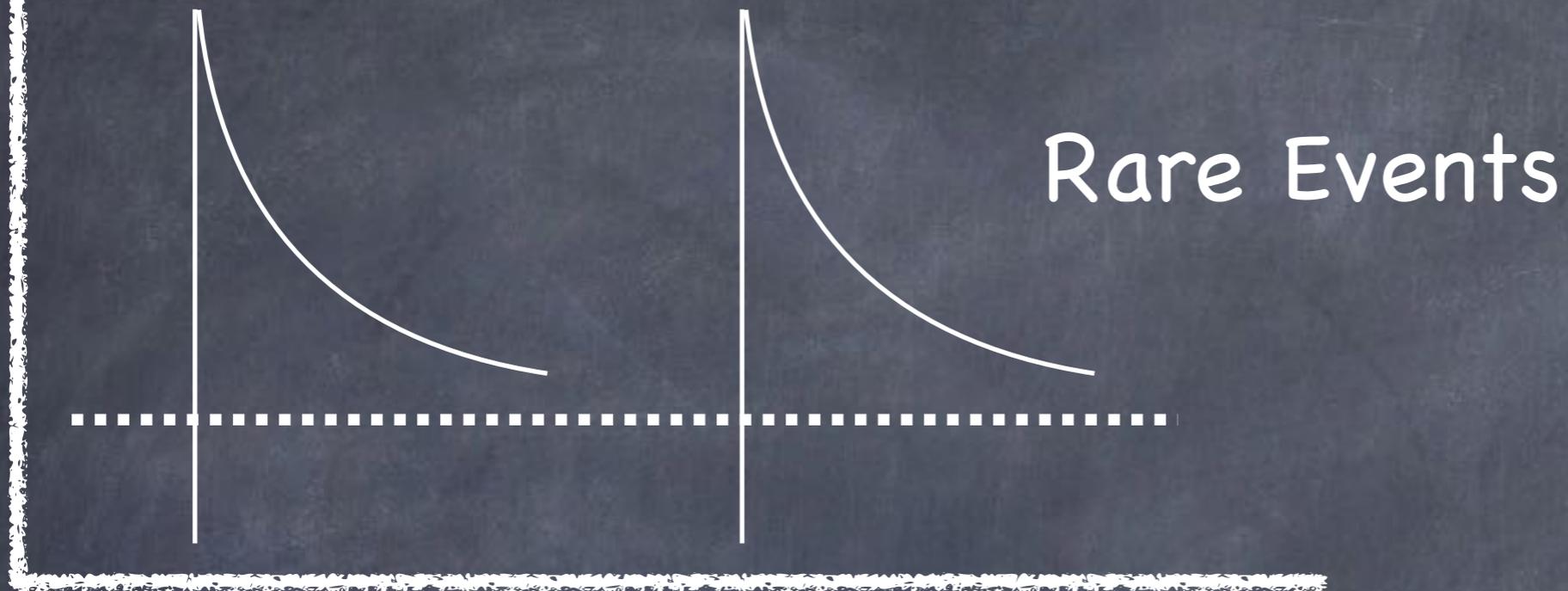
- Irregular production from rare episodes.
=> E.g. a merger within < 50 pc = 150 lyr from the solar system just prior to its formation?



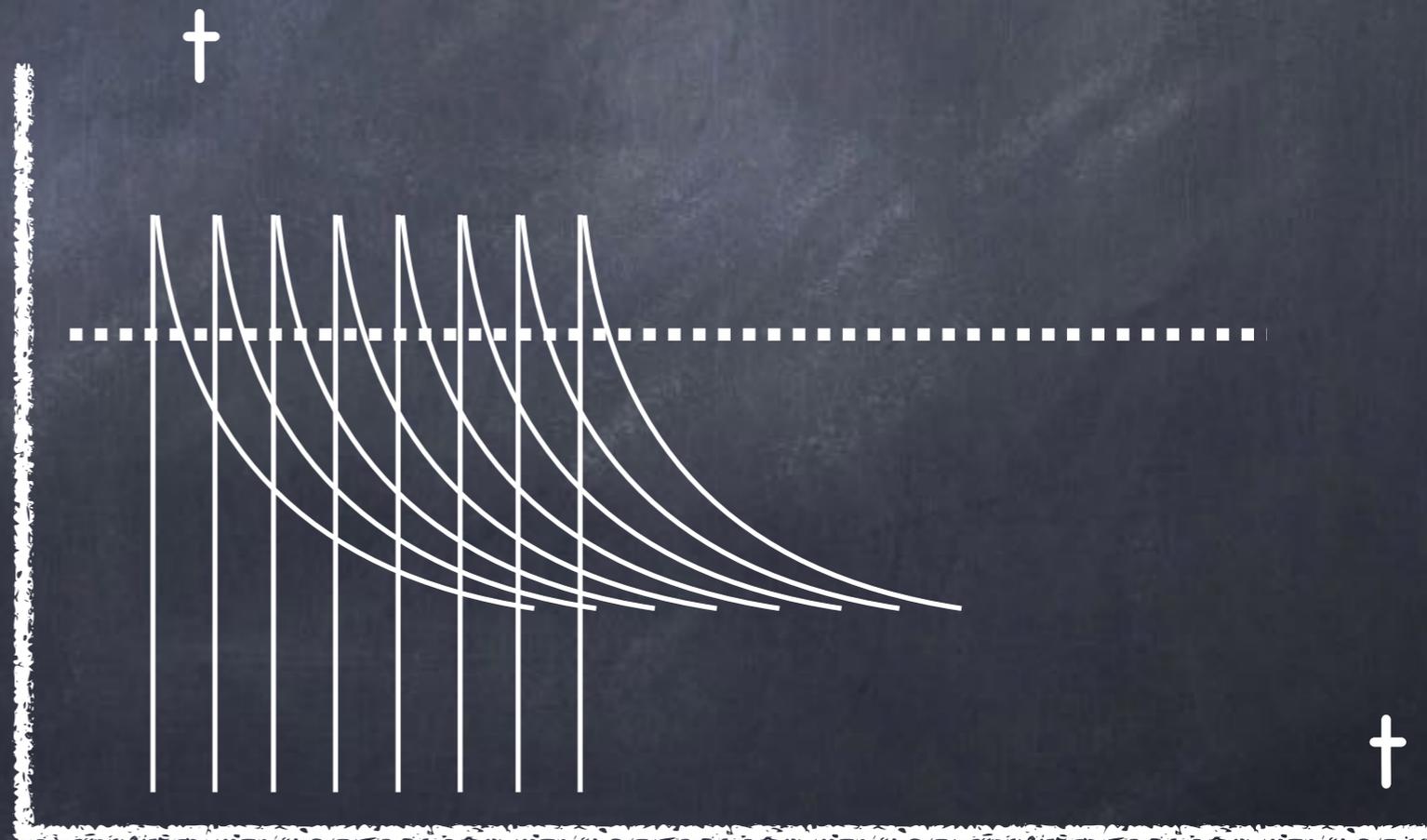
Gerry Wasserburg



Radioactive Elements

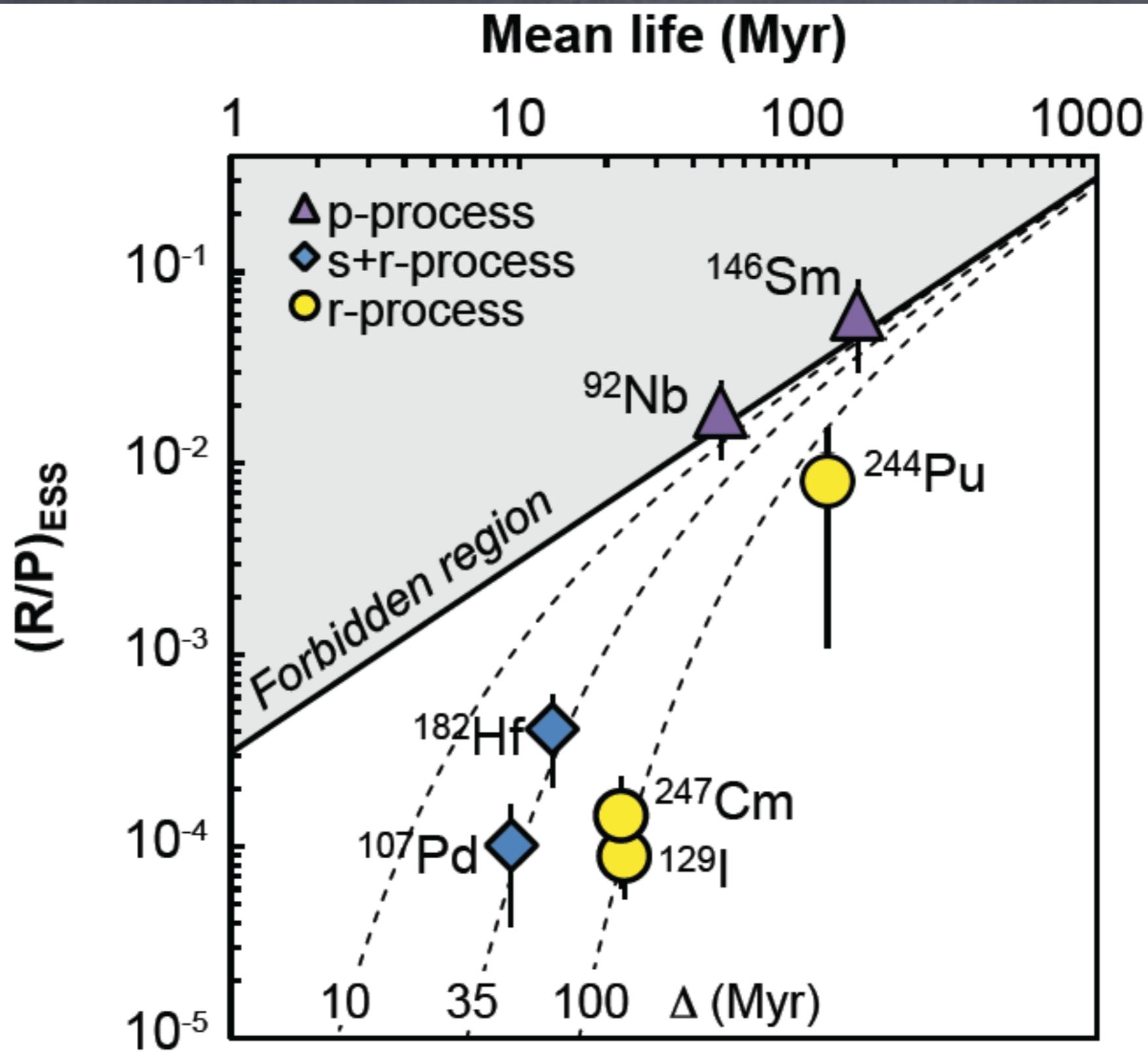


Frequent events



High ^{244}Pu at the early solar system =>

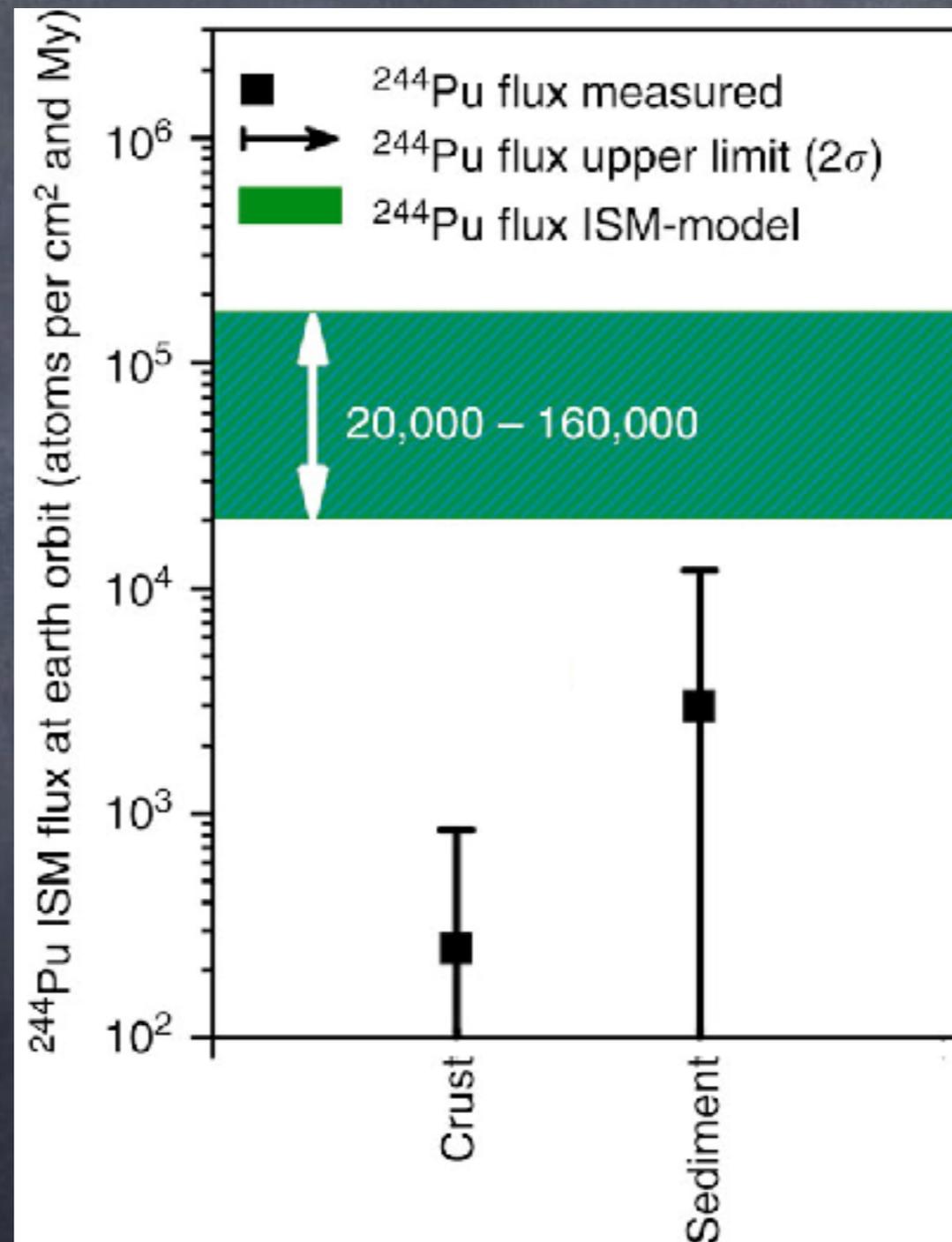
- ^{244}Pu Radioactive decay time ~ 100 Myear
- A nearby event near solar system
- Mixing time < 150 Myr
- Large fluctuations possible => Event rate is low
- Lack of Cu => 10 Myr $<$ Mixing length



Tissot + 16

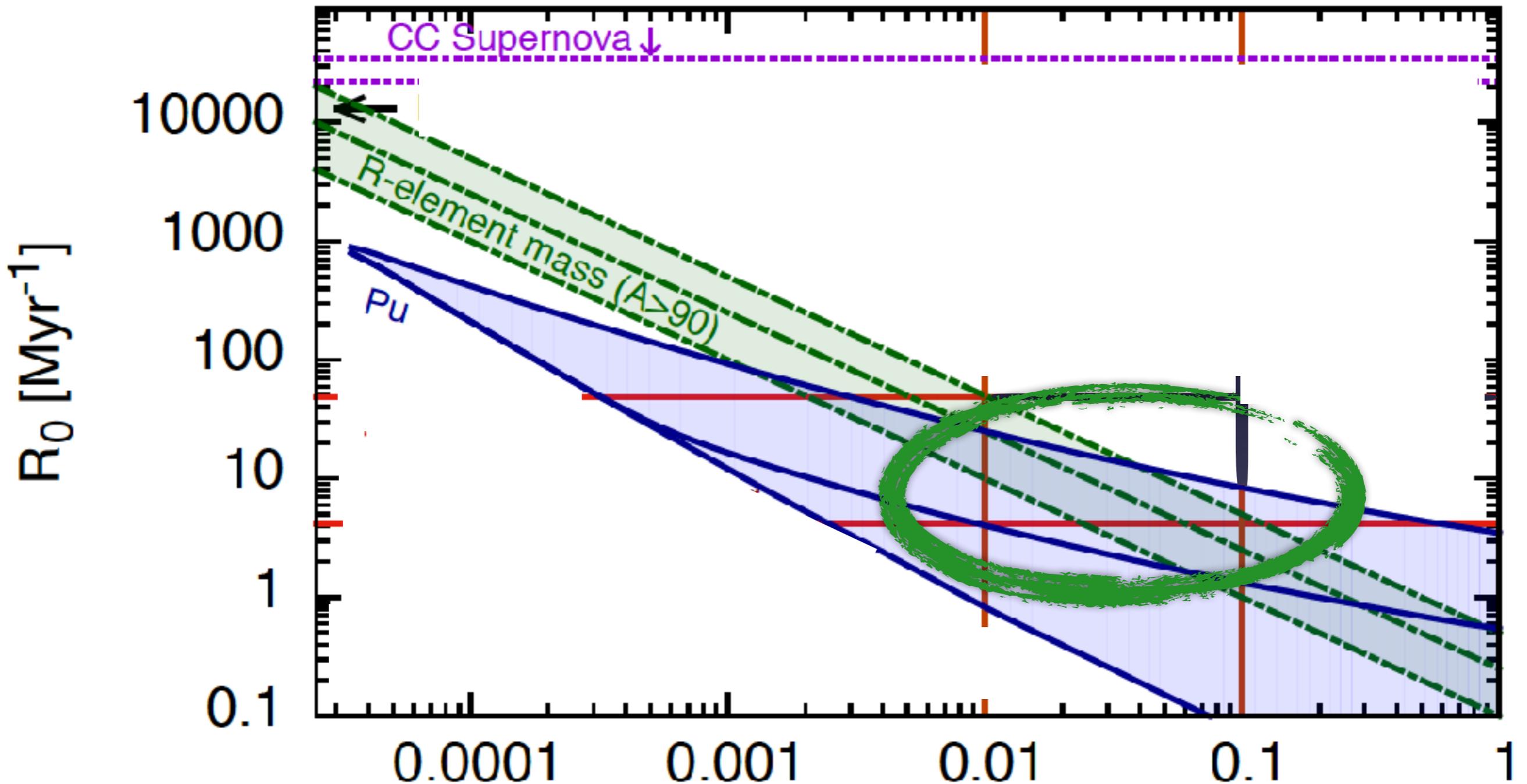
^{244}Pu (half life 81Myr)

The early solar system



Wallner + 14

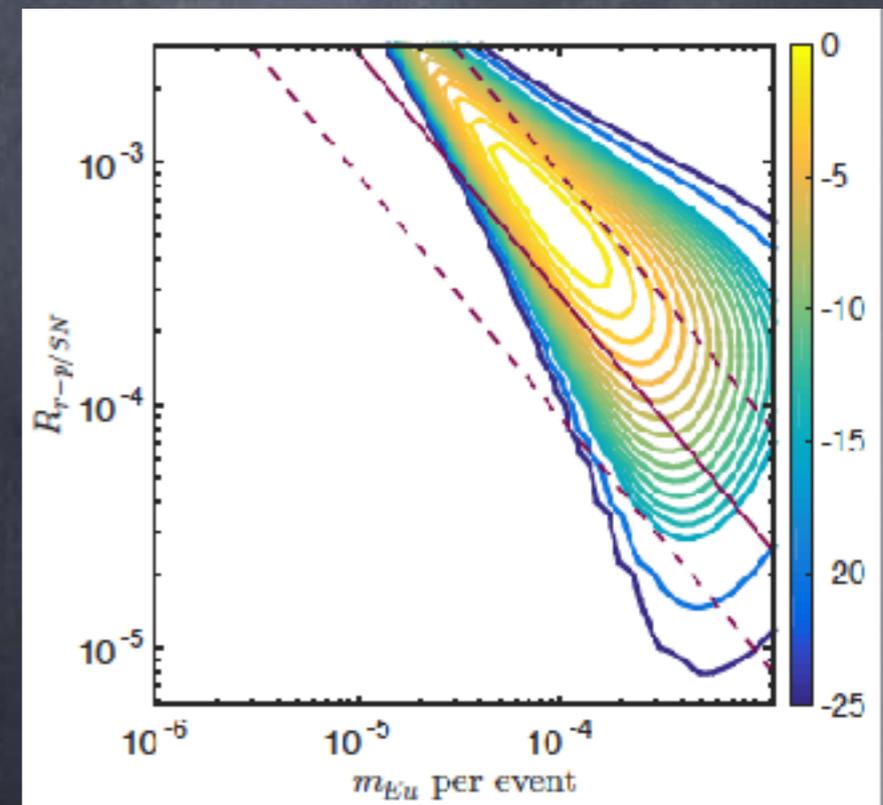
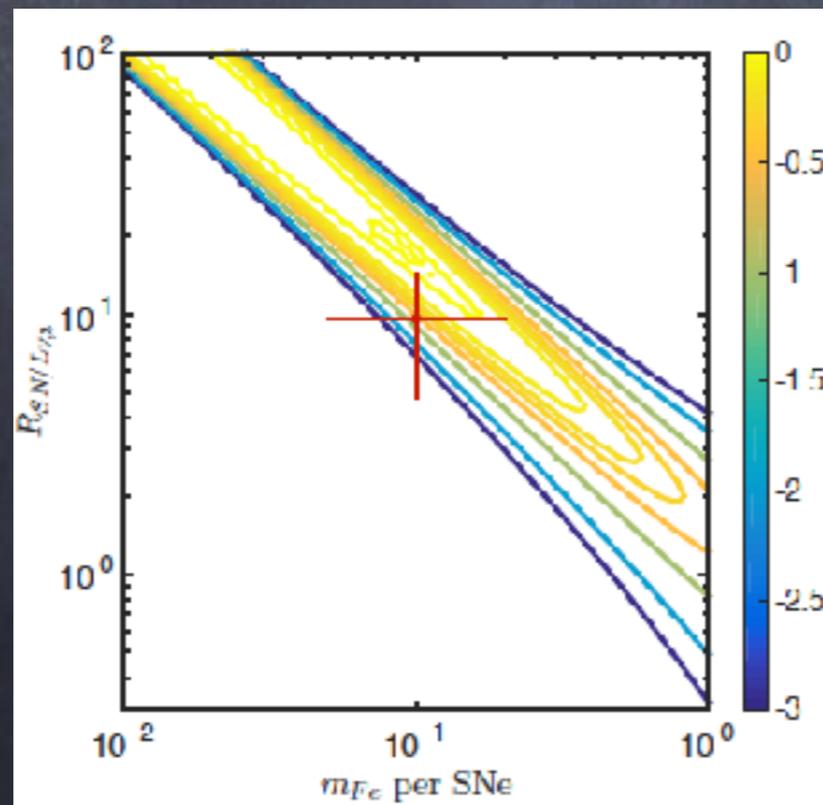
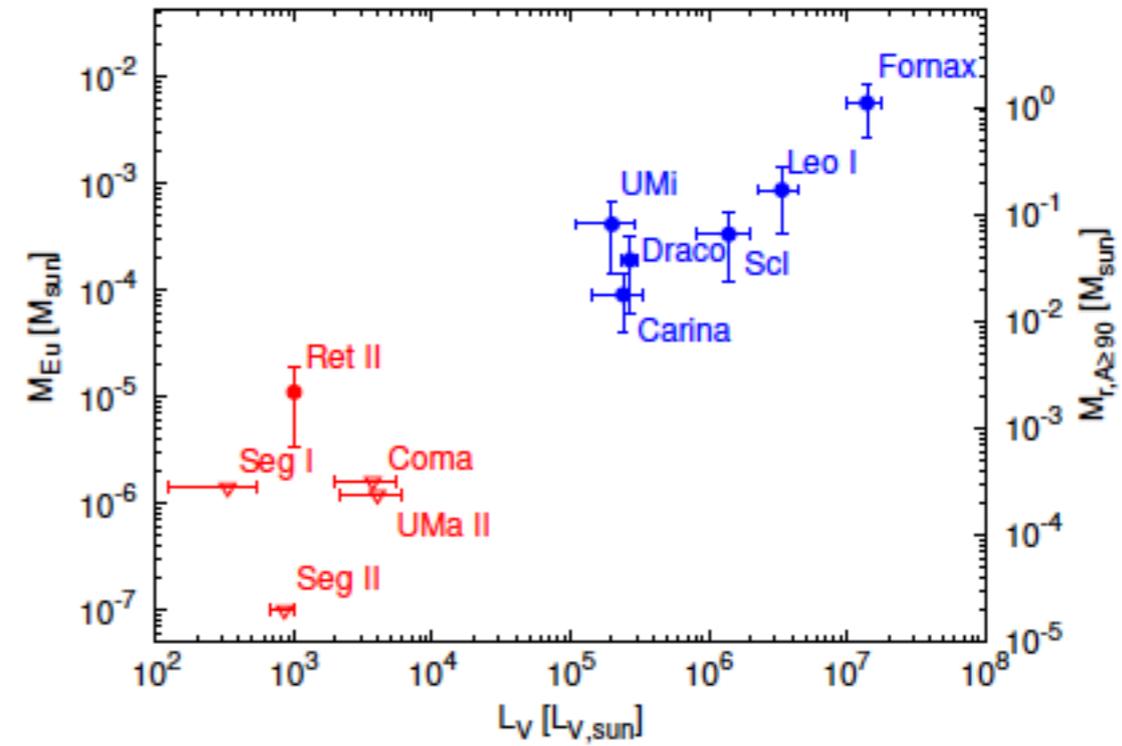
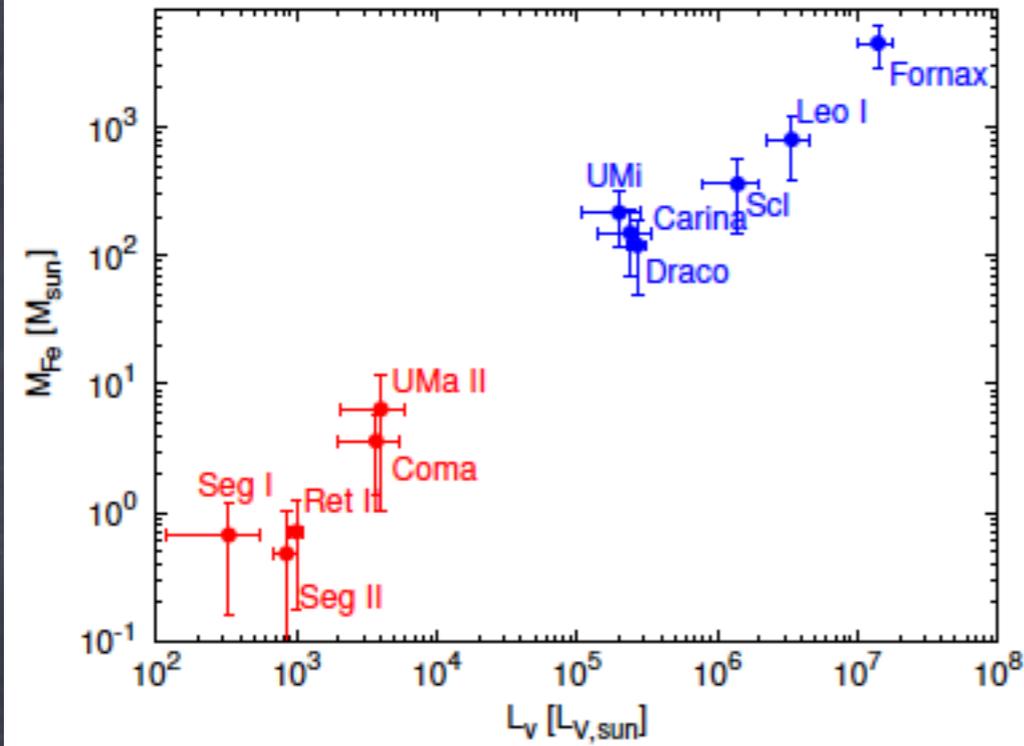
Breaking the degeneracy: ^{244}Pu



Rare and "massive" events
Hotokezaka, TP & Paul, Nature Pays 2015

r-process material in Dwarf Galaxies

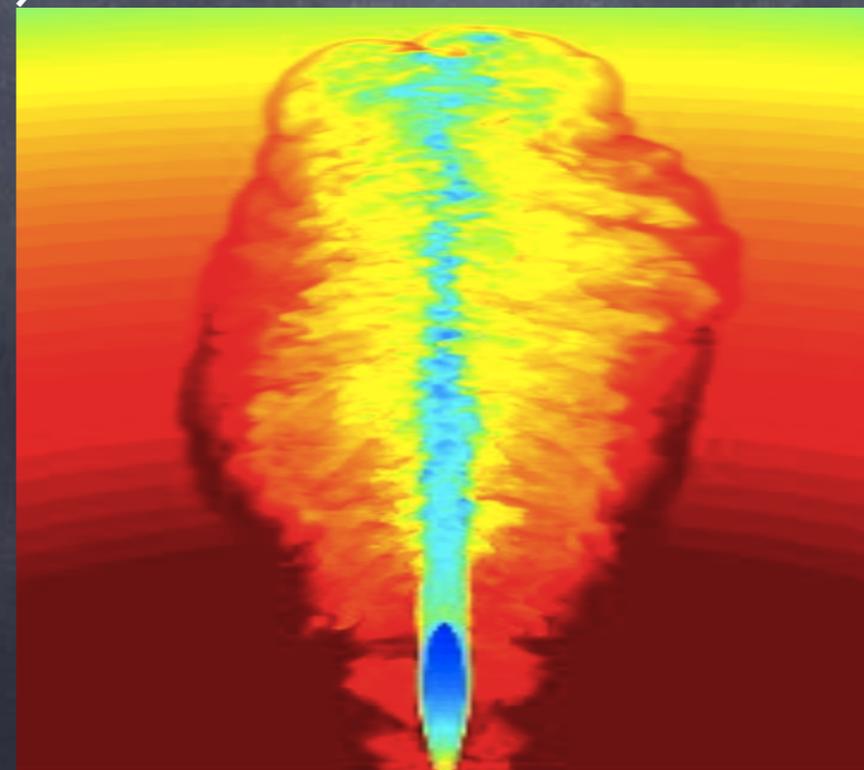
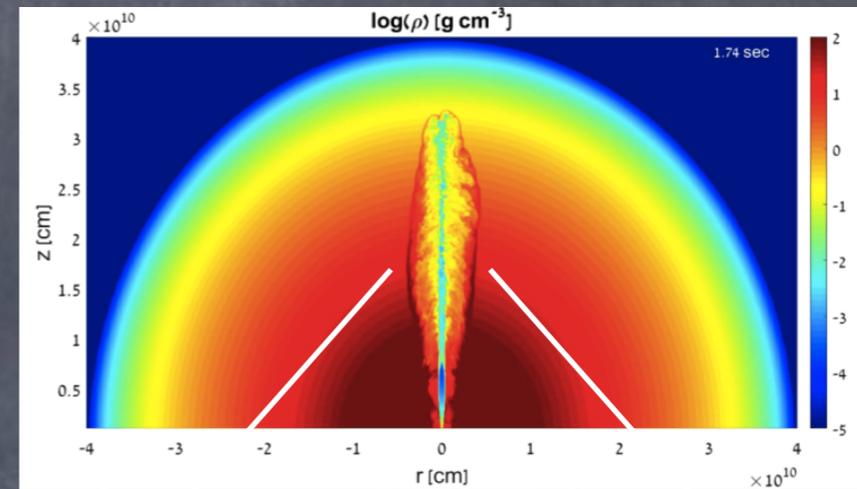
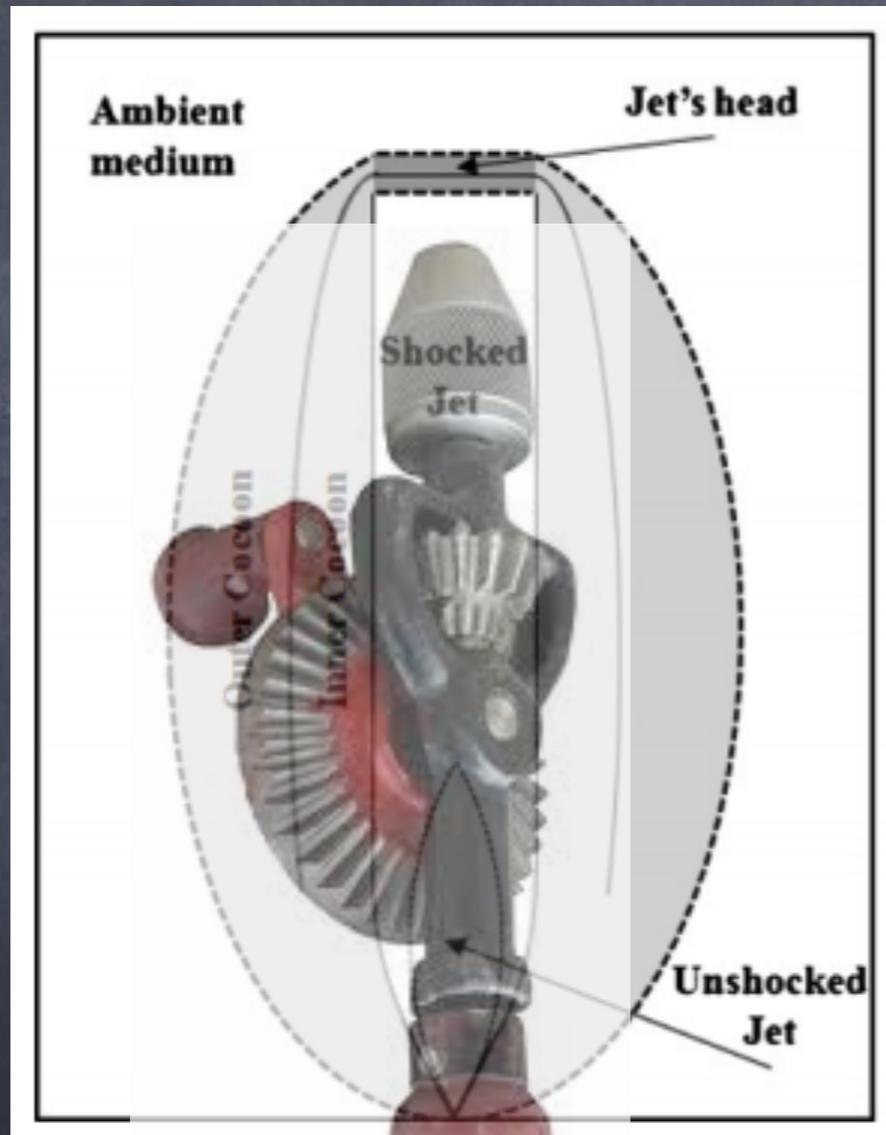
(Beniamini+ 16a,b)



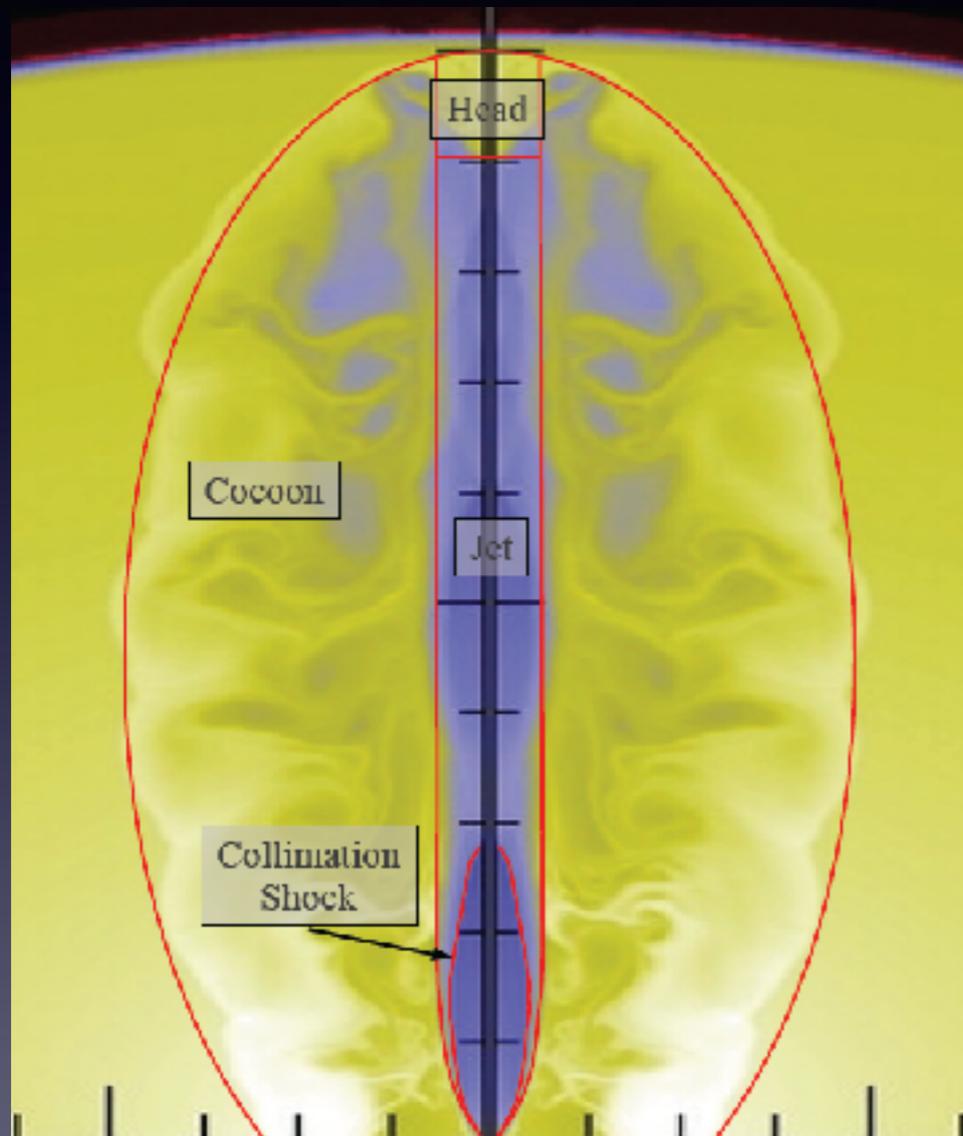
The Origin of Gold

Jet Propagation

(MacFayden & Woosley 1998; Aloy+ 1999; Matzner 2003; Lazzati and Begelman,05; Bromberg + 2011....)



The engine must be active until the jet's head breaks out!



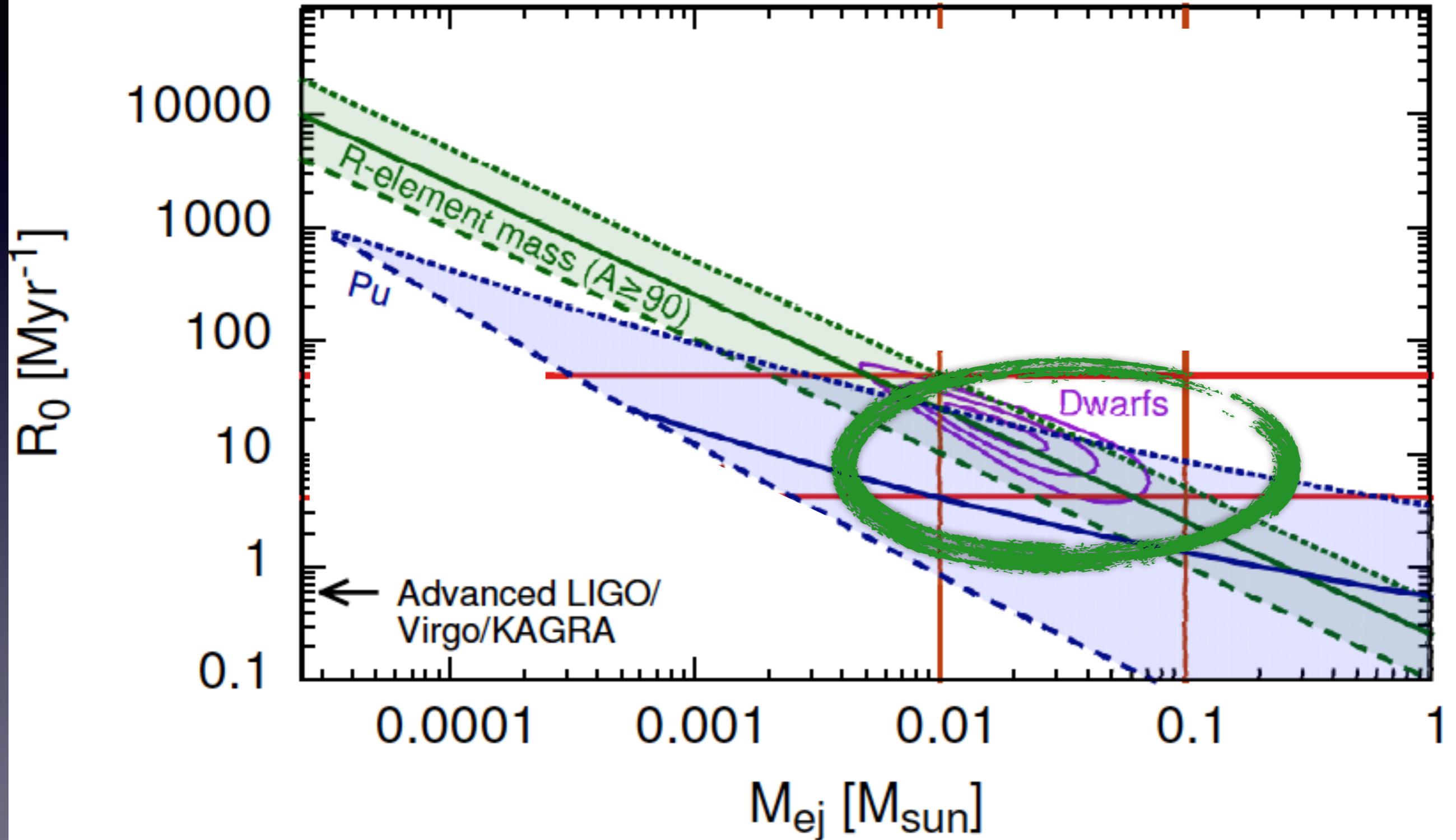
Observed duration

$$T_{90} = T_e - T_B$$

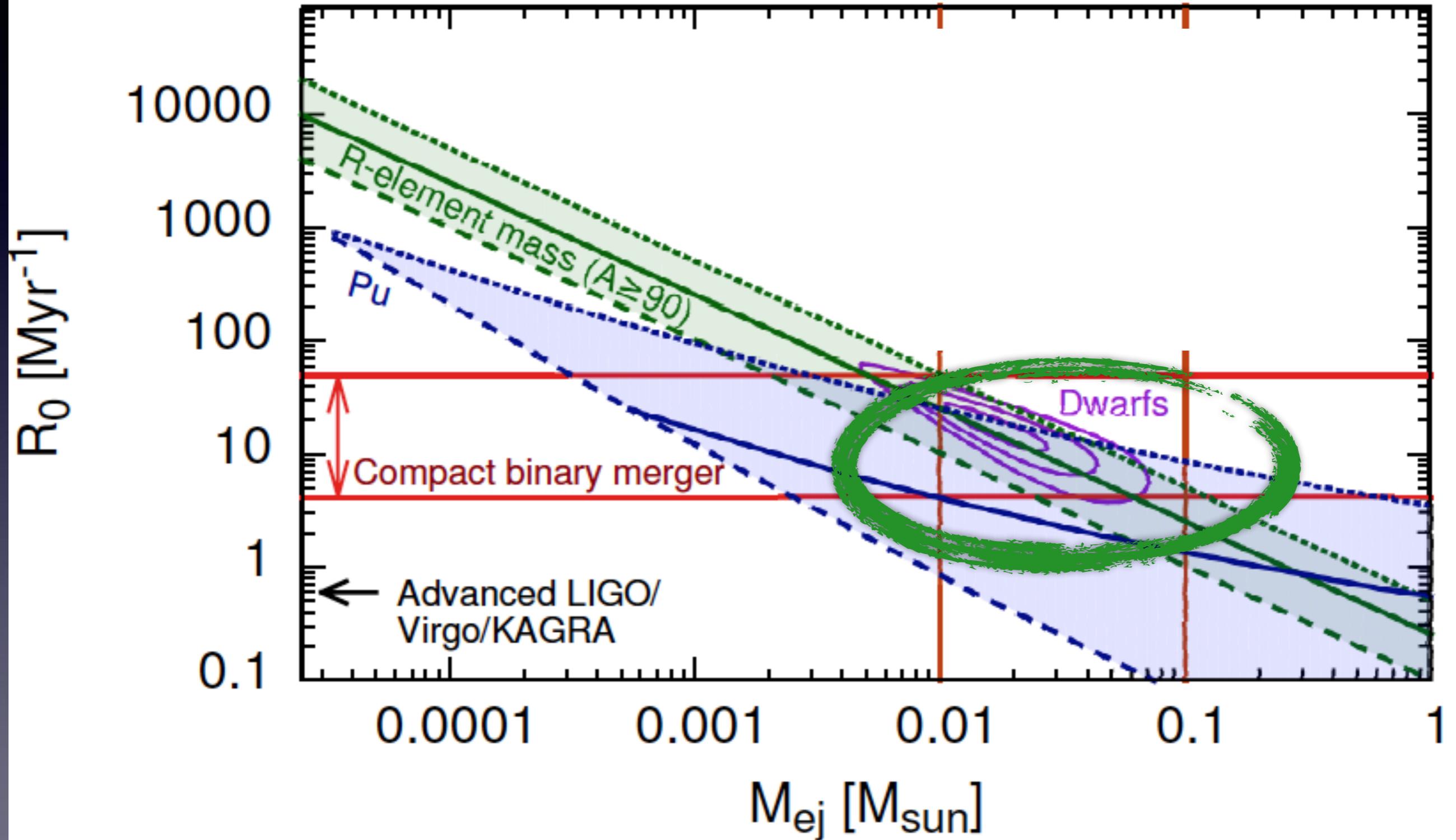
Engine time Break out time

$$t_b \simeq 15 \text{ sec} \cdot \left(\frac{L_{iso}}{10^{51} \text{ erg/sec}} \right)^{-1/3} \left(\frac{\theta}{10^\circ} \right)^{2/3} \left(\frac{R_*}{5R_\odot} \right)^{2/3} \left(\frac{M_*}{15M_\odot} \right)^{1/3}$$

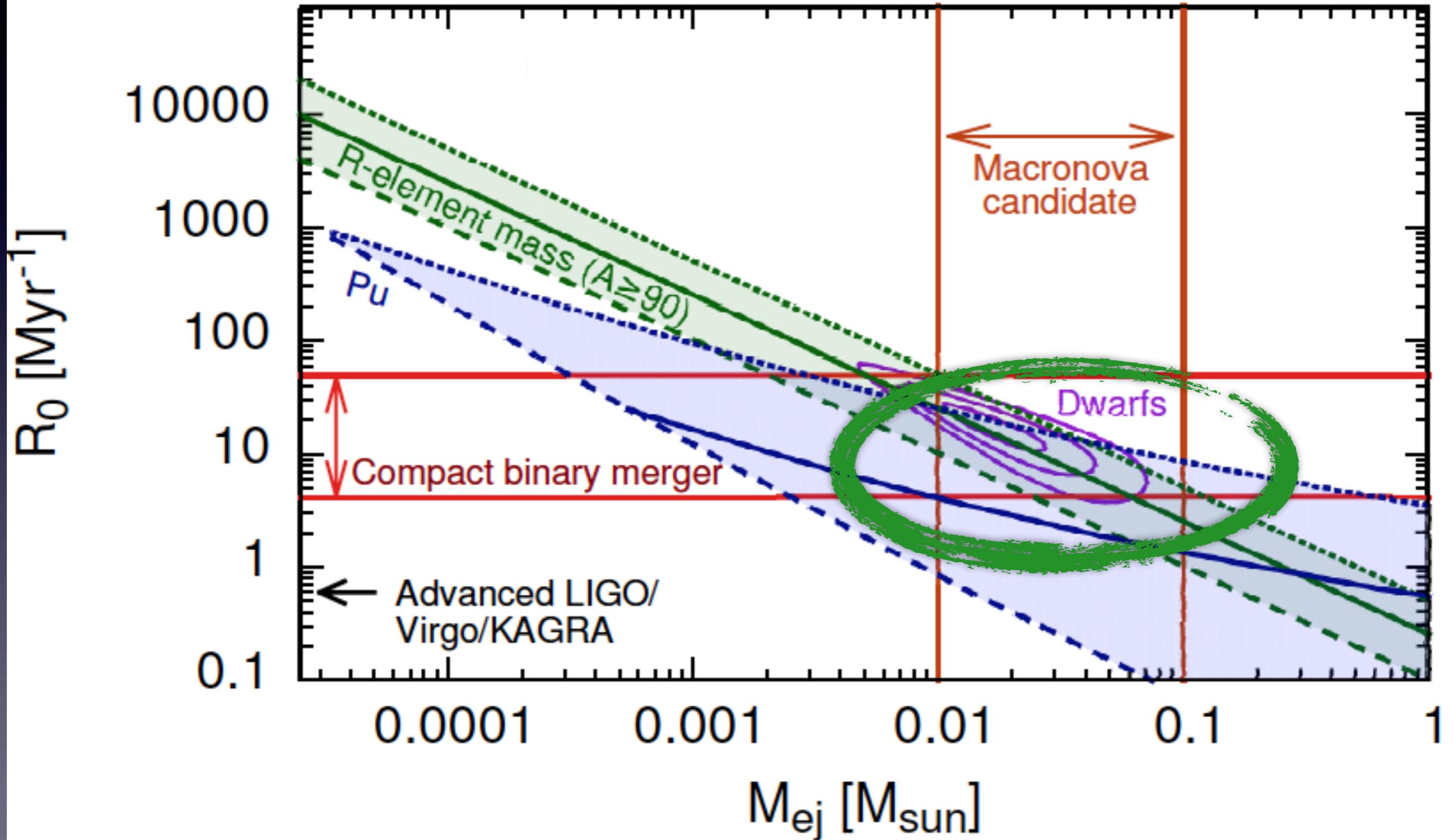
Dwarf Galaxies



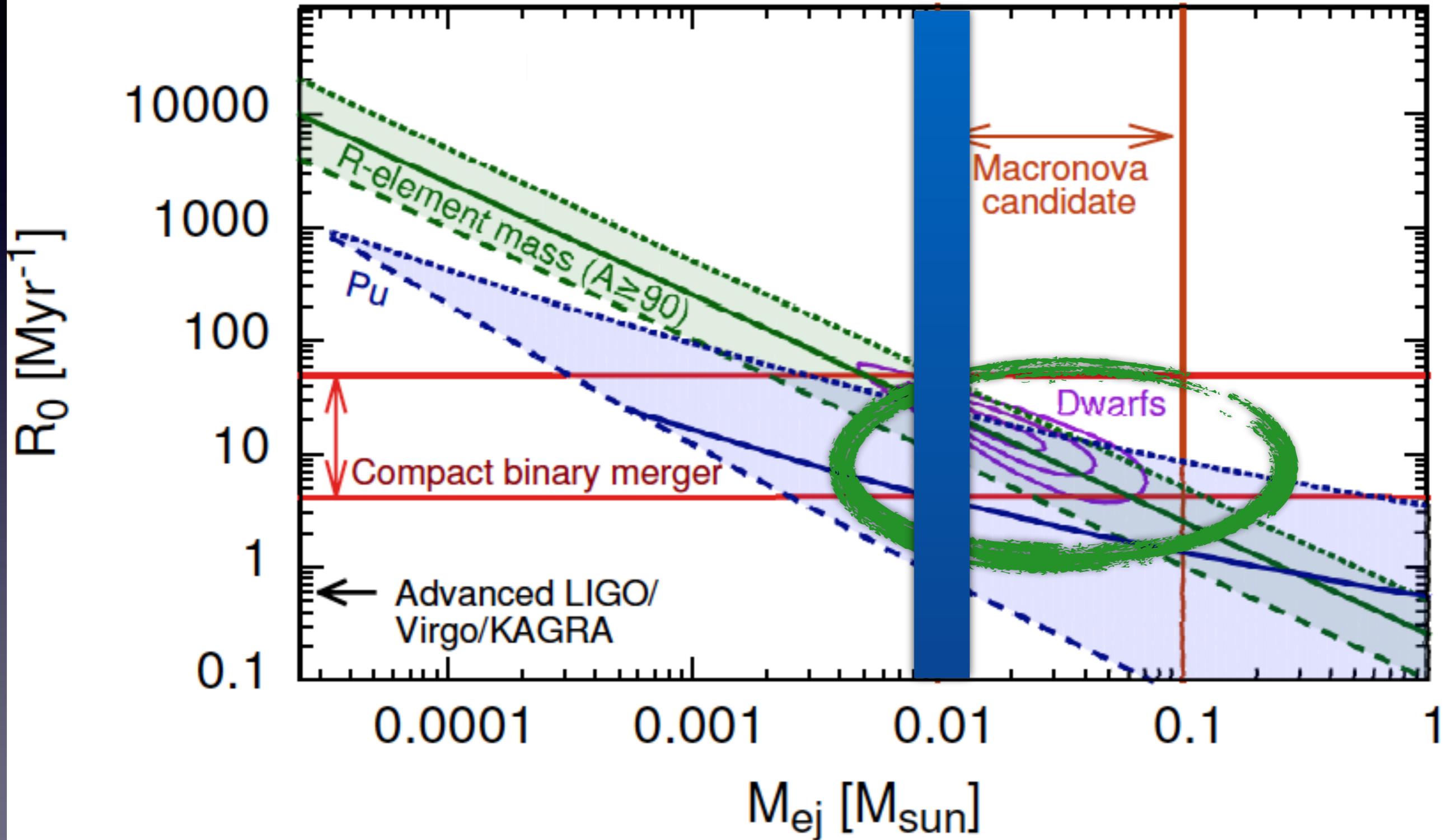
The Merger rate



r-process consistency



From SGRB Plateau



A prediction of the Collapsar model

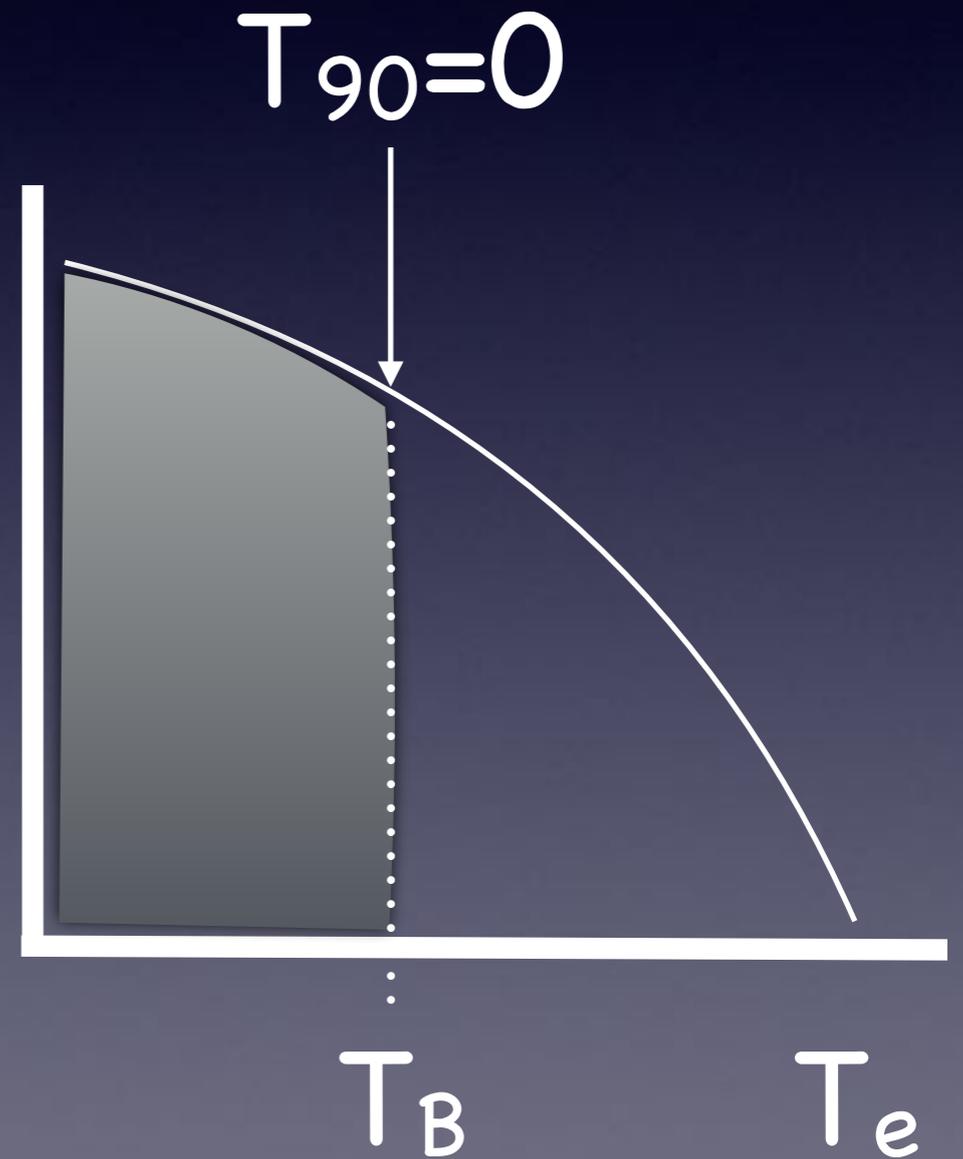
Observed duration

$$T_{90} = T_e - T_B$$

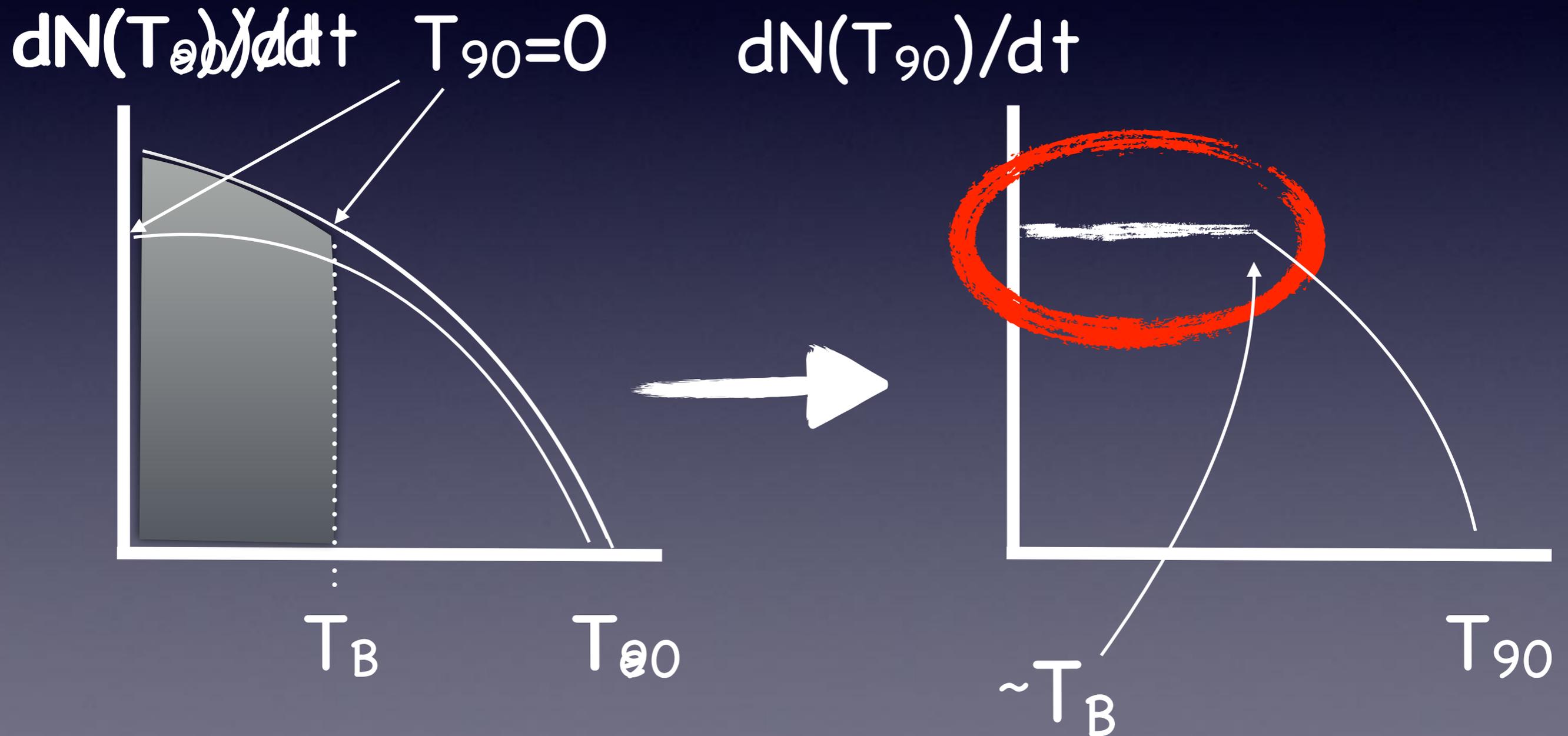
Engine time

Break out time

$dN(T_e)/dt$

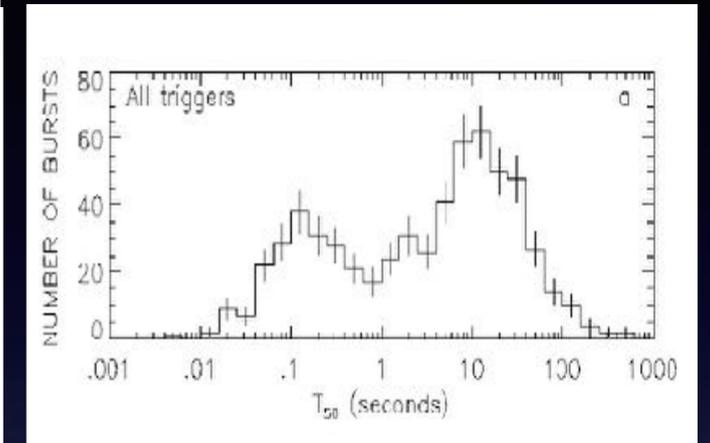
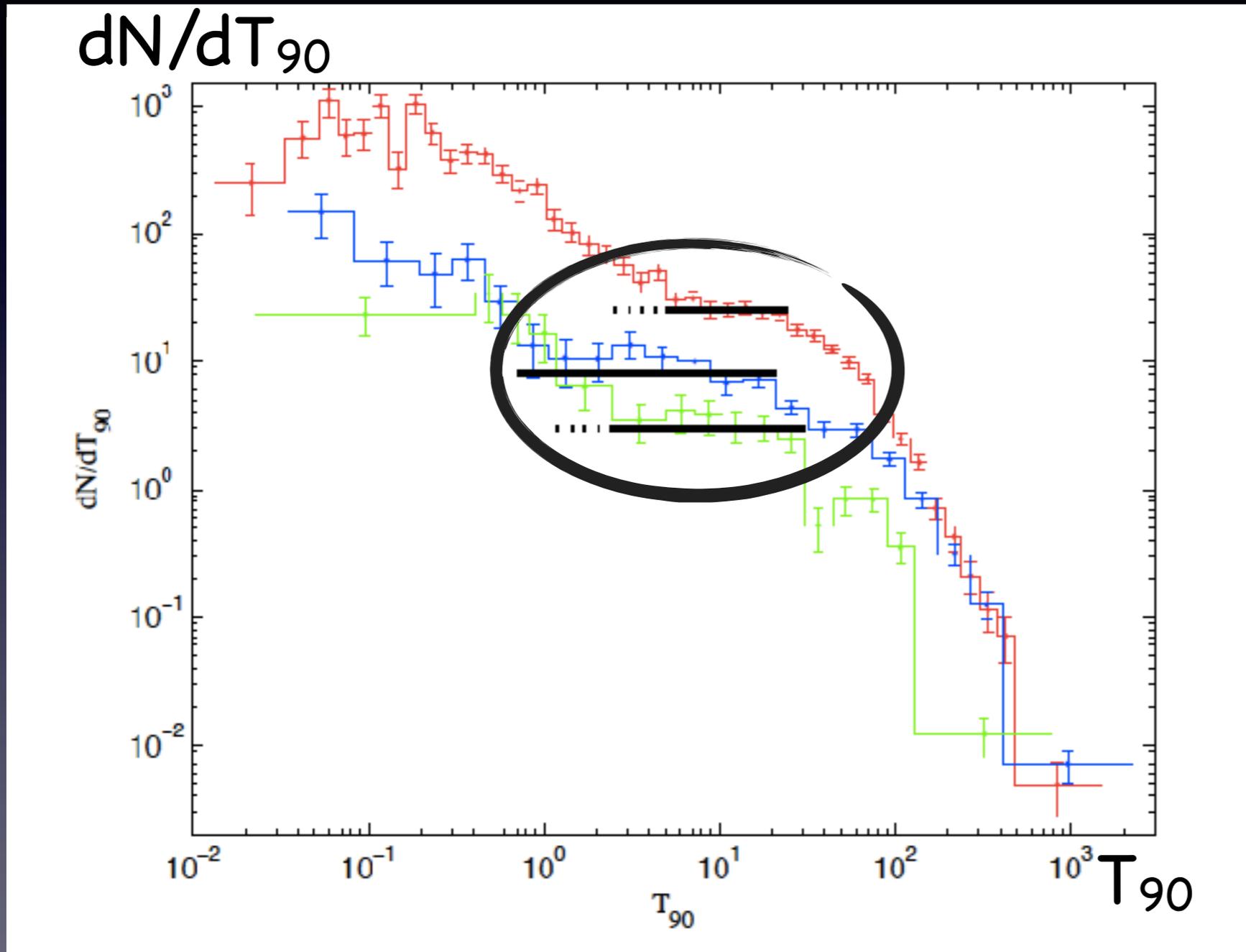


A prediction of the Collapsar model



The duration distribution

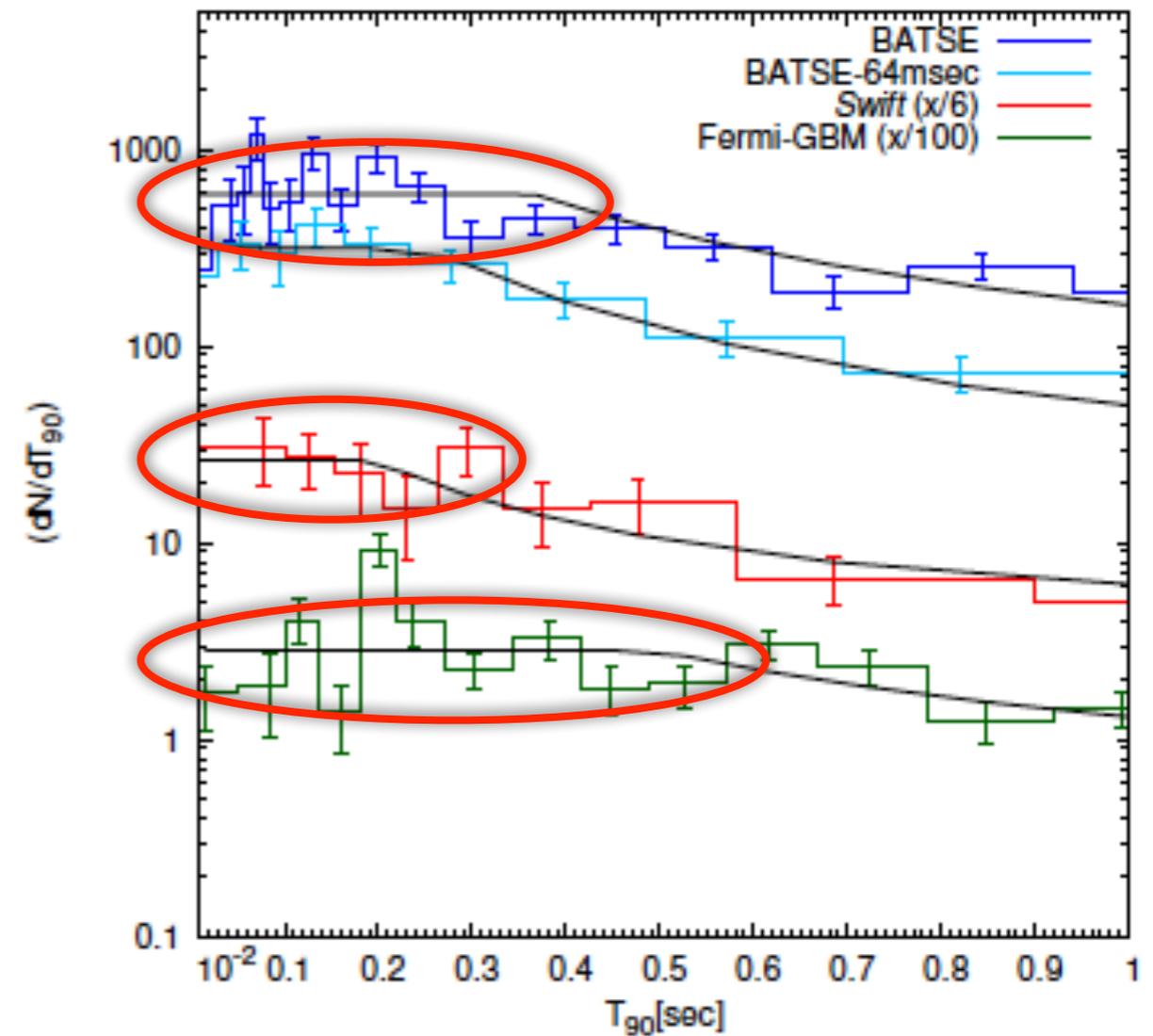
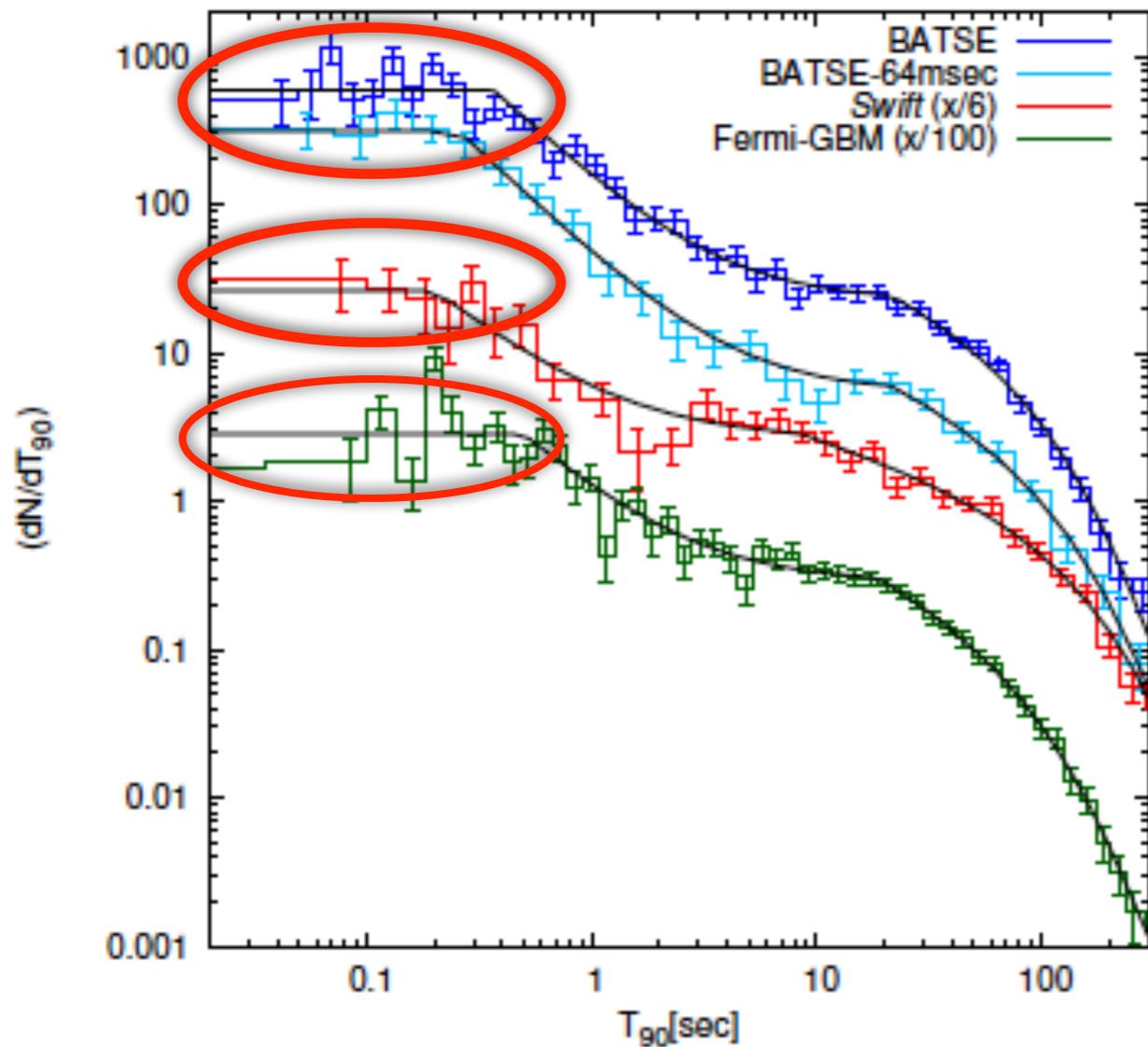
(Bromberg Nakar, TP & Sari, 2011)



A direct observational proof of the Collapsar model.

The “short” plateau

Moharana & TP 17 [arXiv170502598](https://arxiv.org/abs/1705.02598)



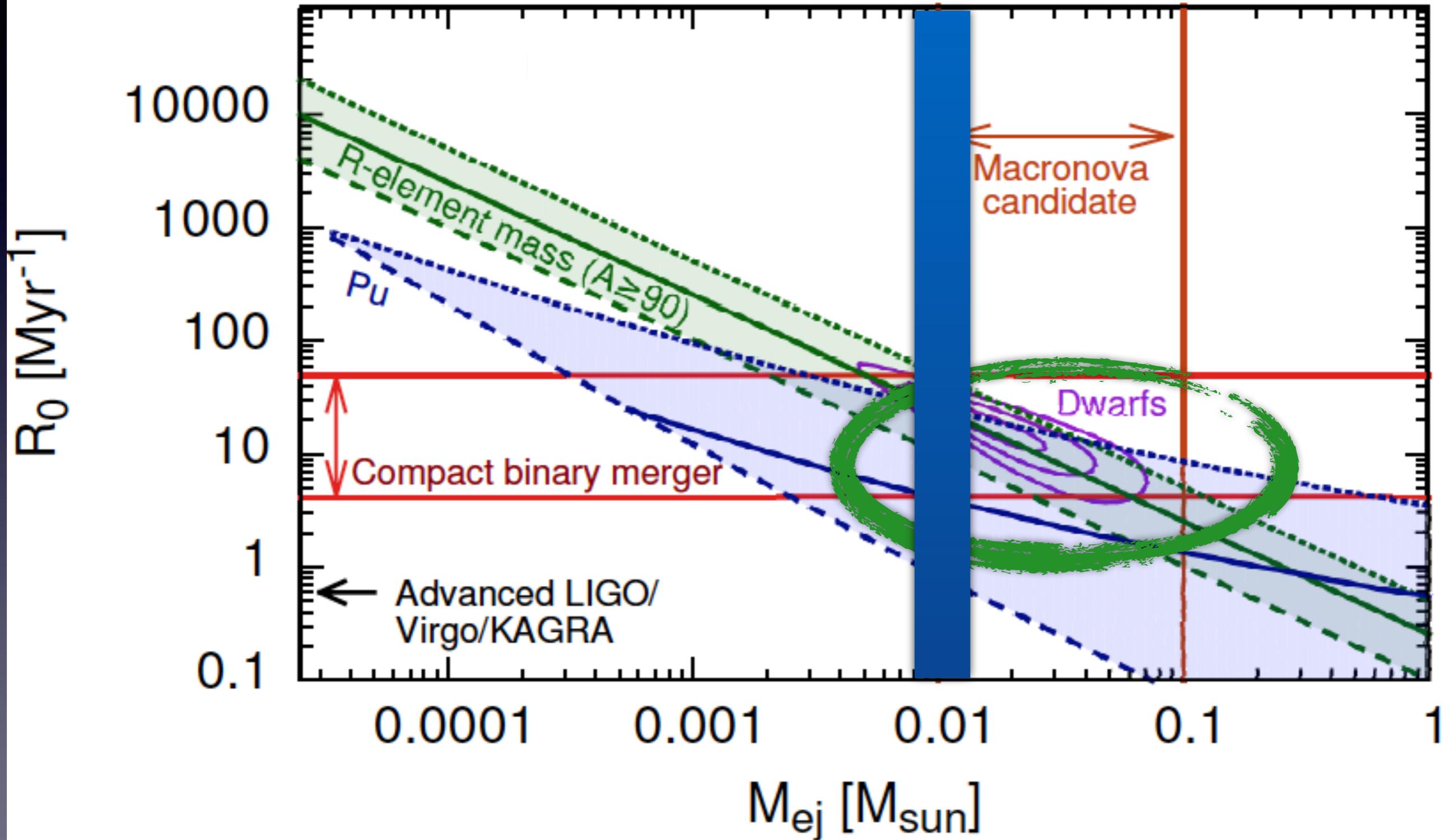
$t_b \sim 0.4$ Sec

$$t_b = 0.4 \text{ sec} \left(\frac{L_{iso,j}}{10^{51} \text{ ergs/sec}} \right)^{-1/3} \left(\frac{\theta_j}{15^\circ} \right)^{2/3} \left(\frac{R_e}{10^9 \text{ cm}} \right)^{2/3} \left(\frac{M_e}{10^{-2} M_\odot} \right)^{1/3}$$

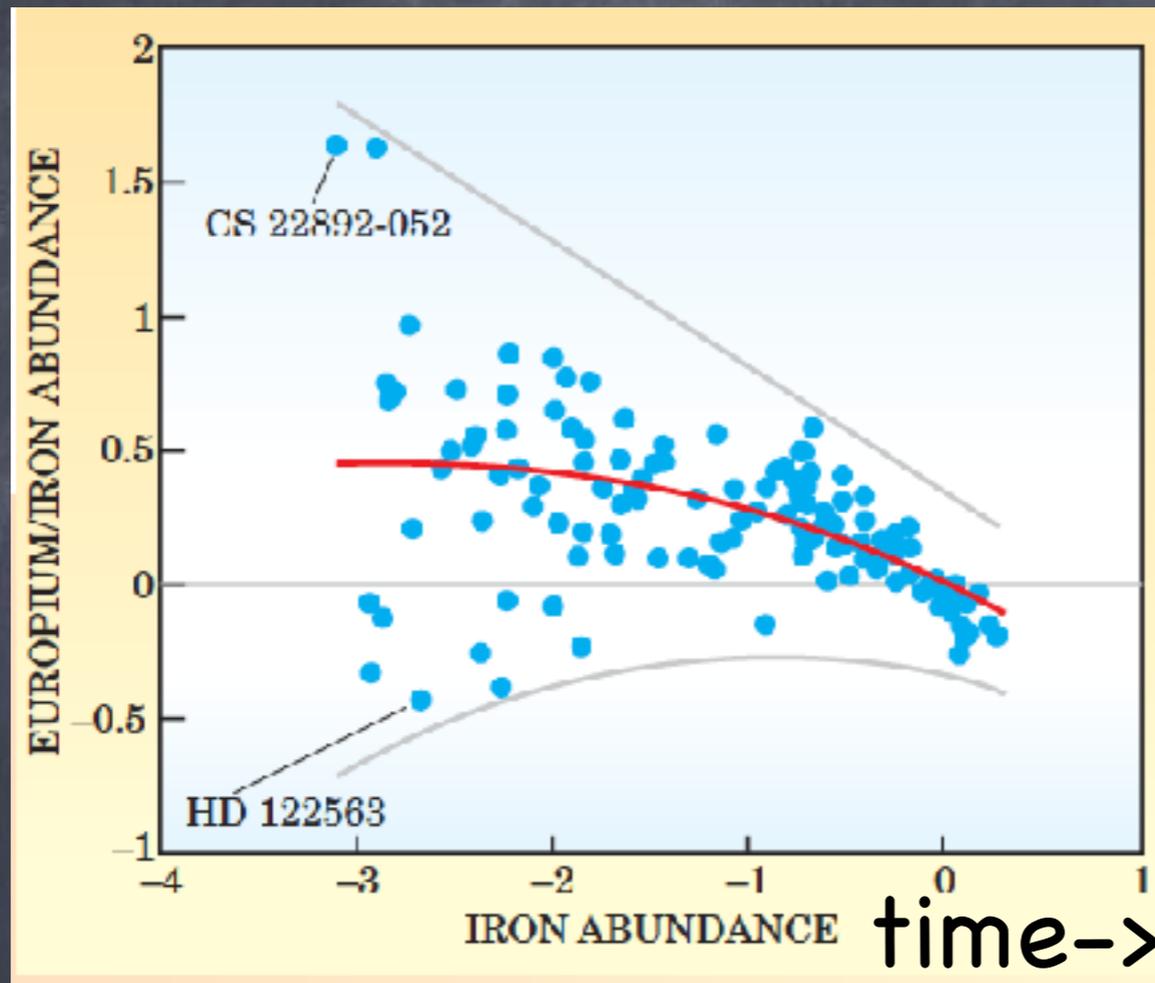
$$\int_0^{T_b} (\beta_h(t) - \beta_{max}) dt = \beta_{max} \Delta t ,$$

There are mergers in which the jet don't break out!

From SGRB Plateau



Early nucleosynthesis – a challenge



A population of fast mergers?

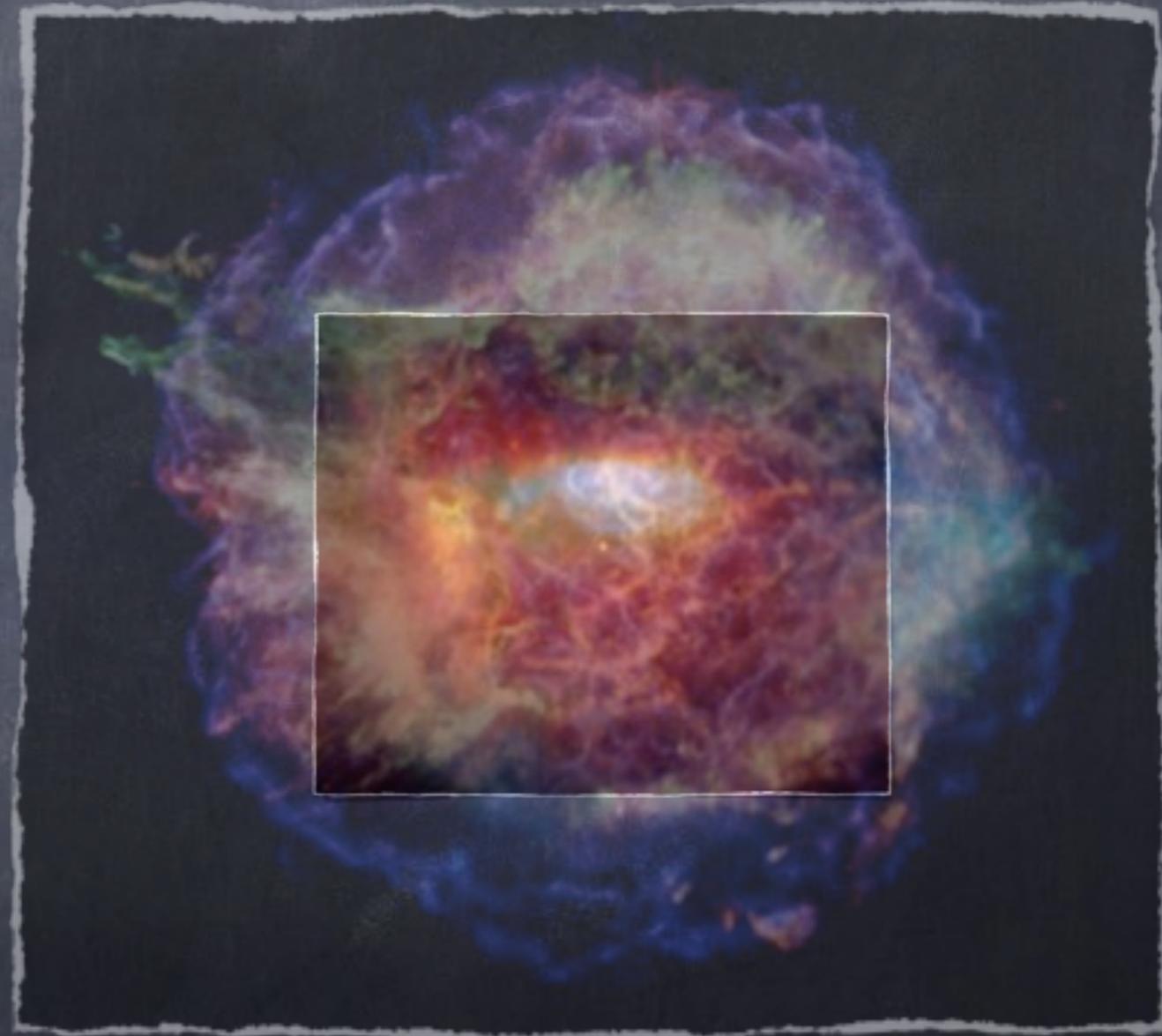
Figure 6. Europium abundance in a large sample of old and young stars, age being inferred from Fe abundance. The halo star HD 122563 is almost as Fe-poor as CS 22892-052, and therefore presumably just about as old, but it has much less Eu, an element made only in the r-process. The red line is a least-square-fit to the data, and the gray flanking curves indicate decreasing scatter in the data with increasing time. Numerical conventions are as in figure 5. Zero on the abscissa means Fe abundance like that of the 4.6-billion-year-old Sun.

From Cowan and Thielemann

The radio – flare (Nakar & Piran 2011)

Testing the Macronova interpretation

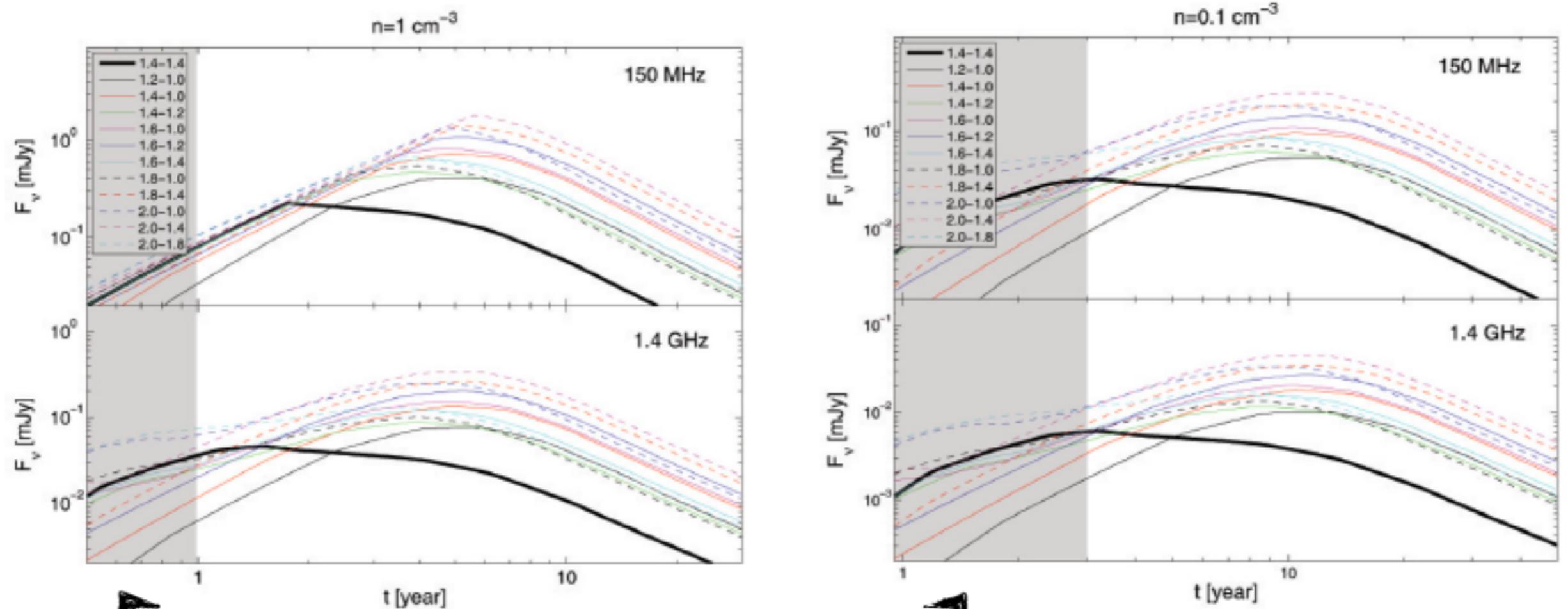
A long lasting radio flare due to the interaction of the ejecta with surrounding matter may follow the macronova.



Supernova → Supernova remnant

Macronova → Radio Flare

Radio flares from neutron star mergers



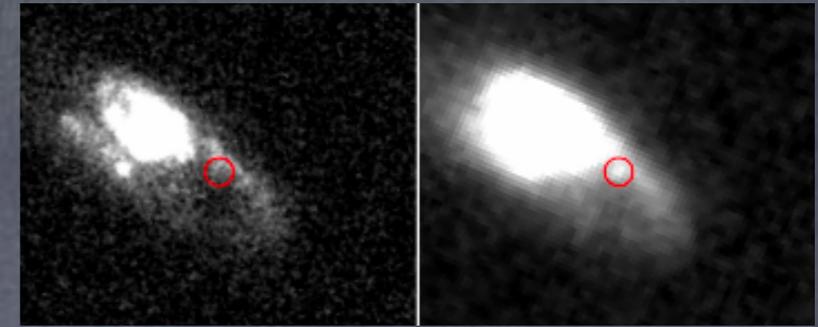
dominated by high velocity ejecta

A flare from GRB 130603B should be easily detected by the EVLA (if external density is not too small)



Summary

- There are a few caveats - But
- The nIR flare that followed the short GRB 130603B could have been a Macronova. If so than:
 - ✓ Short GRBs arise from mergers.
 - ✓ Gold and other $A > 130$ elements are produced in mergers. (But large m_{ej} and short time delay).
- A radio flare may confirm this!
- Another strong well localized short GRB is expected within a year or so.
- A GW signal + Merger + macronova (in 10 years)



The End ?