Search for new physics: Diboson resonance searches in ATLAS Run-2



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Outline

- Motivation
- Reconstruction techniques for high mass resonances
- Recent Run-2 searches results
 - VV resonances
 - VH resonances
 - V_γ resonances
- Conclusions







Lessons from Run-1

 Detectors performed very well in challenging LHC environment

 A Higgs boson was discovered with less luminosity and half the energy

 Remeasuring the Standard Model (SM)

 Nothing beyond the SM yet (not significant ones at least)

| A. | TLAS Exotics S | Search | es* - | 95% | 6 CL | Exclusion | ATL | AS Preliminary |
|---------------------|--|--|---|--|--|---|---|---|
| Sta | atus: March 2015 | | | | | | $\int \mathcal{L} dt = (1.0 - 20.3) \text{fb}^{-1}$ | \sqrt{s} = 7, 8 TeV |
| | Model | <i>ℓ</i> ,γ | Jets | $\mathbf{E}_{\mathrm{T}}^{\mathrm{miss}}$ | ∫£ dt[fb | -1] Mass limit | 5 | Reference |
| Extra dimensions | $\begin{array}{l} \mbox{ADD} \ {\cal G}_{KK} + g/q \\ \mbox{ADD} \ {\rm non-resonant} \ \ell\ell \\ \mbox{ADD} \ {\rm OBH} \ {\rm d} q \\ \mbox{ADD} \ {\rm OBH} \ {\rm d} q \\ \mbox{ADD} \ {\rm BH} \ {\rm high} \ {\cal N}_{t/k} \\ \mbox{ADD} \ {\rm BH} \ {\rm high} \ {\cal N}_{t/k} \\ \mbox{ADD} \ {\rm BH} \ {\rm high} \\ \mbox{ADD} \ {\rm BH} \ {\rm high} \ {\rm high} \ {\rm high} \ {\rm high} \\ \mbox{ADD} \ {\rm BH} \ {\rm high} \\ \mbox{ADD} \ {\rm BH} \ {\rm high} \ {\rm hig$ | $\begin{array}{c} - \\ 2e, \mu \\ 1 e, \mu \\ - \\ 2 \mu (SS) \\ \geq 1 e, \mu \\ - \\ 2 e, \mu \\ 2 \gamma \\ 2 e, \mu \\ 1 e, \mu \\ - \\ 1 e, \mu \\ 2 e, \mu (SS) \end{array}$ | $ \begin{array}{c} \geq 1 j \\ - \\ 1 j \\ 2 j \\ \geq 2 j \\ \geq 2 j \\ - \\ 2 j / 1 J \\ 2 j / 1 J \\ 2 j / 1 J \\ \geq 1 b, \geq 1 J \\ \geq 1 b, \geq 1 b, \geq 1 \end{array} $ | Yes - - - - - Yes j Yes | 20.3 20.3 20.3 20.3 20.3 20.3 20.3 20.3 | Mp 5.25 TeV Ms 4.7 TeV Min 5.2 TeV Min 5.2 TeV Min 5.2 TeV Min 5.2 TeV Min 5.8 TeV Min 5.8 TeV Min 5.8 TeV Girk mass 2.66 TeV Girk mass 740 GeV W'mass 700 GeV Girk mass 590-710 GeV Kir mass 590-710 GeV KK mass 960 GeV | $\begin{array}{l} n=2\\ n=3 \ \text{HLZ}\\ n=6\\ n=6\\ n=6, \ M_0=3 \ \text{TeV}, \ \text{non-rot} \ \text{BH}\\ n=6, \ M_0=3 \ \text{TeV}, \ \text{non-rot} \ \text{BH}\\ n=6, \ M_0=3 \ \text{TeV}, \ \text{non-rot} \ \text{BH}\\ k/\overline{M}_{\mathcal{H}}=0.1\\ k/\overline{M}_{\mathcal{H}}=1.0\\ k/\overline{M}_{\mathcal{H}}=1.0\\ k/\overline{M}_{\mathcal{H}}=1.0\\ \text{BR}=0.925 \end{array}$ | 1502.01518 1407.2410 1311.2006 1407.1376 1308.4075 1405.4254 Preliminary 1405.4123 Preliminary 1409.6190 1503.04677 ATLAS-CONF-2015-009 Preliminary |
| Gauge bosons | $\begin{array}{l} \text{SSM } Z' \rightarrow \ell\ell \\ \text{SSM } Z' \rightarrow \tau\tau \\ \text{SSM } W' \rightarrow b\gamma \\ \text{EGM } W' \rightarrow WZ \rightarrow \delta\gamma \ell' \ell' \\ \text{EGM } W' \rightarrow WZ \rightarrow qq\ell\ell \\ \text{HVT } W' \rightarrow WH \rightarrow \delta\gamma bb \\ \text{LRSM } W_R' \rightarrow t\bar{b} \\ \text{LRSM } W_R' \rightarrow t\bar{b} \end{array}$ | $2 e, \mu 2 \tau 1 e, \mu 3 e, \mu 2 e, \mu 1 e, \mu 1 e, \mu 0 e, \mu$ | - - 2 j / 1 J 2 b 2 b, 0-1 j ≥ 1 b, 1 J | – Yes Yes – Yes J – | 20.3 19.5 20.3 20.3 20.3 20.3 20.3 20.3 20.3 | Z' mass 2.9 TeV Z' mass 2.02 TeV W' mass 3.24 TeV W' mass 1.52 TeV W' mass 1.59 TeV W' mass 1.47 TeV W' mass 1.92 TeV W' mass 1.92 TeV | $g_V = 1$ | 1405.4123 1502.07177 1407.7494 1406.4456 1409.6190 Preliminary 1410.4103 1408.0886 |
| ũ | Cl qqqq Cl qqℓℓ Cl uutt | 2 _ e, μ 2 e, μ (SS) | $\begin{array}{c} 2 \ j \\ - \\ \geq 1 \ b, \geq 1 \end{array}$ | – j Yes | 17.3 20.3 20.3 | Λ Λ Λ 4.35 TeV | 12.0 TeV $\eta_{LL} = -1$ 21.6 TeV $\eta_{LL} = -1$ $ C_{LL} = 1$ | Preliminary 1407.2410 Preliminary |
| MQ | EFT D5 operator (Dirac) EFT D9 operator (Dirac) | 0 e, μ 0 e, μ | $ \geq 1 j \\ 1 J, \leq 1 j $ | Yes Yes | 20.3 20.3 | M, 974 GeV M, 2.4 TeV | at 90% CL for $m(\chi) < 100 \text{ GeV}$ at 90% CL for $m(\chi) < 100 \text{ GeV}$ | 1502.01518 1309.4017 |
| ΓQ | Scalar LQ 1 st gen Scalar LQ 2 nd gen Scalar LQ 3 rd gen | 2 e 2 μ 1 e, μ, 1 τ | ≥2j ≥2j 1b,1j | - - - | 1.0 1.0 4.7 | LO mass 660 GeV LO mass 685 GeV LO mass 534 GeV | $\begin{array}{l} \beta = 1 \\ \beta = 1 \\ \beta = 1 \end{array}$ | 1112.4828 1203.3172 1303.0526 |
| Heavy quarks | $\begin{array}{l} VLQ \ TT \rightarrow Ht + X, \ Wb + X \\ VLQ \ TT \rightarrow Zt + X \\ VLQ \ BB \rightarrow Zb + X \\ VLQ \ BB \rightarrow Wt + X \\ T_{5/3} \rightarrow Wt \end{array}$ | 1 e,μ 2/≥3 e,μ 2/≥3 e,μ 1 e,μ 1 e,μ | $ \begin{array}{l} \geq 1 \ \text{b}, \geq 3 \\ \geq 2l \geq 1 \ \text{b} \\ \geq 2l \geq 1 \ \text{b} \\ \geq 1 \ \text{b}, \geq 5 \\ \geq 1 \ \text{b}, \geq 5 \end{array} $ | j Yes – j Yes j Yes | 20.3 20.3 20.3 20.3 20.3 20.3 | T mass 785 GeV T mass 735 GeV B mass 755 GeV B mass 640 GeV T _{5,3} mass 840 GeV | isospin singlet T in (T,B) doublet B in (B,Y) doublet isospin singlet | ATLAS-CONF-2015-012 1409.5500 1409.5500 Preliminary Preliminary |
| Excited fermions | Excited quark $q^* \rightarrow q\gamma$ Excited quark $q^* \rightarrow qg$ Excited quark $b^* \rightarrow Wt$ Excited lepton $\ell^* \rightarrow \ell\gamma$ Excited lepton $v^* \rightarrow \ell W, vZ$ | $ 1 \gamma - 1 or 2 e, \mu 2 e, \mu, 1 \gamma 3 e, \mu, \tau $ | 1 j 2 j 1 b, 2 j or 1 – – | - - 1 j Yes - - | 20.3 20.3 4.7 13.0 20.3 | q' mass 3.5 TeV q' mass 3.5 TeV b' mass 870 GeV '' mass 2.2 TeV '' mass 1.6 TeV | only u^* and d^* , $\Lambda = m(q^*)$ only u^* and d^* , $\Lambda = m(q^*)$ left-handed coupling $\Lambda = 2.2 \text{ TeV}$ $\Lambda = 1.6 \text{ TeV}$ | 1309.3230 1407.1376 1301.1583 1308.1364 1411.2921 |
| Other | LSTC $a_T \rightarrow W\gamma$ LRSM Majorana ν Higgs triplet $H^{\pm\pm} \rightarrow \ell\ell$ Higgs triplet $H^{\pm\pm} \rightarrow \ell\tau$ Monctop (non-res prod) Multi-charged particles Magnetic monopoles | $ \begin{array}{c} 1 \ e, \mu, 1 \ \gamma \\ 2 \ e, \mu \\ 2 \ e, \mu (SS) \\ 3 \ e, \mu, \tau \\ 1 \ e, \mu \\ - \\ - \\ - \\ - \\ - \\ - \\ - \\ - \\ - \\ -$ | - 2j - 1b - - | Yes Yes | 20.3 2.1 20.3 20.3 20.3 20.3 20.3 2.0 | ar mass 960 GeV N ⁹ mass 1.5 TeV H ^{±±} mass 551 GeV H ^{±±} mass 400 GeV spin-1 invisible particle mass 657 GeV multi-changed particle mass 785 GeV monopole mass 862 GeV | $\begin{split} m(W_{\mathbb{R}}) &= 2 \text{ TeV, no mixing} \\ \mathrm{DY \ production, } BR(H_{L}^{\pm\pm} \to \ell \partial = 1 \\ \mathrm{DY \ production, } BR(H_{L}^{\pm\pm} \to \ell \tau) = 1 \\ a_{0,m-set} &= 0.2 \\ \mathrm{DY \ production, } g &= 5 e \\ \mathrm{DY \ production, } g &= 1 g_{D} \end{split}$ | 1407.8150 1203.5420 1412.0237 1411.2921 1410.5404 Preliminary 1207.6411 |
| | | $\sqrt{s} =$ | 7 TeV | √s = | 8 TeV | 10 ⁻¹ 1 | ¹⁰ Mass scale [TeV | - |

*Only a selection of the available mass limits on new states or phenomena is shown.

LHC Run-2 started in 2015

LHC proton-proton collisions restarted in June 2015 at 13 TeV



ATLAS program for Run-2?

- Is this boson the Higgs boson?
- Increase precision in SM measurements
- Push searches BSM further

Discovery of the Higgs boson

Guided by *clues* from the SM of particle physics



- Everywhere starting with high masses
- Increase of energy → Increase of reach for new phenomena
- Example: production rate of a W' with a mass of 2 TeV increased by 5x times from Run-1 to Run-2
 - You already can see where this is going...
- High mass searches are challenging!
 - Usually require revisiting our object reconstruction and analysis techniques



- Interactions with pairs of electro-weak
 bosons are a fundamental element in
 electroweak symmetry breaking
- These interaction played a key role in the Higgs boson discovery in Run-1

 The high mass is particularly sensitive to a wide range of BSM physics models

Diboson searches important in general



- Clear signature in detector
- Known properties and decay kinematics

High mass diboson searches are well motivated

Several theories predict the existence of new heavy resonance coupling to W/Z/Higgs

- Extra dimensions, compositness, GUT
- Resonance benchmarks you will hear about in this talk
 - Spin 0 Higgs-like scalar singlet
 - Spin-1 HVT (simplified Lagrangian)
 - Model A: Stronger constraints from leptonic searches
 - Model B: Enhanced couplings to dibosons
 - Spin-2 Randram Sundrum graviton (RGS)



- Clear signature in detector
- Known properties and decay kinematics

| | WW | WZ | ZZ | VH |
|--------|--------|----|--------|-------|
| HVT | Ζ' | W' | | Z'/W' |
| Gravi | RSG | | RSG | |
| Scalar | Scalar | | Scalar | |

LHC Run-2 started in 2015

- Up to the summer 2016 ~15/fb of 25 ns data good for physics
- Data quality efficiency ~92%
- Smooth running and excellent trigger performance!
 - ◆ >96% working channels (pixels, cells, …) in each sub-detector



Reconstructing boosted bosons Traditional jet algorithms



Reconstructing boosted bosons



In the boosted regime the decay products are collimated in the direction of the initial particle

- Rule of thumb for angular separation of decay products $\Delta R = \sqrt{\Delta \phi^2 + \Delta \eta^2} \approx \frac{2m}{\rho_T} \quad A W \text{ boson with } p_T \sim 200 \text{ GeV}, \ \Delta R = 0.8$
- From a practical point of view this means:
 - Hadronic decay products merge into a single jet
 - Leptons close to (or even) inside jets
 - (need to modify lepton isolation criteria)
 - At very high p_{τ} the calorimeter granularity is a limitation

Reconstructing boosted bosons

Pile-up conditions make this task more challenging



BooBo/top tagging: what do we need?





B-tagging

- H→bb, Z→bb
- For obvious reasons
- Standard algorithms not adapted for dense environments

Cartoons by A. Martyniuk

ATLAS calorimetry



Good resolution to pick apart the large-R jets and look at its substructure

Grooming techniques



Trimming:

- Reclustering of constituents of large-R jet into small-R jets of size R_{sub}
- Remove subjet *i* if $p_T^i < f_{cut} x p_T^{jet}$
- Default ATLAs groomer (stable against PU)

ATLAS-CONF-2012-065

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Jet variables for tagging

Again...can not cover all! Some of them strongly linear correlated

Mass

 Deduced from 4-momentum sum of all jet constituents

Expected to be small for QCD jets, but closer to the boson/top mass for signal jets

N-subjettiness (JHEP 03 (2011) 015)

Measure of how N or less "subjet"-like a large-R is

Energy correlation variables (*JHEP* 1306 (2013) 108)

• Energy correlations functions (ECFs) construct a complete representation of the jet by combining the p_{T} and angular separation of all jet constituents (ECF1), all pairs of jet constituents (ECF2) and triplets (ECF3)

 Ratios of these are powerful in rejecting jets from multi jet processes

$$D_2^{\beta=1} = E_{\rm CF3} \left(\frac{E_{\rm CF1}}{E_{\rm CF2}}\right)^3$$

ATL-PHYS-PUB-2015-033



W/Z boson tagging for early Run-2

- Huge optimization effort at the end of Run-1: arXiv:1510.05821
- 4 sets of algorithms studied for Run-2, of which this is the most performant:

Mass and energy calibrated anti-kt jets with R = 1.0 Trimmed with fcut = 5% and Rsub = 0.2 Dynamic mass window cut (68%) + p_T dependent D2 cut for jets gives the best rejection (~90%) at 50% signal efficiency

 Uncertainties derived by comparing the measured calorimeter jet energy and mass to the same quantities measured by the tracker in both data and MC, using a double ratio method



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H→bb boson tagging for early Run-2

- **Tag small-radius** (R=0.2) jets made of tracks:
 - Match tracks directly to $PV \rightarrow pileup$ insensitive
 - Smaller radius jets for close-by b-tagging
 - Better resolution w.r.t. b-hadron direction than calo



Mass and energy calibrated anti-kt jets with R = 1.0

Trimmed with fcut = 5% and Rsub = 0.2

Mass window cut + track-jet b-tagging (MV2c20 70% wp) + p_{π} dependent D2 cut (optional)

Uncertainties driven by b-tagging (calibrated in 2015 data using ttbar events)

ATLAS-CONF-2016-039



Boosted boson tagging at work

ATLAS-CONF-2016-055



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How would new resonances manifest themselves?

New particles: resonant excess (bump) over Standard Model background



(~new particle mass)

Some extra from Run-1: an analysis you might remember

Fully hadronic channel observed a 3σ local excess with the Run-1 dataset arXiv:1506.00962v2

Not seen by other channels in ATLAS (arXiv:1512.05099)
 CMS observed smaller deviations around 1.9 TeV (EXOT14010) and at 650 GeV ()



Quick review of the results



Analysis strategy



$VV \rightarrow qqqq$: Selection

Events / 0.15

Entirely dominated by multijet background

- Two highest p₊ jets (>450/200 GeV) boson tagged
 - + Cut on number of tracks associated to the jet
 - 70% bkg rejection, 30% loss in signal
- $|y_1 y_2| < 1.2$, p₁ asymmetry < 0.15
- Veto on leptons and E^{miss}
- 3 partially overlapping signal regions WW,WZ, ZZ





VV → qqqq: Background and signal regions

Background estimated from fit to the data

Validation on simulation and checked against

data in mass sidebands control regions

 $\frac{dn}{dx} = p_1(1-x)^{p_2+\xi p_3} x^{p_3}, x = \frac{m_{jj}}{13TeV}$

10

 10^{3}

10

10

ATLAS Preliminary

s = 13 TeV, 15.5 fb

WZ selection

Events / 0.1 TeV

ATLAS-CONF-2016-055

Data 2015+2016

Fit bkg estimation

Fit exp. stats error

HVT Model A m=1.5 TeV T Model A m=2.4 TeV

WZ signal region



3.5

$VV \rightarrow IIqq/Ivqq/vvqq$: Selection



- High purity with a large-R boson tagged jet
- Low purity with inverse D₂ cut applied to gain sensitivity at high masses

VV → Ilqq/lvqq/vvqq: Backgrounds

- Backgrounds are estimated using Monte Carlo simulation and checked in control regions defined using:
 - The jet mass sidebands for W/Z+jets
 - By asking additional number of b-jets in the event for ttbar
- Control regions are included in the final fit to better constrain the normalisation

Data

1////

Z+iets CR for 2 lepton

2000

Z + jets

SM Diboson

Top Quarks

Stat.

Syst. Uncert.

Pre-fit background

2500

300

m(*llJ*) [GeV]

2 8 10

Events / 10²

10¹

1.0

10

10-

10-3

10

0.5 -500

Data/Pred

ATLAS Preliminary

 $\sqrt{s} = 13$ TeV. 13.2 fb⁻¹

1000

1500

Merged high-purity ZCR, ggF

 $H \rightarrow ZZ \rightarrow \ell \ell q q$



VV → Ilqq/lvqq/vvqq: Signal regions

 \blacksquare Use $m_{_{VV}}$ (m_ $_{_{T,VV}}$ for the vvqq channel) as a discriminant

Larger systematic uncertainties are jet and background modelling related

No significant deviation observed
 Limits set on HVT (spin-1), RSG (spin-2) and Heavy Higgs (spin-0)





VV resonances searches: A bit of perspective



Here we can see which channel is more sensitive and where for the ATLAS searches

Deviation at 650 GeV in the llqq channel not present when adding more data (B2G-16-022)

Current ATLAS Run-2 results highly disfavor the Run-1 2 TeV excess at its observed signal strength

VH → IIbb/Ivbb/vvbb: Selection and backgrounds

- Similar selection to the un-tagged (VV) analyses on leptonic side
- Use $H \rightarrow$ bb boosted tagger on the hadronic side
 - 1 and 2 b-tag categories
- Main backgrounds
 - Top and W/Z+heavy flavour
 - Estimated with simulation and checked in control regions
- No excess over the SM background



JU



2-tag

VH → IIbb/Ivbb/vvbb: Limits results



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Search for VH, V \rightarrow qq and H \rightarrow bb



Complementary to semi-leptonic
 VH searches at high masses

- Main backgrounds: multijet and tt
- ATLAS results with 13 TeV 3.2/fb
- No significant deviation found
 Largest deviation at 3 TeV reported by ATLAS (2.5σ global significance)



Extra thoughts for the future: the challenges we still face

- Excellent performance of ATLAS! Our searches look into many different final states/phase space
- Modeling of complex observables by Monte Carlo simulation is essential
 - Measurements of jet mass and other jet shapes need to be fed back to simulation
- We have many tools to mitigate PU, but the increases we face in Run-2 and beyond will be challenging
- Going to even higher $p_{\tau} \rightarrow$ very very collimated decay products... Limited calorimeter granularity



Looking into the future ATLAS Run-3 boosted objects trigger development

Hadronic decays of high p_T bosons and fermions is a vital part of the ATLAS physics program

- Isolated lepton triggers inefficient for some analyses
 → use jet triggers (e.g. VV → qqqq)
- Many techniques to suppress PU and to identify substructure in jets are implementable in the ATLAS High Level Trigger (HLT)
 - Large-R acceptance in HLT depends on the L1 requirements
 - But adding these in the level-1 hardware-based trigger is more complicated

ATLAS is planning major detector updates in Run-3, like Level-1 trigger (calorimeter) system

- Including the Global Feature Extractor (gFEX)
- Institutions: BNL, UChicago, Indiana, Pittsburgh, Oregon and Stockholm



Nice overview in M. Begel's talk at BOOST15

Global Feature EXtractor (gFEX)

gFEX is a single board that will have access to the information from the whole calorimeter!

Will identify events with large-radius jets

- Improving acceptance for boosted objects
- Jet-level pile-up substraction
- Substructure variables could be used
- Will calculate global event variables:

E₁^{miss}, centrality

 Implemented in a highly parallelized structure (3 large Xilinx Ultrascale FPGAs and Zync System-On-Chip)

Prototype is available. Initial LAr calorimeter-L1Calo link speed communication tests very successful!



acceptance gain for boosted top

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Summarizing

- Search for heavy resonances is one of the most direct ways to find new physics at TeV scale
- Many Run-2 searches are now using these boosted techniques and many more are coming. Wealth of physics encoded inside jets!
 - A very active field
 - Important to build confidence in these tools through its use in more SM measurements
 - Main challenges: pile-up, detector granularity and simulation modeling
- No significant excess observed so far but only a third of the data collected in 2016 has been analyzed so far
- Thinking ahead: gFEX will increase trigger rejection to present acceptance in Run-3

Thanks for your attention!





ATLAS improvements in Run-2

Insertable b-layer (IBL) in place

- Upcoming trigger improvements:
 - Fast TracK trigger (FTK)
 - L1 topological trigger
- Software improvements:
 - New data format for analysis
 - Online (trigger) and offline jet reconstruction are ~the same



Insertable b-layer insertion (2014)

Trigger and data acquisition:

- First step: fast hardware selection
 - Run-1 data taking rate: 75 kHz
 - Run-2 data taking rate: 100 kHz
- Second step: computer farm
 - Run-1 data taking rate: 400 Hz
 - Run-2 data taking rate: 1000 Hz From C. Doglioni



IBL takes first proton-proton data

Grooming techniques

Split-filtering: http://arxiv.org/abs/0802.2470

- Decluster and discard soft junk
- Requiring symmetric splitting
- Repeat until find hard structure
- Small-radius jet reclustering, keeping only the three highest p₁ subjets

Pruning: http://arxiv.org/abs/0912.0033

- Constituents of large-R jet are reclustered with either C/A or kt algorithm
- In each clustering step, large angle and soft clusterings are removed

Trimming: http://arxiv.org/abs/0912.1342

- Reclustering of constituents of large-R jet into small-R jets of size R_{sub}
- Remove subjet *i* if $p_{T}^{i} < f_{cut} x p_{T}^{jet}$
- Default ATLAs groomer (stable against PU)

Can not cover all tools...but these 3 are widely used



global Feature Extraction (gFEX) The big picture

A new level 1 calorimeter trigger system for Run 3 (~2020)



- Entire calorimeter in one single board:
 - Jet substructure in Run 3 and beyond: fat jet reconstruction and jet-level pile-up corrections
 - Global event variables, e.g: E_{T}^{miss} and centrality
- Physics algorithms run within 5 bunch crossings (125 ns), not including data input/output More on algorithms in Walter's talk! ⁴¹

global Feature Extraction (gFEX) The big picture

A new level 1 calorimeter trigger system for Run 3 (~2020)

- One single module with several FPGAs for data processing
 - Inter-communication to avoid environments
- Hybrid FPGA (FPGA+CPU system-onship or Zynq) for control and monitoring
 - Process the event data from processor FPGAs
 - Algorithms to quickly detect calorimeter issues
 - Emulate the feature identification algorithms
 - Histograms interesting quantities



Zynq in gFEX



Zynq functionality/features in gFEX Control and monitoring



Disclaimer: GPIO and MGT counts out of date

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gFEX: area based PU subtraction



Global Feature EXtractor (gFEX)

Successful Link Speed Test @CERN

Unfortunately not all the people that worked hard in the project are in this picture (taken during LAr-gFEX test)



VH → IIbb/Ivbb/vvbb: backgrounds



ATLAS (A Toroidal LHC ApparatuS)

ATLAS consists of a series of concentric sub-detectors around the interaction point

Divided into 4 major parts: the inner detector, the calorimeters, the muon spectrometer and the magnet systems



Search for new diboson resonances in ATLAS: Run-1 Event selection

- Trigger selection: EF_j360_a10tcem (lowest unprescaled jet trigger for 2012)
- Quality:
 - GRL
 - DQ checks from data preparation: coreFlags, LArError, TileError, TileTripReader
 - Bad/ugly jets
 - BCH_CORR_CELL cleaning
- Lepton/MET veto: orthogonal with other diboson searches
- Jets:
 - 2 jets filtered with BDRS-A
 - Mass of dijet-system is required to be above 1.05 TeV¹, in order to avoid region with trigger inefficiency
 - ► Rapidity gap between the two leading jets |∆y₁₂| < 1.2, to reject QCD t-channel dijet production</p>
 - > $p_{\rm T}$ asymmetry, A < 0.15, between two leading jets, to select balanced events
 - $|\eta < 2.0|$ to ensure good overlap with the inner detector
 - Boson tagging criteria applied



WW, WZ, ZZ partially overlapping selections



Full hadronic diboson Run-1 results



Diboson Run-2 results: ATLAS and CMS



Information about benchmarks used in the 13 TeV analyses

Table 1: The resonance width (Γ) and the product of cross-section times branching ratio (BR) for diboson final states, for different values of the mass pole *m* of the resonances predicted by the CP-even scalar model ($\Lambda = 1$ TeV, $c_H = 0.9, c_3 = 1/16\pi^2$), by model B of the HVT parameterisation ($g_V = 3$), and by the graviton model ($k/\overline{M}_{Pl} = 1$).

| | | Scalar | | H | HVT W' and Z' | | | G* | | |
|-------|-------|--------------------|--------------------|-------|-----------------------------|-----------------------------|-------|--------------------|--------------------|--|
| | | WW | ZZ | | WW | WZ | | WW | ZZ | |
| m | Г | $\sigma \times BR$ | $\sigma \times BR$ | Г | $\sigma \times \mathrm{BR}$ | $\sigma \times \mathrm{BR}$ | Г | $\sigma \times BR$ | $\sigma \times BR$ | |
| [TeV] | [GeV] | [fb] | [fb] | [GeV] | [fb] | [fb] | [GeV] | [fb] | [fb] | |
| 0.8 | 4.2 | 730 | 359 | 32 | 682 | 354 | 46 | 301 | 155 | |
| 1.6 | 33 | 7.8 | 3.9 | 51 | 79.3 | 38.5 | 96 | 4.4 | 2.2 | |
| 2.4 | 111 | 0.32 | 0.16 | 74 | 10.5 | 4.87 | 148 | 0.28 | 0.14 | |

Table 2: Generators and PDFs used in the simulation of the various background processes.

| Process | PDF | Generator |
|---|-----------|----------------------------|
| W/Z + jets | CT10 | Sherpa 2.1.1 |
| tī | CT10 | Powheg-BOX v2+Pythia 6.428 |
| Single top (<i>Wt</i> , <i>s</i> -channel) | CT10 | Powheg-BOX v2+Pythia 6.428 |
| Single top (<i>t</i> -channel) | CT10 | Powheg-BOX v1+Pythia 6.428 |
| Diboson (WW , WZ , ZZ) | CT10 | Sherpa 2.1.1 |
| Dijet | NNPDF23LO | Рутніа 8.186 |

Table 4: Channels, signal regions and mass ranges where the channels contribute to the search.

| Channal | Signal ragion | Scalar | HVT W' and Z' | G^* |
|---------|-------------------|------------------|-------------------|------------------|
| Channel | Signal Tegion | mass range [TeV] | mass range [TeV] | mass range [TeV] |
| | WW + ZZ selection | 1.2–3.0 | _ | 1.2–3.0 |
| 9999 | WW + WZ selection | _ | 1.2–3.0 | _ |
| | WZ selection | _ | 0.5–3.0 | _ |
| vvqq | ZZ selection | 0.5–3.0 | _ | 0.5–3.0 |
| luga | WW + WZ selection | _ | 0.5–3.0 | _ |
| ivqq | WW selection | 0.5–3.0 | _ | 0.5-3.0 |
| llaa | WZ selection | _ | 0.5–3.0 | _ |
| uqq | ZZ selection | 0.5–3.0 | _ | 0.5–3.0 |

Information about benchmarks used in the 13 TeV analyses



(c)

Figure 1: Signal acceptance times efficiency for the different analyses contributing to the searches for (a) a scalar decaying to WW and ZZ, (b) HVT decaying to WW and WZ and (c) bulk RS gravitons decaying to WW and ZZ. The branching ratio of the new resonance to diboson is included in the denominator of the calculation. The error bands represent statistical and systematic uncertainties.

Other intriguing results... diphotons at 750 GeV

• Limits set on $\sigma(pp \rightarrow X)^*BR(X \rightarrow Z\gamma)$ assuming scalar X produced in gluon fusion • Observed limits between 295 fb for m_x = 340 GeV and 7.5 fb for m_x = 2.15 TeV



D2 definition

$$e_{2}^{(\beta)} = \frac{1}{p_{TJ}^{2}} \sum_{1 \le i < j \le n_{J}} p_{Ti} p_{Tj} R_{ij}^{\beta},$$

$$e_{3}^{(\beta)} = \frac{1}{p_{TJ}^{3}} \sum_{1 \le i < j < k \le n_{J}} p_{Ti} p_{Tj} p_{Tk} R_{ij}^{\beta} R_{ik}^{\beta} R_{jk}^{\beta},$$

$$D_{2}^{(\beta)} = \frac{e_{3}^{(\beta)}}{(e_{2}^{(\beta)})^{3}}$$
Larkoski et al.arXiV:1409.6298
$$D_{2}: \text{ large for 1-prong jet (e.g. QCD bkg.)}$$

$$(e_{2})^{3} \le e_{3} \le (e_{2})^{2}.$$

$$I-prong (e.g. QCD)$$

$$2-prong (e.g. Hbb)$$
Plots from R. Jacobs

small for 2-prong jet (Higgs signal) $0 < e_3 \ll (e_2)^3$



W/Z boson tagging for early Run-2

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